

Temporally sparse controls for infinite horizon semilinear parabolic equations with norm constraints*

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Dedicated to Jan Sokolowski on the occasion of his 75th birthday

Abstract

In this paper, infinite horizon optimal control problems subject to semilinear parabolic equations are investigated. A finite number of only time-dependent controls localized at disjoint positions in the space domain are considered. The controls are subject to integral constraints and a term is included in the cost functional that promotes control sparsity. The existence of optimal controls is proved, first and second order optimality conditions are derived, and the approximation by finite horizon control problems is addressed.

Keywords: semilinear parabolic equation, sparse optimal control, infinite horizon problems, first and second order optimality conditions

AMS Subject classification: 49K20, 49J52 35K58,

1 Introduction

In this paper, we study the following control problem

$$(P) \begin{cases} \text{Minimize } J(u) := F(u) + \alpha j(u) \\ u \in \mathcal{U}_{ad} \end{cases}$$

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where $\alpha > 0$,

$$\begin{aligned} F(u) &= \int_0^\infty \int_\Omega |y_u(x, t) - y_d(x, t)|^2 dx dt, \\ j : L^1(0, \infty)^m &\longrightarrow \mathbb{R}, \quad j(u) = \sum_{i=1}^m \int_0^\infty |u_i(t)| dt, \\ \mathcal{U}_{ad} &= \{u \in \mathcal{U} : \|u_i\|_{L^p(0, \infty)} \leq \gamma_i, \quad 1 \leq i \leq m\}, \end{aligned}$$

with $y_d \in L^2(Q)$, $\mathcal{U} = L^p(0, \infty)^m \cap L^1(0, \infty)^m$, $m \in \mathbb{N}$, $2 \leq p \leq \infty$, and $0 < \gamma_i < \infty$ for $1 \leq i \leq m$. Here y_u denotes the solution of the equation:

$$\begin{cases} \partial_t y_u + A y_u + f(x, t, y_u) = g(x, t) + B u & \text{in } Q = \Omega \times (0, \infty), \\ \partial_{n_A} y_u = 0 & \text{on } \Sigma = \Gamma \times (0, \infty), \quad y_u(0) = y_0 & \text{in } \Omega, \end{cases} \quad (1.1)$$

where $Bu = \sum_{i=1}^m u_i(t) \psi_i(x)$ for some functions $\{\psi_i\}_{i=1}^m \subset L^\infty(\Omega)$ with $\text{supp}(\psi_i) \cap \text{supp}(\psi_j) = \emptyset$ for $i \neq j$, and A is the linear elliptic operator

$$Ay = - \sum_{i,j=1}^n \partial_{x_j} [a_{ij}(x, t) \partial_{x_i} y] + a_0(x, t)y.$$

Assumptions on the coefficients of A and the functions f , g , and y_0 will be given in the next section.

We observe that \mathcal{U} is continuously embedded in $L^2(0, \infty)^m$, which follows by interpolation between the spaces $L^1(0, \infty)^m$ and $L^p(0, \infty)^m$. As a consequence we also have that $Bu \in L^2(Q) \cap L^p(0, \infty; L^\infty(\Omega))$.

By studying (P) we continue our investigations of nonlinear pde-constrained optimal control problems over infinite time horizons. Such problems have received little attention, although they arise quite naturally, for example, in the context of optimal stabilization, or in the case of modeling with finite horizons, where the length of the horizon is ambiguous and choosing an infinite horizon would be a safe way out. For infinite horizon optimal control problems related to ordinary differential equations we cite the monograph [4], and selected papers [1], [3], [19], where the latter might well be one of the earliest publications on the subject. In our work we treated infinite horizon problems related to semilinear parabolic equations in e.g. [10] and to the Navier Stokes equations in [12]. The specificity of problem (P) which is not considered in our previous work consists in the fact that (P) does not contain a quadratic space-time cost of the control, but rather the sparsifying term in time only, and that the explicit control constraint is of energy type in time, rather than pointwise in time as it was in our previous work. This necessitates a different treatment of the optimality conditions especially the second order necessary and sufficient optimality conditions. The control action enters (1.1) by means of finitely many time dependent controls u_i . The spatial distribution of the controllers is fixed and described by ψ_i . We shall utilize a

feasibility assumption guaranteeing that there exists at least one control vector u satisfying the constraints and rendering the corresponding cost J finite. Such an assumption is justified by stabilizability results which guarantee that for properly chosen m and ψ_i , there exists a control in feedback form for which the associated state decays exponentially, see e.g. [2]. The sparsifying term j in the cost then guarantees that optimal controls will individually shut off and remain shut off from times T_i^* on.

The paper is organised as follows. Section 2 presents properties of the state equation needed for the analysis that follows. In the following section 3 we prove the existence of a solution for (P) and establish the first order necessary conditions. Second order necessary and sufficient optimality conditions for (P) are presented in Section 4, where special attention is paid to the fact that the gap between these two conditions is small. Finally, section 5 is devoted to the approximation of the infinite horizon control problem by finite horizon control problems.

2 Analysis of the state equation

In this section, we are concerned with the existence, uniqueness, and regularity of the solution of (1.1). To this end we make the following assumptions:

Assumption 2.1.

We assume $1 \leq n \leq 3$, $y_0 \in L^\infty(\Omega)$ and $g \in L^r(0, \infty; L^s(\Omega)) \cap L^2(Q)$ with $\frac{1}{r} + \frac{n}{2s} < 1$.

The function $f : Q \times \mathbb{R} \rightarrow \mathbb{R}$ is measurable with respect to $(x, t) \in Q$ and of class C^2 with respect to $y \in \mathbb{R}$ and satisfies the following hypotheses

$$f(x, t, 0) = 0, \quad (2.1)$$

$$\exists \delta_f \in [0, 1) : \frac{\partial f}{\partial y}(x, t, 0) \geq -\delta_f a_0(x, t), \quad (2.2)$$

$$\exists M_f > 0 : f(x, t, y)y \geq 0 \text{ and } \frac{\partial f}{\partial y}(x, t, y) \geq 0 \quad \forall |y| > M_f, \quad (2.3)$$

$$\forall M > 0 \exists C_{f,M} > 0 : \left| \frac{\partial^j f}{\partial y^j}(x, t, y) \right| \leq C_{f,M} \quad \forall |y| \leq M, \quad j = 1, 2, \quad (2.4)$$

$\forall \rho > 0$ and $\forall M > 0 \exists \varepsilon > 0$ such that

$$\left| \frac{\partial^2 f}{\partial y^2}(x, t, y_2) - \frac{\partial^2 f}{\partial y^2}(x, t, y_1) \right| \leq \rho \quad \forall |y_i| \leq M, i = 1, 2, \text{ if } |y_2 - y_1| \leq \varepsilon, \quad (2.5)$$

for almost all $(x, t) \in Q$.

The coefficients of A satisfy $a_{i,j}, a_0 \in L^\infty(Q)$ and

$$\exists \Lambda > 0 : \sum_{i,j=1}^n a_{i,j}(x,t) \xi_i \xi_j \geq \Lambda |\xi|^2 \quad \forall \xi \in \mathbb{R}^n \text{ and for a.a. } (x,t) \in Q, \quad (2.6)$$

$$a_0(x,t) \geq 0 \text{ for a.a. } (x,t) \in Q, \text{ and } a_0 \not\equiv 0. \quad (2.7)$$

We mention that (2.3) and (2.4) imply that

$$\frac{\partial f}{\partial y}(x,t,y) \geq -C_{f,M_f} \quad \forall y \in \mathbb{R} \text{ and for a.a. } (x,t) \in Q. \quad (2.8)$$

Hereafter, we will follow the standard notation

$$W(0,T) = \{y \in L^2(0,T; H^1(\Omega)) : \partial_t y \in L^2(0,T; H^1(\Omega)^*)\} \text{ for } 0 < T \leq \infty.$$

A function y is called solution of (1.1) if $y \in W(0,T) \cap L^\infty(Q_T)$ for all $0 < T < \infty$ and it satisfies the following equation in the variational sense

$$\begin{cases} \partial_t y + Ay + f(x,t,y) = g(x,t) + Bu & \text{in } Q_T = \Omega \times (0,T), \\ \partial_{n_A} y = 0 \text{ on } \Sigma_T = \Gamma \times (0,T), \quad y(0) = y_0 & \text{in } \Omega, \end{cases} \quad (2.9)$$

We know that (2.9) admits a unique solution y_u in $W(0,T) \cap L^\infty(Q_T)$ for every $u \in L^p(0,\infty)$ and all $T < \infty$. Moreover, if $u \in \mathcal{U}$ and $y_u \in L^2(Q)$ then the regularity $y_u \in W(0,\infty) \cap L^\infty(Q)$ holds. Further the following estimates are satisfied:

$$\|y_u\|_Q \leq K_1 \left(\|y_u\|_{L^2(Q)} + \|y_0\|_{L^2(\Omega)} + \|g + Bu\|_{L^2(Q)} \right), \quad (2.10)$$

$$\begin{aligned} \|y_u\|_{L^\infty(Q)} &\leq K_2 \left(\|y_u\|_{L^2(Q)} + \|y_0\|_{L^\infty(\Omega)} + \|g + Bu\|_{L^2(Q)} \right. \\ &\quad \left. + \|g\|_{L^r(0,\infty; L^s(\Omega))} + \|Bu\|_{L^p(0,\infty; L^\infty(\Omega))} + M_f \right), \end{aligned} \quad (2.11)$$

where

$$\|y_u\|_Q = \left(\|y_u\|_{L^2(0,\infty; H^1(\Omega))}^2 + \|y_u\|_{L^\infty(0,\infty; L^2(\Omega))}^2 \right)^{\frac{1}{2}};$$

see [8] for the proof. Using (2.1), (2.4), and the mean value theorem we infer

$$|f(x,t, y_u(x,t))| = \left| \frac{\partial f}{\partial y}(x,t, \theta(x,t) y_u(x,t)) \right| |y_u(x,t)| \leq C_{f,M} |y_u(x,t)|,$$

where $M = \|y_u\|_{L^\infty(Q)}$ and $0 \leq \theta(x,t) \leq 1$. Since $y_u \in L^2(Q) \cap L^\infty(Q)$, the above inequality implies that $f(\cdot, \cdot, y_u) \in L^2(Q) \cap L^\infty(Q)$ and the following estimates hold

$$\|f(\cdot, \cdot, y_u)\|_{L^\infty(Q)} \leq C_{f,M} M \quad \text{and} \quad \|f(\cdot, \cdot, y_u)\|_{L^2(Q)} \leq C_{f,M} \|y_u\|_{L^2(Q)}. \quad (2.12)$$

Using these properties we deduce from the equation (1.1) that $y_u \in W(0, \infty) \cap L^\infty(Q)$. Moreover, from (2.10) and (2.12) we get

$$\|y_u\|_{W(0, \infty)} \leq K_3 \left((1 + C_{f, M}) \|y_u\|_{L^2(Q)} + \|y_0\|_{L^2(\Omega)} + \|g + Bu\|_{L^2(Q)} \right). \quad (2.13)$$

We define the set

$$\mathcal{A} = \{u \in L^2(0, \infty)^m : y_u \in L^2(Q)\}.$$

Now we introduce the mapping $G : \mathcal{A} \rightarrow W(0, \infty) \cap L^\infty(Q)$ define by $G(u) = y_u$. The following theorem was proved in [10, Theorems 2.2 and 3.1].

THEOREM 2.2. *Let us assume that \mathcal{A} is not empty. Then, \mathcal{A} is an open subset of $L^2(0, \infty)^m$ and the mapping G is of class C^2 . Moreover, given $u \in \mathcal{A}$ and $v, v_1, v_2 \in L^2(0, \infty)^m$ we have that $z_v = G'(u)v$ and $z_{v_1, v_2} = G''(u)(v_1, v_2)$ are the unique solutions of the equations*

$$\begin{cases} \partial_t z + Az + \frac{\partial f}{\partial y}(x, t, y_u)z = Bv & \text{in } Q, \\ \partial_{n_A} z = 0 & \text{on } \Sigma, \quad z(0) = 0 & \text{in } \Omega, \end{cases} \quad (2.14)$$

$$\begin{cases} \partial_t z + Az + \frac{\partial f}{\partial y}(x, t, y_u)z = -\frac{\partial^2 f}{\partial y^2}(x, t, y_u)z_{v_1}z_{v_2} & \text{in } Q, \\ \partial_{n_A} z = 0 & \text{on } \Sigma, \quad z(0) = 0 & \text{in } \Omega. \end{cases} \quad (2.15)$$

3 Existence of a solution and first order optimality conditions for (P)

To address the existence of a solution of (P) first we observe that there exists a unique solution y_u of (1.1) for every $u \in \mathcal{U}_{ad}$. However, it could happen that the solution y_u does not belong to $L^2(Q)$ and, consequently, $J(u) = \infty$. In the sequel, we say that u is a feasible control for (P) if $u \in \mathcal{U}_{ad}$ and the associated state y_u belongs to $L^2(Q)$, or equivalently $u \in \mathcal{U}_{ad} \cap \mathcal{A}$. We point out that $\mathcal{U}_{ad} \cap \mathcal{A}$ is not necessarily convex. Hence, regarding the existence of a solution for (P) we have the following result.

THEOREM 3.1. *If there exists a feasible control u_0 for (P), then it has at least one solution.*

Proof. Let $\{u_k\}_{k=1}^\infty$ be a minimizing sequence of (P) formed by feasible controls with associated states $\{y_k\}_{k=1}^\infty$. Since $J(u_k) \rightarrow \inf (P) \leq J(u_0) < \infty$ we deduce that $\{u_k\}_{k=1}^\infty$ and $\{y_k\}_{k=1}^\infty$ are bounded in \mathcal{U} and $L^2(Q)$, respectively. From (2.11) and (2.13) we get that $\{y_k\}_{k=1}^\infty$ is bounded in $W(0, \infty) \cap L^\infty(Q)$. Moreover, the continuous embedding $\mathcal{U} \subset L^2(0, \infty)$ implies that $\{u_k\}_{k=1}^\infty$ is bounded in $L^2(0, \infty)^m$. Therefore, taking a subsequence we obtain that $(u_k, y_k) \xrightarrow{*} (\bar{u}, \bar{y})$ in $L^2(0, \infty)^m \cap L^p(0, \infty)^m \times W(0, \infty) \cap L^\infty(Q)$. This implies that $Bu_k \xrightarrow{*} B\bar{u}$ in

$L^2(Q) \cap L^p(0, \infty; L^\infty(\Omega))$. Using these properties one can pass to the limit in the state equation (1.1) and deduce that \bar{y} is the state associated with \bar{u} ; see [10, Theorem 2.1] for details. Moreover, since $u_k \xrightarrow{*} \bar{u}$ in $L^p(0, \infty)^m$ we deduce that \bar{u} satisfies the control constraints. It remains to prove that $\bar{u} \in L^1(0, \infty)^m$ and $J(\bar{u}) = \inf(P)$. To this end we proceed as follows. For every $T < \infty$, using the compact embedding $W(0, T) \subset L^2(Q_T)$ and the weak convergence $u_k \rightharpoonup \bar{u}$ in $L^1(0, T)^m$ we obtain

$$\begin{aligned} & \frac{1}{2} \int_0^T \int_\Omega |\bar{y} - y_d|^2 dx dt + \alpha \sum_{j=1}^m \int_0^T |\bar{u}_j| dt \\ & \leq \liminf_{k \rightarrow \infty} \left(\frac{1}{2} \int_{Q_T} |y_k - y_d|^2 dx dt + \alpha \sum_{j=1}^m \int_0^T |u_{k,j}| dt \right) \leq \liminf_{k \rightarrow \infty} J(u_k) = \inf(P). \end{aligned}$$

Taking the supremum as $T \rightarrow \infty$ we infer that $\bar{u} \in L^1(0, \infty)^m$ and $J(\bar{u}) \leq \inf(P)$, which concludes the proof. \square

Next we derive the first order optimality conditions satisfied by a local minimizer \bar{u} of (P). If nothing is specifically said, \bar{u} is called a local minimizer of (P) if $J(\bar{u}) < \infty$ and there exists $\varepsilon > 0$ such that $J(\bar{u}) \leq J(u)$ for every $u \in \mathcal{U}_{ad}$ such that $\|u - \bar{u}\|_{\mathcal{U}} \leq \varepsilon$. By interpolation we have that \mathcal{U} is continuously embedded in $L^q(0, \infty)^m$ for every $q \in [1, p]$. Therefore, if \bar{u} is a local minimizer of (P), then it is also a local minimizer in the $L^q(0, \infty)^m$ sense.

To write the optimality conditions satisfied by a local minimizer we need to analyze separately the functions F and j defining the cost functional J . Regarding the functional F we make the following assumption on y_d :

Assumption 3.2. $y_d \in L^2(Q) \cap L^{\hat{r}}(0, \infty; L^{\hat{s}}(\Omega))$ with $\frac{1}{\hat{r}} + \frac{n}{2\hat{s}} < 1$.

As a straightforward consequence of this assumption and Theorem 2.2 we get that $F : \mathcal{A} \rightarrow \mathbb{R}$ is of class C^2 and for all $v, v_1, v_2 \in L^2(0, \infty)^m$ we have the following expressions

$$F'(u)v = \int_Q \varphi_u Bv dx dt = \sum_{i=1}^m \int_0^\infty \phi_{u,i}(t)v_i(t) dt, \quad (3.1)$$

$$F''(u)(v_1, v_2) = \int_Q \left(1 - \frac{\partial^2 f}{\partial y^2}(x, t, y_u) \right) z_{v_1} z_{v_2} dx dt, \quad (3.2)$$

where $\phi_{u,i}(t) = \int_\Omega \varphi_u(t) \psi_i dx$ and $\varphi_u \in W(0, \infty) \cap L^\infty(Q)$ is the adjoint state associated with u satisfying

$$\begin{cases} -\partial_t \varphi + A^* \varphi + \frac{\partial f}{\partial y}(x, t, y_u) \varphi = y_u - y_d & \text{in } Q, \\ \partial_{n_{A^*}} \varphi = 0 & \text{on } \Sigma, \quad \lim_{t \rightarrow \infty} \|\varphi(t)\|_{L^2(\Omega)} = 0 & \text{in } \Omega. \end{cases} \quad (3.3)$$

For the proof of (3.1) and (3.2) the reader is referred to [10, Theorem 2.3 and Corollary 3.1]. The existence and uniqueness of a solution of (3.3) was established

in [10, Theorem A.4]. We observe that the identity $\lim_{t \rightarrow \infty} \|\varphi_u(t)\|_{L^2(\Omega)} = 0$ is a consequence of the fact that $\varphi_u \in W(0, \infty)$; see [9, Theorem 2.4].

The functional $j : L^1(0, \infty)^m \rightarrow \mathbb{R}$ can be written in the form $j(u) = \sum_{i=1}^m j_0(u_i)$, where $j_0 : L^1(0, \infty) \rightarrow \mathbb{R}$ is defined by $j_0(w) = \|w\|_{L^1(0, \infty)}$. It is clear that j_0 is Lipschitz continuous and convex. Hence, the convex subdifferential $\partial j_0(w) \neq \emptyset$ and the directional derivatives $j'_0(w; v)$ exist for all $w, v \in L^1(0, \infty)$. The following properties hold:

$$\lambda \in \partial j_0(w) \text{ iff } \lambda(t) \begin{cases} = +1 & \text{if } w(t) > 0, \\ = -1 & \text{if } w(t) < 0, \\ \in [-1, +1] & \text{if } w(t) = 0, \end{cases} \quad (3.4)$$

$$j'_0(w; v) = \int_{I_w^+} v(t) dt - \int_{I_w^-} v(t) dt + \int_{I_w^0} |v(t)| dt, \quad (3.5)$$

where $I_w^+ = \{t \in (0, \infty) : w(t) > 0\}$, $I_w^- = \{t \in (0, \infty) : w(t) < 0\}$, and $I_w^0 = \{t \in (0, \infty) : w(t) = 0\}$.

We define the sets

$$\mathcal{U}_{ad}^i = \{v \in L^p(0, \infty) \cap L^1(0, \infty) : \|v\|_{L^p(0, \infty)} \leq \gamma_i\}, \quad 1 \leq i \leq m$$

and observe that $\mathcal{U}_{ad} = \prod_{i=1}^m \mathcal{U}_{ad}^i$. We have the following necessary optimality conditions.

THEOREM 3.3. *If \bar{u} is a local minimizer of (P), then there exist $\bar{\lambda}_i \in \partial j_0(\bar{u}_i)$ for $i = 1, \dots, m$ satisfying*

$$\int_0^\infty (\bar{\phi}_i(t) + \alpha \bar{\lambda}_i(t))(u(t) - \bar{u}_i(t)) dt \geq 0 \quad \forall u \in \mathcal{U}_{ad}^i, \quad (3.6)$$

where $\bar{\phi}_i(t) = \int_\Omega \bar{\varphi}(t) \psi_i dx$ and $\bar{\varphi}$ is the adjoint state associated with \bar{u} .

Proof. Using the convexity of \mathcal{U}_{ad} and j we get for every $u \in \mathcal{U}_{ad}$

$$0 \leq \lim_{\rho \rightarrow 0} \frac{J(\bar{u} + \rho(u - \bar{u})) - J(\bar{u})}{\rho} \leq F'(\bar{u})(u - \bar{u}) + \alpha[j(u) - j(\bar{u})]. \quad (3.7)$$

Let us fix $i \in \{1, \dots, m\}$ and take $u \in \mathcal{U}_{ad}$ with $u_k = \bar{u}_k$ if $k \neq i$ and $u_i = v \in \mathcal{U}_{ad}^i$. Then using (3.1) and the definition of $\bar{\phi}_i$ we infer from the above inequality

$$\int_0^\infty \bar{\phi}_i(t)(v(t) - \bar{u}_i(t)) dt + \alpha[j_0(v) - j_0(\bar{u}_i)] \geq 0 \quad \forall v \in \mathcal{U}_{ad}^i.$$

Denoting by $I_{\mathcal{U}_{ad}^i} : L^1(0, \infty) \cap L^p(0, \infty) \rightarrow [0, \infty]$ the indicator function of the convex set \mathcal{U}_{ad}^i we obtain that \bar{u}_i is the minimizer of the convex function

$$\mathcal{J}(v) = \int_0^\infty \bar{\phi}_i(t)v(t) dt + \alpha j_0(v) + I_{\mathcal{U}_{ad}^i}(v).$$

This implies the existence of $\bar{\lambda}_i \in \partial j_0(\bar{u}_i)$ such that

$$0 \in \bar{\phi}_i + \alpha \bar{\lambda}_i + \partial I_{\mathcal{U}_{ad}^i}(\bar{u}_i),$$

which is equivalent to (3.6). \square

COROLLARY 3.4. *Let $\bar{u} = \{\bar{u}_i\}_{i=1}^m$ and $\{(\bar{\lambda}_i, \bar{\phi}_i)\}_{i=1}^m$ be as in Theorem 3.3. Then, $\{(\bar{\lambda}_i, \bar{\phi}_i)\}_{i=1}^m$ are continuous functions in $[0, \infty)$ and the following relations hold for all $t \in [0, \infty)$ and all $i = 1, \dots, m$*

$$\text{If } \|\bar{u}_i\|_{L^p(0, \infty)} < \gamma_i, \text{ then } \bar{\phi}_i(t) + \alpha \bar{\lambda}_i(t) = 0, \quad (3.8)$$

$$\bar{\lambda}_i(t) = \text{Proj}_{[-1, +1]} \left(-\frac{1}{\alpha} \bar{\phi}_i(t) \right). \quad (3.9)$$

If $p \in [2, \infty)$ and $\bar{\phi}_i + \alpha \bar{\lambda}_i \equiv 0$, then for almost every $t \in [0, \infty)$ we have

$$\begin{aligned} \text{if } \bar{u}_i(t) > 0, \text{ then } \bar{\phi}_i(t) &= -\alpha, \\ \text{if } \bar{u}_i(t) < 0, \text{ then } \bar{\phi}_i(t) &= +\alpha, \\ \text{if } |\bar{\phi}_i(t)| < \alpha, \text{ then } \bar{u}_i(t) &= 0. \end{aligned} \quad (3.10)$$

If $p \in [2, \infty)$ and $\bar{\phi}_i + \alpha \bar{\lambda}_i \not\equiv 0$, then \bar{u} is a continuous function and we have for all $t \in [0, \infty)$

$$\begin{aligned} \bar{u}_i(t) > 0 &\text{ iff } \bar{\phi}_i(t) < -\alpha, \\ \bar{u}_i(t) < 0 &\text{ iff } \bar{\phi}_i(t) > +\alpha, \\ \bar{u}_i(t) = 0 &\text{ iff } |\bar{\phi}_i(t)| \leq \alpha. \end{aligned} \quad (3.11)$$

If $p = \infty$, then we have for almost all $t \in [0, \infty)$

$$\begin{aligned} \text{if } \bar{\phi}_i(t) < -\alpha \text{ then } \bar{u}_i(t) &= +\gamma_i, \\ \text{if } \bar{\phi}_i(t) > +\alpha \text{ then } \bar{u}_i(t) &= -\gamma_i, \\ \text{if } |\bar{\phi}_i(t)| < \alpha \text{ then } \bar{u}_i(t) &= 0. \end{aligned} \quad (3.12)$$

Finally, the multipliers $\{\bar{\lambda}_i\}_{i=1}^m$ satisfying (3.6) are unique.

Proof. Since $\bar{\varphi} \in W(0, \infty) \subset C([0, \infty); L^2(\Omega))$, we deduce from the definition of $\bar{\phi}_i$ its continuity in $[0, \infty)$. The continuity of $\bar{\lambda}_i$ is a consequence of the identity (3.9), which will be proved below.

We start assuming that $\|\bar{u}_i\|_{L^p(0, \infty)} < \gamma_i$ and set $\varepsilon_i = \gamma_i - \|\bar{u}_i\|_{L^p(0, \infty)}$. Given $T < \infty$, for every $\hat{v} \in L^p(0, T)$ with $\|\hat{v}\|_{L^p(0, T)} < \varepsilon_i$ we put

$$v(t) = \begin{cases} \hat{v}(t) + \bar{u}_i(t) & \text{if } t \in (0, T), \\ \bar{u}_i(t) & \text{if } t \geq T. \end{cases}$$

Then, it is evident that $v \in \mathcal{U}_{ad}^i$ and, consequently, we deduce from (3.6) that

$$\int_0^T (\bar{\phi}_i + \alpha \bar{\lambda}_i) \hat{v} dt = \int_0^\infty (\bar{\phi}_i + \alpha \bar{\lambda}_i) (v - \bar{u}_i) dt \geq 0 \quad \text{for every } \|\hat{v}\|_{L^p(0, T)} < \varepsilon_i.$$

This implies that $\bar{\phi}_i + \alpha\bar{\lambda}_i = 0$ in $(0, T)$. Since T was arbitrarily large, (3.8) follows.

The relations (3.10) are an immediate consequence of (3.4). Let us prove (3.11). Since $\bar{\phi}_i + \alpha\bar{\lambda}_i \neq 0$, the equality $\|\bar{u}_i\|_{L^p(0, \infty)} = \gamma_i$ follows from (3.8). Then, there exists $T_0 < \infty$ such that $\bar{\phi}_i + \alpha\bar{\lambda}_i \neq 0$ in $[0, T_0]$ and $\|\bar{u}_i\|_{L^p(0, T_0)} > 0$ by (3.6). For every $T > T_0$ we put $\gamma_{i, T} = \|\bar{u}_i\|_{L^p(0, T)} > 0$. For all $v \in L^p(0, T)$ such that $\|v\|_{L^p(0, T)} \leq \gamma_{i, T}$ we define

$$\hat{v}(t) = \begin{cases} v(t) & \text{if } t \in (0, T), \\ \bar{u}_i(t) & \text{if } t \geq T. \end{cases}$$

We observe that $\hat{v} \in L^p(0, \infty) \cap L^1(0, \infty)$ and $\|\hat{v}\|_{L^p(0, \infty)} \leq \|\bar{u}_i\|_{L^p(0, \infty)} = \gamma_i$, and thus $\hat{v} \in \mathcal{U}_{ad}^i$. Then, for $\bar{\eta}_i = \bar{\phi}_i + \alpha\bar{\lambda}_i$, (3.6) leads to

$$\int_0^T \bar{\eta}_i(t)(v(t) - \bar{u}_i(t)) dt \geq 0.$$

This yields

$$-\int_0^T \bar{\eta}_i v dt \leq -\int_0^T \bar{\eta}_i \bar{u}_i dt \leq \|\bar{\eta}_i\|_{L^{p'}(0, T)} \|\bar{u}_i\|_{L^p(0, T)} = \gamma_{i, T} \|\bar{\eta}_i\|_{L^{p'}(0, T)}.$$

Taking the supremum over all elements $v \in L^p(0, T)$ such that $\|v\|_{L^p(0, T)} \leq \gamma_{i, T}$ we obtain

$$\gamma_{i, T} \|\bar{\eta}_i\|_{L^{p'}(0, T)} \leq -\int_0^T \bar{\eta}_i(t) \bar{u}_i(t) dt \leq \gamma_{i, T} \|\bar{\eta}_i\|_{L^{p'}(0, T)}.$$

This implies the representation formula

$$\bar{u}_i(t) = -\gamma_{i, T} \frac{[\bar{\phi}_i(t) + \alpha\bar{\lambda}_i(t)]^{p'-2} (\bar{\phi}_i(t) + \alpha\bar{\lambda}_i(t))}{\|\bar{\phi}_i + \alpha\bar{\lambda}_i\|_{L^{p'}(0, T)}^{p'-1}} \quad \text{a.e. in } (0, T). \quad (3.13)$$

Since $T > T_0$ is arbitrary, from the above formula and (3.4), the relations (3.11) follow. Moreover, if $p \in [2, \infty)$ then (3.9) is deduced from (3.8) and (3.11) by simple computations. The continuity of $\bar{\lambda}_i$ is consequence of the continuity of $\bar{\phi}_i$ and (3.9). Finally, the continuity of \bar{u}_i is consequence of (3.13) and the continuity of $\bar{\lambda}_i$ and $\bar{\phi}_i$.

For the proof of (3.12) and the associated representation formula (3.9) the reader is referred to [5]. The uniqueness of the multipliers $\{\bar{\lambda}_i\}_{i=1}^m$ is an immediate consequence of (3.9). \square

COROLLARY 3.5. *Let $\bar{u} \in \mathcal{U}_{ad}$ satisfy the optimality conditions (3.6), then there exist times $\{T_i^*\}_{i=1}^m \subset (0, \infty)$ such that $\bar{u}_i(t) = \bar{\phi}_i(t) + \alpha\bar{\lambda}_i(t) = 0$ for every $t > T_i^*$, $1 \leq i \leq m$.*

Proof. Since $\lim_{t \rightarrow \infty} \|\bar{\varphi}(t)\|_{L^2(\Omega)} = 0$, we deduce that

$$\lim_{t \rightarrow \infty} |\bar{\phi}_i(t)| \leq \lim_{t \rightarrow \infty} \|\bar{\varphi}(t)\|_{L^2(\Omega)} \|\psi_i\|_{L^2(\Omega)} = 0, \quad 1 \leq i \leq m.$$

This leads the existence of $T_i^* < \infty$ such that $|\bar{\phi}_i(t)| < \alpha$ for every $t > T_i^*$. Then, the equality $\bar{u}_i(t) = 0$ follows from (3.10)–(3.12). Finally, it is enough to use the representation formula (3.13) to deduce that $\bar{\phi}_i(t) + \alpha \bar{\lambda}_i(t) = 0$ for every $t > T_i^*$ as well. \square

COROLLARY 3.6. *Let $\bar{u} \in \mathcal{U}_{ad}$ satisfy the optimality conditions (3.6) and assume that $\bar{\phi}_i + \alpha \bar{\lambda}_i \not\equiv 0$ for some $1 \leq i \leq m$. Then, the following representation formula for \bar{u}_i holds:*

$$\bar{u}_i(t) = -\gamma_i \frac{|\bar{\phi}_i(t) + \alpha \bar{\lambda}_i(t)|^{p'-2} (\bar{\phi}_i(t) + \alpha \bar{\lambda}_i(t))}{\|\bar{\phi}_i + \alpha \bar{\lambda}_i\|_{L^{p'}(0, \infty)}^{p'-1}} \quad \text{for a.e. } t \in [0, \infty). \quad (3.14)$$

Proof. We use (3.13) for $T > T_i^*$. As a consequence of Corollary 3.4 we have that $\gamma_{i,T} = \gamma_i$ and $\|\bar{\phi}_i + \alpha \bar{\lambda}_i\|_{L^{p'}(0,T)} = \|\bar{\phi}_i + \alpha \bar{\lambda}_i\|_{L^{p'}(0, \infty)}$. Hence, (3.13) implies (3.14). \square

4 Second order optimality conditions for (P)

In this section we address the second order optimality conditions for problem (P). We will distinguish the cases $2 \leq p < \infty$ and $p = \infty$.

4.1 Case I: $2 \leq p < \infty$.

We associate the following Lagrangian functions to the control problem (P)

$$\begin{aligned} \mathcal{L}, \mathcal{F} : \mathcal{A} \times \mathbb{R}^m &\longrightarrow \mathbb{R} \\ \mathcal{F}(u, \mu) &= F(u) + \frac{1}{p} \sum_{i=1}^m \frac{\mu_i}{\gamma_i^p} \|u_i\|_{L^p(0, \infty)}^p \quad \text{and} \quad \mathcal{L}(u, \mu) = \mathcal{F}(u, \mu) + \alpha j(u). \end{aligned}$$

According to (3.1) and (3.4) we have the following directional derivative

$$\begin{aligned} \frac{\partial \mathcal{L}}{\partial u}(u, \mu; v) &= \frac{\partial \mathcal{F}}{\partial u}(u, \mu)v + \alpha j'(u; v) \\ &= \sum_{i=1}^m \left(\int_0^\infty [\phi_i(t) + \frac{\mu_i}{\gamma_i^p} |u_i(t)|^{p-2} u_i(t)] v_i(t) dt + \alpha j'_0(u_i; v_i) \right), \quad (4.1) \end{aligned}$$

where $\phi_i(t) = \int_{\Omega} \varphi_u(t) \psi_i dx$. The second derivative of \mathcal{F} with respect to u is given by the expression

$$\begin{aligned} \frac{\partial^2 \mathcal{F}}{\partial u^2}(u, \mu)(v_1, v_2) &= \int_Q \left(1 - \frac{\partial^2 f}{\partial y^2}(x, t, y_u) \varphi_u\right) z_{v_1} z_{v_2} dx dt \\ &+ (p-1) \sum_{i=1}^m \frac{\mu_i}{\gamma_i^p} \int_0^{\infty} |u_i(t)|^{p-2} v_1(t) v_2(t) dt. \end{aligned} \quad (4.2)$$

Let $\bar{u} \in \mathcal{U}_{ad}$ be a control satisfying the first order optimality conditions (3.6). Associated with \bar{u} we define the Lagrange multipliers

$$\bar{\mu} = \{\bar{\mu}_i\}_{i=1}^m \quad \text{with} \quad \bar{\mu}_i = \gamma_i \|\bar{\phi}_i + \alpha \bar{\lambda}_i\|_{L^{p'}(0, \infty)}. \quad (4.3)$$

Since $\bar{\phi}_i, \bar{\lambda}_i \in L^\infty(0, \infty)$ and $\bar{\phi}_i(t) + \alpha \bar{\lambda}_i(t) = 0$ for $t > T_i^*$, we have that $\bar{\phi}_i + \alpha \bar{\lambda}_i \in L^q(0, \infty)$ for all $q \in [1, \infty]$ and every $1 \leq i \leq m$. The choice of $\bar{\mu}$ is justified by the next lemma.

PROPOSITION 4.1. *Let \bar{u} satisfy the first order optimality conditions (3.6) and let $\bar{\mu}$ be defined by (4.3). Then, the following properties hold:*

- i) $\bar{\mu}_i \geq 0$ and $\bar{\mu}_i (\|\bar{u}_i\|_{L^p(0, \infty)} - \gamma_i) = 0$ for $i = 1, \dots, m$,
- ii) $\frac{\partial \mathcal{L}}{\partial u}(\bar{u}, \bar{\mu}; v) \geq 0 \quad \forall v \in \mathcal{U}$,
- iii) $\frac{\partial \mathcal{L}}{\partial u}(\bar{u}, \bar{\mu}; v) = 0$ iff $|v_i(t)| = \bar{\lambda}_i(t) v_i(t)$ a.e. in $I_{\bar{u}_i}^0$ for $i = 1, \dots, m$.

Proof. The statement i) is an immediate consequence of the definition (4.3) of $\bar{\mu}_i$ and (3.8). Let us prove ii). If $\bar{\mu}_i = 0$, then the definition (4.3) along with (3.4) and (3.5) imply

$$\begin{aligned} &\int_0^{\infty} [\bar{\phi}_i(t) + \frac{\bar{\mu}_i}{\gamma_i^p} |\bar{u}_i(t)|^{p-2} \bar{u}_i(t)] v_i(t) dt + \alpha j'_0(\bar{u}_i; v_i) \\ &= -\alpha \int_0^{\infty} \bar{\lambda}_i(t) v_i(t) dt + \alpha j'_0(\bar{u}_i; v_i) = \alpha \int_{I_{\bar{u}_i}^0} [|v_i(t)| - \bar{\lambda}_i(t) v_i(t)] dt \geq 0. \end{aligned} \quad (4.4)$$

Now we consider the case where $\bar{\mu}_i > 0$. Using the representation formula (3.14) we get

$$|\bar{u}_i(t)|^{p-2} \bar{u}_i(t) = -\gamma_i^{p-1} \frac{\bar{\phi}_i(t) + \alpha \bar{\lambda}_i(t)}{\|\bar{\phi}_i + \alpha \bar{\lambda}\|_{L^{p'}(0, \infty)}}.$$

Combining this with the definition of $\bar{\mu}_i$ we obtain

$$\frac{\bar{\mu}_i}{\gamma_i^p} |\bar{u}_i(t)|^{p-2} \bar{u}_i(t) = -(\bar{\phi}_i(t) + \alpha \bar{\lambda}_i(t)).$$

Therefore, the same expressions as in (4.4) apply in this situation. Finally, (4.4) yields iii). \square

Remark 4.2. We observe that it was established in the above proof that

$$\int_0^\infty [\bar{\phi}_i(t) + \frac{\bar{\mu}_i}{\gamma_i^p} |\bar{u}_i(t)|^{p-2} \bar{u}_i(t)] v_i(t) dt + \alpha j_0'(\bar{u}_i; v_i) \geq 0 \quad \forall v \in \mathcal{U} \quad (4.5)$$

and for all $i = 1, \dots, m$.

Now we address the necessary second order conditions. To this end we define the cone of critical directions as follows

$$C_{\bar{u}} = \left\{ v \in \mathcal{U} : J'(\bar{u}; v) = 0 \text{ and } \int_0^\infty |\bar{u}_i(t)|^{p-2} \bar{u}_i(t) v_i(t) dt \leq 0 \quad \forall i \in S_0 \right\},$$

where $S_0 = \{i \in \{1, \dots, m\} : \|\bar{u}_i\|_{L^p(0, \infty)} = \gamma_i\}$. We also define $S_0^+ = \{i \in S_0 : \bar{\mu}_i > 0\}$. We have the following property on $C_{\bar{u}}$.

PROPOSITION 4.3. *If $v \in C_{\bar{u}}$ then we have*

$$\begin{aligned} i) & \int_0^\infty |\bar{u}_i(t)|^{p-2} \bar{u}_i(t) v_i(t) dt = 0 \quad \forall i \in S_0^+, \\ ii) & \frac{\partial \mathcal{L}}{\partial u}(\bar{u}, \bar{\mu}; v) = 0. \end{aligned}$$

Proof. Using Proposition (4.1), the definition of $C_{\bar{u}}$, and (4.1) we have that

$$\begin{aligned} 0 & \leq \frac{\partial \mathcal{L}}{\partial u}(\bar{u}, \bar{\mu}; v) = J'(\bar{u}; v) + \sum_{i=1}^m \bar{\mu}_i \int_0^\infty |\bar{u}_i(t)|^{p-2} \bar{u}_i(t) v_i(t) dt \\ & = \sum_{i \in S_0^+} \bar{\mu}_i \int_0^\infty |\bar{u}_i(t)|^{p-2} \bar{u}_i(t) v_i(t) dt \leq 0, \end{aligned}$$

This implies the statement of the proposition. \square

THEOREM 4.4. *If \bar{u} is a local minimizer of (P), then $\frac{\partial^2 \mathcal{F}}{\partial u^2}(\bar{u}, \bar{\mu}) v^2 \geq 0 \quad \forall v \in C_{\bar{u}}$.*

Proof. First we take an element $v \in C_{\bar{u}} \cap L^\infty(0, \infty)$ satisfying the following property

$$\exists \delta > 0 \text{ such that } v_i(t) = 0 \text{ if } 0 < |\bar{u}_i(t)| < \delta \text{ for } 1 \leq i \leq m. \quad (4.6)$$

We will get rid of these assumptions later. Let us denote

$$E_0 = \left\{ i \in S_0 : \int_0^\infty |\bar{u}_i(t)|^{p-2} \bar{u}_i(t) v_i(t) dt = 0 \right\}.$$

If $i \notin E_0$ we define the mappings $h_i : \mathbb{R} \rightarrow L^p(0, \infty) \cap L^1(0, \infty)$ and $\sigma_i : \mathbb{R} \rightarrow \mathbb{R}$ by $h_i(\rho) = \bar{u}_i + \rho v_i$ and $\sigma_i(\rho) = \|h_i(\rho)\|_{L^p(0, \infty)}^p$. If $i \notin S_0$, then $\sigma_i(0) < \gamma_i^p$ and, consequently, there exists $\varepsilon_i > 0$ such that $\sigma_i(\rho) < \gamma_i^p$ for every $|\rho| < \varepsilon_i$.

If $i \in S_0 \setminus E_0$, then $\sigma_i(0) = \gamma_i^p$ and

$$\sigma_i'(0) = p \int_0^T |\bar{u}_i(t)|^{p-2} \bar{u}_i(t) v_i(t) dt < 0.$$

Again, this implies the existence of $\varepsilon_i > 0$ such that $\sigma_i(\rho) < \gamma_i^p$ for all $\rho \in (0, \varepsilon_i)$. In both cases we have that $h_i(\rho) \in \mathcal{U}_{ad}^i$ for every $\rho \in (0, \varepsilon_i)$. In all cases, we assume that $\varepsilon_i \leq \frac{\delta}{2\|v_i\|_{L^\infty(0,\infty)}}$.

If $i \in E_0$, then there exists $\varepsilon_i > 0$ such that

$$\|\bar{u}_i + \rho v_i\|_{L^p(0,\infty)} \geq \|\bar{u}_i\|_{L^p(0,\infty)} - \rho \|v_i\|_{L^p(0,\infty)} \geq \frac{\gamma_i}{2} \quad \text{if } |\rho| < \varepsilon_i.$$

We define $h_i : (-\varepsilon_i, \varepsilon_i) \rightarrow L^p(0, \infty) \cap L^1(0, \infty)$ by $h_i(\rho) = \gamma_i \frac{\bar{u}_i + \rho v_i}{\|\bar{u}_i + \rho v_i\|_{L^p(0,\infty)}}$.

This choice also implies that $h_i(\rho) \in \mathcal{U}_{ad}^i$. For $i \in E_0$ we define

$$0 < \varepsilon_i \leq \frac{1}{2} \min \left\{ \frac{\gamma_i}{\|v_i\|_{L^p(0,\infty)}}, \frac{\delta}{\|v_i\|_{L^\infty(0,\infty)}} \right\}.$$

For $0 < \varepsilon \leq \min\{\varepsilon_i : 1 \leq i \leq m\}$ the mapping $h : [0, \varepsilon) \rightarrow \mathcal{U}_{ad}$ given by $h(\rho) = (h_i(\rho))_{i=1}^m$ is well defined and of class C^2 . We observe that $h_i(0) = \bar{u}_i$ and $h_i'(0) = v_i$ for every $i = 1, \dots, m$ and, as a consequence, we get that $h(0) = \bar{u}$ and $h'(0) = v$. Associated with this function we set $w : [0, \varepsilon) \rightarrow \mathbb{R}$ by $w(\rho) = J(h(\rho))$. From Propositions 4.3-ii) and 4.1-iii), using (4.6) and the choice of ε it follows that $|h_i(\rho)(t)| = \bar{\lambda}_i(t) h_i(\rho)(t)$ for almost all $t \in (0, \infty)$. Hence we have

$$w(\rho) = F(h(\rho)) + \alpha \sum_{i=1}^m \int_0^\infty \bar{\lambda}_i h_i(\rho) dt.$$

Therefore, w is of class C^2 and satisfies $w(0) = J(\bar{u})$ and

$$w'(0) = F'(\bar{u})h'(0) + \alpha \sum_{i=1}^m \int_0^\infty \bar{\lambda}_i h_i'(0) dt = J'(\bar{u}; v) = 0.$$

Since \bar{u} is a local minimizer of (P), then 0 is a local minimizer of w , hence $w''(0) \geq 0$. Let us compute this derivative. First we observe that

$$w'(\rho) = F'(h(\rho))h'(\rho) + \alpha \sum_{i=1}^m \int_0^\infty \bar{\lambda}_i h_i'(\rho) dt.$$

Derivating this expression we get

$$\begin{aligned} w''(0) &= F''(\bar{u})v^2 + F'(\bar{u})h''(0) + \alpha \sum_{i=1}^m \int_0^\infty \bar{\lambda}_i h_i''(0) dt \\ &= F''(\bar{u})v^2 + \sum_{i=1}^m \int_0^\infty (\bar{\phi}_i + \alpha \bar{\lambda}_i) h_i''(0) dt \\ &= F''(\bar{u})v^2 + \sum_{i \in S_0^+}^m \int_0^\infty (\bar{\phi}_i + \alpha \bar{\lambda}_i) h_i''(0) dt, \end{aligned}$$

where we have used that $\gamma_i \|\bar{\phi}_i + \alpha \bar{\lambda}_i\|_{L^p(0,\infty)} = \bar{\mu}_i = 0$ if $i \notin S_0^+$. Let us compute $h_i''(0)$ for $i \in S_0^+ \subset E_0$. The first derivative is given by

$$h_i'(\rho) = \frac{\gamma_i}{\|\bar{u}_i + \rho v_i\|_{L^p(0,\infty)}} v_i - \frac{\gamma_i \int_0^\infty |\bar{u}_i + \rho v_i|^{p-2} (\bar{u}_i + \rho v_i) v_i dt}{\|\bar{u}_i + \rho v_i\|_{L^p(0,\infty)}^{p+1}} (\bar{u}_i + \rho v_i).$$

Now we use Proposition 4.3-*i*) to deduce

$$h_i''(0) = (p-1) \frac{\int_0^\infty |\bar{u}_i|^{p-2} v_i^2 dt}{\gamma_i^p} \bar{u}_i.$$

Inserting this expression in the obtained formula for $w''(0)$ we infer with Hölder inequality, (4.3), (3.2), and (4.2) that

$$\begin{aligned} 0 \leq w''(0) &= F''(\bar{u})v^2 + (p-1) \sum_{i \in S_0^+} \frac{\int_0^\infty |\bar{u}_i|^{p-2} v_i^2 dt}{\gamma_i^p} \int_0^\infty (\bar{\phi}_i + \alpha \bar{\lambda}_i) \bar{u}_i dt \\ &\leq F''(\bar{u})v^2 + (p-1) \sum_{i \in S_0^+} \frac{\|\bar{\phi}_i + \alpha \bar{\lambda}_i\|_{L^{p'}(0,\infty)}}{\gamma_i^p} \|\bar{u}_i\|_{L^p(0,\infty)} \int_0^\infty |\bar{u}_i|^{p-2} v_i^2 dt \\ &\leq F''(\bar{u})v^2 + (p-1) \sum_{i \in S_0^+} \frac{\bar{\mu}_i}{\gamma_i^p} \int_0^\infty |\bar{u}_i|^{p-2} v_i^2 dt \\ &= F''(\bar{u})v^2 + (p-1) \sum_{i=1}^m \frac{\bar{\mu}_i}{\gamma_i^p} \int_0^\infty |\bar{u}_i|^{p-2} v_i^2 dt = \frac{\partial^2 \mathcal{F}}{\partial u^2}(\bar{u}, \bar{\mu})v^2. \end{aligned}$$

To conclude the proof we prove that all element $v \in C_{\bar{u}}$ can be approximated in the norm of \mathcal{U} by a sequence $\{v_k\}_{k=1}^\infty \subset C_{\bar{u}} \cap L^\infty(0,\infty)$ satisfying (4.6) for appropriate $\delta_k > 0$. By doing this we obtain that

$$\frac{\partial^2 \mathcal{F}}{\partial u^2}(\bar{u}, \bar{\mu})v^2 = \lim_{k \rightarrow \infty} \frac{\partial^2 \mathcal{F}}{\partial u^2}(\bar{u}, \bar{\mu})v_k^2 \geq 0.$$

Let us construct such a sequence $\{v_k\}_{k=1}^\infty$. Given $v \in C_{\bar{u}}$, for every $i = 1, \dots, m$ and every integer $k \geq 1$ we introduce the functions

$$\hat{v}_{i,k} = \begin{cases} 0 & \text{if } 0 < |\bar{u}_i(t)| < \frac{1}{k}, \\ v_i(t) & \text{otherwise.} \end{cases}$$

Now, we set $\hat{S}_0 = \{i \in S_0 : \int_0^\infty |\bar{u}_i|^{p-2} \bar{u}_i v_i dt = 0\}$. For every $i \in \hat{S}_0$ we put

$$\theta_{i,k} = \int_0^\infty |\bar{u}_i|^{p-2} \bar{u}_i \text{Proj}_{[-k,+k]}(\bar{u}_i) dt \text{ and } \varepsilon_{i,k} = \frac{1}{\theta_{i,k}} \int_0^\infty |\bar{u}_i|^{p-2} \bar{u}_i \frac{\hat{v}_{i,k}}{1 + \frac{1}{k} |\hat{v}_{i,k}|} dt.$$

Finally we define

$$v_{i,k}(t) = \begin{cases} \frac{\hat{v}_{i,k}(t)}{1 + \frac{1}{k} |\hat{v}_{i,k}(t)|} - \varepsilon_{i,k} \text{Proj}_{[-k,+k]}(\bar{u}_i(t)) & \text{if } i \in \hat{S}_0, \\ \frac{\hat{v}_{i,k}(t)}{1 + \frac{1}{k} |\hat{v}_{i,k}(t)|} & \text{otherwise.} \end{cases}$$

It is obvious that $\theta_{i,k} \rightarrow \|\bar{u}_i\|_{L^p(0,\infty)}^p = \gamma_i^p$ for all $i \in S_0$. We also have that $\frac{\hat{v}_{i,k}(t)}{1 + \frac{1}{k} |\hat{v}_{i,k}(t)|} \rightarrow v_i(t)$ pointwise in $(0, \infty)$ and the sequence is dominated by $v_i \in L^p(0, \infty) \cap L^1(0, \infty)$. Hence, $\frac{\hat{v}_{i,k}}{1 + \frac{1}{k} |\hat{v}_{i,k}|} \rightarrow v_i$ strongly in $L^p(0, \infty) \cap L^1(0, \infty)$. As a consequence we get

$$\varepsilon_{i,k} \rightarrow \frac{1}{\gamma_i^p} \int_0^\infty |\bar{u}_i|^{p-2} \bar{u}_i v_i dt = 0 \quad \forall i \in \hat{S}_0.$$

Setting $v_k = (v_{i,k})_{k=1}^m$, we deduce from these facts that $v_k \rightarrow v$ strongly in \mathcal{U} and by the choice of $\theta_{i,k}$ and $\varepsilon_{i,k}$ we find for k large enough

$$\int_0^\infty |\bar{u}_i|^{p-2} \bar{u}_i v_{i,k} dt \begin{cases} = 0 & \forall i \in \hat{S}_0, \\ < 0 & \forall i \in S_0 \setminus \hat{S}_0. \end{cases}$$

Further, by construction it is obvious that $\{v_k\}_{k=1}^\infty \subset L^\infty(0, \infty)^m$. It remains to prove that $J'(\bar{u}; v_k) = 0$ to conclude that $\{v_k\}_{k=1}^\infty \subset C_{\bar{u}}$. To this end we first observe that since $v \in C_{\bar{u}}$, Propositions 4.3-*ii*) and 4.1-*iii*) imply that $|v_i(t)| = \bar{\lambda}_i(t) v_i(t)$ for almost all $t \in (0, \infty)$ and all $i = 1, \dots, m$. By construction, it is immediate to check that the property is satisfied by the functions $v_{i,k}$. Using again Proposition 4.1-*iii*) we infer that $\frac{\partial \mathcal{L}}{\partial \bar{u}}(\bar{u}, \bar{\mu}; v_k) = 0$ for every k . This leads to

$$0 = \frac{\partial \mathcal{L}}{\partial \bar{u}}(\bar{u}, \bar{\mu}; v_k) = J'(\bar{u}; v_k) + \sum_{i=1}^m \frac{\bar{\mu}_i}{\gamma_i^p} \int_0^\infty |\bar{u}_i(t)|^{p-2} \bar{u}_i(t) v_{i,k}(t) dt = J'(\bar{u}; v_k),$$

where we have used that the above integral vanishes if $i \in \hat{S}_0$ and $\bar{\mu}_i = 0$ if $i \notin \hat{S}_0$. \square

Now we address the second order sufficient optimality conditions. We limit this study to the case $p = 2$. In infinite dimension optimization, it is well known that we cannot consider the same cone $C_{\bar{u}}$ for the second order necessary and sufficient conditions. In general an extended cone is necessary to deal with the sufficient conditions; see, for instance, [13], [15], [18], or [20]. Given a control $\bar{u} \in \mathcal{U}_{ad}$ satisfying the first order optimality conditions (3.6), we define for every $\tau > 0$ the extended cone

$$C_{\bar{u}}^\tau = \left\{ v \in \mathcal{U} : J'(\bar{u}; v) \leq \tau \|z_v\|_{L^2(Q)} \text{ and } \int_0^\infty \bar{u}_i v_i dt \begin{cases} \leq 0 & \forall i \in S_0 \\ \geq -\tau \|z_v\|_{L^2(Q)} & \forall i \in S_0^+ \end{cases} \right\},$$

where $z_v = G'(\bar{u})v$ is the solution of (2.14) with u replaced by \bar{u} . Using Proposition 4.1-ii we get for every $v \in C_{\bar{u}}^\tau$

$$0 \leq \frac{\partial \mathcal{L}}{\partial u}(\bar{u}, \bar{\mu}; v) = J'(\bar{u}; v) + \sum_{i \in S_0^+} \frac{\bar{\mu}_i}{\gamma_i^2} \int_0^\infty \bar{u}_i(t) v_i(t) dt \leq J'(\bar{u}; v).$$

Thus, for every small $\tau > 0$ and all $v \in C_{\bar{u}}^\tau$ the terms $J'(\bar{u}; v)$ and $\int_0^\infty \bar{u}_i v_i dt$ for $i \in S_0^+$ are not necessarily zero, but they are small. Taking into account Proposition 4.3-i), we observe that $C_{\bar{u}} \subset C_{\bar{u}}^\tau$ for all $\tau > 0$ and $C_{\bar{u}}^\tau$ is a small extension of $C_{\bar{u}}$ if τ is small.

THEOREM 4.5. *Let $\bar{u} \in \mathcal{U}_{ad} \cap \mathcal{A}$ satisfy the first order optimality conditions (3.6) and the following second order condition:*

$$\exists \delta > 0 \text{ and } \exists \tau > 0 : \frac{\partial^2 \mathcal{F}}{\partial u^2}(\bar{u}, \bar{\mu})v^2 \geq \delta \|z_v\|_{L^2(Q)}^2 \quad \forall v \in C_{\bar{u}}^\tau, \quad (4.7)$$

where $z_v = G'(\bar{u})v$. Then, there exist $\varepsilon > 0$ and $\kappa > 0$ such that

$$J(\bar{u}) + \frac{\kappa}{2} \|y_u - \bar{y}\|_{L^2(Q)}^2 \leq J(u) \quad \forall u \in \mathcal{U}_{ad} : \|y_u - \bar{y}\|_{L^2(Q)} + \|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon. \quad (4.8)$$

Before proving this theorem we establish two auxiliary lemmas.

LEMMA 4.6. *Assume that $\bar{u} \in \mathcal{U}_{ad} \cap \mathcal{A}$. Then, there exist $\varepsilon_1 > 0$ and $\bar{M} > 0$ such that for every $u \in \mathcal{U}_{ad}$ with $\|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon_1$ we have that $u \in \mathcal{A}$ and $\|y_u\|_{W(0, \infty)} \leq \bar{M}$. Moreover, the following inequalities hold*

$$\|y_u - (\bar{y} + z_{u-\bar{u}})\|_{L^2(Q)} \leq K_1 \|y_u - \bar{y}\|_{L^\infty(Q)} \|y_u - \bar{y}\|_{L^2(Q)}, \quad (4.9)$$

$$\|y_u - \bar{y}\|_{L^2(Q)} \leq 2 \|z_{u-\bar{u}}\|_{L^2(Q)}, \quad (4.10)$$

$$\|z_{u-\bar{u}}\|_{L^2(Q)} \leq \frac{3}{2} \|y_u - \bar{y}\|_{L^2(Q)}, \quad (4.11)$$

$$\|z_{u,v} - z_v\|_{L^2(Q)} \leq K_2 \|y_u - \bar{y}\|_{L^\infty(Q)} \|z_v\|_{L^2(Q)} \quad \forall v \in L^2(0, \infty)^m, \quad (4.12)$$

$$\|z_{u,v}\|_{L^2(Q)} \leq 2 \|z_v\|_{L^2(Q)} \quad \forall v \in L^2(0, \infty)^m, \quad (4.13)$$

where $z_{u,v} = G'(u)v$, $z_v = G'(\bar{u})v$, and $z_{u-\bar{u}} = G'(\bar{u})(u - \bar{u})$.

Proof. Let u satisfy the assumptions of the lemma and set $w = y_u - \bar{y}$. Subtracting the equations satisfied by y_u and \bar{y} and performing a Taylor expansion of f around \bar{y} we obtain

$$\begin{cases} \partial_t w + Aw + \frac{\partial f}{\partial y}(x, t, \bar{y})w = B(u - \bar{u}) - \frac{1}{2} \frac{\partial^2 f}{\partial y^2}(x, t, y_\theta)(y_u - \bar{y})^2 & \text{in } Q, \\ \partial_{n_A} w = 0 \text{ on } \Sigma, \quad w(0) = 0 & \text{in } \Omega, \end{cases} \quad (4.14)$$

where $0 \leq \theta(x, t) \leq 1$ and $y_\theta = \bar{y} + \theta(y_u - \bar{y})$. Since $\|y_\theta\|_{L^\infty(Q)} \leq M = \|\bar{y}\|_{L^\infty(Q)} + \varepsilon_1$, we infer from assumption (2.4) and [10, Theorem A.3]

$$\|w\|_{W(0,T)} \leq C \left(C_\gamma + C_{f,M} \varepsilon_1 \|y_u - \bar{y}\|_{L^2(Q_T)} \right) \quad \forall T < \infty,$$

where C_γ depends on the parameters $\{\gamma_i\}_{i=1}^m$. This implies that ε_1 can be chosen small enough such that $\|w\|_{W(0,T)} \leq \hat{M}$ for every $T < \infty$ and some constant \hat{M} . Hence, the inequality $\|y_u\|_{W(0,\infty)} \leq \bar{M} = \hat{M} + \|\bar{y}\|_{W(0,\infty)}$ holds and, consequently, we have that $u \in \mathcal{A}$.

Now we set $\hat{w} = y_u - (\bar{y} + z_{u-\bar{u}}) = w - z_{u-\bar{u}}$. Subtracting the equations satisfied by w and $z_{u-\bar{u}}$ it follows

$$\begin{cases} \partial_t \hat{w} + A\hat{w} + \frac{\partial f}{\partial y}(x, t, \bar{y})\hat{w} = -\frac{1}{2} \frac{\partial^2 f}{\partial y^2}(x, t, y_\theta)(y_u - \bar{y})^2 & \text{in } Q, \\ \partial_{n_A} \hat{w} = 0 \text{ on } \Sigma, \quad \hat{w}(0) = 0 & \text{in } \Omega. \end{cases}$$

Arguing as we did for w we deduce (4.9). Now, we redefine $\varepsilon_1 = \min\{\varepsilon_1, \frac{1}{2K_1}\}$. Then, using (4.9) we infer

$$\begin{aligned} \|y_u - \bar{y}\|_{L^2(Q)} &\leq \|y_u - (\bar{y} + z_{u-\bar{u}})\|_{L^2(Q)} + \|z_{u-\bar{u}}\|_{L^2(Q)} \\ &\leq \frac{1}{2} \|y_u - \bar{y}\|_{L^2(Q)} + \|z_{u-\bar{u}}\|_{L^2(Q)}, \end{aligned}$$

which implies (4.10). Inequality (4.11) follows in a similar way:

$$\|z_{u-\bar{u}}\|_{L^2(Q)} \leq \|y_u - (\bar{y} + z_{u-\bar{u}})\|_{L^2(Q)} + \|y_u - \bar{y}\|_{L^2(Q)} \leq \frac{3}{2} \|y_u - \bar{y}\|_{L^2(Q)}.$$

Now, we prove (4.12). Setting $z = z_{u,v} - z_v$, subtracting the equations satisfied by $z_{u,v}$ and z_v , and using the mean value theorem it follows

$$\begin{cases} \partial_t z + Az + \frac{\partial f}{\partial y}(x, t, y_u)z = -\frac{1}{2} \frac{\partial^2 f}{\partial y^2}(x, t, \hat{y}_\theta)(y_u - \bar{y})z_v & \text{in } Q, \\ \partial_{n_A} z = 0 \text{ on } \Sigma, \quad z(0) = 0 & \text{in } \Omega. \end{cases}$$

Then, (4.12) is consequence of [10, Theorem A.3] applied to the above equation and the assumptions on y_u .

Finally, we redefine again $\varepsilon_1 = \min\{\varepsilon_1, \frac{1}{K_2}\}$. Then (4.13) is an immediate consequence of (4.12). \square

LEMMA 4.7. *Let $\bar{u} \in \mathcal{U} \cap \mathcal{A}$ and ε_1 be as in Lemma 4.6. Given $\rho > 0$, there exists $\varepsilon_\rho \in (0, \varepsilon_1]$ such that*

$$F(u) - F(\bar{u}) \geq F'(\bar{u})(u - \bar{u}) + \frac{1}{2} F''(\bar{u})(u - \bar{u})^2 - \frac{\rho}{2} \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \quad (4.15)$$

for every $u \in \mathcal{U}$ such that $\|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon_\rho$, where $z_{u-\bar{u}} = G'(\bar{u})(u - \bar{u})$.

Proof. As in the proof of Lemma 4.6 we set $w = y_u - \bar{y}$. Then, using the adjoint state equation (3.3) associated with \bar{u} and (4.14) we get

$$\begin{aligned}
F(u) - F(\bar{u}) &= \int_Q (\bar{y} - y_d)(y_u - \bar{y}) \, dx \, dt + \frac{1}{2} \int_Q (\bar{y} - y_u)^2 \, dx \, dt \\
&= \int_Q \left(-\partial_t \bar{\varphi} + A^* \bar{\varphi} + \frac{\partial f}{\partial y}(x, t, y_u) \bar{\varphi} \right) (y_u - \bar{y}) \, dx \, dt + \frac{1}{2} \int_Q (\bar{y} - y_u)^2 \, dx \, dt \\
&\quad \int_Q \bar{\varphi} \left(\partial_t w + Aw + \frac{\partial f}{\partial y}(x, t, \bar{y}) w \right) \, dx \, dt + \frac{1}{2} \int_Q (\bar{y} - y_u)^2 \, dx \, dt \\
&= \int_Q \bar{\varphi} B(u - \bar{u}) \, dx \, dt - \frac{1}{2} \int_Q \frac{\partial^2 f}{\partial y^2}(x, t, y_\theta) \bar{\varphi} (y_u - \bar{y})^2 \, dx \, dt + \frac{1}{2} \int_Q (\bar{y} - y_u)^2 \, dx \, dt \\
&= F'(\bar{u})(u - \bar{u}) + \frac{1}{2} \int_Q \left[1 - \frac{\partial^2 f}{\partial y^2}(x, t, y_\theta) \bar{\varphi} \right] (y_u - \bar{y})^2 \, dx \, dt. \tag{4.16}
\end{aligned}$$

From here and (3.2) we deduce

$$\begin{aligned}
F(u) - F(\bar{u}) &= F'(\bar{u})(u - \bar{u}) + \frac{1}{2} F''(\bar{u})(u - \bar{u})^2 \\
&\quad - \frac{1}{2} \left(\int_Q \left[1 - \frac{\partial^2 f}{\partial y^2}(x, t, \bar{y}) \bar{\varphi} \right] z_{u-\bar{u}}^2 \, dx \, dt - \int_Q \left[1 - \frac{\partial^2 f}{\partial y^2}(x, t, y_\theta) \bar{\varphi} \right] (y_u - \bar{y})^2 \, dx \, dt \right).
\end{aligned}$$

To prove (4.15) we have to estimate the difference of the last two integrals. To this end we proceed as follows:

$$\begin{aligned}
&\left| \int_Q \left[1 - \frac{\partial^2 f}{\partial y^2}(x, t, \bar{y}) \bar{\varphi} \right] z_{u-\bar{u}}^2 \, dx \, dt - \int_Q \left[1 - \frac{\partial^2 f}{\partial y^2}(x, t, y_\theta) \bar{\varphi} \right] (y_u - \bar{y})^2 \, dx \, dt \right| \\
&\leq \int_Q \left| z_{u-\bar{u}}^2 - (y_u - \bar{y})^2 \right| \, dx \, dt + \int_Q \left| \frac{\partial^2 f}{\partial y^2}(x, t, y_\theta) - \frac{\partial^2 f}{\partial y^2}(x, t, \bar{y}) \right| |\bar{\varphi}| z_{u-\bar{u}}^2 \, dx \, dt \\
&\quad + \int_Q \left| \frac{\partial^2 f}{\partial y^2}(x, t, y_\theta) \bar{\varphi} \right| \left| z_{u-\bar{u}}^2 - (y_u - \bar{y})^2 \right| \, dx \, dt = I_1 + I_2 + I_3.
\end{aligned}$$

For the first term we have with (4.9) and (4.10)

$$\begin{aligned}
|I_1| &\leq \|y_u - (\bar{y} + z_{u-\bar{u}})\|_{L^2(Q)} (\|z_{u-\bar{u}}\|_{L^2(Q)} + \|y_u - \bar{y}\|_{L^2(Q)}) \\
&\leq 6K_1 \|y_u - \bar{y}\|_{L^\infty(Q)} \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \leq \frac{\rho}{3} \|z_{u-\bar{u}}\|_{L^2(Q)}^2
\end{aligned}$$

if $\|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon_{\rho,1} = \min\{\varepsilon_1, \frac{\rho}{18K_1}\}$.

To estimate I_2 we use (2.5) and the fact that $\|y_\theta - \bar{y}\|_{L^\infty(Q)} \leq \|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon_1$. Hence, we deduce the existence of $\varepsilon_{\rho,2} \in (0, \varepsilon_1]$ such that for $\|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon_{\rho,2}$ we have

$$|I_2| \leq \left\| \frac{\partial^2 f}{\partial y^2}(x, t, y_\theta) - \frac{\partial^2 f}{\partial y^2}(x, t, \bar{y}) \right\|_{L^\infty(Q)} \|\bar{\varphi}\|_{L^\infty(Q)} \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \leq \frac{\rho}{3} \|z_{u-\bar{u}}\|_{L^2(Q)}^2.$$

The term I_3 is estimated almost as I_1

$$\begin{aligned} |I_3| &\leq C_{f,M} \|\bar{\varphi}\|_{L^\infty(Q)} \int_Q \left| z_{u-\bar{u}}^2 - (y_u - \bar{y})^2 \right| dx dt \\ &\leq 6C_{f,M} \|\bar{\varphi}\|_{L^\infty(Q)} K_1 \|y_u - \bar{y}\|_{L^\infty(Q)} \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \leq \frac{\rho}{3} \end{aligned}$$

if $\|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon_{\rho,3} = \min\{\varepsilon_1, \frac{\rho}{6C_{f,M} \|\bar{\varphi}\|_{L^\infty(Q)} K_1}\}$. Then, it is enough to take $\varepsilon_\rho = \min\{\varepsilon_{\rho,1}, \varepsilon_{\rho,2}, \varepsilon_{\rho,3}\}$ to deduce (4.15). \square

Proof of Theorem 4.5. Let ε_1 be the number given in Lemma 4.6. From (4.16) and (4.10) we infer

$$\begin{aligned} F(u) - F(\bar{u}) &\geq F'(\bar{u})(u - \bar{u}) - \frac{1}{2}(1 + C_{f,M} \|\bar{\varphi}\|_{L^\infty(Q)}) \|y_u - \bar{y}\|_{L^2(Q)}^2 \\ &\geq F'(\bar{u})(u - \bar{u}) - (1 + C_{f,M} \|\bar{\varphi}\|_{L^\infty(Q)}) \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \\ &= F'(\bar{u})(u - \bar{u}) - K_3 \|z_{u-\bar{u}}\|_{L^2(Q)}^2. \end{aligned} \quad (4.17)$$

Let $i_0 \in S_0^+$ satisfy $\frac{\bar{\mu}_{i_0}}{\gamma_{i_0}^2} = \min_{i \in S_0^+} \frac{\bar{\mu}_i}{\gamma_i^2}$. We put $\nu = \min\{1, \frac{\bar{\mu}_{i_0}}{\gamma_{i_0}^2}\}$ and $\tau_0 = \nu\tau$.

We take $\varepsilon \in (0, \varepsilon_1]$ such that $\frac{2\tau_0}{3\varepsilon} - K_3 \geq \frac{\delta}{4}$, where δ satisfies (4.7). We also assume that $\varepsilon \leq \varepsilon_\rho$ and $\rho = \frac{\delta}{2}$, where ε_ρ was introduced in Lemma 4.7.

Given an element $u \in \mathcal{U}_{ad}$ such that $\|y_u - \bar{y}\|_{L^2(Q)} + \|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon$, we distinguish two cases to prove (4.8).

Case I: $J'(\bar{u})(u - \bar{u}) > \tau_0 \|z_{u-\bar{u}}\|_{L^2(Q)}$.

Using the convexity of j , (4.17), (4.10), and the assumption $\|y_u - \bar{y}\|_{L^2(Q)} \leq \varepsilon$ we obtain

$$\begin{aligned} J(u) - J(\bar{u}) &\geq J'(\bar{u}; u - \bar{u}) - K_3 \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \geq \tau_0 \|z_{u-\bar{u}}\|_{L^2(Q)} - K_3 \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \\ &\geq \left(\frac{2\tau_0}{3\|y_u - \bar{y}\|_{L^2(Q)}} - K_3 \right) \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \geq \frac{\delta}{4} \|z_{u-\bar{u}}\|_{L^2(Q)}^2. \end{aligned}$$

Case II: $J'(\bar{u})(u - \bar{u}) \leq \tau_0 \|z_{u-\bar{u}}\|_{L^2(Q)}$.

Under this assumption we have that $u - \bar{u} \in C_{\bar{u}}^\tau$. Indeed, since $\tau_0 \leq \tau$, we have that $J'(\bar{u})(u - \bar{u}) \leq \tau \|z_{u-\bar{u}}\|_{L^2(Q)}$. Moreover, since $u \in \mathcal{U}_{ad}$ we have for every $i \in S_0$

$$\begin{aligned} \int_0^\infty \bar{u}_i(t)(u_i(t) - \bar{u}_i(t)) dt &\leq \|\bar{u}_i\|_{L^2(0,\infty)} \|u_i\|_{L^2(0,\infty)} - \|\bar{u}_i\|_{L^2(0,\infty)}^2 \\ &= \gamma_i \left(\|u_i\|_{L^2(0,\infty)} - \gamma_i \right) \leq 0. \end{aligned}$$

Now we check the last condition to prove that $u - \bar{u}$ belongs to the critical cone $C_{\bar{u}}^\tau$. From Proposition 4.1-ii we get

$$0 \leq \frac{\partial \mathcal{L}}{\partial u}(\bar{u}, \bar{\mu}; u - \bar{u}) = J'(\bar{u}; u - \bar{u}) + \sum_{i \in S_0^+} \frac{\bar{\mu}_i}{\gamma_i^2} \int_0^\infty \bar{u}_i(t)(u_i(t) - \bar{u}_i(t)) dt,$$

which implies with the definition of ν

$$\int_0^\infty \bar{u}_i(t)(u_i(t) - \bar{u}_i(t)) \, dt \geq -\frac{\gamma_i^2}{\mu_i} J'(\bar{u}; u - \bar{u}) \geq -\frac{\tau_0}{\nu} \|z_{u-\bar{u}}\|_{L^2(Q)} = -\tau \|z_{u-\bar{u}}\|_{L^2(Q)}$$

for every $i \in S_0^+$. Thus $u - \bar{u} \in C_{\bar{u}}^\tau$ holds. Using that $\|u_i\|_{L^2(0,\infty)} - \|\bar{u}_i\|_{L^2(0,\infty)} \leq 0$ for all $i \in S_0^+$, the convexity of j , Proposition 4.1-ii, (4.7), and (4.15) with $\rho = \frac{\delta}{2}$ obtain

$$\begin{aligned} J(u) - J(\bar{u}) &\geq \mathcal{L}(u, \bar{\mu}) - \mathcal{L}(\bar{u}, \bar{\mu}) \geq \frac{\partial \mathcal{L}}{\partial u}(\bar{u}, \bar{\mu}; u - \bar{u}) \\ &\quad + \frac{1}{2} \mathcal{F}''(\bar{u})(u - \bar{u})^2 - \frac{\delta}{4} \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \geq \frac{\delta}{4} \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \end{aligned}$$

Finally, using (4.10) we infer for both cases that

$$J(u) - J(\bar{u}) \geq \frac{\delta}{4} \|z_{u-\bar{u}}\|_{L^2(Q)}^2 \geq \frac{\delta}{16} \|y_u - \bar{y}\|_{L^2(Q)}^2. \quad \square$$

Remark 4.8. The inequality (4.15) is a key issue in the proof of Theorem 4.5. The way in which (4.15) is proved here is different from the usual procedure; see [16] or [14]. Here, the main difficulty to follow the approach of [16] or [14] is that we cannot perform a Taylor expansion of $F(u)$ around \bar{u} for arbitrary $u \in \mathcal{A}$ since it is not known whether $u_\theta = \bar{u} + \theta(u - \bar{u})$ is an element of \mathcal{A} . Despite the fact that $y_u, \bar{y} \in L^\infty(Q)$, we do not know if y_{u_θ} belongs to $L^\infty(Q)$.

COROLLARY 4.9. Under the assumptions of Theorem 4.5 there exist $\hat{\varepsilon} > 0$ and $\delta > 0$ such that

$$J(\bar{u}) + \frac{\kappa}{2} \|y_u - \bar{y}\|_{L^2(Q)}^2 \leq J(u) \quad \forall u \in \mathcal{U}_{ad} \cap B_{\hat{\varepsilon}}(\bar{u}),$$

where $B_{\hat{\varepsilon}}(\bar{u}) \subset \mathcal{A} \subset L^2(0, \infty)^m$ is the ball centered at \bar{u} and radius $\hat{\varepsilon}$.

This corollary is an immediate consequence of Theorem 4.5 and the continuity of the mapping $G : \mathcal{A} \rightarrow W(0, \infty) \cap L^\infty(Q)$.

4.2 Case II: $p = \infty$.

In this case, the control constraints are linear, consequently the second order analysis is simpler. We start by establishing the second order necessary conditions for optimality. Assuming that $\bar{u} \in \mathcal{U}_{ad}$ satisfies the first order optimality conditions (3.6), we define the associated cone of critical directions as follows:

$$C_{\bar{u}} = \{v \in \mathcal{U} : J'(\bar{u}; v) = 0 \text{ and } v \text{ satisfies (4.18)}\}$$

with

$$v_i(t) \begin{cases} \geq 0 & \text{if } \bar{u}_i(t) = -\gamma_i, \\ \leq 0 & \text{if } \bar{u}_i(t) = +\gamma_i, \end{cases} \quad 1 \leq i \leq m. \quad (4.18)$$

The proof of the following lemma can be found in [7].

LEMMA 4.10. *The following properties hold:*

i) $J'(\bar{u}; v) \geq 0$ for every v satisfying the sign conditions (4.18).

ii) For every $v \in C_{\bar{u}}$ we have

$$\int_0^T (\bar{\phi}_i(t) + \alpha \bar{\lambda}_i(t)) v_i(t) dt = 0 \quad \text{and} \quad j'_0(\bar{u}_i, v_i) = \int_0^T \bar{\lambda}_i(t) v_i(t) dt.$$

iii) $C_{\bar{u}}$ is a closed, convex cone in \mathcal{U} .

Since $u - \bar{u}$ satisfies the sign conditions (4.18) for every $u \in \mathcal{U}_{ad}$, the statement i) implies that $J'(u; u - \bar{u}) \geq 0$ for all $u \in \mathcal{U}_{ad}$.

Now, we formulate the second order necessary optimality conditions.

THEOREM 4.11. *Let $\bar{u} \in \mathcal{U}_{ad} \cap \mathcal{A}$ be a local minimizer of (P), then $F''(\bar{u})v^2 \geq 0$ for every $v \in C_{\bar{u}}$.*

Proof. Given $v \in C_{\bar{u}}$ we define for every integer $k \geq 1$

$$v_{i,k}(t) = \begin{cases} 0 & \text{if } \gamma_i - \frac{1}{k} < |\bar{u}_i(t)| < \gamma_i, \\ \text{Proj}_{[-k, +k]}(v_i(t)) & \text{otherwise,} \end{cases} \quad 1 \leq i \leq m.$$

It is immediate that $v_k \rightarrow v$ strongly in \mathcal{U} and there exists $\rho_k > 0$ such that $\bar{u} + \rho v_k \in \mathcal{A}$ for every $\rho \in (0, \rho_k)$. Following the steps of the proof of [6, Theorem 3.7], we obtain that $F''(\bar{u})v_k^2 \geq 0$. Then, we pass to the limit as $k \rightarrow \infty$ and get the desired result. \square

As we did for $p = 2$, we need to extend the critical cone $C_{\bar{u}}$ to formulate the second order sufficient optimality conditions. For every $\tau > 0$ we define

$$C_{\bar{u}}^\tau = \{v \in \mathcal{U} : J'(\bar{u}; v) \leq \tau \|z_v\|_{L^2(Q)} \text{ and } v \text{ satisfies (4.18)}\}.$$

Then we have the following result.

THEOREM 4.12. *Let $\bar{u} \in \mathcal{U}_{ad} \cap \mathcal{A}$ satisfy the first order optimality conditions (3.6) and the following second order condition:*

$$\exists \delta > 0 \text{ and } \exists \tau > 0 : \frac{\partial^2 F}{\partial u^2}(\bar{u}, \bar{\mu})v^2 \geq \delta \|z_v\|_{L^2(Q)}^2 \quad \forall v \in C_{\bar{u}}^\tau, \quad (4.19)$$

where $z_v = G'(\bar{u})v$. Then, there exist $\varepsilon > 0$ and $\kappa > 0$ such that

$$J(\bar{u}) + \frac{\kappa}{2} \|y_u - \bar{y}\|_{L^2(Q)}^2 \leq J(u) \quad \forall u \in \mathcal{U}_{ad} : \|y_u - \bar{y}\|_{L^2(Q)} + \|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon. \quad (4.20)$$

The proof of this theorem follows the same steps as the one of Theorem 4.5 with some simplifications. Given $u \in \mathcal{U}_{ad} \cap \mathcal{A}$, we know that $u - \bar{u}$ satisfies the sign conditions (4.18). Hence, $u - \bar{u} \in C_{\bar{u}}^\tau$ holds if and only if $J'(u; u - \bar{u}) \leq \tau \|z_{u-\bar{u}}\|_{L^2(Q)}$. Then, we distinguish two cases as in the proof of Theorem 4.5, but with $\tau_0 = \tau$. Then, the proof is the same just replacing \mathcal{L} by J and \mathcal{F} by F .

5 Approximation by finite horizon problems

In this section we consider the approximation of (P) by finite horizon optimal control problems and provide error estimates for these approximations. For every $0 < T < \infty$ we consider the control problem

$$(P_T) \quad \min_{u \in \mathcal{U}_{T,ad}} J_T(u),$$

where $\mathcal{U}_{T,ad} = \{u \in L^p(0, T)^m : \|u_i\|_{L^p(0, T)} \leq \gamma_i\}$,

$$J_T(u) = F_T(u) + \alpha j_T(u) = \frac{1}{2} \int_{Q_T} (y_{T,u} - y_d)^2 dx dt + \alpha \sum_{i=1}^m \int_0^T |u_i(t)| dt,$$

and $y_{T,u}$ is the solution of

$$\begin{cases} \frac{\partial y}{\partial t} + Ay + ay + f(x, t, y) = g + Bu \text{ in } Q_T = \Omega \times (0, T), \\ \frac{\partial y}{\partial n_A} = 0 \text{ on } \Sigma_T = \Gamma \times (0, T), \quad y(0) = y_0 \text{ in } \Omega. \end{cases} \quad (5.1)$$

For every control $u \in L^2(0, T)^m$ with associated state $y_{T,u}$ and adjoint state $\varphi_{T,u}$ we define extensions to $(0, \infty)$ and Q , denoted by \hat{u} , $\hat{y}_{T,u}$, and $\hat{\varphi}_{T,u}$, by setting $(\hat{u}, \hat{\varphi}_{T,u})(x, t) = (0, 0)$ if $t > T$ and $\hat{y}_{T,u}$ is the solution of (1.1) for $u = \hat{u}_T$. It is obvious that if $u \in \mathcal{U}_{T,ad}$, then $\hat{u} \in \mathcal{U}_{ad}$ holds. Given a local minimizer u_T of (P_T) , we denote by y_T and φ_T its associated state and adjoint state, respectively. Arguing as in the proof of Theorem 3.3, we have that u_T satisfies the following optimality conditions

$$\int_0^T (\phi_{T,i}(t) + \alpha \lambda_{T,i}(t))(u(t) - u_{T,i}(t)) dt \geq 0 \quad \forall u \in \mathcal{U}_{T,ad}^i, \quad (5.2)$$

where $\phi_{T,i}(t) = \int_{\Omega} \varphi_T(t) \psi_i dx$ and $\lambda_{T,i} \in \partial j_{T,0}(u_{T,i})$, $1 \leq i \leq m$. Hereafter, $j_{T,0} : L^1(0, T) \rightarrow \mathbb{R}$ denotes the mapping $j_{T,0}(u) = \|u\|_{L^1(0, T)}$.

As a consequence of (5.2), Corollary 3.4 is also satisfied with $(0, \infty)$ and $(\bar{u}, \bar{\lambda}, \bar{\phi})$ replaced by $(0, T)$ and (u_T, λ_T, ϕ_T) .

The next two theorems establish the convergence of the approximating problems (P_T) to (P) as $T \rightarrow \infty$.

THEOREM 5.1. *For every $T > 0$ the control problem (P_T) has at least one solution u_T . If (P) has a feasible control u_0 , then the extensions $\{\hat{u}_T\}_{T>0}$ of any family of solutions are bounded in $L^p(0, \infty)^m$. Every weak limit \bar{u} in $L^p(0, \infty)^m$ of a sequence $\{\hat{u}_{T_k}\}_{k=1}^{\infty}$ with $T_k \rightarrow \infty$ as $k \rightarrow \infty$ is a solution of (P). Moreover, the weak convergence $\hat{u}_{T_k} \rightharpoonup \bar{u}$ in $L^q(0, \infty)^m$ for $q \in (1, p]$ and the strong convergence $\hat{y}_{T_k} \rightarrow \bar{y}$ in $L^2(Q) \cap L^\infty(Q)$ hold.*

Before proving this theorem, we establish the following lemma.

LEMMA 5.2. Let $\bar{u} \in \mathcal{U} \cap \mathcal{A}$ satisfy that such that $\bar{u}(t) = 0$ for $t \geq T^*$ with $T^* \in (0, \infty)$. Let us denote by $\bar{y} \in W(0, \infty) \cap L^\infty(Q)$ its associated state. Then, for every $T \in [T^*, \infty)$ there exists $\varepsilon > 0$ such that for all $\phi \in L^\infty(\Omega)$ with $\|\phi\|_{L^\infty(\Omega)} < \varepsilon$ the problem

$$\begin{cases} \partial_t y + Ay + f(x, t, y) = g(x, t) & \text{in } Q^T = \Omega \times (T, \infty), \\ \partial_{n_A} y = 0 & \text{on } \Sigma^T = \Gamma \times (T, \infty), \quad y(T) = \bar{y}(T) + \phi & \text{in } \Omega \end{cases} \quad (5.3)$$

has a unique solution $y \in W(T, \infty) \cap L^\infty(Q^T)$. Moreover, there exists a constant C independent of ϕ such that

$$\begin{aligned} \|y\|_{W(T, \infty)} + \|y\|_{L^\infty(Q^T)} &\leq C \left(\|y\|_{L^2(Q^T)} + \|\bar{y}(T) + \phi\|_{L^\infty(\Omega)} \right. \\ &\quad \left. + \|g\|_{L^2(Q^T)} + \|g\|_{L^r(T, \infty; L^s(\Omega))} + M_f \right). \end{aligned} \quad (5.4)$$

Proof. We are going to deduce this result by applying the implicit function theorem. First, we define the space

$$Y = \{y \in W(T, \infty) \cap L^\infty(Q^T) : \partial_t y + Ay \in L^r(T, \infty; L^s(\Omega)) \cap L^2(Q^T)\}$$

and the function $\mathcal{G} : Y \times L^\infty(\Omega) \rightarrow [L^r(T, \infty; L^s(\Omega)) \cap L^2(Q^T)] \times L^\infty(\Omega)$ by

$$\mathcal{G}(y, \phi) = \left(\partial_t y + Ay + f(\cdot, \cdot, y) - g, y(T) - (\phi + \bar{y}(T)) \right).$$

Endowed with the graph norm, Y is a Banach space and \mathcal{G} is of class C^1 . For every $z \in Y$ we have

$$\frac{\partial \mathcal{G}}{\partial y}(y, \phi)z = (\partial_t z + Az + \frac{\partial f}{\partial y}(\cdot, \cdot, y)z, z(T)).$$

Obviously, we have that $\mathcal{G}(\bar{y}, 0) = 0$ and $\frac{\partial \mathcal{G}}{\partial y}(\bar{y}, 0) : Y \rightarrow [L^r(T, \infty; L^s(\Omega)) \cap L^2(Q^T)] \times L^\infty(\Omega)$ is a continuous linear mapping. To prove that it is an isomorphism we have to check that the equation

$$\begin{cases} \partial_t z + Az + \frac{\partial f}{\partial y}(x, t, y)z = h & \text{in } Q^T, \\ \partial_{n_A} z = 0 & \text{on } \Sigma^T, \quad z(T) = z_T & \text{in } \Omega \end{cases}$$

has a unique solution $z \in Y$ for every $(h, z_T) \in [L^r(T, \infty; L^s(\Omega)) \cap L^2(Q^T)] \times L^\infty(\Omega)$. This follows from [10, Theorem A.3] and [8]. Then, the statement of the lemma follows from the implicit function theorem. \square

Proof of Theorem 5.1. Since $\mathcal{U}_{T, ad}$ is not empty, the existence of solution for (P_T) is a classical result. Actually, one can easily adapt the existence proof of solution for (P) to (P_T) . Let y^0 be the solution of (1.1) corresponding to u_0 . By

definition of feasible control we have that $J(u_0) < \infty$. Using the optimality of u_T we obtain

$$J_T(u_T) \leq J_T(u_0) \leq J(u_0) \quad \forall T > 0.$$

This proves the boundedness of $\{\hat{u}_T\}_{T>0}$ in $L^1(0, \infty)^m$ and the existence of a constant K such that $\|y_T\|_{L^2(Q_T)} \leq K$ for every T . Moreover, from the fact that $\{\hat{u}_T\}_{T>0} \subset \mathcal{U}_{ad}$ we deduce the boundedness of $\{\hat{u}_T\}_{T>0}$ in $L^p(0, \infty)^m$. By interpolation between the spaces $L^1(0, \infty)^m$ and $L^p(0, \infty)^m$ we infer that $\{\hat{u}_T\}_{T>0}$ is a bounded sequence in $L^q(0, \infty)^m$ for every $q \in [1, p]$. Let $\{(\hat{u}_{T_k}, y_{T_k} \chi_{(0, T_k)})\}_{k=1}^\infty$ be a sequence with $T_k \rightarrow \infty$ as $k \rightarrow \infty$ converging weakly to (\bar{u}, \bar{y}) in $L^p(0, \infty)^m \times L^2(Q)$. This implies the weak convergence of $\{\hat{u}_{T_k}\}_{k=1}^\infty$ in $L^q(0, \infty)^m$ for every $q \in (1, p]$. Since $\{\hat{u}_{T_k}\}_{k=1}^\infty \subset \mathcal{U}_{ad}$ and \mathcal{U}_{ad} is closed in $L^p(0, \infty)^m$ and convex, we infer that $\bar{u} \in \mathcal{U}_{ad}$. Moreover, we can apply [10, Theorem A1] to the equation (5.1) and deduce the existence of a constant M_1 such that for all $k \geq 1$

$$\begin{aligned} \|y_{T_k}\|_{L^2(0, T_k; H^1(\Omega))} + \|y_{T_k}\|_{L^\infty(Q_{T_k})} &\leq M_1 = C \left(\|g + B\hat{u}_{T_k}\|_{L^2(Q)} \right. \\ &\left. + \|g\|_{L^r(0, \infty; L^s(\Omega))} + \|B\hat{u}_{T_k}\|_{L^p(0, \infty; L^\infty(\Omega))} + \|y_0\|_{L^\infty(\Omega)} + K + M_f \right). \end{aligned}$$

From this estimate and (2.12) we get the existence of a constant M_2 such that

$$\|f(\cdot, \cdot, y_{T_k})\|_{L^2(Q_{T_k})} + \|f(\cdot, \cdot, y_{T_k})\|_{L^\infty(Q_{T_k})} \leq M_2 \quad \forall k \geq 1.$$

The two above estimates and (5.1) imply that

$$\|y_{T_k}\|_{W(0, T_k)} + \|y_{T_k}\|_{L^\infty(Q_{T_k})} \leq M_3 \quad \forall k \geq 1 \quad (5.5)$$

for a new constant M_3 . Using the convergence of $y_{T_k} \rightharpoonup \bar{y}$ in $L^2(Q_T)$ for every $T < \infty$, the compactness of the embedding $W(0, T) \subset L^2(Q_T)$, and the above estimate, it is straightforward to pass to the limit in the equation

$$\begin{cases} \frac{\partial y_{T_k}}{\partial t} + Ay_{T_k} + ay_{T_k} + f(x, t, y_{T_k}) = g + Bu_{T_k} & \text{in } Q_T, \\ \partial_{n_A} y = 0 & \text{on } \Sigma_T, \quad y_{T_k}(0) = y_0 & \text{in } \Omega \end{cases} \quad (5.6)$$

for each $T_k \geq T$, and to deduce that \bar{y} is the solution of (5.1) associated to \bar{u} for arbitrary $0 < T < \infty$. This proves that \bar{y} is the solution of (1.1) corresponding to \bar{u} and $\|\bar{y}\|_{W(0, \infty)} + \|\bar{y}\|_{L^\infty(Q)} \leq M_3$. This implies that $\bar{u} \in \mathcal{A}$. Let us prove that \bar{u} is a solution of (P). Using the convergence $u_{T_k} \rightharpoonup \bar{u}$ in $L^1(Q_T)^m$ for every $T < \infty$, we get for every feasible control u of (P)

$$\begin{aligned} J_T(\bar{u}) &\leq \liminf_{k \rightarrow \infty} J_T(u_{T_k}) \leq \liminf_{k \rightarrow \infty} J_{T_k}(u_{T_k}) \\ &\leq \limsup_{k \rightarrow \infty} J_{T_k}(u_{T_k}) \leq \limsup_{k \rightarrow \infty} J_{T_k}(u) = J(u). \end{aligned}$$

Hence, the inequality $J(\bar{u}) = \sup_{T \rightarrow \infty} J_T(\bar{u}) \leq J(u)$ holds, which proves that \bar{u} is a solution of (P). Moreover, replacing u by \bar{u} in the above inequalities we infer

$$\lim_{k \rightarrow \infty} \left(\frac{1}{2} \int_{Q_{T_k}} (y_{T_k} - y_d)^2 dx dt + \alpha j_{T_k}(u_{T_k}) \right) = \frac{1}{2} \int_Q (\bar{y} - y_d)^2 dx dt + \alpha j(\bar{u}).$$

This is equivalent to the identity

$$\lim_{k \rightarrow \infty} \left(\frac{1}{2} \int_Q (y_{T_k} - y_d)^2 \chi_{(0, T_k)} dx dt + \alpha j(\hat{u}_{T_k}) \right) = \frac{1}{2} \int_Q (\bar{y} - y_d)^2 dx dt + \alpha j(\bar{u}).$$

Once more, using the convergence $u_{T_k} \rightarrow \bar{u}$ in $L^1(Q_T)^m$ for every $T < \infty$ we obtain

$$j_T(\bar{u}) \leq \liminf_{k \rightarrow \infty} j_T(u_{T_k}) \leq \liminf_{k \rightarrow \infty} j(\hat{u}_{T_k}).$$

Taking the supremum in T we deduce $j(\bar{u}) \leq \liminf_{k \rightarrow \infty} j(\hat{u}_{T_k})$. This convergence along with the weak convergence $y_{T_k} \chi_{(0, T_k)} \rightarrow \bar{y}$ in $L^2(Q)$ and the above equality yield the strong convergence $\lim_{k \rightarrow \infty} \|y_{T_k} - \bar{y}\|_{L^2(Q_{T_k})} = 0$; see [11, Lemma 5.2]. It remains to prove that $\hat{y}_{T_k} \rightarrow \bar{y}$ in $L^2(Q) \cap L^\infty(Q)$. The proof of this convergence is split in several steps.

Step I.- $\lim_{k \rightarrow \infty} \|\hat{y}_{T_k} - \bar{y}\|_{L^\infty(Q_T)} = 0$ for every $T < \infty$. Let us set $w_k = \hat{y}_{T_k} - \bar{y}$. Then, we have for every $T_k \geq T$

$$\begin{cases} \frac{\partial w_k}{\partial t} + A w_k + \frac{\partial f}{\partial y}(x, t, y_{\theta_k}) w_k = B(u_{T_k} - \bar{u}) \text{ in } Q_T, \\ \partial_{n_A} w_k = 0 \text{ on } \Sigma_T, \quad w_k(0) = 0 \text{ in } \Omega, \end{cases} \quad (5.7)$$

where $y_{\theta_k} = \bar{y} + \theta_k(\hat{y}_{T_k} - \bar{y})$ and $\theta_k : Q \rightarrow [0, 1]$ is a measurable function. Since $\{w_k\}_{k=1}^\infty$ is bounded in $L^2(Q_T) \cap L^\infty(Q_T)$, we get with (2.4) that $B(u_{T_k} - \bar{u}) - \frac{\partial f}{\partial y}(x, t, y_{\theta_k}) w_k$ is bounded in $L^p(0, T; L^\infty(\Omega))$. Then, we deduce from [17] the boundedness of $\{w_k\}_{k=1}^\infty$ in $C^{0, \beta}(\bar{Q}_T)$ for some $\beta \in (0, 1)$. Using the compactness of the embedding $C^{0, \beta}(\bar{Q}_T) \subset C(\bar{Q}_T)$ along with the strong convergence $y_{T_k} \rightarrow \bar{y}$ in $L^2(Q_T)$, we infer the strong convergence $w_k \rightarrow 0$ in $C(\bar{Q}_T)$ and $\hat{y}_{T_k} \rightarrow \bar{y}$ in $L^\infty(Q_T)$ as $k \rightarrow \infty$ for every $T < \infty$.

Step II.- There exist $T^* < \infty$ and $k^* \geq 1$ such that $\hat{u}_{T_k}(t) = 0$ for all $k \geq k^*$ and almost all $t > T^*$. Indeed, we have that the adjoint states $\hat{\varphi}_{T_k}$ satisfy the adjoint state equations

$$\begin{cases} -\partial_t \hat{\varphi}_{T_k} + A^* \hat{\varphi}_{T_k} + \frac{\partial f}{\partial y}(x, t, \hat{y}_{T_k} \chi_{(0, T_k)}) \hat{\varphi}_{T_k} = (\hat{y}_{T_k} - y_d) \chi_{(0, T_k)} \text{ in } Q, \\ \partial_{n_{A^*}} \hat{\varphi}_{T_k} = 0 \text{ on } \Sigma, \quad \lim_{t \rightarrow \infty} \|\hat{\varphi}_{T_k}(t)\|_{L^2(\Omega)} = 0 \text{ in } \Omega. \end{cases} \quad (5.8)$$

Given $\varepsilon > 0$, the convergence $\lim_{k \rightarrow \infty} \|y_{T_k} - \bar{y}\|_{L^2(Q_{T_k})} = 0$ and the fact that $\bar{y} - y_d \in L^2(Q)$ imply the existence of k_ε and T_ε such that for $k > k_\varepsilon$ and $T_k > T_\varepsilon$

$$\|(\hat{y}_{T_k} - y_d) \chi_{(0, T_k)}\|_{L^2(T_\varepsilon, \infty; L^2(\Omega))} \leq \|\hat{y}_{T_k} - \bar{y}\|_{L^2(Q_{T_k})} + \|\bar{y} - y_d\|_{L^2(T_\varepsilon, \infty; L^2(\Omega))} < \varepsilon.$$

Using this, we deduce from (5.8) and [10, Theorem A.4] that $\|\hat{\varphi}_{T_k}\|_{L^\infty(T_\varepsilon, \infty; L^2(\Omega))} \leq C\varepsilon$ for every $k \geq k_\varepsilon$. For every $1 \leq i \leq m$, this yields

$$|\phi_{T_k, i}(t)| \leq \|\hat{\varphi}_{T_k}(t)\|_{L^2(\Omega)} \|\psi_i\|_{L^2(\Omega)} \leq C \|\psi_i\|_{L^2(\Omega)} \varepsilon \quad \forall k \geq k_\varepsilon \text{ and for a.a. } t \geq T_\varepsilon.$$

Selecting $\varepsilon > 0$ such that $C \|\psi_i\|_{L^2(\Omega)} \varepsilon < \alpha$ and applying Corollary 3.4 we infer that $\hat{u}_{T_k}(t) = 0$ for every $k \geq k_\varepsilon$ and almost all $t > T_\varepsilon$.

Step III.- $\exists k_0 \geq 1$ such that $\{\hat{y}_{T_k}\}_{k \geq k_0} \subset W(0, \infty) \cap L^\infty(Q)$ and the convergence $\lim_{k \rightarrow \infty} (\|\hat{y}_{T_k} - \bar{y}\|_{L^2(Q)} + \|\hat{y}_{T_k} - \bar{y}\|_{L^\infty(Q)}) = 0$ holds. Without loss of generality we can assume that $T^* > T_i^*$ for $1 \leq i \leq m$, where $\{T_i^*\}_{i=1}^m$ are given in Corollary 3.5. Thus, we have that $B(u_{T_k} - \bar{u})(t) = 0 \forall k \geq k^*$ and for almost all $t \geq T^*$.

We take $T > T^*$. Since $w_k \rightarrow 0$ in $C(\bar{Q}_T)$, we have that $w_k(T) \rightarrow 0$ in $C(\bar{\Omega})$. Then, applying Lemma 5.2 we infer the existence of k_0 such that $\{\hat{y}_{T_k}\}_{k \geq k_0} \subset W(T, \infty) \cap L^\infty(Q^T)$ and it is uniformly bounded in this space. Combining this with (5.5) we infer that $\{\hat{y}_{T_k}\}_{k \geq k_0} \subset W(0, \infty) \cap L^\infty(Q)$. Moreover, applying [10, Theorem A.3] along with [8] to the equation

$$\begin{cases} \frac{\partial w_k}{\partial t} + Aw_k + aw_k + \frac{\partial f}{\partial y}(x, t, y_{\theta_k})w_k = 0 \text{ in } Q^T, \\ \partial_{n_A} w_k = 0 \text{ on } \Sigma^T, \end{cases}$$

we obtain

$$\|w_k\|_{W(T, \infty)} + \|w_k\|_{L^\infty(Q^T)} \leq C\|w_k(T)\|_{L^\infty(\Omega)} \rightarrow 0 \text{ as } k \rightarrow \infty.$$

Combining this with *Step I* we get the desired convergence. \square

Now we address a kind of converse theorem for strong local minimizers. We say that \bar{u} is a strong local minimizer of (P) if there exists $\varepsilon > 0$ such that

$$J(\bar{u}) \leq J(u) \quad \forall u \in \mathcal{U}_{ad} \cap \mathcal{A} \text{ satisfying } \|y_u - \bar{y}\|_{L^2(Q)} + \|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon. \quad (5.9)$$

If the above inequality is strict for $u \neq \bar{u}$, then we say that \bar{u} is a strict strong local minimizer.

THEOREM 5.3. *Let \bar{u} be a strict strong local minimizer of (P). Then, there exist $T_0 \in (0, \infty)$ and a family $\{u_T\}_{T > T_0}$ of strong local minimizers to (P_T) such that the weak convergence $\hat{u}_T \rightharpoonup \bar{u}$ in $L^q(0, \infty)^m$ for all $q \in (1, p]$ and the strong convergence $\hat{y}_T \rightarrow \bar{y}$ in $L^2(Q) \cap L^\infty(Q)$ hold as $T \rightarrow \infty$.*

Proof. Let us \bar{u} satisfy (5.9). We consider the control problems

$$(P_\varepsilon) \quad \min_{u \in \mathcal{U}_{ad}^\varepsilon} J(u) \quad \text{and} \quad (P_{T, \varepsilon}) \quad \min_{u \in \mathcal{U}_{T, ad}^\varepsilon} J_T(u),$$

where

$$\begin{aligned} \mathcal{U}_{ad}^\varepsilon &= \{u \in \mathcal{U}_{ad} \cap \mathcal{A} : \|y_u - \bar{y}\|_{L^2(Q)} + \|y_u - \bar{y}\|_{L^\infty(Q)} \leq \varepsilon\}, \\ \mathcal{U}_{T, ad}^\varepsilon &= \{u \in \mathcal{U}_{T, ad} : \|y_u - \bar{y}\|_{L^2(Q_T)} + \|y_u - \bar{y}\|_{L^\infty(Q_T)} \leq \varepsilon\}. \end{aligned}$$

Obviously \bar{u} is the unique solution of (P_ε) . Regarding the problem $(P_{T, \varepsilon})$ we first observe that $\bar{u}|_{(0, T)} \in \mathcal{U}_{T, ad}^\varepsilon$. Moreover, it is easy to check that if $\{u_k\}_{k=1}^\infty \subset \mathcal{U}_{T, ad}^\varepsilon$ and $u_k \rightharpoonup u$ in $L^p(0, T)^m$, then $y_{u_k} \xrightarrow{*} y_u$ in $L^\infty(Q_T)$. Hence, $\mathcal{U}_{T, ad}^\varepsilon$ is nonempty,

bounded, and sequentially weakly closed in $L^p(0, T)^m$. Then, for every T the existence of a solution u_T of $(P_{T, \varepsilon})$ can be proved as usual by taking a minimizing sequence. Now, arguing as in the proof of Theorem 5.1 and using the uniqueness of the solution of (P_ε) , we deduce the convergence $\hat{u}_T \rightharpoonup \bar{u}$ in $L^p(0, \infty)^m$ as $T \rightarrow \infty$ and $\hat{y}_T \rightarrow \bar{y}$ in $L^2(Q) \cap L^\infty(Q)$. This implies the existence of T_0 such that $\|\hat{y}_T - \bar{y}\|_{L^2(Q_T)} + \|\hat{y}_T - \bar{y}\|_{L^\infty(Q_T)} < \varepsilon$ for all $T > T_0$. Hence, u_T is also a strong local minimizer of (P_T) for $T > T_0$. Indeed, let us set $\varepsilon_T = \|\hat{y}_T - \bar{y}\|_{L^2(Q_T)} + \|\hat{y}_T - \bar{y}\|_{L^\infty(Q_T)}$. Then, for every $u \in \mathcal{U}_{T, ad}$ with $\|y_T - y_{T, u}\|_{L^2(Q_T)} + \|y_T - y_{T, u}\|_{L^\infty(Q_T)} \leq \varepsilon - \varepsilon_T$ we have

$$\begin{aligned} \|y_{T, u} - \bar{y}\|_{L^2(Q_T)} + \|y_{T, u} - \bar{y}\|_{L^\infty(Q_T)} &\leq \|y_T - y_{T, u}\|_{L^2(Q_T)} + \|y_T - y_{T, u}\|_{L^\infty(Q_T)} \\ &\quad + \|\hat{y}_T - \bar{y}\|_{L^2(Q_T)} + \|\hat{y}_T - \bar{y}\|_{L^\infty(Q_T)} < \varepsilon. \end{aligned}$$

Since u_T is a minimizer of $(P_{T, \varepsilon})$ and u is a feasible control for $(P_{T, \varepsilon})$, the inequality $J_T(u_T) \leq J_T(u)$ follows. \square

In the previous theorem we proved the existence of strong local minimizers $\{u_T\}_{T > T_0}$ of problems (P_T) weakly converging to \bar{u} , assuming that \bar{u} is a strict strong local minimizer of (P) . Moreover, the strong convergence of the associated states $\hat{y}_T \rightarrow \bar{y}$ in $L^2(Q) \cap L^\infty(Q)$ was established. In addition, the inequality $J_T(u_T) \leq J_T(\bar{u})$ holds for every $T > T_0$. In the next theorem we provide an estimate for the difference of the corresponding states.

THEOREM 5.4. *Suppose that $p = 2$ or $p = \infty$ and that \bar{u} is a strong local minimizer of (P) satisfying the second order sufficient optimality condition. We assume that $\frac{\partial^2 f}{\partial y^2}(x, t, y) \geq 0$ holds for all $y \in \mathbb{R}$ and almost all $(x, t) \in Q$. Let $\{u_T\}_{T > T_0}$ be a sequence of local minimizers of problems (P_T) such that $\hat{u}_T \rightharpoonup \bar{u}$ in $L^q(0, \infty)^m$ $\forall q \in (1, p]$, $\hat{y}_T \rightarrow \bar{y}$ in $L^2(Q) \cap L^\infty(Q)$, and $J_T(\bar{u}) \leq J_T(u_T)$. Then, there exist $T^* \in [T_0, \infty)$ and a constant C such that for every $T \geq T^*$*

$$\|\hat{y}_T - \bar{y}\|_{L^2(Q)} \leq C \left(\|y_T(T)\|_{L^2(\Omega)} + \|y_d\|_{L^2(T, \infty; L^2(\Omega))} + \|g\|_{L^2(T, \infty; L^2(\Omega))} \right). \quad (5.10)$$

Proof. We use the inequalities (4.8) or (4.20). For this purpose, we take $T^* \in [T_0, \infty)$ such that $\|\hat{y}_T - \bar{y}\|_{L^2(Q)} + \|\hat{y}_T - \bar{y}\|_{L^\infty(Q)} < \varepsilon$ for all $T \geq T^*$. Then, we have

$$\begin{aligned} \frac{\kappa}{2} \|\hat{y}_T - \bar{y}\|_{L^2(Q)}^2 &\leq J(\hat{u}_T) - J(\bar{u}) \leq J_T(u_T) - J_T(\bar{u}) \\ &\quad + \frac{1}{2} \int_T^\infty \|\hat{y}_T(t) - y_d(t)\|_{L^2(\Omega)}^2 dt \leq \frac{1}{2} \int_T^\infty \|\hat{y}_T(t) - y_d(t)\|_{L^2(\Omega)}^2 dt, \end{aligned}$$

which leads to

$$\|\hat{y}_T - \bar{y}\|_{L^2(Q)} \leq \frac{1}{\sqrt{\kappa}} \|\hat{y}_T - y_d\|_{L^2(T, \infty; L^2(\Omega))}. \quad (5.11)$$

To prove (5.10) we observe that \hat{y}_T satisfies the equation

$$\begin{cases} \frac{\partial \hat{y}_T}{\partial t} + A\hat{y}_T + f(x, t, \hat{y}_T) = g \text{ in } \Omega \times (T, \infty), \\ \partial_{n_A} \hat{y}_T = 0 \text{ on } \Gamma \times (T, \infty), \hat{y}_T(T) = y_T(T) \text{ in } \Omega. \end{cases}$$

Testing this equation with \hat{y}_T , and using that $f(x, t, \hat{y}_T)\hat{y}_T \geq 0$ due to the monotonicity of f with respect to y and (2.1), it follows that

$$\frac{1}{2} \|\hat{y}_T(t)\|_{L^2(\Omega)}^2 + \int_T^\infty \langle A\hat{y}_T, y_T \rangle dt \leq \frac{1}{2} \|y_T(T)\|_{L^2(\Omega)}^2 + \int_T^\infty \int_\Omega g \hat{y}_T dx dt.$$

From this inequality we infer

$$\|\hat{y}_T\|_{L^2(T, \infty; L^2(\Omega))} \leq C' \left(\|y_T(T)\|_{L^2(\Omega)} + \|g\|_{L^2(T, \infty; L^2(\Omega))} \right).$$

This inequality and (5.9) imply (5.8). \square

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