### Gauge invariance and Observables in Particle Physics

#### **Axel Maas**

19<sup>th</sup> of December 2023 Zurich Switzerland





- Why an invariant formulation?
  - Path integral formulation and symmetries

- Why an invariant formulation?
  - Path integral formulation and symmetries
- Brout-Englert-Higgs Physics

- Why an invariant formulation?
  - Path integral formulation and symmetries
- Brout-Englert-Higgs Physics
- Fröhlich-Morchio-Strocchi mechanism

- Why an invariant formulation?
  - Path integral formulation and symmetries
- Brout-Englert-Higgs Physics
- Fröhlich-Morchio-Strocchi mechanism
  - Standard Model
    - Experimental signatures

- Why an invariant formulation?
  - Path integral formulation and symmetries
- Brout-Englert-Higgs Physics
- Fröhlich-Morchio-Strocchi mechanism
  - Standard Model
    - Experimental signatures
  - Beyond the Standard Model
    - Qualitative changes

- Why an invariant formulation?
  - Path integral formulation and symmetries
- Brout-Englert-Higgs Physics
- Fröhlich-Morchio-Strocchi mechanism
  - Standard Model
    - Experimental signatures
  - Beyond the Standard Model
    - Qualitative changes
  - Quantum (super)gravity

### What's the deal? -Gauge symmetry

 $Z = \int_{\Omega} D \phi e^{iS[\phi]}$ 

Integral over all space-time histories of the universe  $Z = \int_{\Omega} D \phi e^{iS[\phi]}$ 









 $\langle \phi(x) \dots \phi(z) \rangle = \int_{\Omega} D \phi \phi(x) \dots \phi(z) e^{iS[\phi]}$ 

Expectation values are weighted averages over space-time histories



Expectation values are weighted averages over space-time histories

[Review: Maas'17]

 $Z = \int_{\Omega} D \phi^a e^{iS[\phi]}$ 

[Review: Maas'17]



 $Z = \int_{\Omega}$ 

[Review: Maas'17]



- no anomalies

Action is invariant  $S[\phi] = S[G\phi]$ 

iS[  $\phi$ 

 $D\phi$ 

[Review: Maas'17]

### $Z = \int_{\Omega} D \phi^a e^{iS[\phi]}$

Integration range - contains all orbits  $G \phi$ 

[Review: Maas'17]

### $\langle \phi^b(x) \rangle = \int_{\Omega} D \phi^a e^{iS[\phi]} \phi^b(x)$

[Review: Maas'17]

$$\langle \phi^b(x) \rangle = \int_{\Omega} D \phi^a e^{iS[\phi]} \phi^b(x)$$

- There is no preferred point on the group orbit
  - There is no absolute orientation/frame in the internal space
  - Does not change when averaging over position
  - There is no absolute charge

[Review: Maas'17]

$$\langle \phi^b(x) \rangle = \int_{\Omega} D \phi^a e^{iS[\phi]} \phi^b(x) = 0$$

- There is no preferred point on the group orbit
  - There is no absolute orientation/frame in the internal space
  - Does not change when averaging over position
  - There is no absolute charge

[Review: Maas'17]

$$\langle \phi^b(x)\phi^c(y)\rangle = \int_{\Omega} D\phi^a e^{iS[\phi]}\phi^b(x)\phi^c(y)$$

• Relative charge measurement averaged over all possible starting points

[Review: Maas'17]

$$\langle \phi^b(x)\phi^c(y)\rangle = \int_{\Omega} D\phi^a e^{iS[\phi]}\phi^b(x)\phi^c(y) = 0$$

- Relative charge measurement averaged over all possible starting point
  - Vanishes because no preferred absolute starting point

[Review: Maas'17]

# $\left\langle \delta_{bc} \phi^{b}(x) \phi^{c}(y) \right\rangle$ $= \int_{\Omega} D \phi^{a} e^{iS[\phi]} \delta_{bc} \phi^{b}(x) \phi^{c}(y)$

- Group-invariant quantity
  - Measures relative orientation
  - Created from an invariant tensor  $\,\delta_{ab}\,$

[Review: Maas'17]

# $\left\langle \delta_{bc} \phi^{b}(x) \phi^{c}(y) \right\rangle$ $= \int_{\Omega} D \phi^{a} e^{iS[\phi]} \delta_{bc} \phi^{b}(x) \phi^{c}(y) \neq 0$

- Group-invariant quantity
  - Measures relative orientation
  - Created from an invariant tensor  $\,\delta_{ab}\,$

[Review: Maas'17]

### $Z = \int_{\Omega} D \phi^a e^{iS[\phi]}$

[Review: Maas'17]

Field – transforms locally under a group  $\phi^a(x) \rightarrow G^{ab}(x) \phi^b(x)$ 

 $Z = \int_{\Omega} D(\phi^{a}) e^{iS[\phi]}$ 

 $Z = \int_{\Omega}$ 

[Review: Maas'17]



- no anomalies

Action is invariant  $S[\phi] = S[G\phi]$ 

iS[*φ* 

 $D\phi$ 

[Review: Maas'17]

# $Z = \int_{\Omega} D \phi^a e^{iS[\phi]}$

Integration range - contains all orbits  $G \phi$ 

[Review: Maas'17]

## $\left\langle \delta_{bc} \phi^{b}(x) \phi^{c}(y) \right\rangle$ $= \int_{\Omega} D \phi^{a} e^{iS[\phi]} \delta_{bc} \phi^{b}(x) \phi^{c}(y)$

[Review: Maas'17]

# $\left\langle \delta_{bc} \phi^{b}(x) \phi^{c}(y) \right\rangle$ $= \int_{\Omega} D \phi^{a} e^{iS[\phi]} \delta_{bc} \phi^{b}(x) \phi^{c}(y) = 0$

- No longer invariant under gauge transformations
  - Vanishes just as any other non-invariant quantity

[Review: Maas'17]



•Transporter compensates gauge transformations

Implemented by gauge fields

[Review: Maas'17]

### $\langle \phi^{b}(x) U^{bc}(x, y) \phi^{c}(y) \rangle$ = $\int_{\Omega} D \phi^{a} D U e^{iS[\phi]} \phi^{b}(x) U^{bc}(x, y) \phi^{c}(y) \neq 0$

•Transporter compensates gauge transformations

Implemented by gauge fields

[Review: Maas'17]

Reduced integration range

# $= \int_{\Omega_c} \Phi^a DU W(U, \phi) e^{iS[\phi]} \phi^b(x) \phi^c(y) \neq 0$

- Reduction of integration region by gauge fixing
  - Arbitrary choice of coordinates
  - Weight factor to keep gauge-invariant quantities the same
# Brout-Englert-Higgs Physics -The Standard Model

#### A toy model

• Consider an SU(2) with a fundamental scalar

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$

- Ws  $W^a_{\mu}$  W
- Coupling g and some numbers  $f^{abc}$

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + gf^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Ws  $W^a_{\mu}$  W
- Higgs  $h_i$  (h)
- Coupling g and some numbers  $f^{abc}$  and  $t_a^{ij}$

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j}) + D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Ws  $W^a_{\mu}$  W
- Higgs  $h_i$  (h)
- Couplings g, v,  $\lambda$  and some numbers  $f^{abc}$  and  $t_a^{ij}$

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Ws  $W^a_{\mu}$  W
- Higgs  $h_i$  (h)
- Couplings  $g, v, \lambda$  and some numbers  $f^{abc}$  and  $t_a^{ij}$
- Parameters selected for a BEH effect

### A toy model: Symmetries

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

## A toy model: Symmetries

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

• Local SU(2) gauge symmetry  $W^{a}_{\mu} \rightarrow W^{a}_{\mu} + (\delta^{a}_{b}\partial_{\mu} - gf^{a}_{bc}W^{c}_{\mu})\phi^{b}$   $h_{i} \rightarrow h_{i} + gt^{ij}_{a}\phi^{a}h_{j}$ 

# A toy model: Symmetries

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Local SU(2) gauge symmetry  $W^a_{\mu} \rightarrow W^a_{\mu} + (\delta^a_b \partial_{\mu} - g f^a_{bc} W^c_{\mu}) \Phi^b$   $h_i \rightarrow h_i + g t^{ij}_a \Phi^a h_j$
- Global SU(2) custodial (flavor) symmetry
  - Acts as (right-)transformation on the scalar field only  $W^a_{\mu} \rightarrow W^a_{\mu}$   $h \rightarrow h \Omega$

 Choose parameters to get a Brout-Englert-Higgs effect

- Choose parameters to get a Brout-Englert-Higgs effect
- Minimize the classical action

- Choose parameters to get a Brout-Englert-Higgs effect
- Minimize the classical action
- Choose a suitable gauge and obtain 'spontaenous gauge symmetry breaking':  $SU(2) \rightarrow 1$

- Choose parameters to get a Brout-Englert-Higgs effect
- Minimize the classical action
- Choose a suitable gauge and obtain 'spontaenous gauge symmetry breaking': SU(2) → 1
- Get masses and degeneracies at treelevel

- Choose parameters to get a Brout-Englert-Higgs effect
- Minimize the classical action
- Choose a suitable gauge and obtain 'spontaenous gauge symmetry breaking': SU(2) → 1
- Get masses and degeneracies at treelevel
- Perform perturbation theory

#### **Physical spectrum**

Perturbation theory



 $\bigcirc$ 

## **Physical spectrum**

Perturbation theory Scalar fixed charge

• Custodial singlet

Mass

## **Physical spectrum**



Both custodial singlets

 $\bigcirc$ 

[Fröhlich et al.'80, Banks et al.'79]

• Elementary fields are gauge-dependent

[Fröhlich et al.'80, Banks et al.'79]

- Elementary fields are gauge-dependent
  - Change under a gauge transformation

- Elementary fields are gauge-dependent
  - Change under a gauge transformation
  - Gauge transformations are a human choice

- Elementary fields are gauge-dependent
  - Change under a gauge transformation
  - Gauge transformations are a human choice...
  - ...and gauge-symmetry breaking is not there [Elitzur'75, Osterwalder & Seiler'77, Fradkin & Shenker'78]

- Elementary fields are gauge-dependent
  - Change under a gauge transformation
  - Gauge transformations are a human choice...
  - ...and gauge-symmetry breaking is not there [Elitzur'75, Osterwalder & Seiler'77, Fradkin & Shenker'78]
    - Just a figure of speech
    - Actually just ordinary gauge-fixing

- Elementary fields are gauge-dependent
  - Change under a gauge transformation
  - Gauge transformations are a human choice...
  - ...and gauge-symmetry breaking is not there [Elitzur'75, Osterwalder & Seiler'77, Fradkin & Shenker'78]
    - Just a figure of speech
    - Actually just ordinary gauge-fixing
- Physics has to be expressed in terms of manifestly gauge-invariant quantities

- Elementary fields are gauge-dependent
  - Change under a gauge transformation
  - Gauge transformations are a human choice...
  - ...and gauge-symmetry breaking is not there [Elitzur'75, Osterwalder & Seiler'77, Fradkin & Shenker'78]
    - Just a figure of speech
    - Actually just ordinary gauge-fixing
- Physics has to be expressed in terms of manifestly gauge-invariant quantities
  - And this includes non-perturbative aspects

- Elementary fields are gauge-dependent
  - Change under a gauge transformation
  - Gauge transformations are a human choice...
  - ...and gauge-symmetry breaking is not there [Elitzur'75, Osterwalder & Seiler'77, Fradkin & Shenker'78]
    - Just a figure of speech
    - Actually just ordinary gauge-fixing
- Physics has to be expressed in terms of manifestly gauge-invariant quantities
  - And this includes non-perturbative aspects...
  - ...even at weak coupling [Gribov'78,Singer'78,Fujikawa'82]

[Fröhlich et al.'80, Banks et al.'79]

• Need physical, gauge-invariant particles

- Need physical, gauge-invariant particles
  - Cannot be the elementary particles
  - Non-Abelian nature is relevant

- Need physical, gauge-invariant particles
  - Cannot be the elementary particles
  - Non-Abelian nature is relevant
- Need more than one particle: Composite particles

- Need physical, gauge-invariant particles
  - Cannot be the elementary particles
  - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
  - Higgs-Higgs



- Need physical, gauge-invariant particles
  - Cannot be the elementary particles
  - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
  - Higgs-Higgs, W-W



- Need physical, gauge-invariant particles
  - Cannot be the elementary particles
  - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
  - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Need physical, gauge-invariant particles
  - Cannot be the elementary particles
  - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
  - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



Has nothing to do with weak coupling

- Need physical, gauge-invariant particles
  - Cannot be the elementary particles
  - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
  - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Has nothing to do with weak coupling
  - Think QED (hydrogen atom!)

- Need physical, gauge-invariant particles
  - Cannot be the elementary particles
  - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
  - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Has nothing to do with weak coupling
  - Think QED (hydrogen atom!)
- Can this matter?
[Fröhlich et al.'80,'81, Maas & Törek'16,'18, Maas, Sondenheimer & Törek'17]

•  $J^{\text{PC}}$  and custodial charge only quantum numbers

- J<sup>PC</sup> and custodial charge only quantum numbers
  - Different from perturbation theory
    - Operators limited to asymptotic, elementary, gauge-dependent states

- J<sup>PC</sup> and custodial charge only quantum numbers
  - Different from perturbation theory
    - Operators limited to asymptotic, elementary, gauge-dependent states
- Formulate gauge-invariant, composite operators

- J<sup>PC</sup> and custodial charge only quantum numbers
  - Different from perturbation theory
    - Operators limited to asymptotic, elementary, gauge-dependent states
- Formulate gauge-invariant, composite operators
  - Bound state structure

- J<sup>PC</sup> and custodial charge only quantum numbers
  - Different from perturbation theory
    - Operators limited to asymptotic, elementary, gauge-dependent states
- Formulate gauge-invariant, composite operators
  - Bound state structure non-perturbative methods!

- J<sup>PC</sup> and custodial charge only quantum numbers
  - Different from perturbation theory
    - Operators limited to asymptotic, elementary, gauge-dependent states
- Formulate gauge-invariant, composite operators
  - Bound state structure non-perturbative methods! - Lattice



Both custodial singlets

Experiment tells that somehow the left is correct



Experiment tells that somehow the left is correct Theory say the right is correct



Experiment tells that somehow the left is correct Theory say the right is correct There must exist a relation that both are correct



Gauge-invariant

Scalar singlet

Both custodial singlets

$$h(x) + h(x)$$





Both custodial singlets

 $\square$ 



**Custodial singlet** 



#### • Both custodial singlets Custodial singlet



Both custodial singlets

Custodial singlet

$$tr t^{a} \frac{h^{+}}{\sqrt{h^{+} h}} D_{\mu} \frac{h}{\sqrt{h^{+} h}}$$





Both custodial singlets

#### Custodial singlet

Triplet

$$tr \frac{h^{+}}{\sqrt{h^{+} h}} D_{\mu} \frac{h}{\sqrt{h^{+} h}}$$





#### • Both custodial singlets Custodial singlet Triplet

# A microscopic origin -Fröhlich-Morchio-Strocchi mechanism

[Fröhlich et al.'80,'81, Maas & Törek'16,'18, Maas, Sondenheimer & Törek'17 Maas & Sondenheimer '20]

- J<sup>PC</sup> and custodial charge only quantum numbers
  - Different from perturbation theory
    - Operators limited to asymptotic, elementary, gauge-dependent states
- Formulate gauge-invariant, composite operators
  - Bound state structure non-perturbative methods?

[Fröhlich et al.'80,'81, Maas & Törek'16,'18, Maas, Sondenheimer & Törek'17 Maas & Sondenheimer '20]

- J<sup>PC</sup> and custodial charge only quantum numbers
  - Different from perturbation theory
    - Operators limited to asymptotic, elementary, gauge-dependent states
- Formulate gauge-invariant, composite operators
  - Bound state structure non-perturbative methods?
  - But coupling is still weak and there is a BEH

[Fröhlich et al.'80,'81, Maas & Törek'16,'18, Maas, Sondenheimer & Törek'17 Maas & Sondenheimer '20]

- $J^{\text{PC}}$  and custodial charge only quantum numbers
  - Different from perturbation theory
    - Operators limited to asymptotic, elementary, gauge-dependent states
- Formulate gauge-invariant, composite operators
  - Bound state structure non-perturbative methods?
  - But coupling is still weak and there is a BEH
  - Perform double expansion [Fröhlich et al.'80, Maas'12]
    - Vacuum expectation value (FMS mechanism)
    - Standard expansion in couplings
    - Together: Augmented perturbation theory

[Fröhlich et al.'80,'81 Maas'12,'17]

1) Formulate gauge-invariant operator

[Fröhlich et al.'80,'81 Maas'12,'17]

- 1) Formulate gauge-invariant operator
  - **0**<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$

Higgs field

[Fröhlich et al.'80,'81 Maas'12,'17]

1) Formulate gauge-invariant operator

0<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$ 

[Fröhlich et al.'80,'81 Maas'12,'17]

1) Formulate gauge-invariant operator

0<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$ 

2) Expand Higgs field around fluctuations  $h=v+\eta$ 

[Fröhlich et al.'80,'81 Maas'12,'17]

1) Formulate gauge-invariant operator

0<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$ 

2) Expand Higgs field around fluctuations  $h=v+\eta$ 

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle \\ + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

[Fröhlich et al.'80,'81 Maas'12,'17]

1) Formulate gauge-invariant operator

0<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$ 

2) Expand Higgs field around fluctuations  $h=v+\eta$ 

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle \\ + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

3) Standard perturbation theory

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle + \langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$$

[Fröhlich et al.'80,'81 Maas'12,'17]

1) Formulate gauge-invariant operator

0<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$ 

2) Expand Higgs field around fluctuations  $h=v+\eta$ 

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle \\ + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

3) Standard perturbation theory

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle + \langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$$

[Fröhlich et al.'80,'81 Maas'12,'17]

1) Formulate gauge-invariant operator

0<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$ 

2) Expand Higgs field around fluctuations  $h=v+\eta$ 

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle \\ + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

3) Standard perturbation theory

Bound state  $\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle$ mass  $+ \langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$ 

```
[Fröhlich et al.'80,'81
Maas'12,'17]
```

1) Formulate gauge-invariant operator

0<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$ 

2) Expand Higgs field around fluctuations  $h=v+\eta$ 

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle \\ + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

3) Standard perturbation theory

Bound state  $\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle$ mass  $\langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$ 

Trivial two-particle state

```
[Fröhlich et al.'80,'81
Maas'12,'17]
```

1) Formulate gauge-invariant operator

0<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$ 

2) Expand Higgs field around fluctuations  $h=v+\eta$ 

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle \\ + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

3) Standard perturbation theory

Bound state  $\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle$ mass  $+ \langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$  Higgs mass

[Fröhlich et al.'80,'81 Maas'12,'17]

1) Formulate gauge-invariant operator

0<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$ 

2) Expand Higgs field around fluctuations  $h=v+\eta$ 

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle$$
  
+  $v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$ 

3) Standard perturbation theory

Standard Perturbation Theory

Bound state  $\langle (h^+ h)(x)(h^+ h)(y) = v^2 \eta^+ (x)\eta(y) \rangle$ mass  $+ \langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$ 

[Fröhlich et al.'80,'81 Maas'12,'17 Maas & Sondenheimer'20]

1) Formulate gauge-invariant operator

0<sup>+</sup> singlet:  $\langle (h^+ h)(x)(h^+ h)(y) \rangle$ 

2) Expand Higgs field around fluctuations  $h=v+\eta$ 



$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle + \langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$$

[Maas'12,'17 Maas & Sondenheimer'20]

#### $\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle$ $+ v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$

[Maas'12,'17 Maas & Sondenheimer'20]



 $\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle$  $+ v \langle \eta^{+} \eta^{2} + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^{2} \rangle$ 

[Maas'12,'17 Maas & Sondenheimer'20]



 $\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle$ + $v\langle \eta^{+} \eta^{2} + \eta^{+2} \eta \rangle$  +  $\langle \eta^{+2} \eta^{2} \rangle$ 

[Maas'12,'17 Maas & Sondenheimer'20]



$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle \\ + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$


$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle$$
  
+  $v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$ 









[Maas'12,'17 Maas & Sondenheimer'20, Maas et al. unpublished]



[Fröhlich et al.'80,'81 Maas'12]

1) Formulate gauge-invariant operator 1<sup>-</sup> triplet:  $\langle (\tau^i h^+ D_\mu h)(x)(\tau^j h^+ D_\mu h)(y) \rangle$ 

- 1) Formulate gauge-invariant operator
  - **1**<sup>-</sup> triplet:  $\langle (\tau^i h^+ D_\mu h)(x)(\tau^j h^+ D_\mu h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$

- 1) Formulate gauge-invariant operator 1<sup>-</sup> triplet:  $\langle (\tau^i h^+ D_\mu h)(x)(\tau^j h^+ D_\mu h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$ 
  - $\langle (\tau^i h^+ D_{\mu} h)(x)(\tau^j h^+ D_{\mu} h)(y) \rangle = v^2 c_{ij}^{ab} \langle W^a_{\mu}(x) W^b(y)^{\mu} \rangle + \dots$

- 1) Formulate gauge-invariant operator 1<sup>-</sup> triplet:  $\langle (\tau^i h^+ D_\mu h)(x)(\tau^j h^+ D_\mu h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$ 
  - $\langle (\tau^i h^+ D_{\mu} h)(x)(\tau^j h^+ D_{\mu} h)(y) \rangle = v^2 c_{ij}^{ab} \langle W^a_{\mu}(x) W^b(y)^{\mu} \rangle + \dots$

Matrix from group structure

- 1) Formulate gauge-invariant operator 1<sup>-</sup> triplet:  $\langle (\tau^i h^+ D_\mu h)(x)(\tau^j h^+ D_\mu h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$

$$\langle (\tau^i h^+ D_{\mu} h)(x)(\tau^j h^+ D_{\mu} h)(y) \rangle = v^2 c_{ij}^{ab} \langle W^a_{\mu}(x) W^b(y)^{\mu} \rangle + \dots$$
$$= v^2 \langle W^i_{\mu} W^j_{\mu} \rangle + \dots$$

Matrix from group structure

- **1)** Formulate gauge-invariant operator  $1^{-}$  triplet:  $\langle (\tau^{i}h^{+}D_{\mu}h)(x)(\tau^{j}h^{+}D_{\mu}h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$ 
  - $\langle (\tau^{i}h^{+}D_{\mu}h)(x)(\tau^{j}h^{+}D_{\mu}h)(y)\rangle = v^{2}c_{ij}^{ab}\langle W_{\mu}^{a}(x)W^{b}(y)^{\mu}\rangle + \dots$  $= v^{2}\langle W_{\mu}^{i}W_{\mu}^{j}\rangle + \dots$

Matrix from group structure

*c* projects custodial states to gauge states

- 1) Formulate gauge-invariant operator 1<sup>-</sup> triplet:  $\langle (\tau^i h^+ D_{\mu} h)(x)(\tau^j h^+ D_{\mu} h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$ 
  - $\langle (\tau^{i}h^{+}D_{\mu}h)(x)(\tau^{j}h^{+}D_{\mu}h)(y)\rangle = v^{2}c_{ij}^{ab}\langle W_{\mu}^{a}(x)W^{b}(y)^{\mu}\rangle + \dots$  $= v^{2}\langle W_{\mu}^{i}W_{\mu}^{j}\rangle + \dots$ Matrix from group structure

*c* projects custodial states to gauge states

Exactly one gauge boson for every physical state

# Phenomenological Implications

Can we measure this?

 Two possibilities to measure extension

- Two possibilities to measure extension
  - Form factor
    - Difficult
      - Higgs and Z need to be both produced in the same process



- Two possibilities to measure extension
  - Form factor
    - Difficult
      - Higgs and Z need to be both produced in the same process
  - Elastic scattering
    - Standard vector boson scattering process at low energies
    - Use this one



- Elastic region:  $160/180 \, GeV \leq \sqrt{s} \leq 250 \, GeV$ 
  - s is the CMS energy in the initial/final ZZ/WW system

- Elastic region:  $160/180 \, GeV \leq \sqrt{s} \leq 250 \, GeV$ 
  - s is the CMS energy in the initial/final ZZ/WW system
- Requires a partial wave analysis

- Elastic region:  $160/180 \, GeV \leq \sqrt{s} \leq 250 \, GeV$ 
  - s is the CMS energy in the initial/final ZZ/WW system
- Requires a partial wave analysis

$$\frac{d\sigma}{d\Omega} = \frac{1}{64 \, \pi^2 s} |M|^2$$

- Elastic region:  $160/180 \, GeV \leq \sqrt{s} \leq 250 \, GeV$ 
  - s is the CMS energy in the initial/final ZZ/WW system
- Requires a partial wave analysis



Matrix element

- Elastic region:  $160/180 \, GeV \leq \sqrt{s} \leq 250 \, GeV$ 
  - s is the CMS energy in the initial/final ZZ/WW system
- Requires a partial wave analysis

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 s} |M|^2$$
  
M(s,  $\Omega$ ) = 16 $\pi \sum_{J} (2J+1) f_J(s) P_J(\cos\theta)$ 

- Elastic region:  $160/180 \, GeV \leq \sqrt{s} \leq 250 \, GeV$ 
  - s is the CMS energy in the initial/final ZZ/WW system
- Requires a partial wave analysis

Matrix element  

$$\frac{d \sigma}{d \Omega} = \frac{1}{64 \pi^2 s} |M|^2 \quad \text{Partial wave amplitude}$$

$$M(s, \Omega) = 16 \pi \sum_{J} (2J+1) f_{J}(s) P_{J}(\cos \theta)$$
Legendre polynom

- Elastic region:  $160/180 \, GeV \leq \sqrt{s} \leq 250 \, GeV$ 
  - s is the CMS energy in the initial/final ZZ/WW system
- Requires a partial wave analysis

$$\frac{d \sigma}{d \Omega} = \frac{1}{64 \pi^2 s} |M|^2$$
$$M(s, \Omega) = 16 \pi \sum_J (2J+1) f_J(s) P_J(\cos \theta)$$
Partial wave  $f_J(s) = e^{i \delta_J(s)} \sin(\delta_J(s))$ 

- Elastic region:  $160/180 \, GeV \leq \sqrt{s} \leq 250 \, GeV$ 
  - s is the CMS energy in the initial/final ZZ/WW system
- Requires a partial wave analysis

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 s} |M|^2$$

$$M(s,\Omega) = 16\pi \sum_J (2J+1)f_J(s)P_J(\cos\theta)$$
Partial wave  $f_J(s) = e^{i\delta_J(s)}\sin(\delta_J(s))$ 
amplitude

Phase shift

- Elastic region:  $160/180 \, GeV \leq \sqrt{s} \leq 250 \, GeV$ 
  - s is the CMS energy in the initial/final ZZ/WW system
- Requires a partial wave analysis

$$\frac{d \sigma}{d \Omega} = \frac{1}{64 \pi^2 s} |M|^2$$

$$M(s, \Omega) = 16 \pi \sum_J (2J+1) f_J(s) P_J(\cos \theta)$$

$$f_J(s) = e^{i \delta_J(s)} \sin(\delta_J(s))$$

$$s \to 4m_W^2$$

$$a_0 = \tan(\delta_J) / \sqrt{s - 4m_W^2}$$
Phase shift

- Elastic region:  $160/180 \, GeV \leq \sqrt{s} \leq 250 \, GeV$ 
  - s is the CMS energy in the initial/final ZZ/WW system
- Requires a partial wave analysis

$$\frac{d \sigma}{d \Omega} = \frac{1}{64 \pi^2 s} |M|^2$$

$$M(s, \Omega) = 16 \pi \sum_J (2J+1) f_J(s) P_J(\cos \theta)$$

$$f_J(s) = e^{i \delta_J(s)} \sin(\delta_J(s))$$

$$s \to 4m_W^2 \tan(\delta_J) / \sqrt{s - 4m_W^2}$$
Scattering length~"size" Phase shift

• Consider the Higgs: *J*=0



• Consider the Higgs: J=0



- Consider the Higgs: *J*=0
- Mock-up effect
  - Scattering length 1/(40 GeV)



- Consider the Higgs: *J*=0
- Mock-up effect
  - Scattering length 1/(40 GeV)

## Impact on the radius of the Higgs

- Reduced SM: Only W/Z and the Higgs
  - Parameters slightly different
    - Higgs too heavy (145 GeV) and too strong weak coupling
  - Qualitatively but not quantitatively

## Impact on the radius of the Higgs



- Reduced SM: Only W/Z and the Higgs
  - Parameters slightly different
    - Higgs too heavy (145 GeV) and too strong weak coupling
  - Qualitatively but not quantitatively

# Impact on the radius of the Higgs



- Reduced SM: Only W/Z and the Higgs
  - Parameters slightly different
    - Higgs too heavy (145 GeV) and too strong weak coupling
  - Qualitatively but not quantitatively


- Reduced SM: Only W/Z and the Higgs
  - Parameters slightly different
    - Higgs too heavy (145 GeV) and too strong weak coupling
  - Qualitatively but not quantitatively



- Reduced SM: Only W/Z and the Higgs
  - Parameters slightly different
    - Higgs too heavy (145 GeV) and too strong weak coupling
  - Qualitatively but not quantitatively



- Reduced SM: Only W/Z and the Higgs
  - Parameters slightly different
    - Higgs too heavy (145 GeV) and too strong weak coupling
  - Qualitatively but not quantitatively



- Reduced SM: Only W/Z and the Higgs
  - Parameters slightly different
    - Higgs too heavy (145 GeV) and too strong weak coupling
  - Qualitatively but not quantitatively

[Jenny, Maas, Riederer'22, Maas'23]



- Reduced SM: Only W/Z and the Higgs
  - Parameters slightly different
    - Higgs too heavy (145 GeV) and too strong weak coupling
  - Qualitatively but not quantitatively

[Jenny, Maas, Riederer'22, Maas'23]



- Reduced SM: Only W/Z and the Higgs
  - Parameters slightly different
    - Higgs too heavy (145 GeV) and too strong weak coupling
  - Qualitatively but not quantitatively

[Jenny, Maas, Riederer'22, Maas'23]













## **Physical states**

- Need physical, gauge-invariant particles
  - Cannot be the elementary particles
  - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
  - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Has nothing to do with weak coupling
  - Think QED (hydrogen atom!)
- Can this matter?

## **Physical states**

- Need physical, gauge-invariant particles
  - Cannot be the elementary particles
  - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
  - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Has nothing to do with weak coupling
  - Think QED (hydrogen atom!)
- Can this matter?

[Fröhlich et al.'80, Egger, Maas, Sondenheimer'17]

- Flavor has two components
  - Global SU(3) generation
  - Local SU(2) weak gauge (up/down distinction)

- Flavor has two components
  - Global SU(3) generation
  - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable

- Flavor has two components
  - Global SU(3) generation
  - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state

$$\begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} _i (\mathbf{x})$$

- Flavor has two components
  - Global SU(3) generation
  - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state



• Gauge-invariant state

- Flavor has two components
  - Global SU(3) generation
  - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state

$$\begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} (x)$$

• Gauge-invariant state, but custodial doublet

- Flavor has two components
  - Global SU(3) generation
  - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state

$$\left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} | \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_i (x)^+ \left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} | \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_j (y) \right|$$

• Gauge-invariant state, but custodial doublet

- Flavor has two components
  - Global SU(3) generation
  - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state FMS applicable

$$\left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_i (x)^+ \left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_j (y) \right| \stackrel{h=\mathbf{v}+\mathbf{n}}{\approx} \mathbf{v}^2 \left| \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_i (x)^+ \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} (y) + O(\mathbf{\eta})$$

• Gauge-invariant state, but custodial doublet

- Flavor has two components
  - Global SU(3) generation
  - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state FMS applicable

$$\left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} | \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_i (x)^* \left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} | \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_j (y) \right|^{h=\mathbf{v}+\eta} \approx \mathbf{v}^2 \left| \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} |_i (x)^* \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} |_j (y) \right|^{+O(\eta)}$$

- Gauge-invariant state, but custodial doublet
- Yukawa terms break custodial symmetry
  - Different masses for doublet members

- Flavor has two components
  - Global SU(3) generation
  - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state FMS applicable

$$\left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} | \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_i (x)^* \left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} | \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_j (y) \right|^{h=\mathbf{v}+\eta} \approx \mathbf{v}^2 \left| \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} |_i (x)^* \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} |_j (y) \right|^{+O(\eta)}$$

- Gauge-invariant state, but custodial doublet
- Yukawa terms break custodial symmetry
  - Different masses for doublet members
- Extends non-trivially to hadrons

- Only mock-up standard model
  - Compressed mass scales
  - One generation
  - Degenerate leptons and neutrinos
  - Dirac fermions: left/righthanded non-degenerate
  - Quenched

- Only mock-up standard model
  - Compressed mass scales
  - One generation
  - Degenerate leptons and neutrinos
  - Dirac fermions: left/righthanded non-degenerate
  - Quenched
- Same qualitative outcome
  - FMS construction
  - Mass defect
  - Flavor and custodial symmetry patterns

- Only mock-up standard model
  - Compressed mass scales
  - One generation
  - Degenerate leptons and neutrinos
  - Dirac fermions: left/righthanded non-degenerate
  - Quenched
- Same qualitative outcome
  - FMS construction
  - Mass defect
  - Flavor and custodial symmetry patterns



- Only mock-up standard model
  - Compressed mass scales
  - One generation
  - Degenerate leptons and neutrinos
  - Dirac fermions: left/righthanded non-degenerate
  - Quenched
- Same qualitative outcome
  - FMS construction
  - Mass defect
  - Flavor and custodial symmetry patterns



- Only mock-up standard model
  - Compressed mass scales
  - One generation
  - Degenerate leptons and neutrinos
  - Dirac fermions: left/righthanded non-degenerate
  - Quenched
- Same qualitative outcome
  - FMS construction
  - Mass defect
  - Flavor and custodial symmetry patterns



- Only mock-up standard model
  - Compressed mass scales
  - One generation
  - Degenerate leptons and neutrinos
  - Dirac fermions: left/righthanded non-degenerate
  - Quenched
- Same qualitative outcome
  - FMS construction
  - Mass defect
  - Flavor and custodial symmetry patterns
- Supports FMS prediction



# New physics -Qualitative changes

[Maas'15 Maas, Sondenheimer, Törek'17]

Standard model is special

- Standard model is special
  - Mapping of custodial symmetry to gauge symmetry
  - Fits perfectly degrees of freedom

- Standard model is special
  - Mapping of custodial symmetry to gauge symmetry
  - Fits perfectly degrees of freedom
- Is this generally true?

- Standard model is special
  - Mapping of custodial symmetry to gauge symmetry
  - Fits perfectly degrees of freedom
- Is this generally true?
  - No: Depends on gauge group, representations, and custodial groups

- Standard model is special
  - Mapping of custodial symmetry to gauge symmetry
  - Fits perfectly degrees of freedom
- Is this generally true?
  - No: Depends on gauge group, representations, and custodial groups
  - Can work sometimes (2HDM,MSSM) [Maas,Pedro'16, Maas,Schreiner'23]

- Standard model is special
  - Mapping of custodial symmetry to gauge symmetry
  - Fits perfectly degrees of freedom
- Is this generally true?
  - No: Depends on gauge group, representations, and custodial groups
  - Can work sometimes (2HDM,MSSM) [Maas,Pedro'16, Maas,Schreiner'23]
  - Generally qualitative differences
• Consider an SU(3) with a single fundamental scalar

- Consider an SU(3) with a single fundamental scalar
- Looks very similar to the standard model Higgs

$$L = -\frac{1}{4} W^a_{\mu\nu} W^{\mu\nu}_a$$
$$W^a_{\mu\nu} = \partial_\mu W^a_\nu - \partial_\nu W^a_\mu + g f^a_{bc} W^b_\mu W^c_\nu$$

- Ws  $W^a_{\mu}$  W
- Coupling g and some numbers  $f^{abc}$

- Consider an SU(3) with a single fundamental scalar
- Looks very similar to the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Ws  $W^a_{\mu}$  W
- Higgs  $h_i$  (h)
- Coupling g and some numbers  $f^{abc}$  and  $t_a^{ij}$

- Consider an SU(3) with a single fundamental scalar
- Looks very similar to the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j}) + D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Ws  $W^a_{\mu}$  W
- Higgs  $h_i$  (h)
- Couplings g, v,  $\lambda$  and some numbers  $f^{abc}$  and  $t_a^{ij}$

- Consider an SU(3) with a single fundamental scalar
- Looks very similar to the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j}) + D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Ws  $W^a_{\mu}$  W
- Higgs  $h_i$  (h)
- Couplings  $g, v, \lambda$  and some numbers  $f^{abc}$  and  $t_a^{ij}$
- Parameters selected for a BEH effect

- Consider an SU(3) with a single fundamental scalar
- Looks very similar to the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Consider an SU(3) with a single fundamental scalar
- Looks very similar to the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h^{+}_{a} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

• Local SU(3) gauge symmetry  $W^{a}_{\mu} \rightarrow W^{a}_{\mu} + (\delta^{a}_{b}\partial_{\mu} - gf^{a}_{bc}W^{c}_{\mu})\phi^{b}$   $h_{i} \rightarrow h_{i} + gt^{ij}_{a}\phi^{a}h_{j}$ 

- Consider an SU(3) with a single fundamental scalar
- Looks very similar to the standard model Higgs

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Local SU(3) gauge symmetry  $W^a_\mu \rightarrow W^a_\mu + (\delta^a_b \partial_\mu - g f^a_{bc} W^c_\mu) \Phi^b$   $h_i \rightarrow h_i + g t^{ij}_a \Phi^a h_j$
- Global U(1) custodial (flavor) symmetry
  - Acts as (right-)transformation on the scalar field only  $W^a_{\mu} \rightarrow W^a_{\mu}$   $h \rightarrow \exp(ia)h$

Gauge-dependent Vector



 $SU(3) \rightarrow SU(2)'$ 





Confirmed in gauge-fixed lattice calculations [Maas et al.'16]

[Maas & Törek'16,'18 Maas, Sondenheimer & Törek'17]





- 1) Formulate gauge-invariant operator
  - 1<sup>-</sup> singlet

1) Formulate gauge-invariant operator 1<sup>-</sup> singlet:  $\langle (h^+ D_\mu h)(x)(h^+ D_\mu h)(y) \rangle$ 

- 1) Formulate gauge-invariant operator
  - **1**<sup>-</sup> singlet:  $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$

- 1) Formulate gauge-invariant operator 1<sup>-</sup> singlet:  $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$ 
  - $\langle (h + D_{\mu}h)(x)(h + D_{\mu}h)(y)\rangle = v^2 c^{ab} \langle W^a_{\mu}(x)W^b(y)^{\mu}\rangle + \dots$

- 1) Formulate gauge-invariant operator
  - **1**<sup>-</sup> singlet:  $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$ 
  - $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y)\rangle = v^2 c^{ab} \langle W^a_{\mu}(x)W^b(y)^{\mu}\rangle + \dots$

Matrix from group structure

- 1) Formulate gauge-invariant operator
  - **1**<sup>-</sup> singlet:  $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$ 
  - $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y)\rangle = v^2 c^{ab} \langle W^a_{\mu}(x)W^b(y)^{\mu}\rangle + \dots$  $= v^2 \langle W^8_{\mu}W^8_{\mu}\rangle + \dots$

Matrix from group structure

*c*<sup>*ab*</sup> projects out only one field

- 1) Formulate gauge-invariant operator
  - **1**<sup>-</sup> singlet:  $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y) \rangle$
- 2) Expand Higgs field around fluctuations  $h=v+\eta$ 
  - $\langle (h^{+} D_{\mu}h)(x)(h^{+} D_{\mu}h)(y)\rangle = v^{2}c^{ab}\langle W_{\mu}^{a}(x)W^{b}(y)^{\mu}\rangle + \dots$  $= v^{2}\langle W_{\mu}^{8}W_{\mu}^{8}\rangle + \dots$ Matrix from group structure

*c*<sup>*ab*</sup> projects out only one field

Only one state remains in the spectrum at mass of gauge boson 8 (heavy singlet)

## A different kind of states

- Group theory forced same gauge multiplets and custodial multiples for SU(2)
  - Because Higgs is bifundamental
  - Remainder is bound state/resonance or not

# A different kind of states

- Group theory forced same gauge multiplets and custodial multiples for SU(2)
  - Because Higgs is bifundamental
  - Remainder is bound state/resonance or not
- Now: Elementary states without analouge
  - No global symmetry to provide multiplet structure

# A different kind of states

- Group theory forced same gauge multiplets and custodial multiples for SU(2)
  - Because Higgs is bifundamental
  - Remainder is bound state/resonance or not
- Now: Elementary states without analouge
  - No global symmetry to provide multiplet structure
- Now: States without elementary analouge
  - Gauge-invariant states from 3 Higgs fields
  - Baryon analogue U(1) acts as baryon number
  - Lightest must exist and be absolutely stable

## **Possible new states**

• Quantum numbers are J<sup>PC</sup><sub>Custodial</sub>

## **Possible new states**

- Quantum numbers are J<sup>PC</sup><sub>Custodial</sub>
  - Simpelst non-trivial state operator:  $0^{++}_{1}$  $\epsilon_{abc} \phi^a D_\mu \phi^b D_\nu D^\nu D^\mu \phi^c$

## **Possible new states**

- Quantum numbers are J<sup>PC</sup><sub>Custodial</sub>
  - Simpelst non-trivial state operator:  $0^{++}_{1}$  $\epsilon_{abc} \phi^a D_\mu \phi^b D_\nu D^\nu D^\mu \phi^c$
- What is the lightest state?
  - Prediction with constituent model







- Qualitatively different spectrum
- No mass gap!

## **Possible states**

- Quantum numbers are J<sup>PC</sup><sub>Custodial</sub>
  - Simpelst non-trivial state operator: 0<sup>++</sup>
  - $\epsilon_{abc} \phi^a D_\mu \phi^b D_\nu D^\nu D^\mu \phi^c$
- What is the lightest state?
  - Prediction with constituent model
  - Lattice calulations

## **Possible states**

- Quantum numbers are J<sup>PC</sup><sub>Custodial</sub>
  - Simpelst non-trivial state operator: 0<sup>++</sup>
  - $\epsilon_{abc} \phi^a D_\mu \phi^b D_\nu D^\nu D^\mu \phi^c$
- What is the lightest state?
  - Prediction with constituent model
  - Lattice calulations
  - All channels: J<3
  - Aim: Ground state for each channel
    - Characterization through scattering states
















































### Experimental consequences [Maas & 7 Maas'17]

[Maas & Törek'18

Add fundamental fermions

[Maas & Törek'18 Maas'17]

- Add fundamental fermions
- Bhabha scattering

[Maas & Törek'18 Maas'17]



- Add fundamental fermions
- Bhabha scattering

[Maas & Törek'18 Maas'17]



- Add fundamental fermions
- Bhabha scattering

## Experimental consequences [Maas & Törek'18 Maas'17]



- Add fundamental fermions
- Bhabha scattering

## Experimental consequences [Maas & Törek'18 Maas'17]



- Add fundamental fermions
- Bhabha scattering

Toy models representative for mechanisms

- Toy models representative for mechanisms
- All results so far inconsistent with perturbation theory

- Toy models representative for mechanisms
- All results so far inconsistent with perturbation theory
- Consistent with FMS construction

- Toy models representative for mechanisms
- All results so far inconsistent with perturbation theory
- Consistent with FMS construction
  - Different spectrum: Different phenomenology

- Toy models representative for mechanisms
- All results so far inconsistent with perturbation theory
- Consistent with FMS construction
  - Different spectrum: Different phenomenology
- Application of FMS to GUT candidates

- Toy models representative for mechanisms
- All results so far inconsistent with perturbation theory
- Consistent with FMS construction
  - Different spectrum: Different phenomenology
- Application of FMS to GUT candidates: <u>All checked failed</u> [Maas et al.'17,Sondenheimer'19]
  - SU(5), SO(10), Pati-Salam,...

- Toy models representative for mechanisms
- All results so far inconsistent with perturbation theory
- Consistent with FMS construction
  - Different spectrum: Different phenomenology
- Application of FMS to GUT candidates: <u>All checked failed</u> [Maas et al.'17,Sondenheimer'19]
  - SU(5), SO(10), Pati-Salam,...
    - None found so far that works

- Toy models representative for mechanisms
- All results so far inconsistent with perturbation theory
- Consistent with FMS construction
  - Different spectrum: Different phenomenology
- Application of FMS to GUT candidates: <u>All checked failed</u> [Maas et al.'17,Sondenheimer'19]
  - SU(5), SO(10), Pati-Salam,...
    - None found so far that works
  - Depends on Higgs sector

- Toy models representative for mechanisms
- All results so far inconsistent with perturbation theory
- Consistent with FMS construction
  - Different spectrum: Different phenomenology
- Application of FMS to GUT candidates: <u>All checked failed</u> [Maas et al.'17,Sondenheimer'19]
  - SU(5), SO(10), Pati-Salam,...
    - None found so far that works
  - Depends on Higgs sector
    - May still be possible

- Toy models representative for mechanisms
- All results so far inconsistent with perturbation theory
- Consistent with FMS construction
  - Different spectrum: Different phenomenology
- Application of FMS to GUT candidates: <u>All checked failed</u> [Maas et al.'17,Sondenheimer'19]
  - SU(5), SO(10), Pati-Salam,...
    - None found so far that works
  - Depends on Higgs sector
    - May still be possible (hopefully?)

[Maas'19, Maas, Markl, Müller'22]

Quantum gravity is a gauge theory

- Quantum gravity is a gauge theory
- Empirically dominated by a field configuration
  - De Sitter/FLRW metric

- Quantum gravity is a gauge theory
- Empirically dominated by a field configuration
  - De Sitter/FLRW metric
- Observables need to be fully invariant
  - Diffeomorphism and local Lorentz

- Quantum gravity is a gauge theory
- Empirically dominated by a field configuration
  - De Sitter/FLRW metric
- Observables need to be fully invariant
  - Diffeomorphism and local Lorentz
- FMS mechanism applicable
  - A 'BEH effect' for gravity
  - Technically much more involved
  - First predictions agree with dynamical triangulation results [Dai et al.'21]

- Gravity and supersymmetry imply supergravity
  - Supersymmetry becomes a local gauge symmetry

- Gravity and supersymmetry imply supergravity
  - Supersymmetry becomes a local gauge symmetry
- Same reasoning: Observables need to be gauge invariant

- Gravity and supersymmetry imply supergravity
  - Supersymmetry becomes a local gauge symmetry
- Same reasoning: Observables need to be gauge invariant
  - Observables cannot show supersymmetry

- Gravity and supersymmetry imply supergravity
  - Supersymmetry becomes a local gauge symmetry
- Same reasoning: Observables need to be gauge invariant
  - Observables cannot show supersymmetry
  - Could explain absence of supersymmetry in experiment

- Gravity and supersymmetry imply supergravity
  - Supersymmetry becomes a local gauge symmetry
- Same reasoning: Observables need to be gauge invariant
  - Observables cannot show supersymmetry
  - Could explain absence of supersymmetry in experiment
- FMS mechanism as applicable as to quantum gravity

### Summary

• Full invariance necessary for physical observables in path integrals

Review: 1712.04721 Update: 2305.01960
## Summary

 Full invariance necessary for physical observables in path integrals

 FMS mechanism allows estimates of quantum effects in a systematic expansion

Review: 1712.04721 Update: 2305.01960

## Summary

 Full invariance necessary for physical observables in path integrals

 FMS mechanism allows estimates of quantum effects in a systematic expansion

 Gives a new perspective on particle physics and quantum gravity

> Philosophy of physics perspective: 2110.00616 Review: 1712.04721 Update: 2305.01960

## **Come to Graz!**

- Upcoming jobs advertisments:
- Postdoc in this talk's topic in January '24
- Full professorship in particle physics in January'24
- Upcoming workshops:
- Parton Shower and Resummation in July '24
- Philosophical Reflections on Gauge Symmetries in July'24

