Monte Carlo methods

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Overview

- **1** Examples: Ising model etc.
- **2** Integration vs Sampling
- 3 Pseudo random number generators
- 4 Generation of random numbers with a given given distribution
- 5 Percolation, Random Walks
- **6** Improtance Sampling, Markov chains, Metropolis Alg.
- Statistics, error estimates: Jackknife, bootstrap

Random Walks

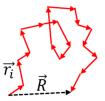
Conformation of a polymer chain in a solutions How far are the ends?

N monomers $ightarrow R \sim N^{
u}$

Simplest model: step in a random direction for each i Also models a drunkard's movement \rightarrow "Random walk"

 $R \sim t^{1/2}$ disagrees with experiment (of polimer chains)





1D random walk

We walk on a line, every step p = 0.5 for left, p = 0.5 for right.

Step at time *i*: Z_i , i.e. $Z_i \in \{-1, 1\}$, after *i* steps: $R_i = \sum_{i=1}^n Z_i$ each Z_i is independent $\implies \langle R_i \rangle = \sum \langle Z_i \rangle = 0$

What's the typical distance?

$$\langle R_n^2 \rangle = \sum_{i=1}^n Z_i^2 + 2 \sum_{i < j} \langle Z_i Z_j \rangle = n \tag{1}$$

If *n* is large R_n is the sum of many independent variables \rightarrow Gaussian distribution.

Generalizations, relations:

- random walk in many dimensions: each direction has 1/(2d) chance
- Gaussian random walk: stepsize is a Gaussian random variable. Used in finance
- Random walk without lattice: step in a random direction (sometimes also at random times)
- Wiener Process: limit of random walk with very small stepsizes (both space and time)

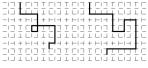
Restricted Random walks

monomers repulse each other \rightarrow Polymer chain cannot cross itself $\rightarrow R \sim N^{1/2}$ is wrong

No backtrack random walk The next step cannot immediately turn back. In the first step there are 2d possibilities, later only 2d - 1 Crossings are still allowed.

Self avoiding random walk (SAW) Visitng the same site twice is not allowed (monomers repulse) Defining the ensemble: What should be the weight of a random walk?

RW: all walks of length N are equally probable. **SAW**: all walks of length N are equally probable, as long as they are SAWs.





 \implies naive algorithm: Generate random walks. Throw result away if it intersects.

Works if N is small, otherwise most walks get thrown away: "attrition problem"

Biased sampling

Check at each step how many steps are available: l_i , choose one among them with equal probabilities. If $l_i = 0$: throw away and start from the beginning. Now each walk has a weight:

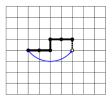
$$w = \prod_{i=1}^{N} \frac{l_i}{2d - 1}$$
(2)

use these weights when calculating averages.

Reptation algorithm

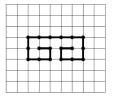
Generating a "chain" of walks: $\ldots \rightarrow W_i \rightarrow W_{i+1} \rightarrow \ldots$

- **1** Assume we have a walk of length N: W_i
- 2 delete a point from a random end
- **B** add one point to a random end with random direction. if it is intersecting: $W_{i+1} = W_i$ (change rejected) else W_{i+1} = new walk (change accepted)



Reptation

Can we get to every walk? No, there are trapped configurations:



Reptation can not change this or arrive here \implies process is non ergodic, sampling is biased (i.e. not uniform).

Such configurations are so rare that there absence in the sampling is not noticeable on the results.

Pivot algorithm

- I Pick a random point on the random walk, this divides walk to two sides
- 2 choose a random side
- Itransform the choosen side using lattice symmetries (rotations, mirroring)
- 4 accept the new walk if it is SAW, otherwise keep the old walk

Random walks conclusion

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Results form B. Li et. al. J. Stat. Phys. 80, 661 (1995)

$$\nu = 3/4 \quad \text{in } d = 2, \tag{3}$$

$$\nu = 0.59... \quad \text{in } d = 3,$$

$$\nu = 1/2 \quad \text{in } d = 4$$

Further generalizations of the polymer model: monomer interactions (itroducing a potential) interaction with solvent Thermal effects

Importance Sampling

Usual question (thermal average):

$$\langle F \rangle = \frac{1}{Z} \int F(p,q) e^{-\beta H(p,q)} d^N p d^N q =?$$
(4)

Can be also written as (assume F depends on energy E or do an average at constant energy otherwise)

$$\langle F \rangle = \int F(E) \rho(E) e^{-\beta E} dE$$
 (5)

For $N \gg 1$ the energy is sharply peaked. A typical uniform random configuration will miss the peak \implies random sampling very unefficient

If we could sample configurations with probability $p(C_i) \sim e^{-\beta H(p_i,q_i)}$ Averages calculated simply as

$$\langle F \rangle = \frac{1}{N} \sum_{i=1}^{N} F(p_i, q_i)$$
 (6)

This is called Importance sampling

Markov chains

Consider a sequence of "configurations" (e.g. phase space points of some system) governed by some random process

 $\ldots \rightarrow \Phi_i \rightarrow \Phi_{i+1} \rightarrow \ldots$

It's called a **Markov chain** if $T(\Phi'_{i+1}|\Phi_i)$, the probability to go from Φ does not depend on *i* and only depends on Φ_i , not on older Φ_j with j < i (no memory)

Notation: transition probabilities: $T(\Phi_k | \Phi_j) = T_{kj}$ At step *i* the system is in state *k* with probability $P_k^{(i)}$

conservation of probability \implies :

$$\sum_{j} T_{jk} = 1 \tag{7}$$

at step i + 1 we have

$$P_{k}^{(i+1)} = T_{kj} P_{j}^{(i)}$$
(8)

Change of probability

How does $P^{(i)}$ change? We had $P_k^{(i+1)} = T_{kj}P_j^{(i)}$

$$P_{j}^{(i+1)} - P_{j}^{(i)} = \sum_{k} T_{jk} P_{k}^{(n)} - P_{j}^{(n)} = \underbrace{\sum_{k} T_{jk} P_{k}^{(n)}}_{\text{inflow}} - \underbrace{\sum_{k} T_{kj} P_{j}^{(n)}}_{\text{outflow}}$$
(9)

Similar to the master equation

$$\frac{\partial P(\Phi,t)}{\partial t} = \sum_{\Phi'} W(\Phi \leftarrow \Phi') P(\Phi') - \sum_{\Phi'} W(\Phi' \leftarrow \Phi) P(\Phi)$$

We are looking for a stationary solution Π_k $\implies T_{ik}\Pi_k = \Pi_i$ eigenvector with eigenvalue $\lambda = 1$ **Balance equation**:

$$\sum_{k} T_{jk} \Pi_{k}^{(n)} = \sum_{k} T_{kj} \Pi_{j}^{(n)}$$
(10)

Detailed balance: $T_{jk}\Pi_k^{(n)} = T_{kj}\Pi_j^{(n)}$

Convergence to equilibrium

Suppose we choose T_{jk} such that the detailed balance is satisfied for some Π_k Consider the evolution of $P: P^{(1)} \rightarrow P^{(2)} \rightarrow \dots$ Does it converge to Π_k ? Let's define $\sigma_n = \sum_i |P_i^{(n)} - \Pi_i|$

$$\sigma_{n+1} = \sum_{i} \left| \sum_{j} T_{ij} P_j^{(n)} - \Pi_i \right| = \sum_{i} \left| \sum_{j} (T_{ij} P_j^{(n)} - T_{ji} \Pi_i) \right| =$$
(11)

using detailed balance

$$=\sum_{i}\left|\sum_{j}(T_{ij}P_{j}^{(n)}-T_{ij}\Pi_{j})\right|=\sum_{i}\left|\sum_{j}T_{ij}(P_{j}^{(n)}-\Pi_{j})\right|\leq\sum_{i,j}T_{ij}\left|(P_{j}^{(n)}-\Pi_{j})\right|=\sigma_{n}$$

Since $\sum_{j} (P_{j}^{(n)} - \Pi_{j}) = 0$, there is equality only if $\sigma_{n} = 0$, if we have $T_{ij} > 0$ (strong ergodicity)

 \implies convergence to unique equilibrium Π_k

Convergence rate is given by second largest eigenvalue of T_{jk}

Metropolis Algorithm

If we can set up a process that satisfies detailed balance for some Π_k , the process will converge to Π_k .

Metropolis algorithm achieves this.

Starting from configuration Φ_n

1 proposal with probability $T_0(\Phi'|\Phi_{n-1})$

2 Accepted with probability $T_A(\Phi'|\Phi_{n-1})$ if accepted: $\Phi_{n+1} = \Phi'$ else: $\Phi_{n+1} = \Phi_n$

To satisfy detailed balance, we must have

$$T_0(\Phi'|\Phi)T_A(\Phi'|\Phi)\Pi(\Phi) = T_0(\Phi|\Phi')T_A(\Phi|\Phi')\Pi(\Phi')$$
(12)

Metropolis

detailed balance:

$$\frac{T_A(\Phi|\Phi')}{T_A(\Phi'|\Phi)} = \frac{\Pi(\Phi)}{\Pi(\Phi')} \frac{T_0(\Phi'|\Phi)}{T_0(\Phi|\Phi')} = F(\Phi|\Phi')$$
(13)

Note that $F(\Phi|\Phi') = 1/F(\Phi'|\Phi)$ We choose e.g.

$$T_{A}(\Phi'|\Phi) = \begin{cases} F(\Phi'|\Phi) \text{ if } F(\Phi'|\Phi) < 1, \\ 1 \text{ otherwise} \end{cases}$$
(14)

or

$$T_A(\Phi'|\Phi) = \frac{F(\Phi'|\Phi)}{1 + F(\Phi'|\Phi)}$$
(15)

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Metropolis

Usually we choose the symmetric (aka. reversible) proposal: $T_0(\Phi'|\Phi) = T_0(\Phi|\Phi')$ Equilibrium is given by some action: $\Pi(\Phi) = e^{-S[\Phi]}$ In this case:

$$F(\Phi'|\Phi) = \frac{\Pi(\Phi')}{\Pi(\Phi)} = e^{-S(\Phi') + S(\Phi)} = e^{-\Delta S}$$
(16)

In this case the acceptance probability is:

$$T_{A}(\Phi'|\Phi) = \begin{cases} e^{-\Delta S} \text{ if } \Delta S > 0, \\ 1 \quad \text{if } \Delta S < 0 \end{cases}$$
(17)

How to choose proposal probability? change too big \implies acceptance rate will be small change too small \implies exploration of configuration space too slow Should be optimized

Running a Monte Carlo simulation

- 1 Initialize configuration
- 2 Choose parameters for proposal step (if applicable)
- **3** Generate Φ' for Φ_n
- 4 calculate $T_A(\Phi'|\Phi_n)$
- **5** generate random uniform $r \in [0, 1]$
- 6 accept if $r < T_A \rightarrow \Phi_{n+1} = \Phi'$ else $\Phi_{n+1} = \Phi_n$
- goto 3 if n_{max} not reached

The sequence of configurations can be seen as dependent on "MC time". This is not the same as the real dynamical evolution of the system. One measures time in MC hits/lattice site to make it volume independent

Example: Ising

$$H = -J \sum_{neighbors} s_i s_j - h \sum_i s_i$$
(18)

Update step: take a spin, and flip it.

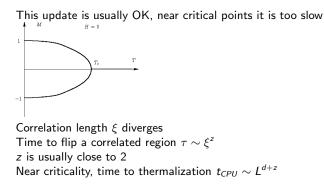
Either go through spins over the lattice sequentally, or always choose a random spin. Ergodic (any configuration can be reached) and reversible.

Since energy is local, only the neighboring spins are needed to calculate energy change ΔE , which depends on the number of antiparallel spins. Can calculate probabilities beforehand.

e.g. in d = 2 and h = 0 (there are more possibilities for $h \neq 0$)

# of antipar. spins	0	1	2	3	4
$\Delta E/J$	8	4	0	-4	-8
$\ln(T_A)$	$-8\beta J$	$-4\beta J$	1	1	1

Critical slowing down



Efficient update strategies

z is dependent on the update algorithm. If we flip a correlated cluster of spins at once, we eliminate z **Cluster algorithm** Another fast algorithm: **Worm algorithm**

Statistics

Suppose we have a MC simulation, results: x_i , i = 1...NIn equilibrium, they all have the same average and variance:

$$\langle x_i \rangle = \langle x \rangle, \quad \sigma_{x_i}^2 = \left\langle \left(x_i - \langle x_i \rangle \right)^2 \right\rangle = \sigma_x^2$$
 (19)

Unbiased estimators:

$$\hat{x} = \frac{1}{N} \sum x_i, \qquad \hat{\sigma}_x^2 = \frac{1}{N-1} \sum (x_i - \hat{x})^2$$
 (20)

 $\langle \hat{x} \rangle = \langle x \rangle$, but how close is it? Look at its variance:

$$\sigma_{\hat{x}}^{2} = \left\langle (\hat{x} - \langle x \rangle)^{2} \right\rangle = \left\langle \left(\frac{1}{N} \sum (x_{i} - \langle x \rangle) \right)^{2} \right\rangle =$$
(21)
$$= \frac{1}{N^{2}} \left\langle \sum_{i,j} (x_{i} - \langle x \rangle) (x_{j} - \langle x \rangle) \right\rangle = \frac{1}{N} \langle x^{2} \rangle + \frac{1}{N^{2}} \sum_{i \neq j} \langle x_{i} x_{j} \rangle + \langle x \rangle^{2}$$

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Statistics

For uncorrelated configurations, we have $\langle x_i x_j \rangle = \langle x_i \rangle \langle x_j \rangle$

$$\sigma_{\hat{x}}^2 = \frac{1}{N} \hat{\sigma}_x^2 \tag{22}$$

In this case we quote the results of the simulation as

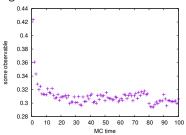
$$\hat{x} \pm \sigma$$
 (23)

with $\sigma = \sigma_{\hat{x}} = \hat{\sigma}_x / \sqrt{N}$

MC results

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In a Monte Carlo simulation, measuring observables as a function of n one gets e.g.:



exponential thermalization to thermal averages Initial thermalization should be discarded from averages Thermalization rate can be observable dependent! Averages are easy:

$$\langle O \rangle = \frac{1}{n_{max} - n_{therm}} \sum_{i > n_{therm}} O_i$$
 (24)

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How do we calculate errors?

Correlations

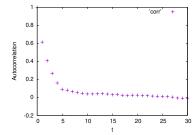
Since we generate Φ_{n+1} from Φ_n , they are correlated ($\langle A \rangle = \mu_A$).

$$C(A,B) = \frac{\langle (A - \mu_A)(B - \mu_B) \rangle}{\sigma_A \sigma_B} = \frac{\langle AB \rangle - \langle A \rangle \langle B \rangle}{\sigma_A \sigma_B}$$
(25)

Suppose we have some observable A_n at MC time n. We define **Autocorrelation**:

$$C_A(t) = C(A_i, A_{i+t}) = \frac{\langle A_i A_{i+t} \rangle - \langle A \rangle^2}{\langle A^2 \rangle - \langle A \rangle^2}$$
(26)

(In equilibrium correlation depends only on t)



 $C_A(t) \sim exp(-t/ au_A)$ au_A : autorrelation time

We need $\tau = \sup \tau_A$

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Statistics

In the correlated case, we have to be more careful estimating the variance. Using $D_A(t) = \langle A_i A_{i+t} \rangle - \langle A \rangle^2$

$$\begin{aligned} \sigma_{\hat{x}}^{2} &= \frac{1}{N^{2}} \left\langle \sum_{i,j} (x_{i} - \langle x \rangle) (x_{j} - \langle x \rangle) \right\rangle = \frac{1}{N^{2}} \sum_{i,j} D_{A}(|i - j|) = \\ &= \frac{1}{N^{2}} \sum_{t = -(N-1)}^{N-1} \sum_{k=1}^{N-|t|} D_{A}(|t|) = \sigma_{A}^{2} \sum_{t = -(N-1)}^{N-1} \frac{N - |t|}{N^{2}} C_{A}(t) \\ &= \frac{\sigma_{A}^{2}}{N} \sum_{t = -(N-1)}^{N-1} C_{A}(t) \left(1 - \frac{|t|}{N}\right) \approx \frac{\sigma_{A}^{2}}{N} 2 \left(\frac{1}{2} + \sum_{t=1}^{N} C_{A}(t)\right) = \frac{\sigma_{A}^{2}}{N} 2 \tau_{A,int} \end{aligned}$$

we introduced integrated autocorrelation time:

$$\tau_{A,int} = \frac{1}{2} + \sum_{t=1}^{N} C_A(t) \approx \int_0^\infty e^{-t/\tau} dt = \tau$$
(27)

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Autocorrelation

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This can be understood as haveing

$$N_{indep} = \frac{N}{2\tau_{A,int}}$$
(28)

independent configurations. Quote results as :

$$\hat{x} \pm \sqrt{\frac{1}{N}} 2\tau_{A,int} \sigma_A^2 \tag{29}$$

Getting a good estimate for the autocorrelation requires a large dataset

Blocking

Divide dateset into consequent dataset of size N_B . If uncorrelated, their variance should decrease as $\sim 1/N_B$. \implies repeat blocking for different sizes, and consider the results statistically independent as $\sim 1/N_B$ dependence sets in.

Problem: Some observables require "global quantities".

example 1: do a fit on the data as a function of some parameter. What is the statistical error of the fit parameters? one data point might give very noisy data such that no senseful fit can be performed

example 2: What is the error of the variance estimator?

How to calculate errors for such qunatities? Using standard error propagation might be cumbersome or too difficult

(standard error propagation for 1 variable: $\sigma_{f(A)} = \frac{df}{dA}\sigma_A$)

Jackknife

Consider an initial sample set of N uncorrelated datapoints x_i . The original observable from the dataset:

$$\hat{\Theta} = f(\{x_i\}) \tag{30}$$

Now construct N smaller datasets by removing the *n*-th datapoint from the set. For each of this smaller datasets, we determine our observable Θ_n . Estimator of the variance:

$$\sigma_{\hat{\Theta}}^2 = \frac{N-1}{N} \sum_{n=1}^{N} \left(\Theta_n - \hat{\Theta}\right)^2 \tag{31}$$

We can also remove a bias from the estimator (or at least the 1/N part), by calculating

$$\tilde{\Theta} = \frac{1}{N} \sum_{k=1}^{N} \Theta_k \tag{32}$$

and finally

$$\langle \Theta \rangle = \hat{\Theta} - (N - 1)(\tilde{\Theta} - \hat{\Theta}) \pm \sigma_{\hat{\Theta}}$$
(33)

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Usually one needs blocking to get rid of autocorrelation. divide the set x_i into N_{jk} separate "streams" of data Consider the blocks as the jackknife "points" e.g. x_i , i = 1...1000, and take $N_{jk} = 10$, than we have

$$\hat{\Theta} = f(\{x_{i=1...1000}\}),$$

$$\Theta_1 = f(\{x_{i=101...1000}\}),$$

$$\Theta_2 = f(\{x_{i=1...100,201...1000}\}),$$
and so on

$$(34)$$

and finally use the formulas on prev. slide for N = 10

Bootstrap

This is a resampling method. Consider an initial sample set of N uncorrelated datapoints x_i

The original observable from the dataset:

$$\hat{\Theta} = f(\{x_i\}) \tag{35}$$

now create K new datasets S_k by drawing N random elements from the set $\{x_i\}$, with repetitions allowed. for each set we calculate $\Theta_k = f(S_k)$. We than calculate

$$\tilde{\Theta} = \frac{1}{K} \sum_{k=1}^{K} \Theta_k, \qquad \sigma_{\tilde{\Theta}}^2 = \frac{1}{K} \sum_{k=1}^{K} \left(\Theta_k - \tilde{\Theta} \right)^2$$
(36)

We can use this as estimators of the average and the variance (these are biased, i.e. one needs a large K (check $\hat{\Theta} - \tilde{\Theta}$)):

$$\langle \Theta \rangle = \tilde{\Theta} \pm \sigma_{\tilde{\Theta}} \tag{37}$$