

Brout-Englert-Higgs physics: From foundations to phenomenology

Axel Maas

with

Larissa Egger, Leonardo Pedro, and Pascal Törek

15th of November 2016

Vienna

Austria



NAWI Graz

Natural Sciences

FWF

Der Wissenschaftsfonds

What is this talk about?

- Why it is not obvious that the Higgs and W/Z are physical particles?

What is this talk about?

- Why it is not obvious that the Higgs and W/Z are physical particles?
- Does it matter in the standard model?

What is this talk about?

- Why it is not obvious that the Higgs and W/Z are physical particles?
- Does it matter in the standard model?
 - No. And why.

What is this talk about?

- Why it is not obvious that the Higgs and W/Z are physical particles?
- Does it matter in the standard model?
 - No. And why.
- Why it can matter beyond the standard model

What is this talk about?

- Why it is not obvious that the Higgs and W/Z are physical particles?
- Does it matter in the standard model?
 - No. And why.
- Why it can matter beyond the standard model
- How this can be treated
 - Introducing gauge-invariant perturbation theory
 - Checking its validity

What is non-perturbative?

- Strong interactions are non-perturbative
 - Like QCD
 - But not always: Asymptotic freedom

What is non-perturbative?

- Strong interactions are non-perturbative
 - Like QCD
 - But not always: Asymptotic freedom
- Weak interactions can be non-perturbative
 - QED is weakly interacting, but has non-perturbative features like atoms, molecules, matter with phase structure,...
 - Bound states, phase transitions,...

What is non-perturbative?

- Strong interactions are non-perturbative
 - Like QCD
 - But not always: Asymptotic freedom
- Weak interactions can be non-perturbative
 - QED is weakly interacting, but has non-perturbative features like atoms, molecules, matter with phase structure,...
 - Bound states, phase transitions,...
- Are there (relevant) non-perturbative effects in the weak interactions and the Higgs?

Why it is not obvious that the Higgs and W/Z are physical particles

Or: What states can be gauge-invariant

The Problem

- Consider the Higgs sector of the standard model

The Problem

- Consider the Higgs sector of the standard model
- The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu}$$

$$W_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + gf_{bc}^a W_\mu^b W_\nu^c$$

- W_s W_μ^a 

- Coupling g and some numbers f^{abc}



The Problem

- Consider the Higgs sector of the standard model
- The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^\dagger D_{ik}^\mu h_k$$

$$W_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + gf_{bc}^a W_\mu^b W_\nu^c$$

$$D_\mu^{ij} = \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij}$$

- **Ws** W_μ^a 
- **Higgs** h_i 

- Coupling g and some numbers f^{abc} and t_a^{ij}



The Problem

- Consider the Higgs sector of the standard model
- The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^\dagger D_{ik}^\mu h_k + \lambda (h^a h_a^\dagger - v^2)^2$$

$$W_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + gf_{bc}^a W_\mu^b W_\nu^c$$

$$D_\mu^{ij} = \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij}$$

- **Ws** W_μ^a 
- **Higgs** h_i 
- No QED: Ws and Zs are degenerate
- Couplings g, v, λ and some numbers f^{abc} and t_a^{ij}

The Problem

- Consider the Higgs sector of the standard model
- The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^\dagger D_{ik}^\mu h_k + \lambda (h^a h_a^\dagger - v^2)^2$$

$$W_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + gf_{bc}^a W_\mu^b W_\nu^c$$

$$D_\mu^{ij} = \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij}$$

The Problem

- Consider the Higgs sector of the standard model
- The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^\dagger D_{ik}^\mu h_k + \lambda (h^a h_a^\dagger - v^2)^2$$

$$W_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + g f_{bc}^a W_\mu^b W_\nu^c$$

$$D_\mu^{ij} = \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij}$$

- Local SU(2) gauge symmetry

$$W_\mu^a \rightarrow W_\mu^a + (\delta_b^a \partial_\mu - g f_{bc}^a W_\mu^c) \phi^b$$

$$h_i \rightarrow h_i + g t_a^{ij} \phi^a h_j$$

The Problem

- Consider the Higgs sector of the standard model
- The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^\dagger D_{ik}^\mu h_k + \lambda (h^a h_a^\dagger - v^2)^2$$

$$W_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + g f_{bc}^a W_\mu^b W_\nu^c$$

$$D_\mu^{ij} = \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij}$$

- Local SU(2) gauge symmetry

$$W_\mu^a \rightarrow W_\mu^a + (\delta_b^a \partial_\mu - g f_{bc}^a W_\mu^c) \phi^b \qquad h_i \rightarrow h_i + g t_a^{ij} \phi^a h_j$$

- Global SU(2) Higgs custodial (flavor) symmetry

- Acts as right-transformation on the Higgs field only

$$W_\mu^a \rightarrow W_\mu^a \qquad h_i \rightarrow h_i + a^{ij} h_j + b^{ij} h_j^*$$

Standard approach

[Bohm et al. 2001]

- Minimize action classically
 - Yields $hh^+ = v^2$ - Higgs vev
 - Assume quantum corrections to this are small

Standard approach

[Bohm et al. 2001]

- Minimize action classically
 - Yields $hh^+ = v^2$ - Higgs vev
 - Assume quantum corrections to this are small
- Perform global gauge transformation such that

$$h(x) = v + \eta(x) = \begin{pmatrix} \phi^1(x) + i\phi^2(x) \\ v + \omega(x) + i\phi^3(x) \end{pmatrix} \Rightarrow \langle h \rangle = \begin{pmatrix} 0 \\ v \end{pmatrix}$$

- η mass depends at tree-level on v

Standard approach

[Bohm et al. 2001]

- Minimize action classically
 - Yields $hh^+ = v^2$ - Higgs vev
 - Assume quantum corrections to this are small
- Perform global gauge transformation such that

$$h(x) = v + \eta(x) = \begin{pmatrix} \phi^1(x) + i\phi^2(x) \\ v + \omega(x) + i\phi^3(x) \end{pmatrix} \Rightarrow \langle h \rangle = \begin{pmatrix} 0 \\ v \end{pmatrix}$$

- η mass depends at tree-level on v
- Perform perturbation theory

Implications of global transformation

- Not all charge directions equal
 - This is not physical, but merely a choice of gauge

Implications of global transformation

- Not all charge directions equal
 - This is not physical, but merely a choice of gauge
 - “Spontaneous gauge symmetry breaking”
 - Broken by the transformation, not by the dynamics
 - Dynamics only affect the length of the Higgs field
 - Local symmetry intact and cannot be broken

[Elitzur'75]

Implications of global transformation

- Not all charge directions equal
 - This is not physical, but merely a choice of gauge
 - “Spontaneous gauge symmetry breaking”
 - Broken by the transformation, not by the dynamics
 - Dynamics only affect the length of the Higgs field
 - Local symmetry intact and cannot be broken
- There are gauges where the vev always vanishes
 - Perturbation theory not sensible

[Elitzur'75]

[Maas'13]

[Lee et al.'72]

Implications of global transformation

- Not all charge directions equal
 - This is not physical, but merely a choice of gauge
 - “Spontaneous gauge symmetry breaking”
 - Broken by the transformation, not by the dynamics
 - Dynamics only affect the length of the Higgs field
 - Local symmetry intact and cannot be broken
- There are gauges where the vev always vanishes
 - Perturbation theory not sensible
- Consequence: Symmetry in charge space not manifest (hidden)
 - Symmetry expressed in STIs/WTIs

[Elitzur'75]

[Maas'13]

[Lee et al.'72]

Physical states

[Fröhlich et al.'80,
't Hooft'80,
Bank et al.'79]

- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections

Physical states

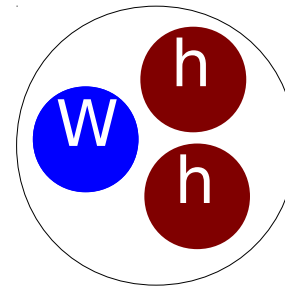
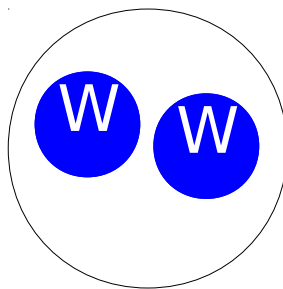
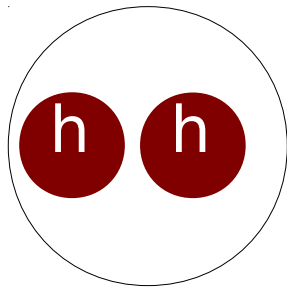
[Fröhlich et al.'80,
't Hooft'80,
Bank et al.'79]

- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable

Physical states

[Fröhlich et al.'80,
't Hooft'80,
Bank et al.'79]

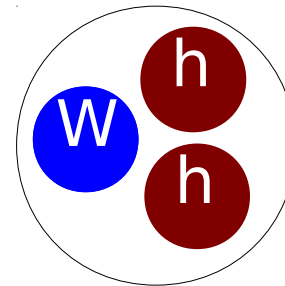
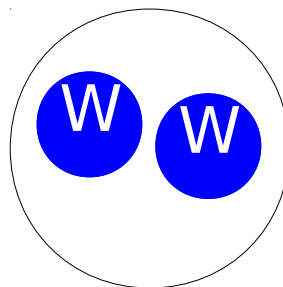
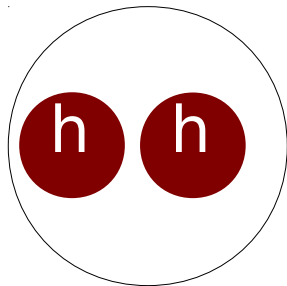
- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable
- Gauge-invariant states are composite
 - Not asymptotic states in perturbation theory
 - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



Physical states

[Fröhlich et al.'80,
't Hooft'80,
Bank et al.'79]

- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable
- Gauge-invariant states are composite
 - Not asymptotic states in perturbation theory
 - Higgs-Higgs, W-W, Higgs-Higgs-W etc.

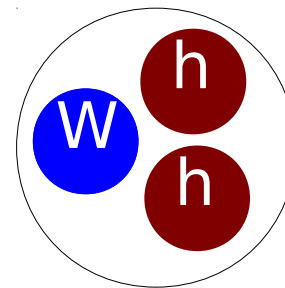
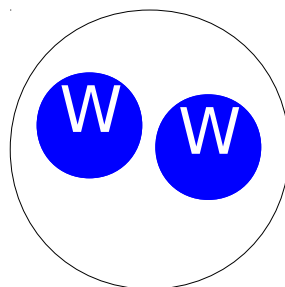
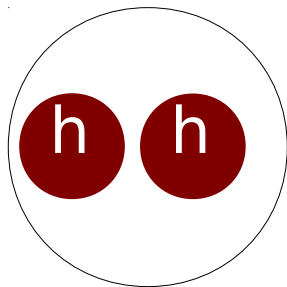


- Why does perturbation theory work?

Physical states

[Fröhlich et al.'80,
't Hooft'80,
Bank et al.'79]

- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable
- Gauge-invariant states are composite
 - Not asymptotic states in perturbation theory
 - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Why does perturbation theory work?
- Mass spectrum?

Why it does not matter in the standard model

Introducing gauge-invariant perturbation theory

Masses from propagators

- Masses are determined by poles of propagators

Masses from propagators

- Masses are determined by poles of propagators
- 2 propagators
 - W/Z $D_{\mu\nu}^{ab}(x-y) = \langle W_{\mu}^a(x) W_{\nu}^b(y) \rangle$
 - Degenerate without QED
 - Scalar $D_H^{ij}(x-y) = \langle \eta^i(x) \eta^{j+}(y) \rangle$

Masses from propagators

- Masses are determined by poles of propagators
- 2 propagators
 - W/Z $D_{\mu\nu}^{ab}(x-y) = \langle W_{\mu}^a(x) W_{\nu}^b(y) \rangle$
 - Degenerate without QED
 - Scalar $D_H^{ij}(x-y) = \langle \eta^i(x) \eta^{j+}(y) \rangle$
- (Tree-level/perturbative) poles of Higgs and W

Masses from propagators

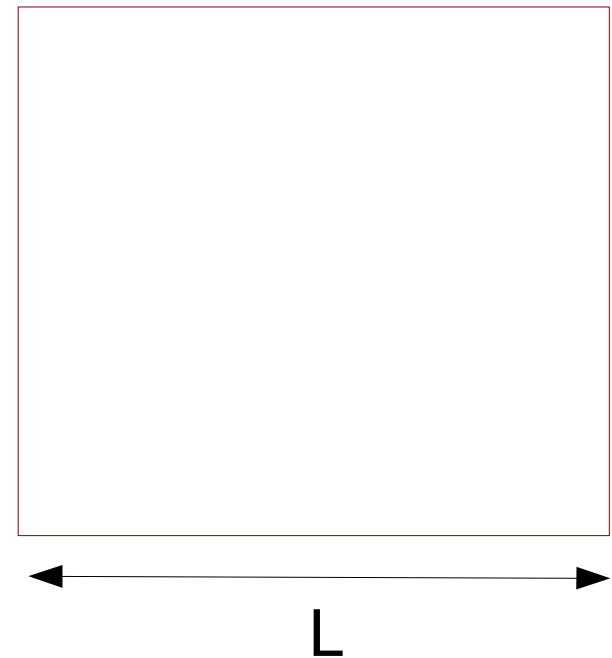
- Masses are determined by poles of propagators
- 2 propagators
 - W/Z $D_{\mu\nu}^{ab}(x-y) = \langle W_\mu^a(x) W_\nu^b(y) \rangle$
 - Degenerate without QED
 - Scalar $D_H^{ij}(x-y) = \langle \eta^i(x) \eta^{j+}(y) \rangle$
- (Tree-level/perturbative) poles of Higgs and W
 - But only in a fixed gauge
 - Elementary fields are gauge-dependent
 - Without gauge fixing propagators are $\sim \delta(x-y)$

Masses from propagators

- Masses are determined by poles of propagators
- 2 propagators
 - W/Z $D_{\mu\nu}^{ab}(x-y) = \langle W_{\mu}^a(x) W_{\nu}^b(y) \rangle$
 - Degenerate without QED
 - Scalar $D_H^{ij}(x-y) = \langle \eta^i(x) \eta^{j+}(y) \rangle$
- (Tree-level/perturbative) poles of Higgs and W
 - But only in a fixed gauge
 - Elementary fields are gauge-dependent
 - Without gauge fixing propagators are $\sim \delta(x-y)$
- Gauge-invariant: Non-perturbative method

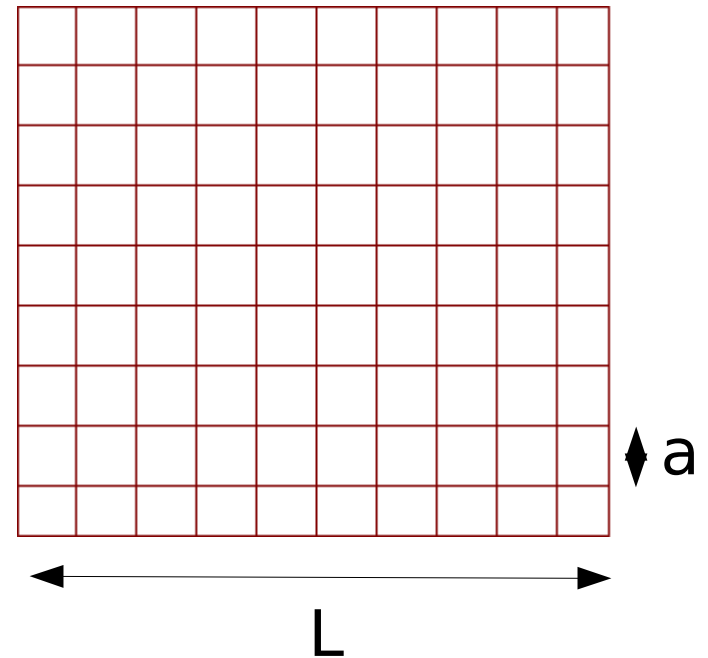
Lattice calculations

- Take a **finite volume** - usually a hypercube



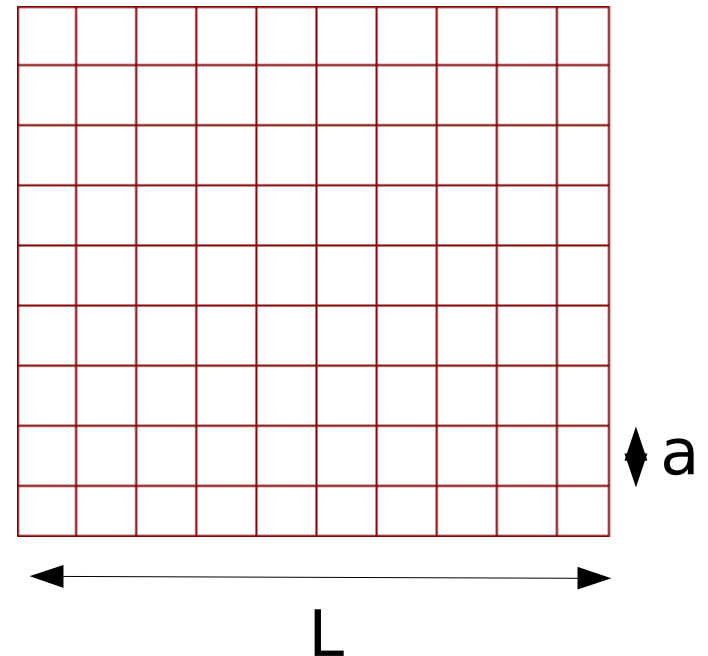
Lattice calculations

- Take a **finite volume** - usually a hypercube
- **Discretize it**, and get a finite, hypercubic lattice



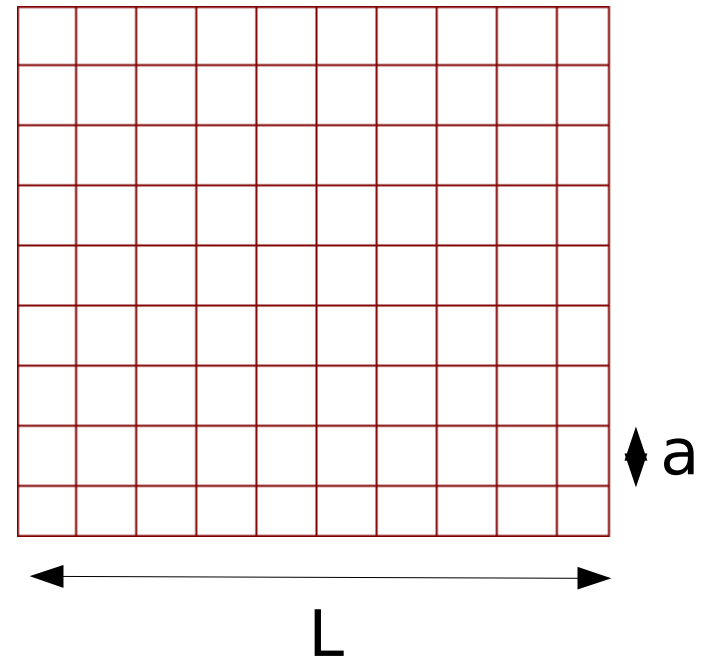
Lattice calculations

- Take a **finite volume** - usually a hypercube
- **Discretize it**, and get a finite, hypercubic lattice
- Calculate observables using path integral
 - Can be done numerically
 - Uses Monte-Carlo methods



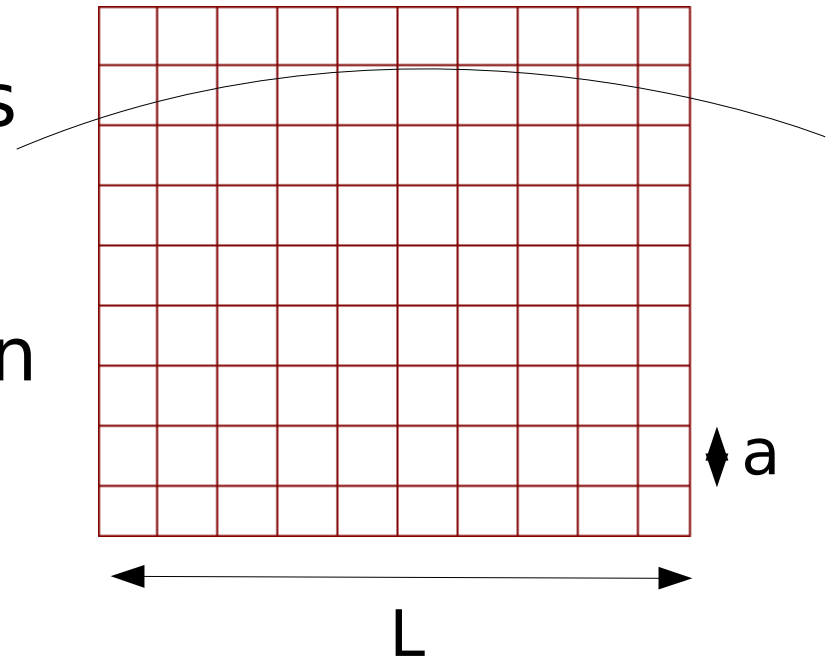
Lattice calculations

- Take a **finite volume** - usually a hypercube
- **Discretize it**, and get a finite, hypercubic lattice
- Calculate observables using path integral
 - Can be done numerically
 - Uses Monte-Carlo methods
- **Artifacts**
 - Finite volume/discretization



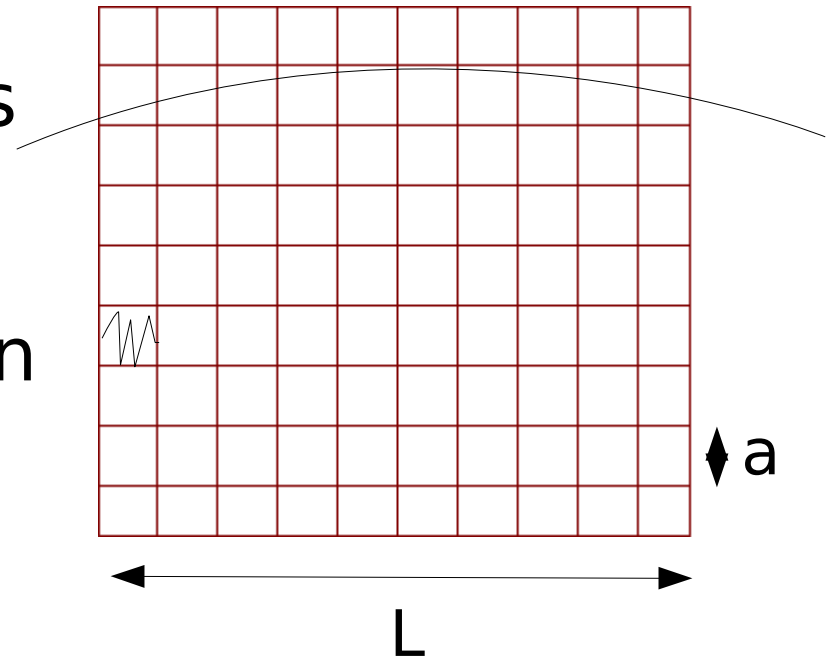
Lattice calculations

- Take a **finite volume** - usually a hypercube
- **Discretize it**, and get a finite, hypercubic lattice
- Calculate observables using path integral
 - Can be done numerically
 - Uses Monte-Carlo methods
- Artifacts
 - Finite volume/discretization
 - Masses vs. wave-lengths



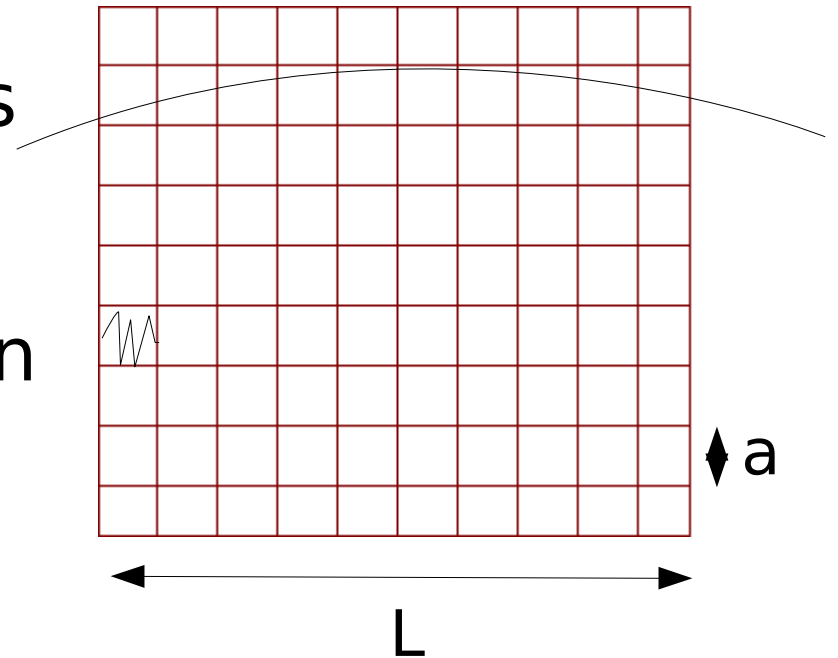
Lattice calculations

- Take a **finite volume** - usually a hypercube
- **Discretize it**, and get a finite, hypercubic lattice
- Calculate observables using path integral
 - Can be done numerically
 - Uses Monte-Carlo methods
- Artifacts
 - Finite volume/discretization
 - Masses vs. wave-lengths



Lattice calculations

- Take a **finite volume** - usually a hypercube
- **Discretize it**, and get a finite, hypercubic lattice
- Calculate observables using path integral
 - Can be done numerically
 - Uses Monte-Carlo methods
- Artifacts
 - Finite volume/discretization
 - Masses vs. wave-lengths
 - Euclidean formulation



Masses from Euclidean propagators

Masses from Euclidean propagators

$$D(p) = \langle O^+(p) O(-p) \rangle$$

- Masses can be inferred from propagators

Masses from Euclidean propagators

$$D(p) = \langle O^+(p) O(-p) \rangle \sim \frac{1}{p^2 + m^2}$$

- Masses can be inferred from propagators

Masses from Euclidean propagators

$$D(p) = \langle O^+(p) O(-p) \rangle \sim \frac{1}{p^2 + m^2}$$

$$C(t) = \langle O^+(x) O(y) \rangle \sim \exp(-m \Delta t)$$

- Masses can be inferred from propagators

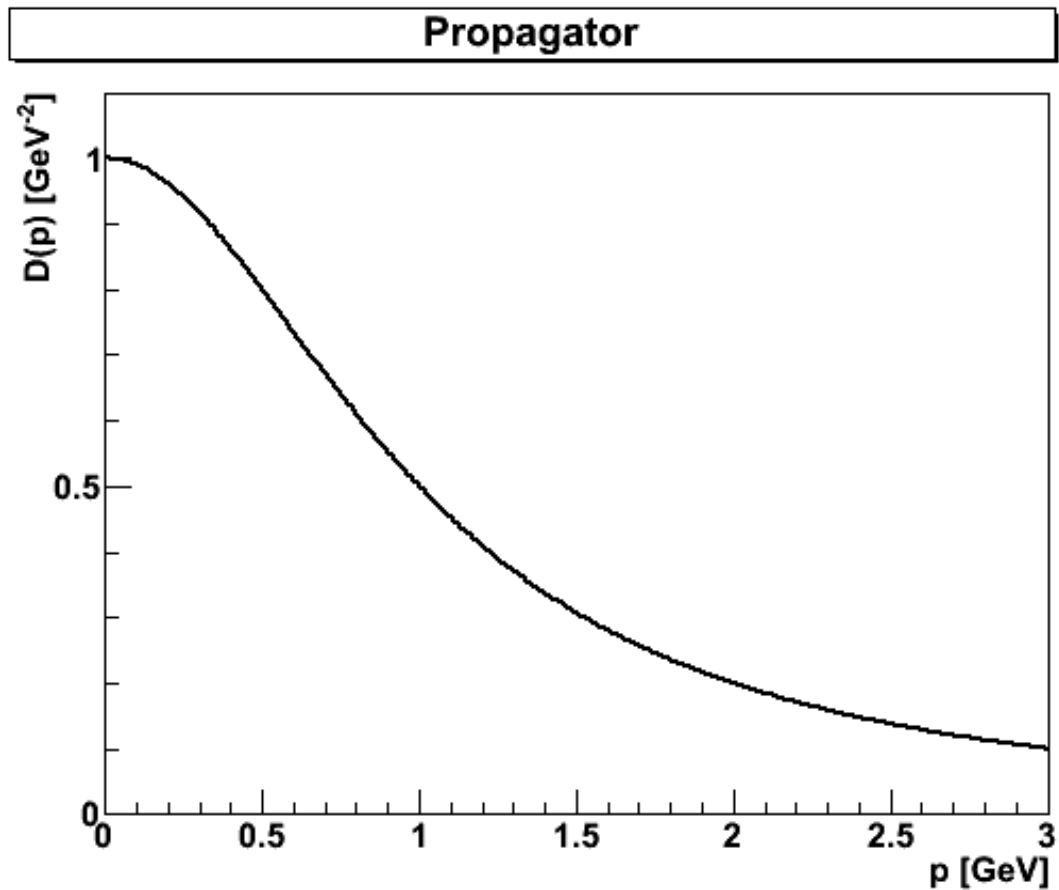
Masses from Euclidean propagators

$$D(p) = \langle O^+(p) O(-p) \rangle \sim \sum \frac{a_i}{p^2 + m_i^2}$$

$$C(t) = \langle O^+(x) O(y) \rangle \sim \sum a_i \exp(-m_i \Delta t)$$
$$\sum a_i = 1 \wedge m_0 < m_1 < \dots$$

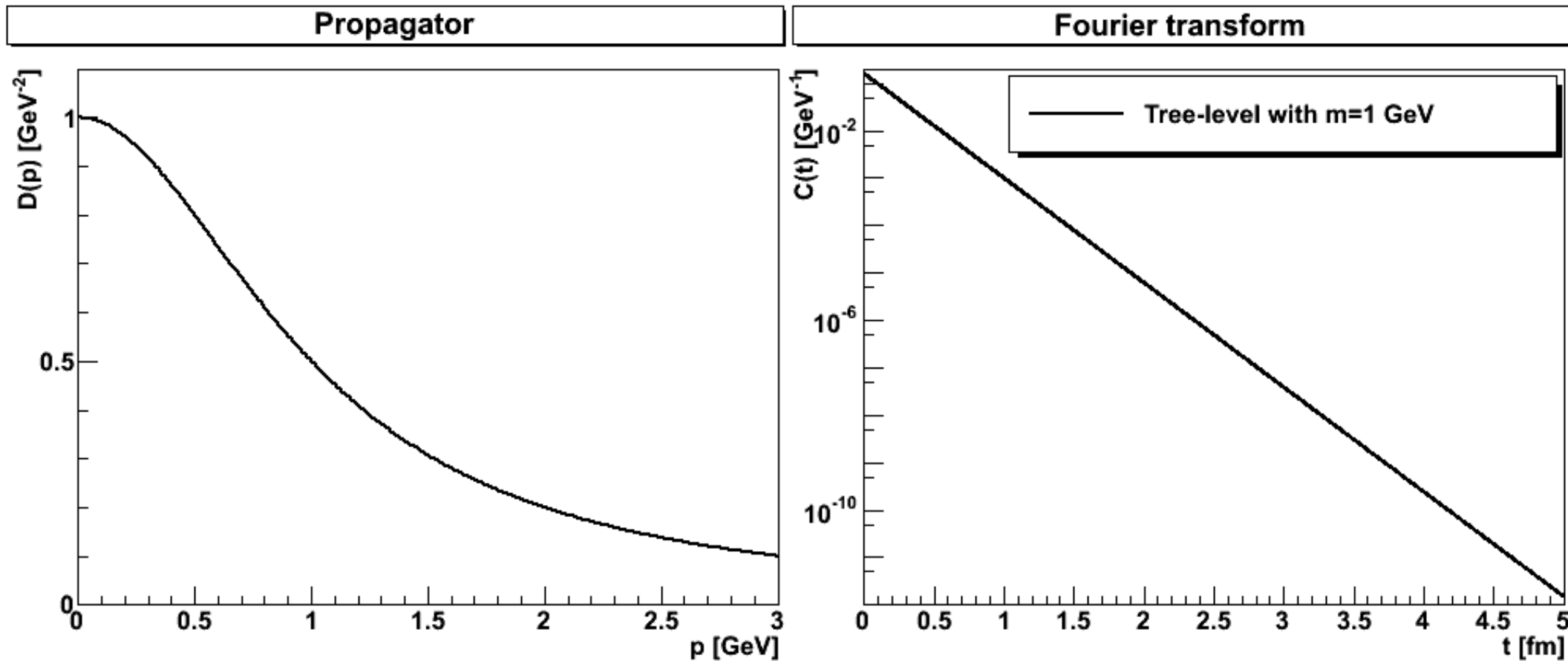
- Masses can be inferred from propagators
- Long-time behavior relevant
 - No exact results on time-like momenta

Masses from Euclidean propagators



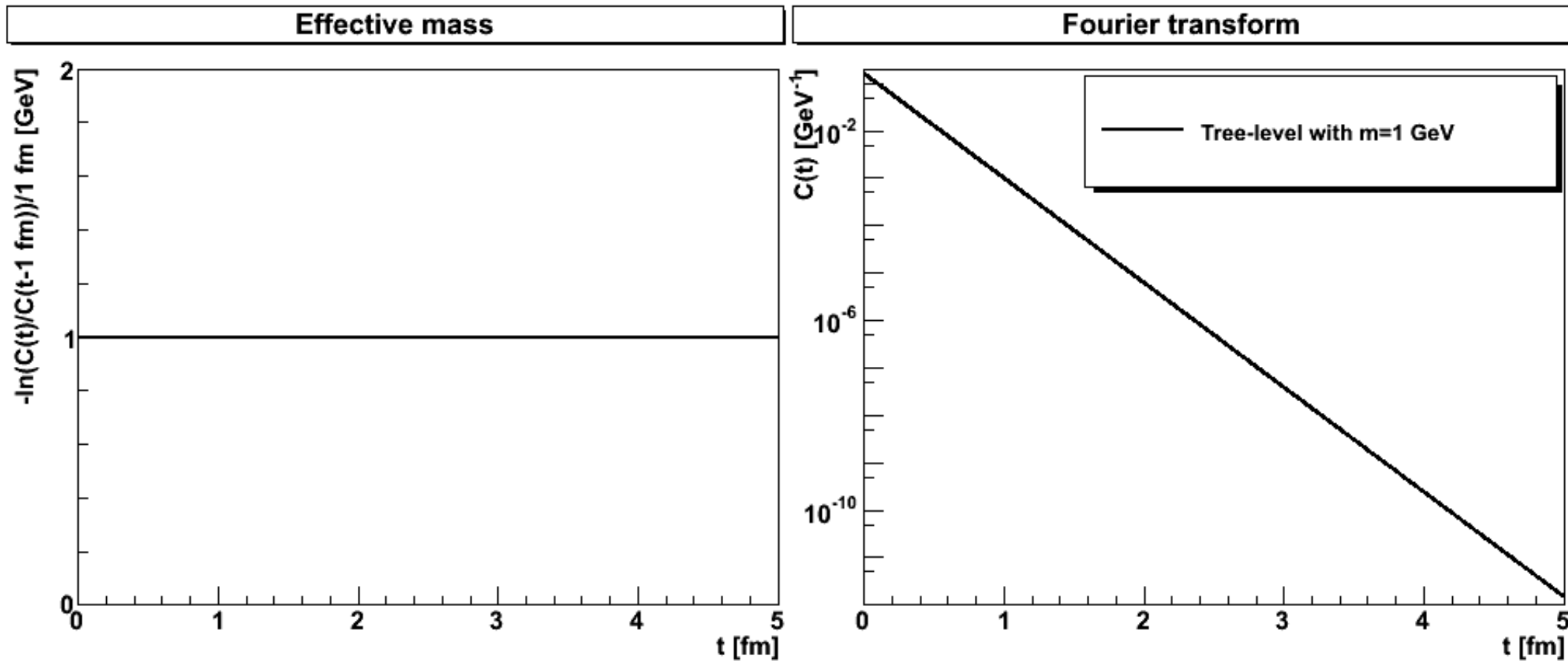
- Masses can be inferred from propagators
- Long-time behavior relevant
 - No exact results on time-like momenta

Masses from Euclidean propagators



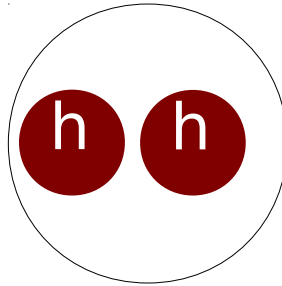
- Masses can be inferred from propagators
- Long-time behavior relevant
 - No exact results on time-like momenta

Masses from Euclidean propagators



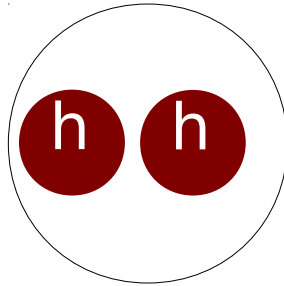
- Masses can be inferred from propagators
- Long-time behavior relevant
 - No exact results on time-like momenta

Higgsonium



- Simplest 0^+ bound state $h^+(x)h(x)$

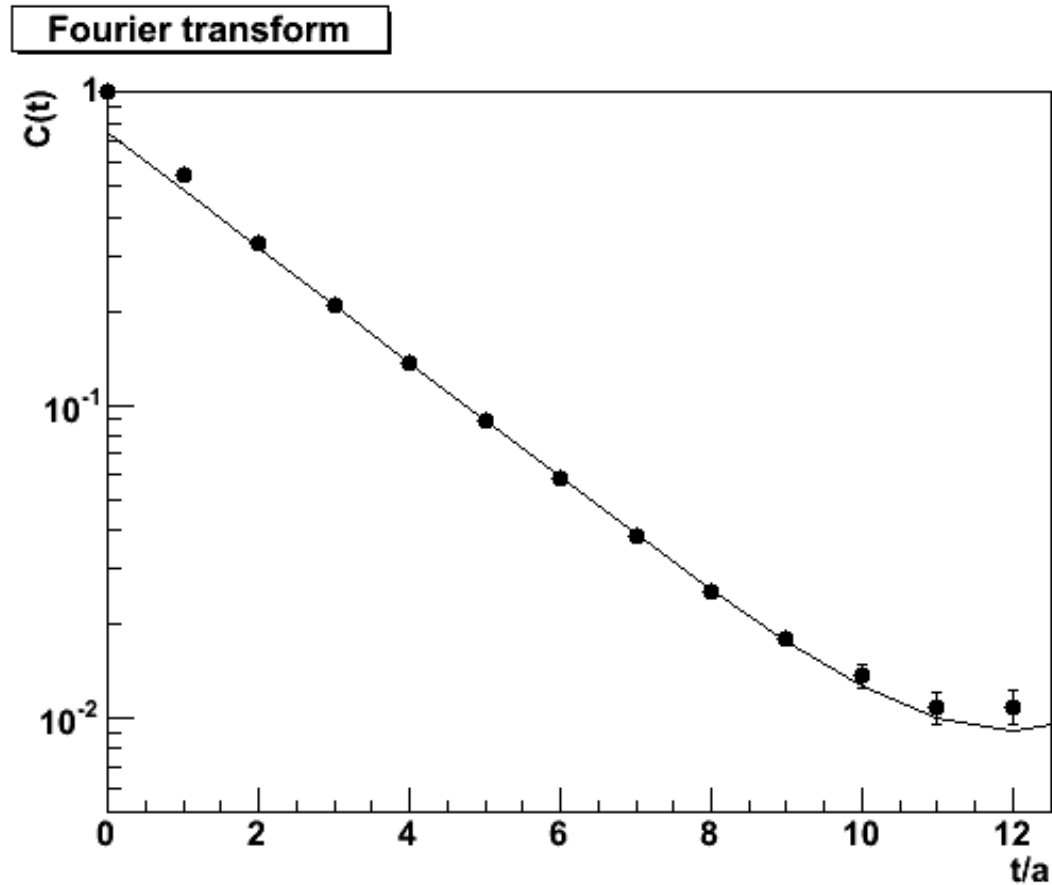
Higgsonium



- Simplest 0^+ bound state $h^+(x)h(x)$
 - Same quantum numbers as the Higgs
 - No weak or flavor charge

Higgsonium

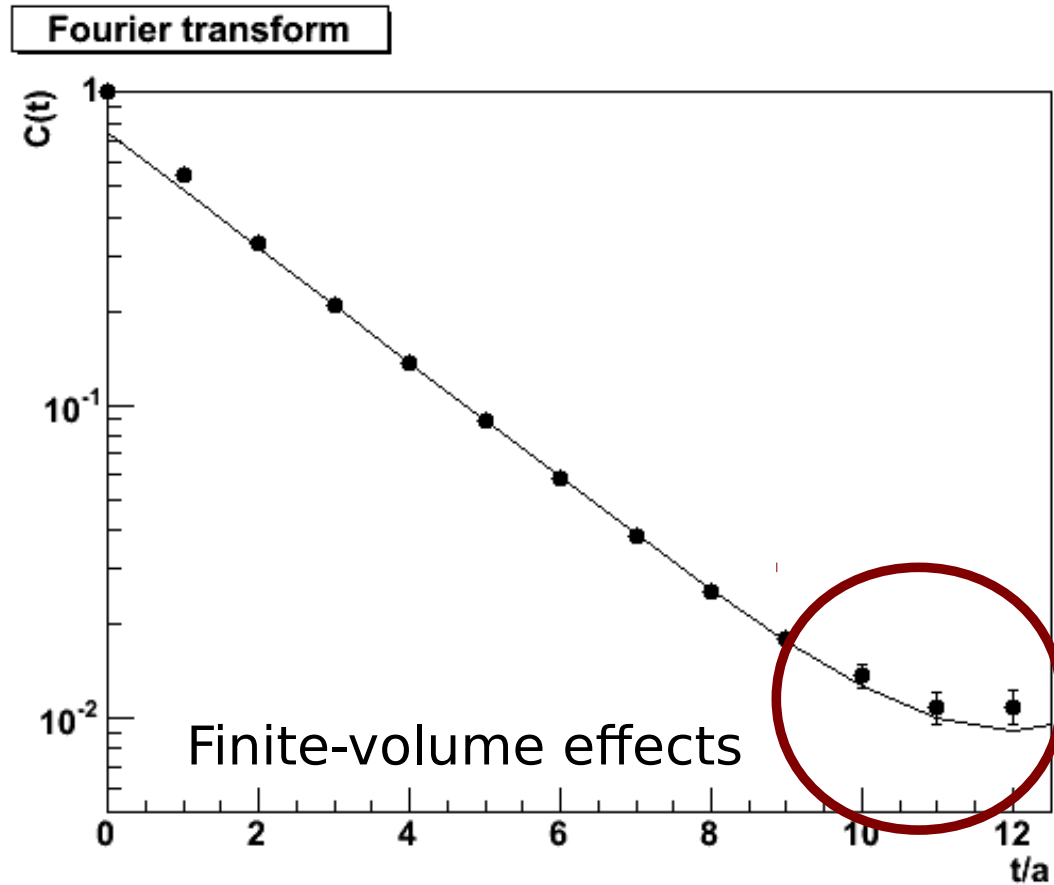
[Maas et al. '13]



- Simplest 0^+ bound state $h^+(x)h(x)$
 - Same quantum numbers as the Higgs
 - No weak or flavor charge

Higgsonium

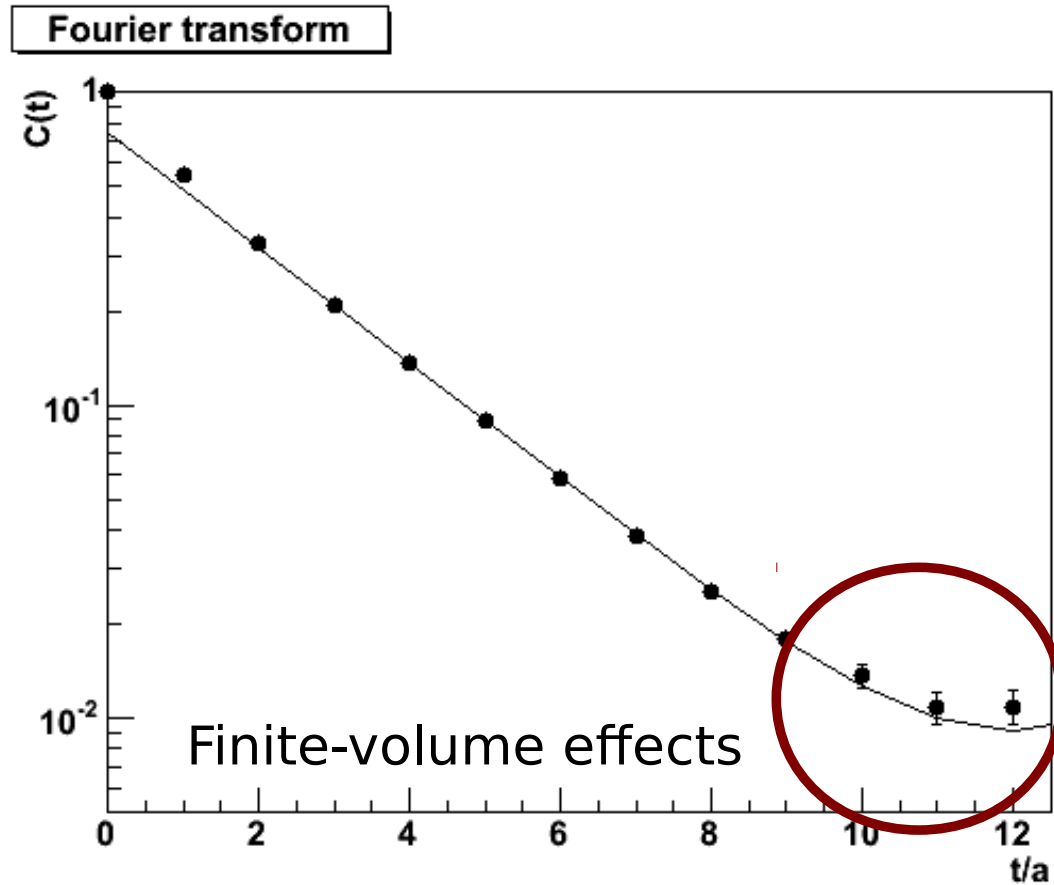
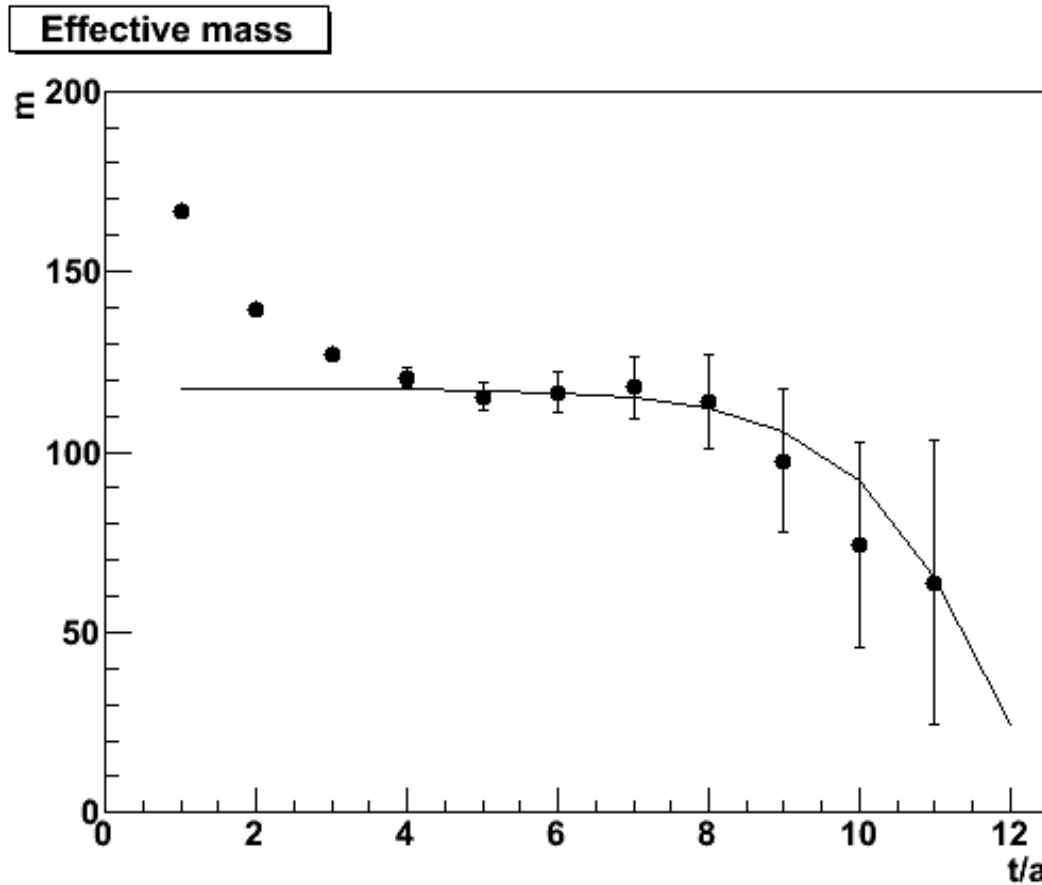
[Maas et al. '13]



- Simplest 0^+ bound state $h^+(x)h(x)$
 - Same quantum numbers as the Higgs
 - No weak or flavor charge

Higgsonium

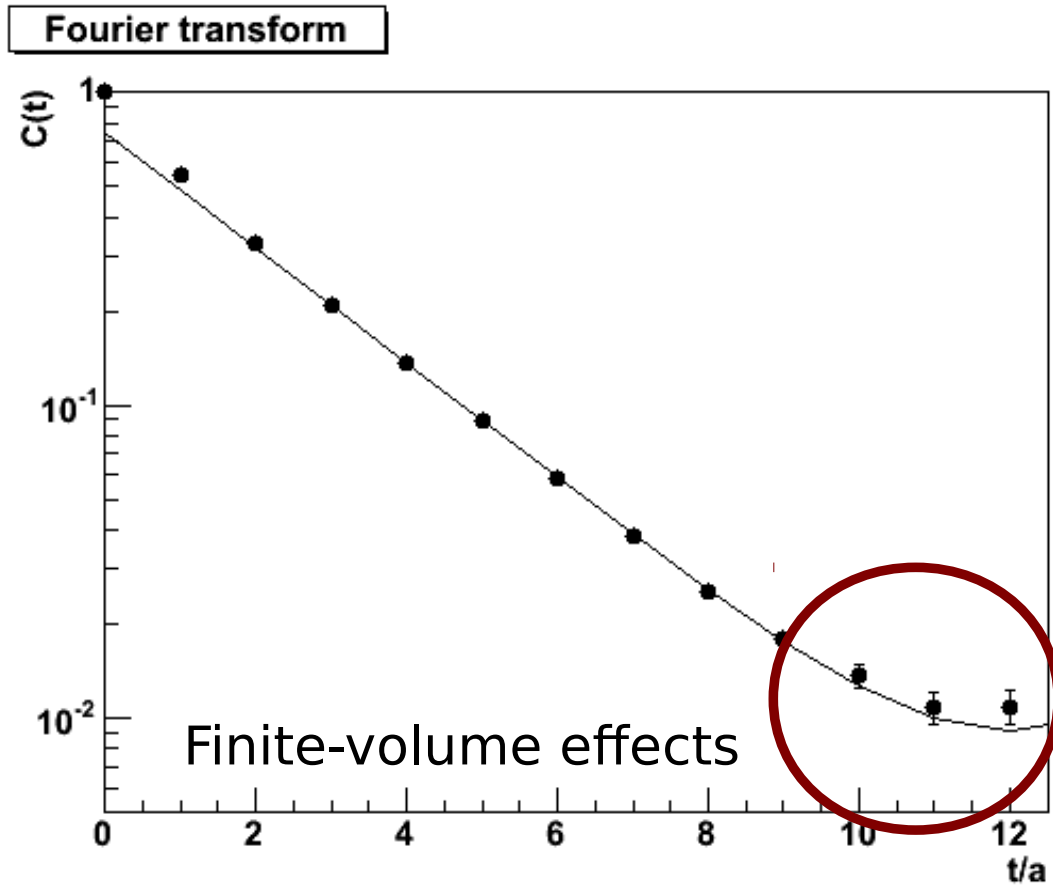
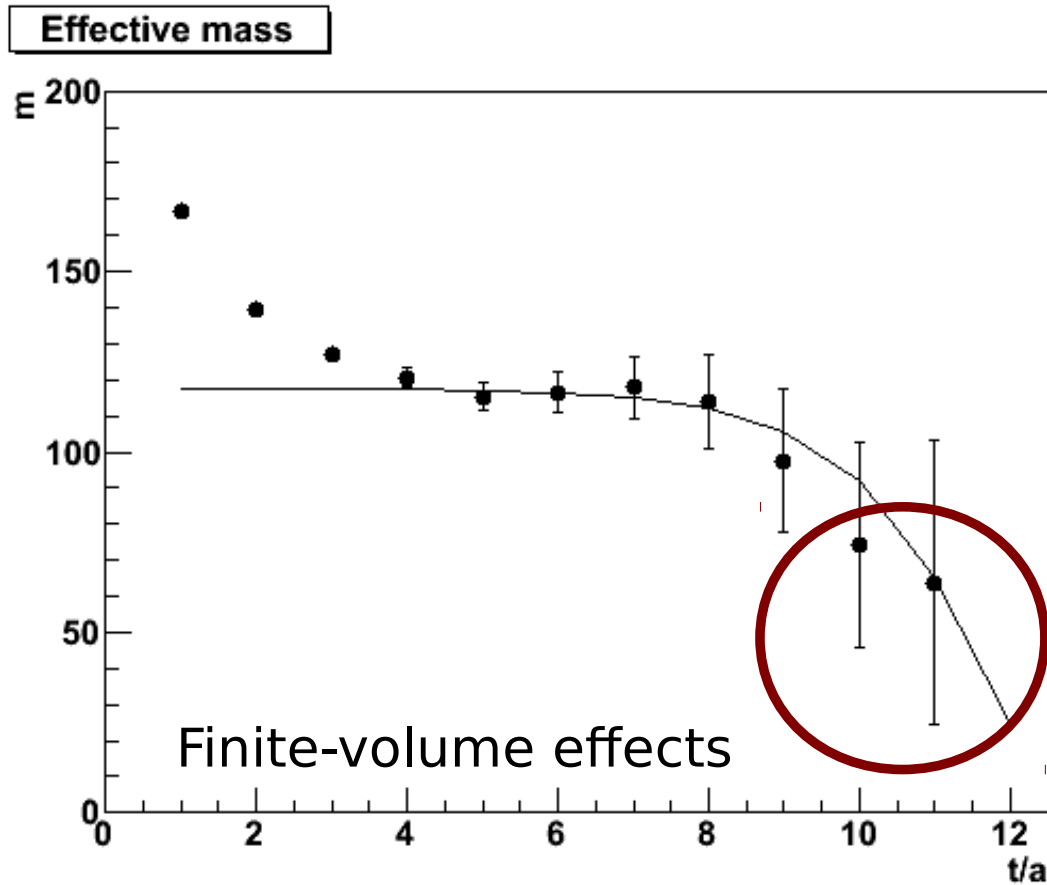
[Maas et al. '13]



- Simplest 0^+ bound state $h^+(x)h(x)$
 - Same quantum numbers as the Higgs
 - No weak or flavor charge
 - Mass is about 120 GeV

Higgsonium

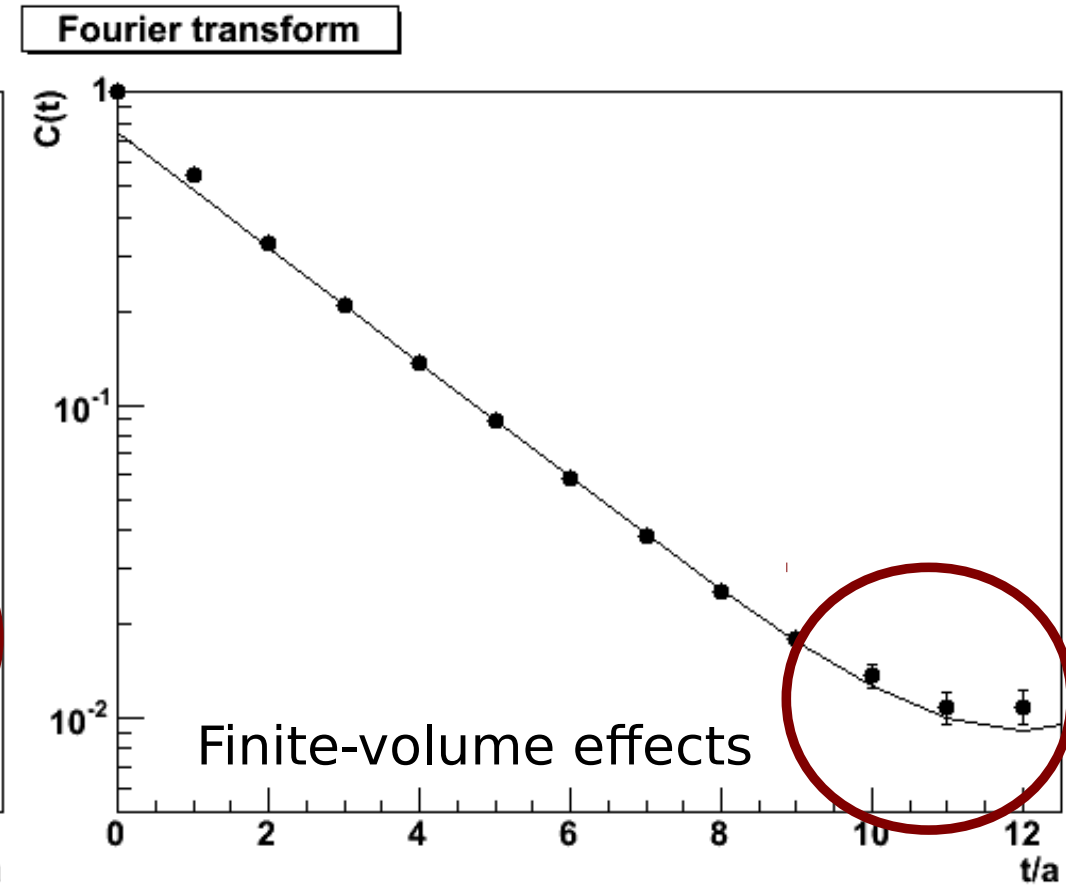
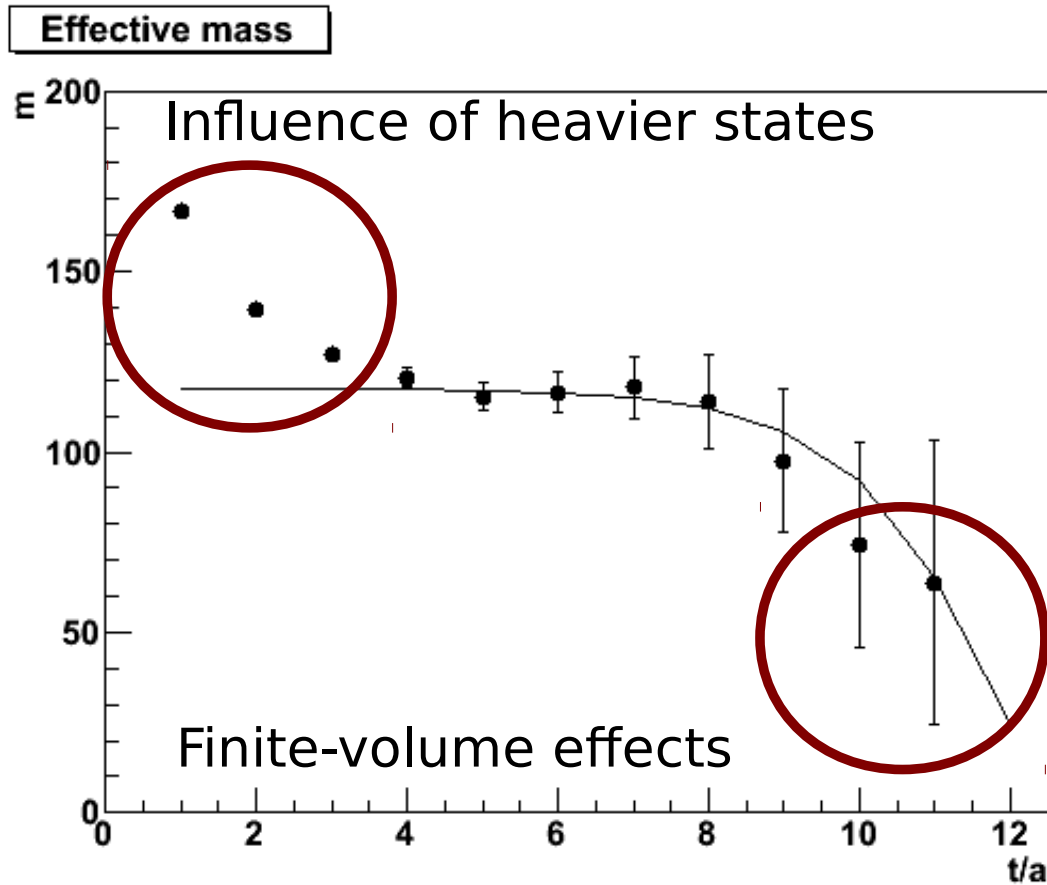
[Maas et al. '13]



- Simplest 0^+ bound state $h^+(x)h(x)$
 - Same quantum numbers as the Higgs
 - No weak or flavor charge
 - Mass is about 120 GeV

Higgsonium

[Maas et al. '13]



- Simplest 0^+ bound state $h^+(x)h(x)$
 - Same quantum numbers as the Higgs
 - No weak or flavor charge
 - Mass is about 120 GeV

Mass relation - Higgs

[Fröhlich et al.'80
Maas'12, Maas & Mufti'13]

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV

Mass relation - Higgs

[Fröhlich et al.'80
Maas'12, Maas & Mufti'13]

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV
- Coincidence?

Mass relation - Higgs

[Fröhlich et al.'80
Maas'12, Maas & Mufti'13]

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV
- Coincidence? No.

Gauge-invariant perturbation theory

Gauge-invariant perturbation theory

[Fröhlich et al. PLB 80
Maas'12, Törek & Maas'16]

1) Formulate gauge-invariant operator

Gauge-invariant perturbation theory

[Fröhlich et al. PLB 80
Maas'12, Törek & Maas'16]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^+ h)(x)(h^+ h)(y) \rangle$$

Gauge-invariant perturbation theory

[Fröhlich et al. PLB 80
Maas'12, Tórek & Maas'16]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle$$

2) Expand Higgs field around fluctuations

$$\begin{aligned} \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle &= c + v^2 \langle \eta^\dagger(x) \eta(y) \rangle \\ &+ v \langle \eta^\dagger \eta^2 + \eta^{\dagger 2} \eta \rangle + \langle \eta^{\dagger 2} \eta^2 \rangle \end{aligned}$$

Gauge-invariant perturbation theory

[Fröhlich et al. PLB 80
Maas'12, Tórek & Maas'16]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle$$

2) Expand Higgs field around fluctuations

$$\begin{aligned} \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle &= c + v^2 \langle \eta^\dagger(x) \eta(y) \rangle \\ &\quad + v \langle \eta^\dagger \eta^2 + \eta^{\dagger 2} \eta \rangle + \langle \eta^{\dagger 2} \eta^2 \rangle \end{aligned}$$

3) Standard perturbation theory

$$\begin{aligned} \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle &= c + v^2 \langle \eta^\dagger(x) \eta(y) \rangle \\ &\quad + \langle \eta^\dagger(x) \eta(y) \rangle \langle \eta^\dagger(x) \eta(y) \rangle + O(g, \lambda) \end{aligned}$$

Gauge-invariant perturbation theory

[Fröhlich et al. PLB 80
Maas'12, Törek & Maas'16]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle$$

2) Expand Higgs field around fluctuations

$$\begin{aligned} \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle &= c + v^2 \langle \eta^\dagger(x) \eta(y) \rangle \\ &+ v \langle \eta^\dagger \eta^2 + \eta^{\dagger 2} \eta \rangle + \langle \eta^{\dagger 2} \eta^2 \rangle \end{aligned}$$

3) Standard perturbation theory

$$\begin{aligned} \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle &= c + v^2 \langle \eta^\dagger(x) \eta(y) \rangle \\ &+ \langle \eta^\dagger(x) \eta(y) \rangle \langle \eta^\dagger(x) \eta(y) \rangle + O(g, \lambda) \end{aligned}$$

4) Compare poles on both sides

Gauge-invariant perturbation theory

[Fröhlich et al. PLB 80
Maas'12, Tórek & Maas'16]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^+ h)(x)(h^+ h)(y) \rangle$$

2) Expand Higgs field around fluctuations

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= c + v^2 \langle \eta^+(x) \eta(y) \rangle \\ &+ v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle \end{aligned}$$

3) Standard perturbation theory

Bound
state
mass

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= c + v^2 \langle \eta^+(x) \eta(y) \rangle \\ &+ \langle \eta^+(x) \eta(y) \rangle \langle \eta^+(x) \eta(y) \rangle + O(g, \lambda) \end{aligned}$$

4) Compare poles on both sides

Gauge-invariant perturbation theory

[Fröhlich et al. PLB 80
Maas'12, Tórek & Maas'16]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle$$

2) Expand Higgs field around fluctuations

$$\begin{aligned} \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle &= c + v^2 \langle \eta^\dagger(x) \eta(y) \rangle \\ &+ v \langle \eta^\dagger \eta^2 + \eta^{\dagger 2} \eta \rangle + \langle \eta^{\dagger 2} \eta^2 \rangle \end{aligned}$$

3) Standard perturbation theory

Bound
state
mass

$$\begin{aligned} \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle &= c + v^2 \langle \eta^\dagger(x) \eta(y) \rangle \\ &+ \langle \eta^\dagger(x) \eta(y) \rangle \langle \eta^\dagger(x) \eta(y) \rangle + O(g, \lambda) \end{aligned}$$

2 x Higgs mass:
Scattering state

4) Compare poles on both sides

Gauge-invariant perturbation theory

[Fröhlich et al. PLB 80
Maas'12, Tórek & Maas'16]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle$$

2) Expand Higgs field around fluctuations

$$\begin{aligned} \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle &= c + v^2 \langle \eta^\dagger(x) \eta(y) \rangle \\ &+ v \langle \eta^\dagger \eta^2 + \eta^{\dagger 2} \eta \rangle + \langle \eta^{\dagger 2} \eta^2 \rangle \end{aligned}$$

3) Standard perturbation theory

Bound
state
mass

$$\begin{aligned} \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle &= c + v^2 \langle \eta^\dagger(x) \eta(y) \rangle \\ &+ \langle \eta^\dagger(x) \eta(y) \rangle \langle \eta^\dagger(x) \eta(y) \rangle + O(g, \lambda) \end{aligned}$$

Higgs
mass

2 x Higgs mass:
Scattering state

4) Compare poles on both sides

Mass relation - Higgs

[Fröhlich et al.'80
Maas'12, Maas & Mufti'13]

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV
- Coincidence? No.
 - Duality between elementary states and bound states

[Fröhlich et al.'80]

$$\langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle \stackrel{h=v+\eta}{\approx} \text{const.} + \langle \eta^\dagger(x)\eta(y) \rangle + O(\eta^3)$$

- Same poles to leading order

Mass relation - Higgs

[Fröhlich et al.'80
Maas'12, Maas & Mufti'13]

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV
- Coincidence? No.
 - Duality between elementary states and bound states

[Fröhlich et al.'80]

$$\langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle \stackrel{h=v+\eta}{\approx} \text{const.} + \langle \eta^\dagger(x)\eta(y) \rangle + O(\eta^3)$$

- Same poles to leading order
- Fröhlich-Morchio-Strocchi (FMS) mechanism

Mass relation - Higgs

[Fröhlich et al.'80
Maas'12, Maas & Mufti'13]

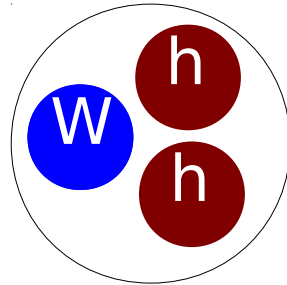
- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV
- Coincidence? No.
 - Duality between elementary states and bound states

[Fröhlich et al.'80]

$$\langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle \stackrel{h=v+\eta}{\approx} \text{const.} + \langle \eta^\dagger(x)\eta(y) \rangle + O(\eta^3)$$

- Same poles to leading order
- Fröhlich-Morchio-Strocchi (FMS) mechanism
- Deeply-bound relativistic state
 - Mass defect \sim constituent mass
 - Cannot describe with quantum mechanics
 - Very different from QCD bound states

Isovector-vector state

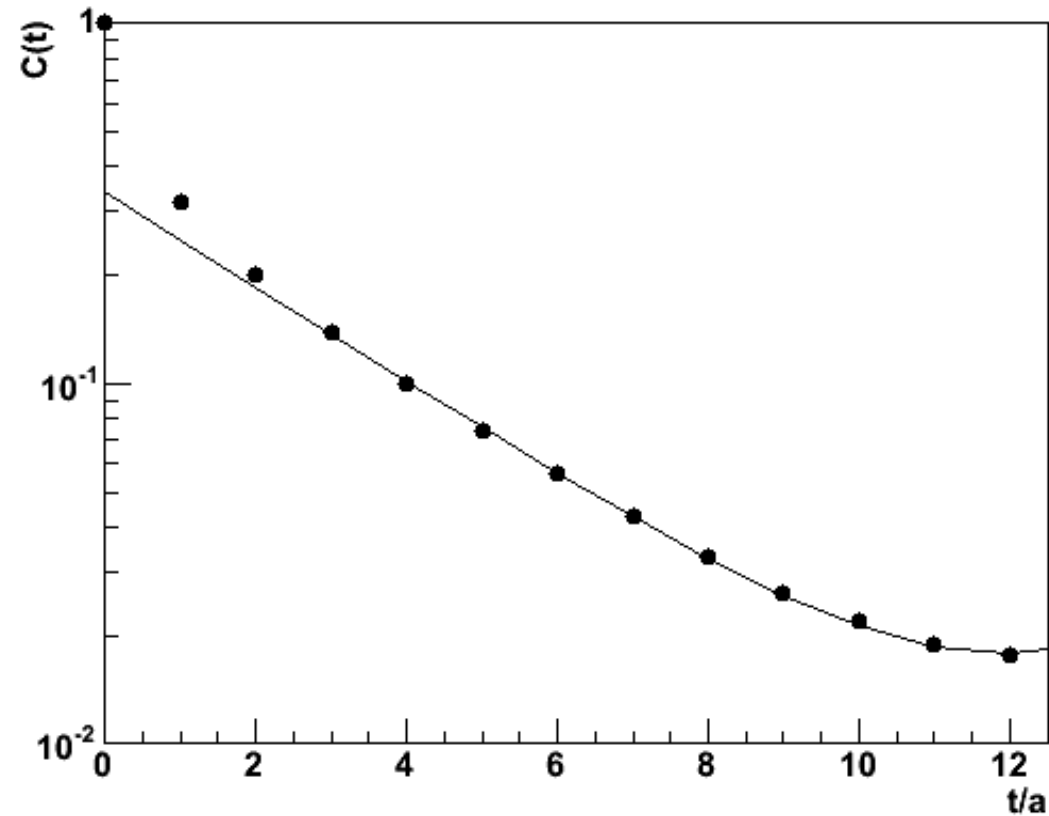


- Vector state 1^- with operator $tr t^a \frac{h^+}{\sqrt{h^+ h}} D_\mu \frac{h}{\sqrt{h^+ h}}$
 - Only in a Higgs phase close to a simple particle
 - Custodial triplet, instead of gauge triplet

Isvector-vector state

[Maas et al. '13]

Fourier transform

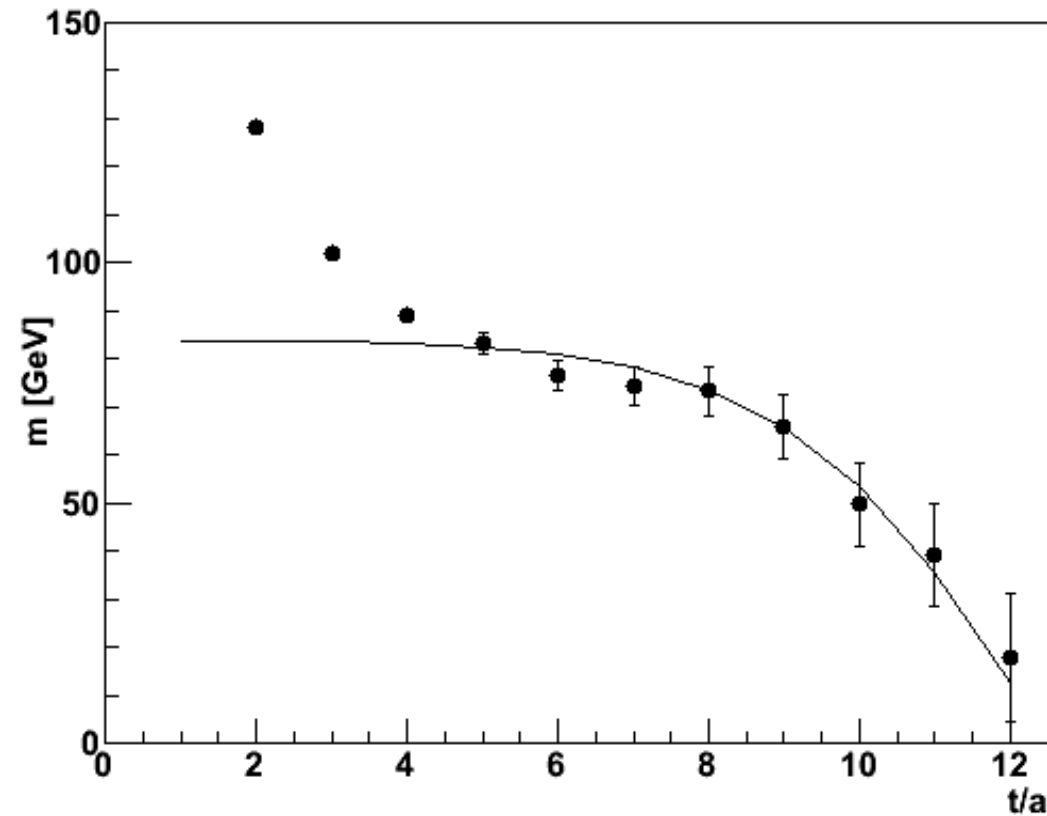


- Vector state 1^- with operator $tr t^a \frac{h^+}{\sqrt{h^+ h}} D_\mu \frac{h}{\sqrt{h^+ h}}$
 - Only in a Higgs phase close to a simple particle
 - Custodial triplet, instead of gauge triplet

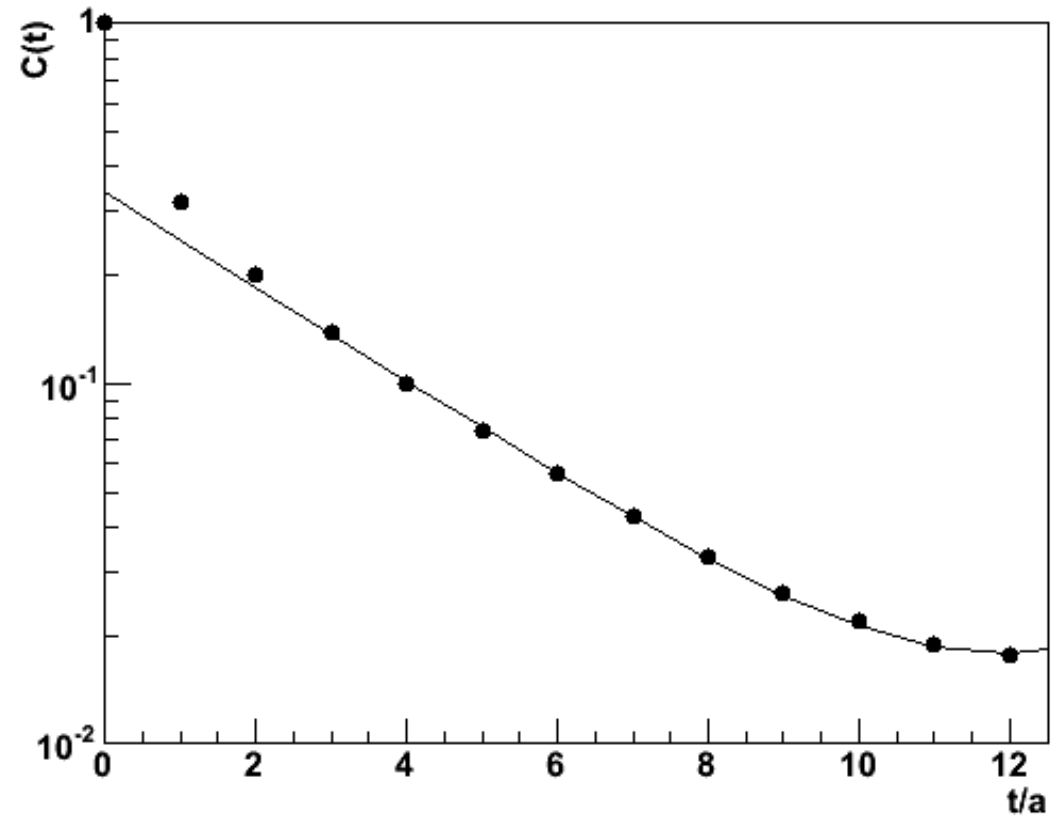
Isovector-vector state

[Maas et al. '13]

Effective mass



Fourier transform



- Vector state 1^- with operator $tr t^a \frac{h^+}{\sqrt{h^+ h}} D_\mu \frac{h}{\sqrt{h^+ h}}$
 - Only in a Higgs phase close to a simple particle
 - Custodial triplet, instead of gauge triplet
 - Mass about 80 GeV

Mass relation - W

[Fröhlich et al.'80
Maas'12]

- Vector state: 80 GeV
- W at tree-level: 80 GeV
 - W not scale or scheme dependent

Mass relation - W

[Fröhlich et al.'80
Maas'12]

- Vector state: 80 GeV
- W at tree-level: 80 GeV
 - W not scale or scheme dependent
- Same mechanism

$$\langle (h^\dagger D_\mu h)(x) (h^\dagger D_\mu h)(y) \rangle$$

Mass relation - W

[Fröhlich et al.'80
Maas'12]

- Vector state: 80 GeV
- W at tree-level: 80 GeV
 - W not scale or scheme dependent
- Same mechanism

$$\langle (h^\dagger D_\mu h)(x) (h^\dagger D_\mu h)(y) \rangle$$
$$h = v + \eta$$
$$\approx \text{const.} + \langle W_\mu(x) W_\mu(y) \rangle + O(\eta^3)$$
$$\partial v = 0$$

Mass relation - W

[Fröhlich et al.'80
Maas'12]

- Vector state: 80 GeV
- W at tree-level: 80 GeV
 - W not scale or scheme dependent
- Same mechanism

$$\langle (h^\dagger D_\mu h)(x) (h^\dagger D_\mu h)(y) \rangle$$

$$h = v + \eta$$
$$\approx \text{const.} + \langle W_\mu(x) W_\mu(y) \rangle + O(\eta^3)$$

$$\partial v = 0$$

- Same poles at leading order
 - Remains true beyond leading order

Mass relation - W

[Fröhlich et al.'80
Maas'12]

- Vector state: 80 GeV
- W at tree-level: 80 GeV
 - W not scale or scheme dependent
- Same mechanism

$$\langle (h^\dagger D_\mu h)(x) (h^\dagger D_\mu h)(y) \rangle$$

$$h = v + \eta$$

$$\approx \text{const.} + \langle W_\mu(x) W_\mu(y) \rangle + O(\eta^3)$$

$$\partial v = 0$$

- Same poles at leading order
 - Remains true beyond leading order
 - Exchanges a gauge for a custodial triplet

What about the rest?

[Fröhlich et al.'80]

- Quarks and gluons
 - Anyhow bound by confinement in bound states
 - Top subtle, but same principle

What about the rest?

[Fröhlich et al.'80,
Egger et al., unpublished]

- Quarks and gluons
 - Anyhow bound by confinement in bound states
 - Top subtle, but same principle
 - But open flavor needs a Higgs

What about the rest?

[Fröhlich et al.'80,
Egger et al., unpublished]

- Quarks and gluons
 - Anyhow bound by confinement in bound states
 - Top subtle, but same principle
 - But open flavor needs a Higgs - $qqqh$

What about the rest?

[Fröhlich et al.'80,
Egger et al., unpublished]

- Quarks and gluons
 - Anyhow bound by confinement in bound states
 - Top subtle, but same principle
 - But open flavor needs a Higgs - $qqqh$
- Leptons
 - Actually Higgs-lepton bound-states
 - Enormous mass defects
 - Requires confirmation
 - Except for right-handed (Dirac) neutrino

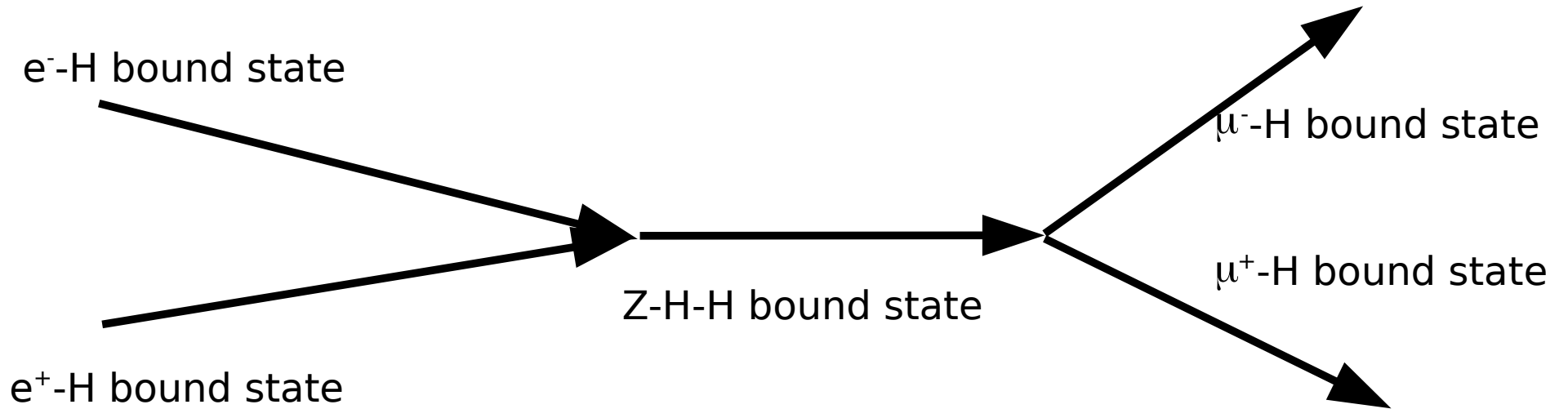
What about the rest?

[Fröhlich et al.'80,
Egger et al., unpublished]

- Quarks and gluons
 - Anyhow bound by confinement in bound states
 - Top subtle, but same principle
 - But open flavor needs a Higgs - $qqqh$
- Leptons
 - Actually Higgs-lepton bound-states
 - Enormous mass defects
 - Requires confirmation
 - Except for right-handed (Dirac) neutrino
- Photons
 - QED similar but simpler

How events looks like (LEP/ILC)

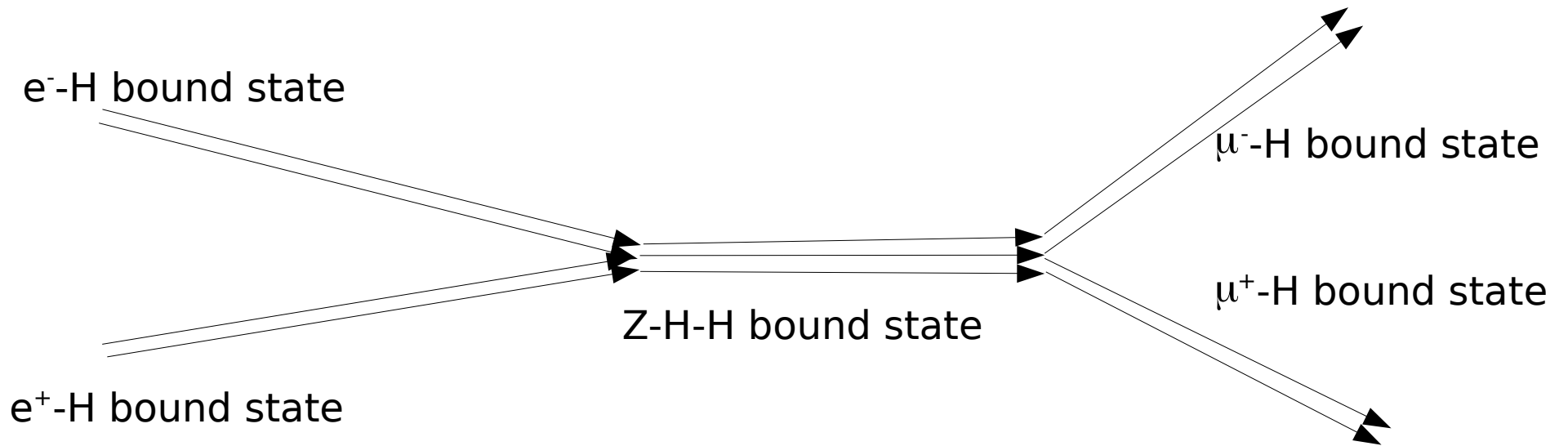
[Maas'12]



- Collision of bound states

How events looks like (LEP/ILC)

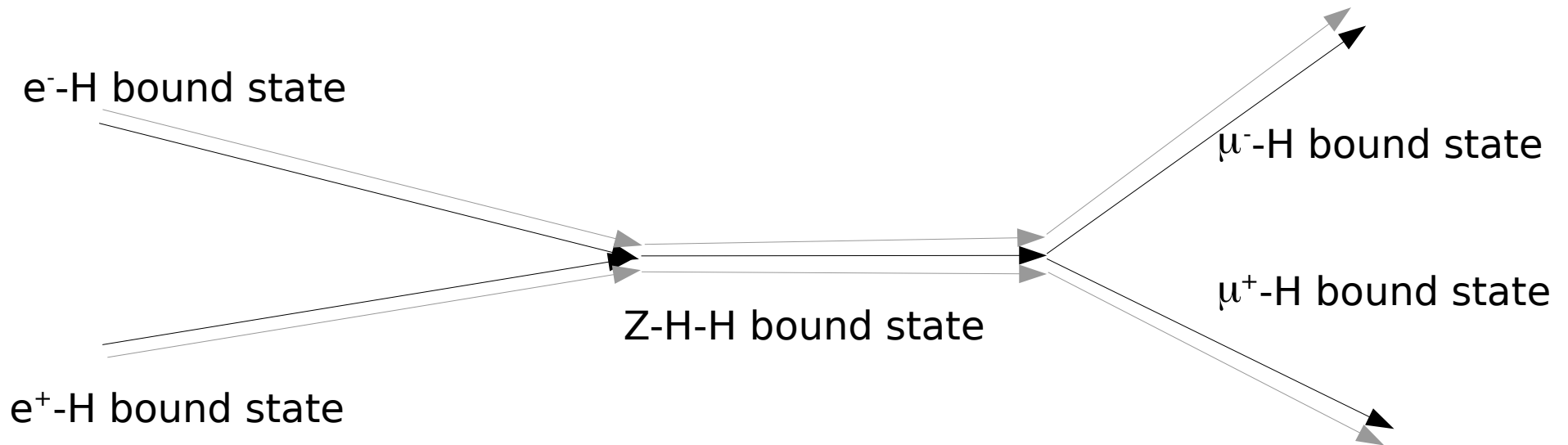
[Maas'12]



- Collision of bound states - 'constituent' particles

How events looks like (LEP/ILC)

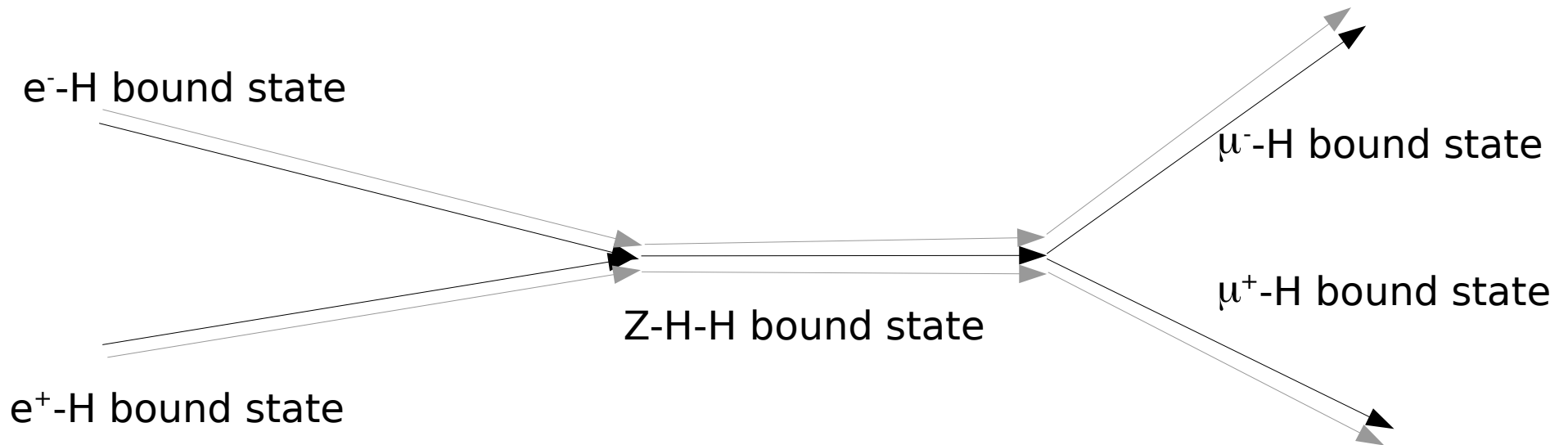
[Maas'12]



- Collision of bound states - 'constituent' particles
- Higgs partners just spectators
 - Similar to pp collisions

How events looks like (LEP/ILC)

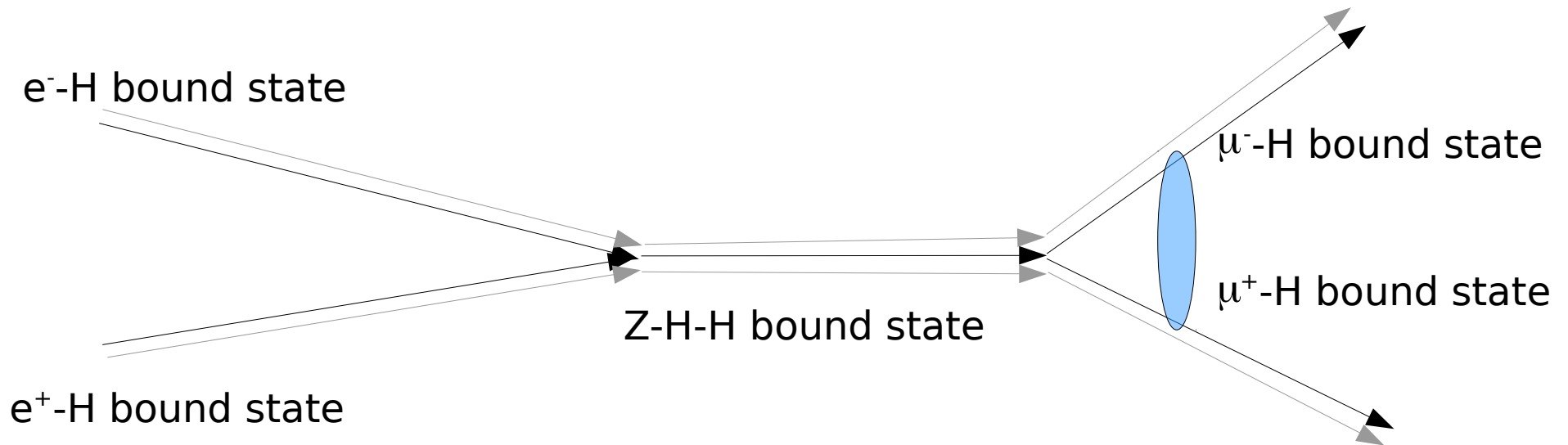
[Maas'12]



- Collision of bound states - 'constituent' particles
- Higgs partners just spectators
 - Similar to pp collisions
- Sub-leading contributions

How events looks like (LEP/ILC)

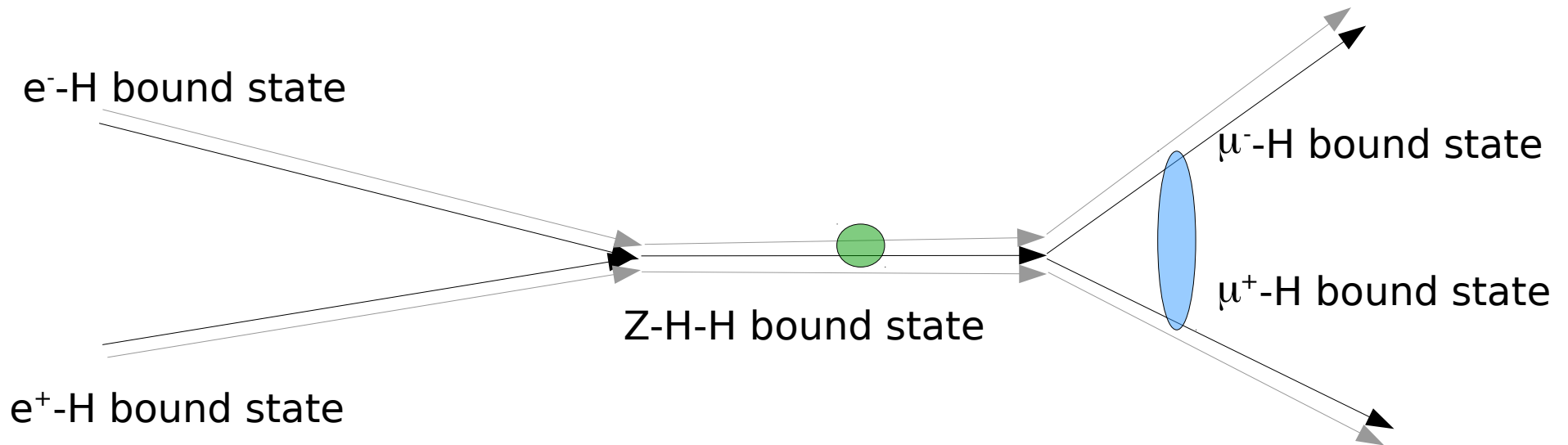
[Maas'12]



- Collision of bound states - 'constituent' particles
- Higgs partners just spectators
 - Similar to pp collisions
- Sub-leading contributions
 - **Ordinary ones**: Large and detected

How events looks like (LEP/ILC)

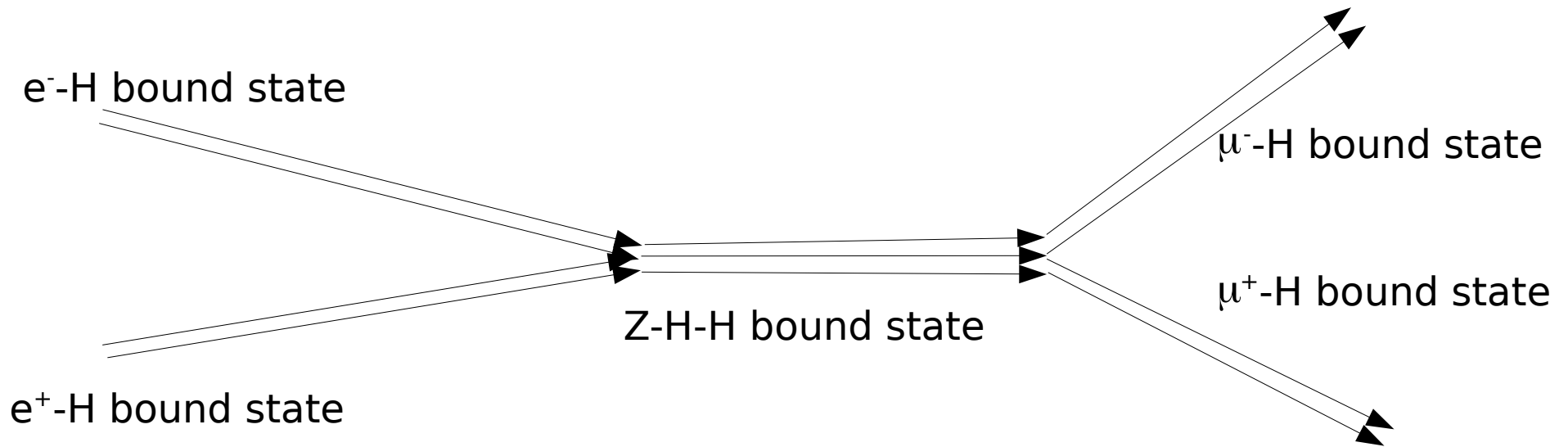
[Maas'12]



- Collision of bound states - 'constituent' particles
- Higgs partners just spectators
 - Similar to pp collisions
- Sub-leading contributions
 - **Ordinary ones**: Large and detected
 - **New ones**: **Small**, require more sensitivity

How events looks like (LEP/ILC)

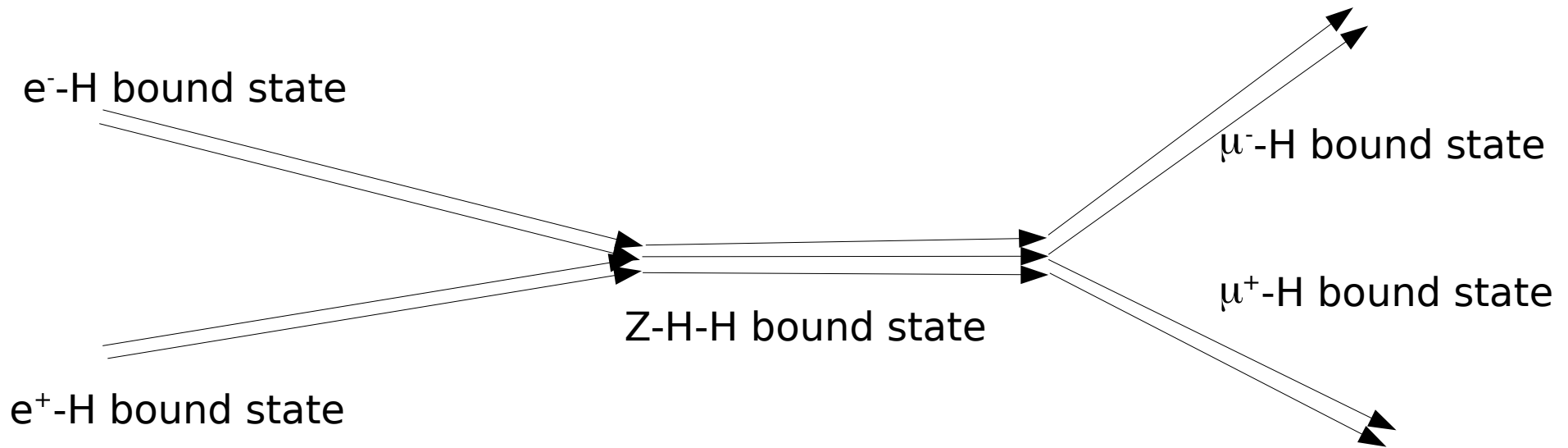
[Maas'12,
Egger et al., unpublished]



- Description of impact?

How events looks like (LEP/ILC)

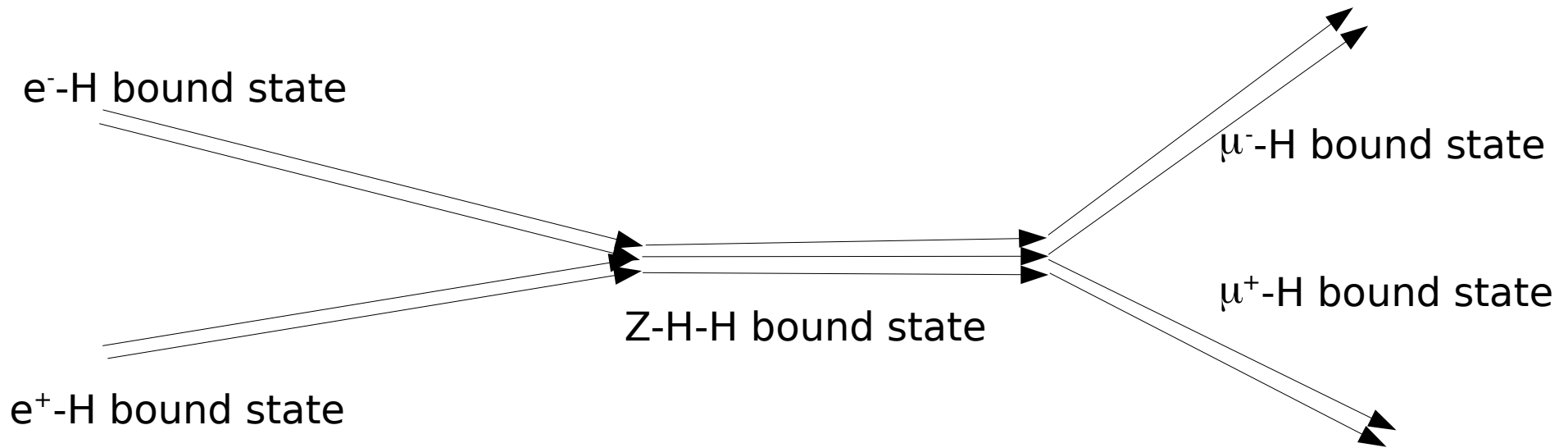
[Maas'12,
Egger et al., unpublished]



- Description of impact? Gauge-invariant perturbation theory!

How events looks like (LEP/ILC)

[Maas'12,
Egger et al., unpublished]

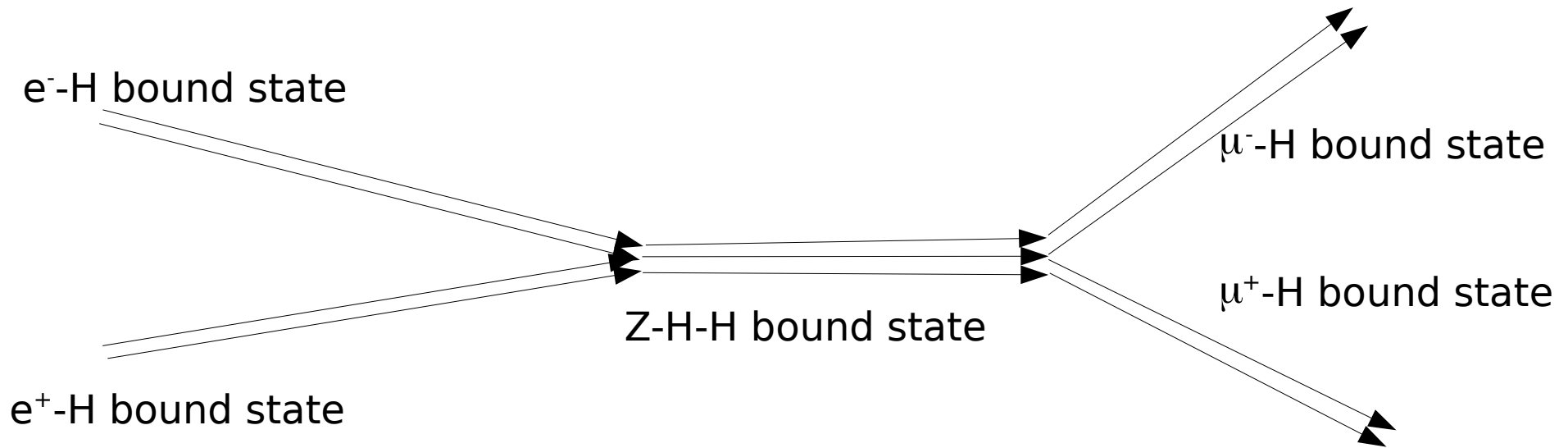


- Description of impact? Gauge-invariant perturbation theory!

$$\langle hehe|h\mu h\mu \rangle$$

How events looks like (LEP/ILC)

[Maas'12,
Egger et al., unpublished]



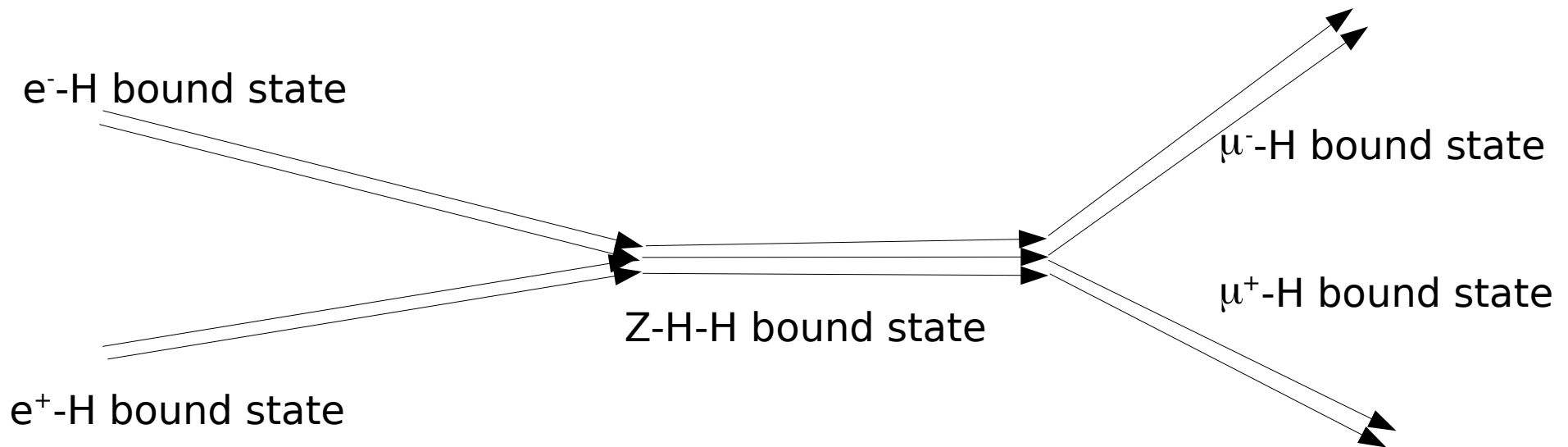
- Description of impact? Gauge-invariant perturbation theory!

$$\langle hehe|h\mu h\mu \rangle = \langle ee|\mu\mu \rangle$$

- Ordinary contribution

How events looks like (LEP/ILC)

[Maas'12,
Egger et al., unpublished]



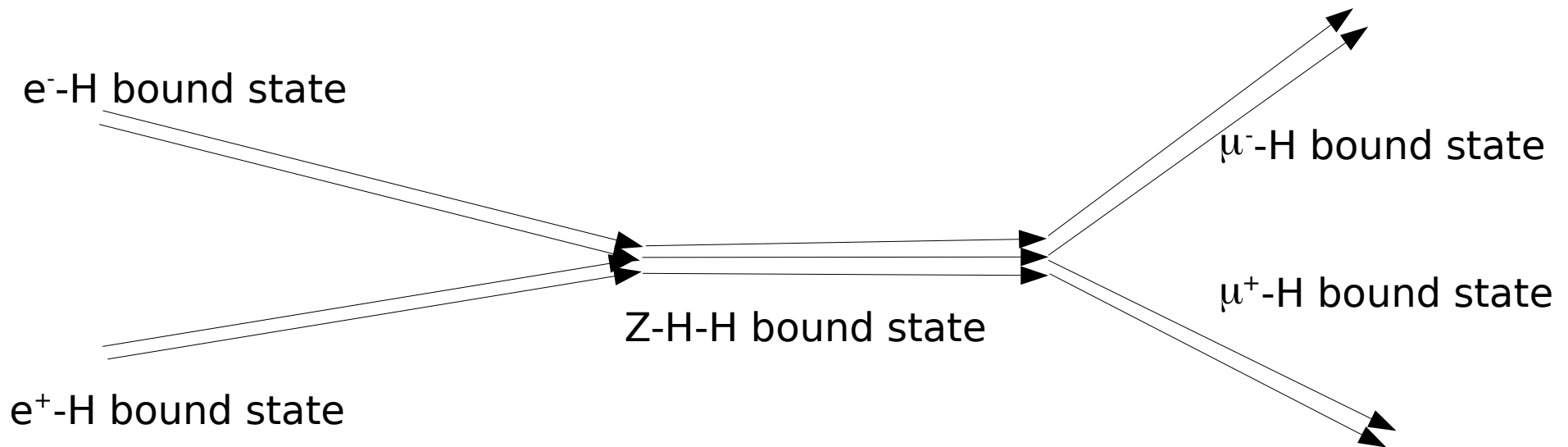
- Description of impact? Gauge-invariant perturbation theory!

$$\langle hehe|h\mu h\mu \rangle = \langle ee|\mu\mu \rangle + \langle \eta\eta \rangle \langle ee|\mu\mu \rangle$$

- Ordinary contribution
- Modification of ordinary contribution

How events looks like (LEP/ILC)

[Maas'12,
Egger et al., unpublished]



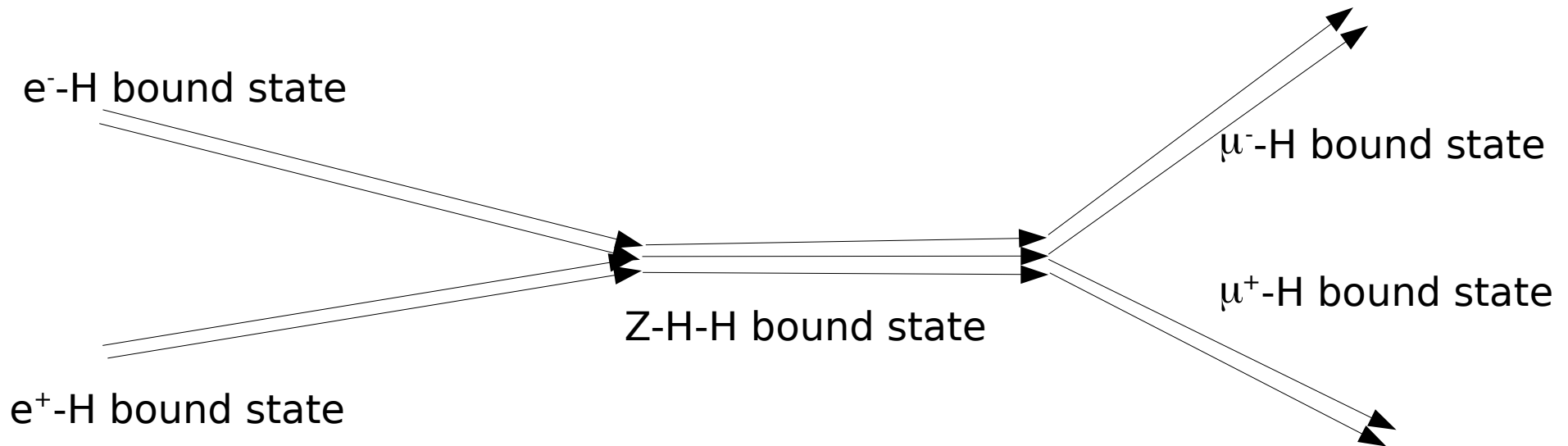
- Description of impact? Gauge-invariant perturbation theory!

$$\langle hehe | h\mu h\mu \rangle = \langle ee | \mu\mu \rangle + \langle \eta\eta \rangle \langle ee | \mu\mu \rangle + \langle ee \rangle \langle \eta\eta | \mu\mu \rangle$$

- Ordinary contribution
- Modification of ordinary contribution
- Higgs as initial state

How events looks like (LEP/ILC)

[Maas'12,
Egger et al., unpublished]



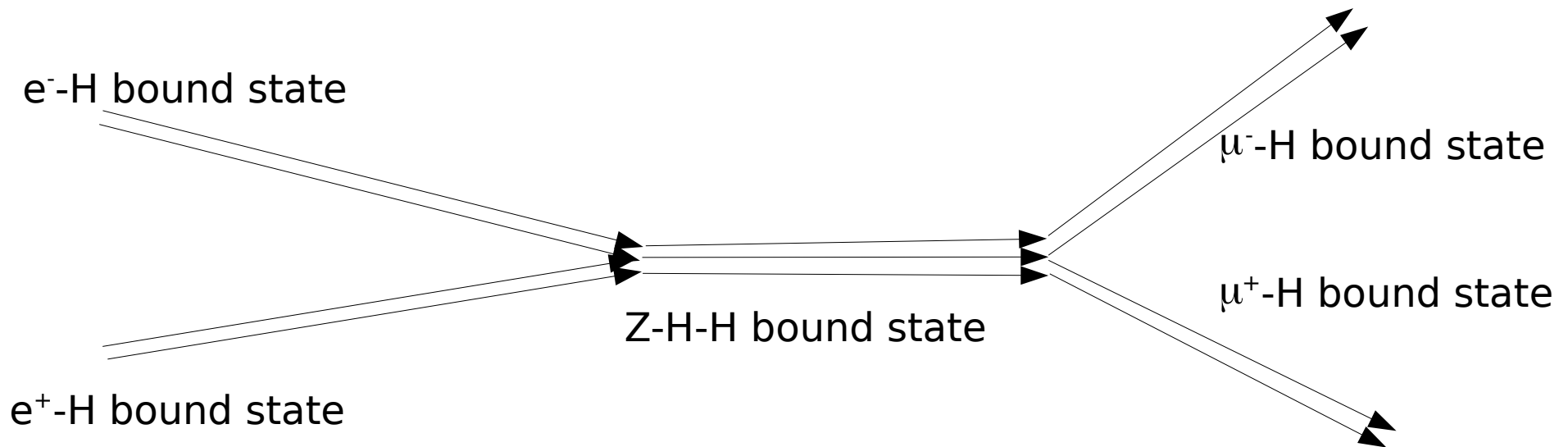
- Description of impact? Gauge-invariant perturbation theory!

$$\langle hehe|h\mu h\mu \rangle = \langle ee|\mu\mu \rangle + \langle \eta\eta \rangle \langle ee|\mu\mu \rangle + \langle ee \rangle \langle \eta\eta|\mu\mu \rangle + \dots$$

- Ordinary contribution
- Modification of ordinary contribution
- Higgs as initial state
- More contributions...

How events looks like (LEP/ILC)

[Maas'12,
Egger et al., unpublished]



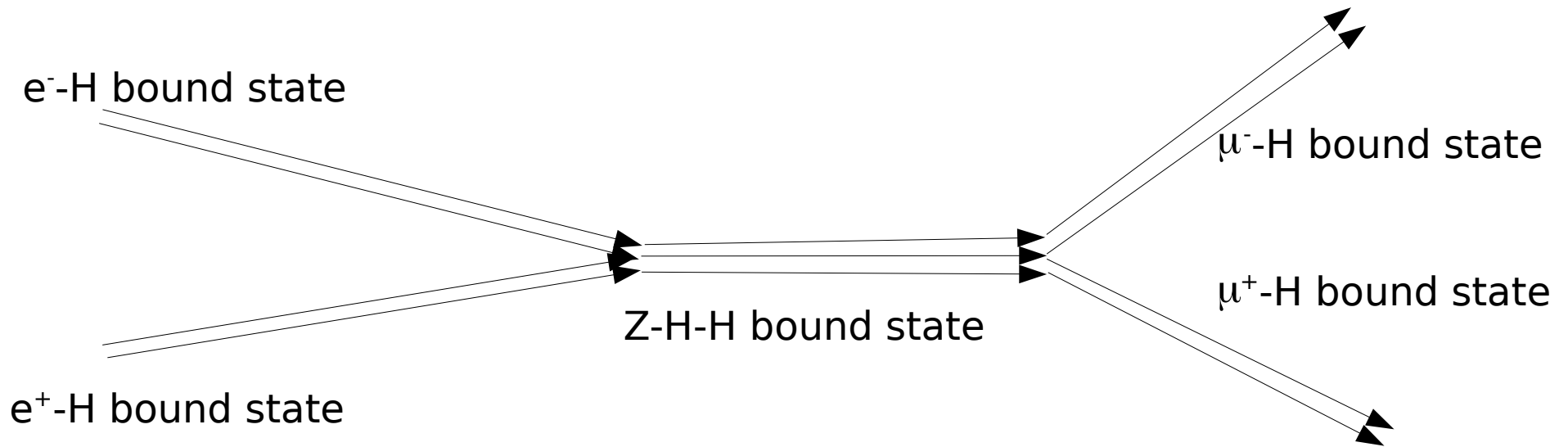
- Description of impact? Gauge-invariant perturbation theory!

$$\langle hehe|h\mu h\mu \rangle = \langle ee|\mu\mu \rangle + \langle \eta\eta \rangle \langle ee|\mu\mu \rangle + \langle ee \rangle \langle \eta\eta|\mu\mu \rangle + \dots$$

- Ordinary contribution
- Modification of ordinary contribution
- Higgs as initial state
- More contributions...complicated

How events looks like (LEP/ILC)

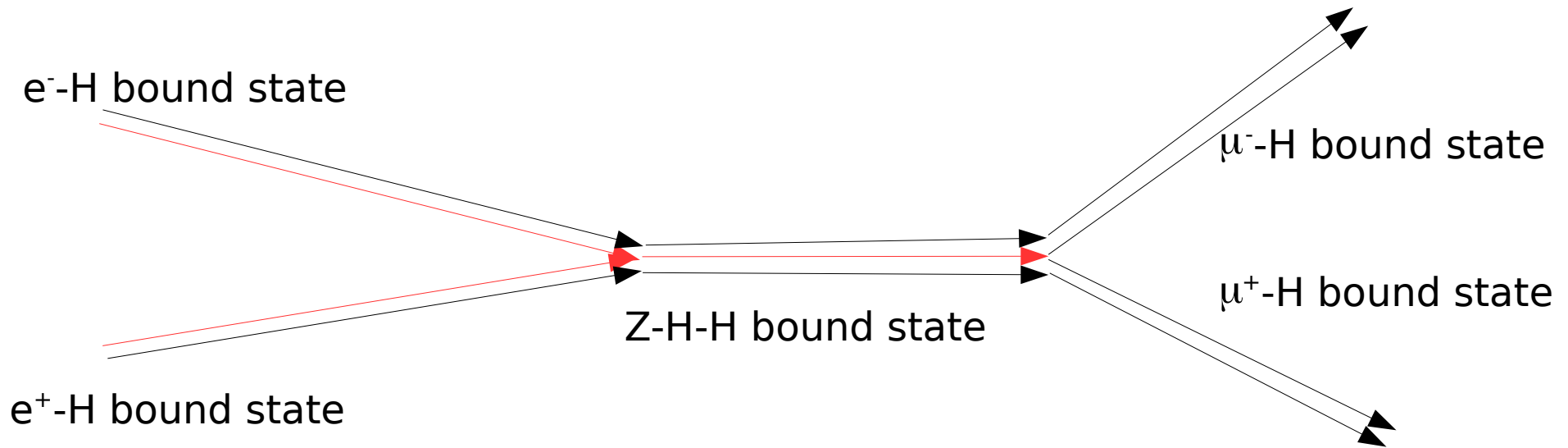
[Maas'12,
Egger et al., unpublished]



- Description of impact? PDF-type language!

How events looks like (LEP/ILC)

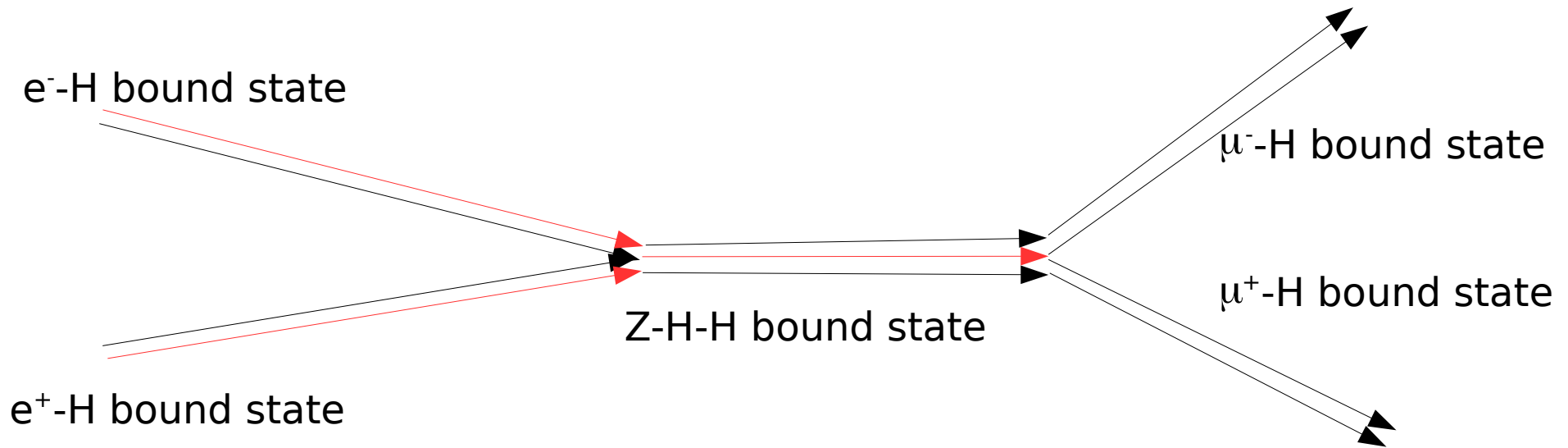
[Maas'12,
Egger et al., unpublished]



- Description of impact? PDF-type language!
- Interacting particles either electrons

How events looks like (LEP/ILC)

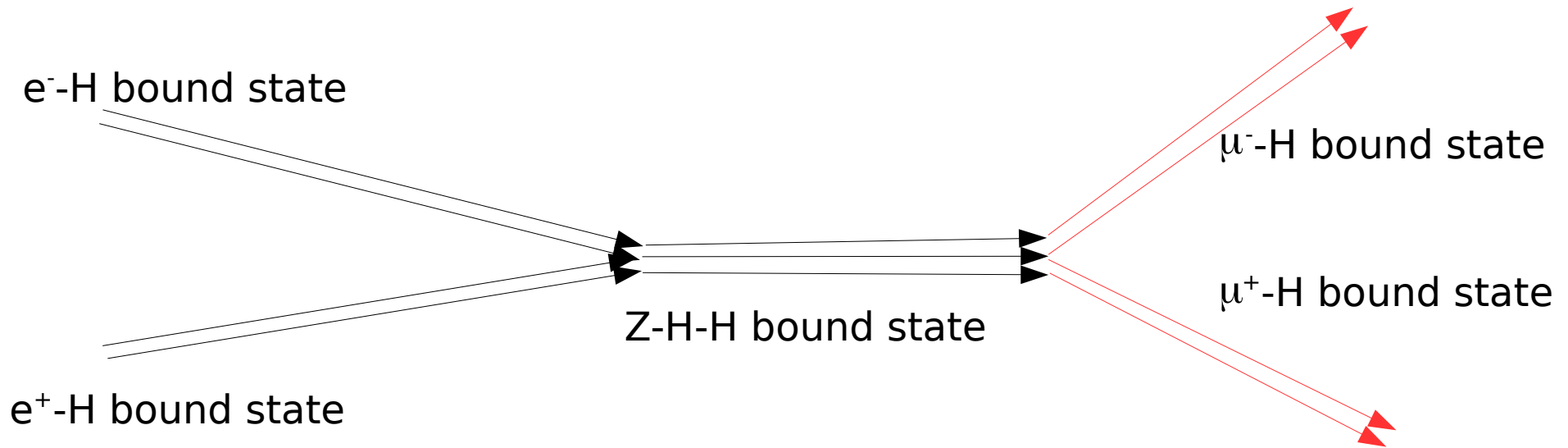
[Maas'12,
Egger et al., unpublished]



- Description of impact? PDF-type language!
- Interacting particles either electrons or Higgs

How events looks like (LEP/ILC)

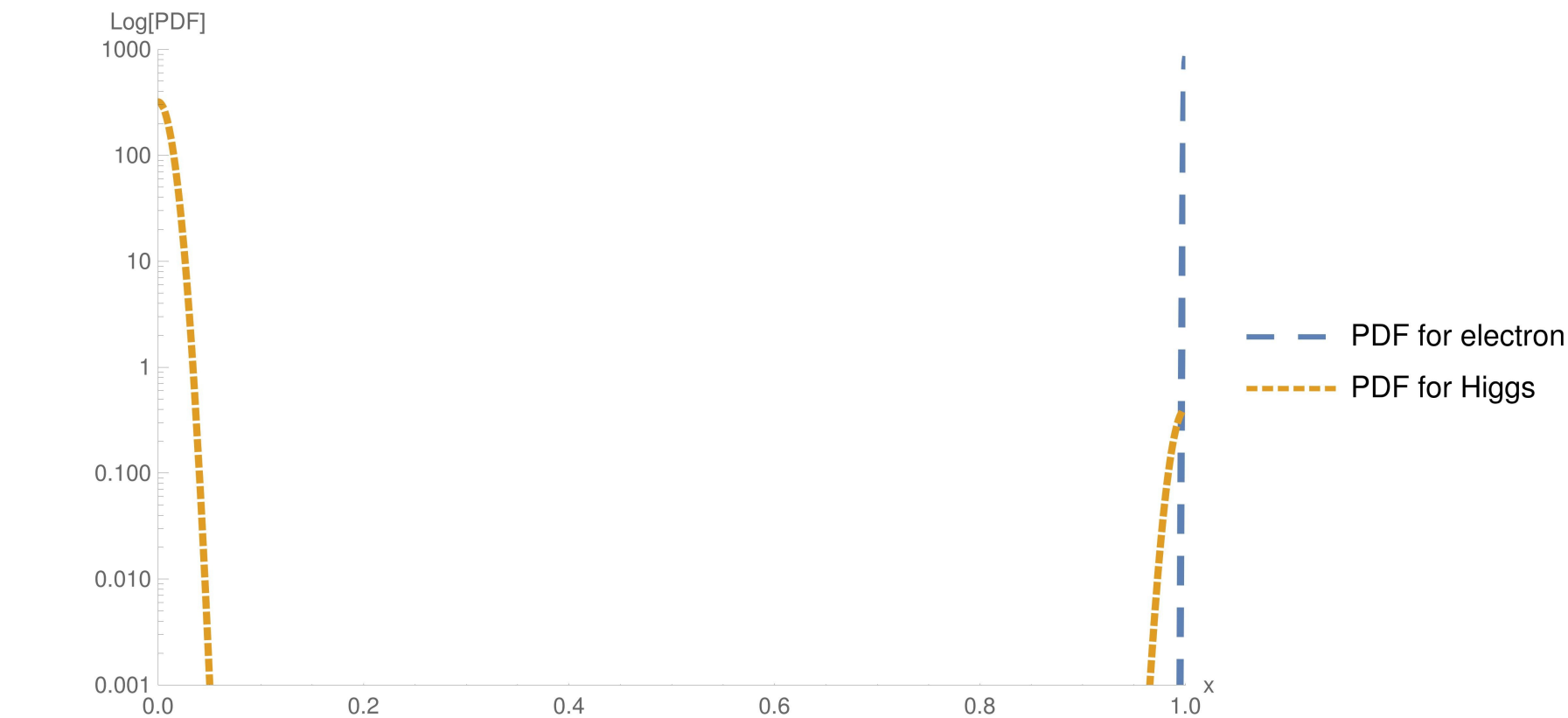
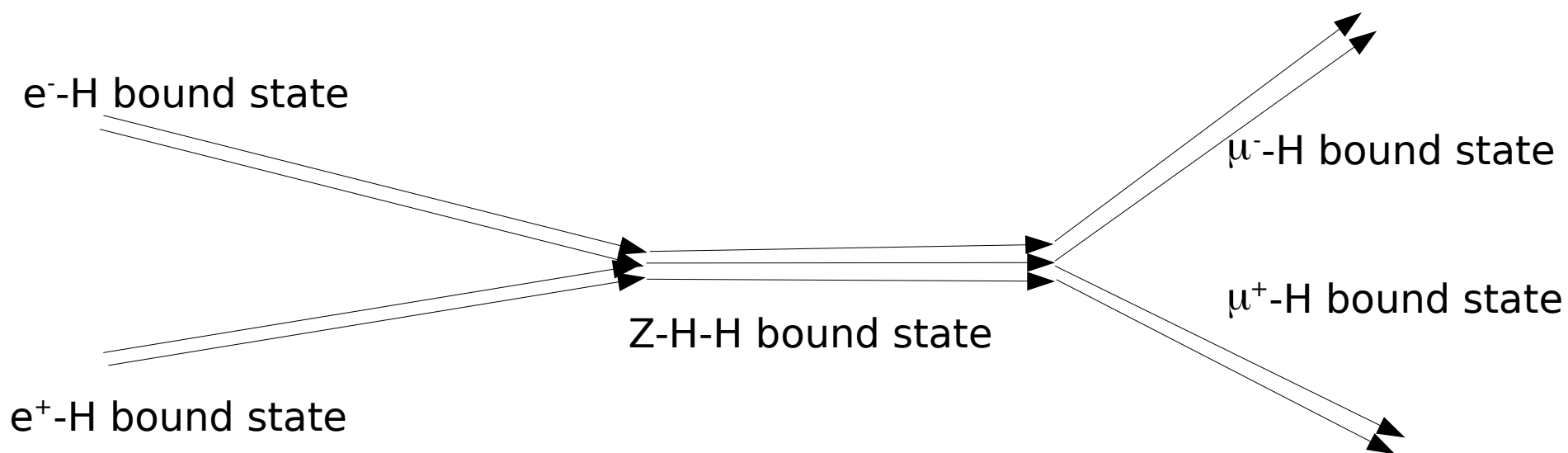
[Maas'12,
Egger et al., unpublished]



- Description of impact? PDF-type language!
- Interacting particles either electrons or Higgs
- Fragmentation 100% efficient - like for quarks

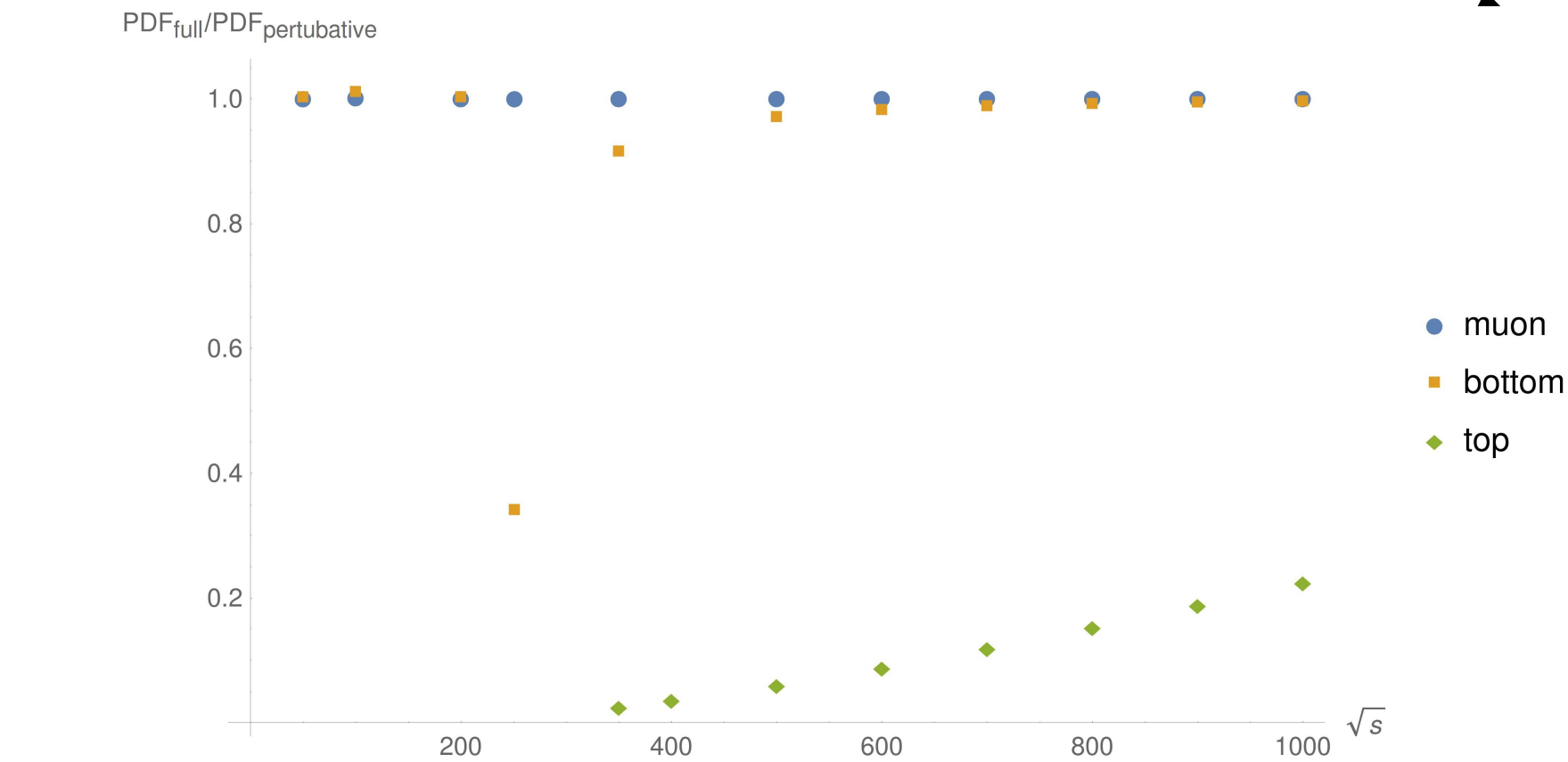
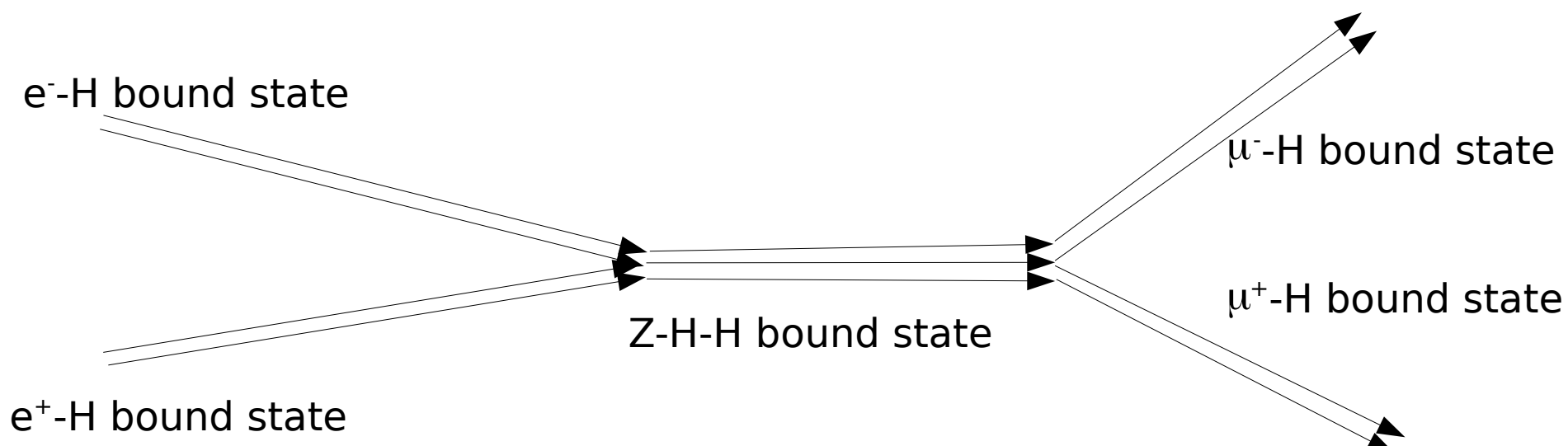
How events looks like (LEP/ILC)

[Maas'12,
Egger et al., unpublished]



How events looks like (LEP/ILC)

[Maas'12,
Egger et al., unpublished]



Why it can matter beyond the standard model

And when this can be dealt with using gauge-invariant perturbation theory

Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory

Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true?

Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]

Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]
 - Fluctuations can invalidate it

Lines of constant physics

- Lattice simulations have an intrinsic cutoff – the lattice spacing a

Lines of constant physics

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem

Lines of constant physics

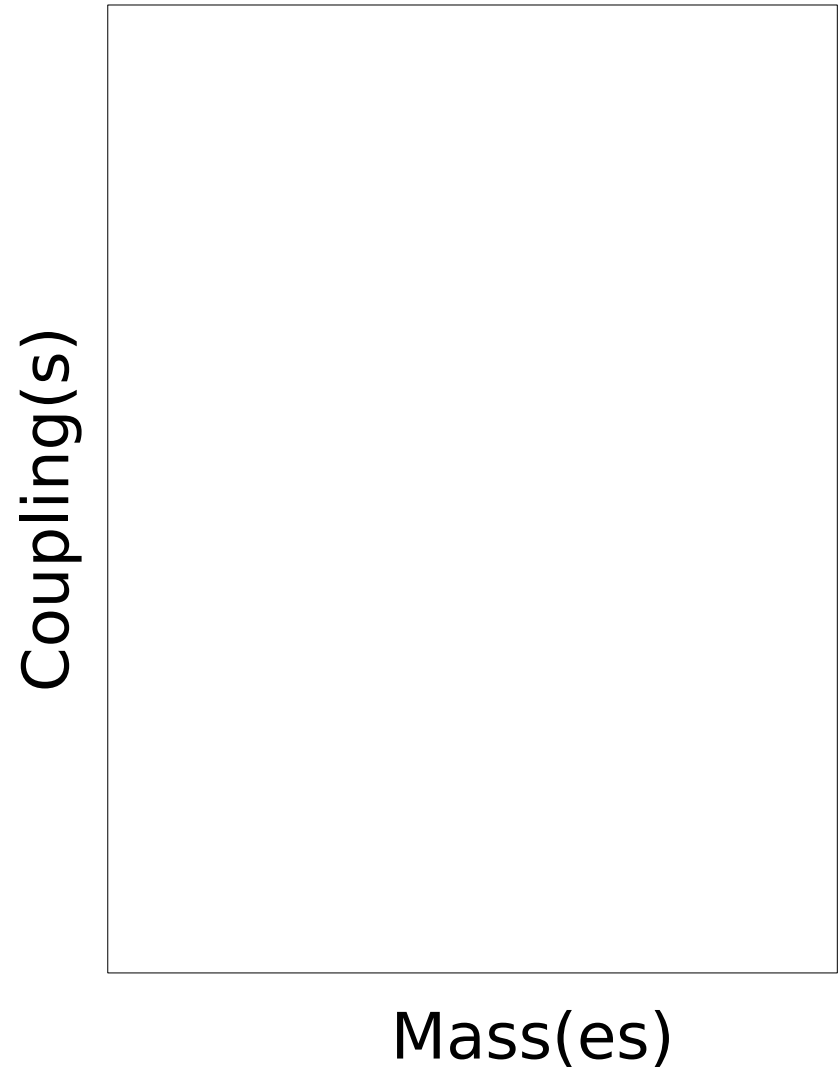
- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale

Lines of constant physics

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
Calculate for several a with all independent observables fixed
- “Lines of constant physics”

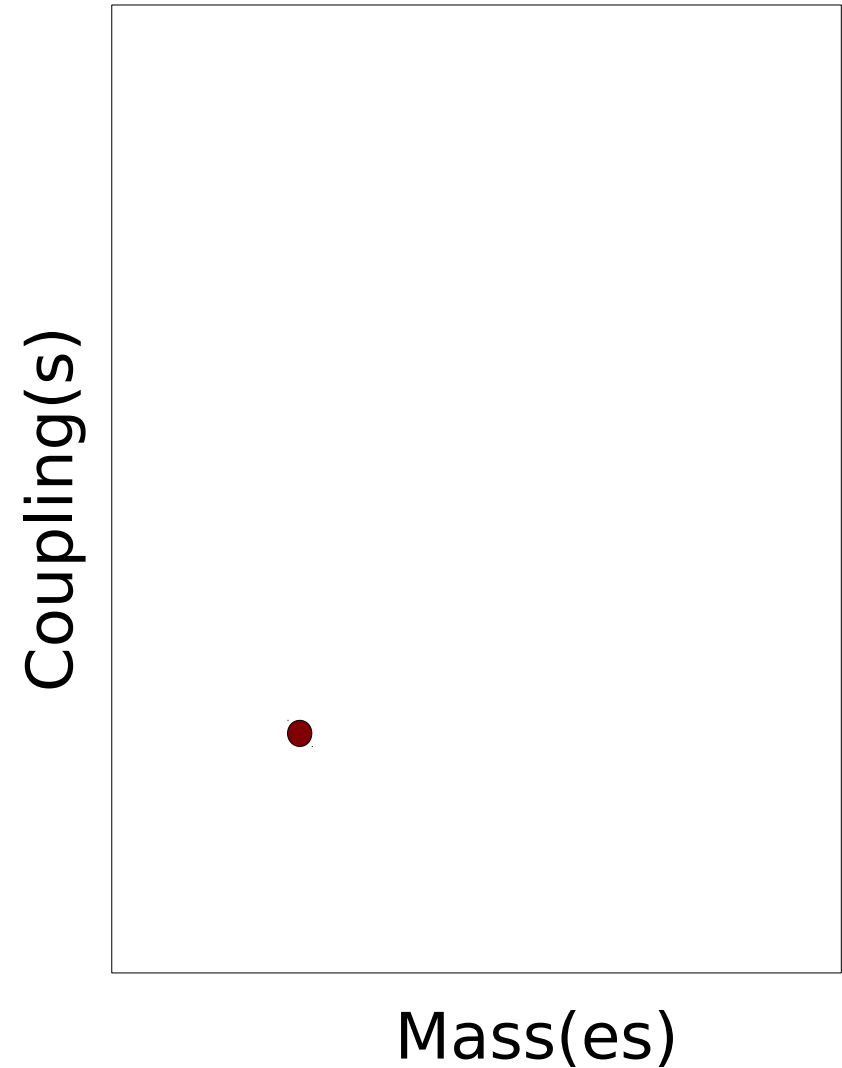
Lines of constant physics

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
Calculate for several a with all independent observables fixed
- “Lines of constant physics”



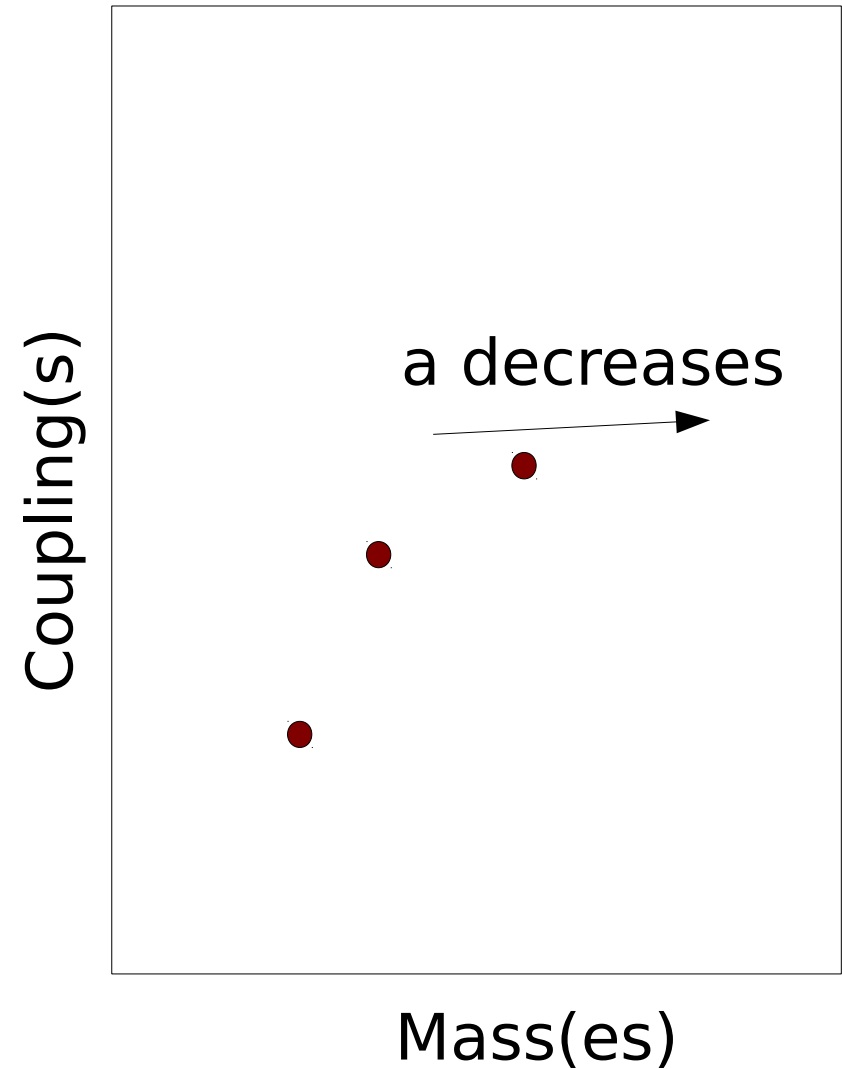
Lines of constant physics

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
Calculate for several a with all independent observables fixed
- “Lines of constant physics”



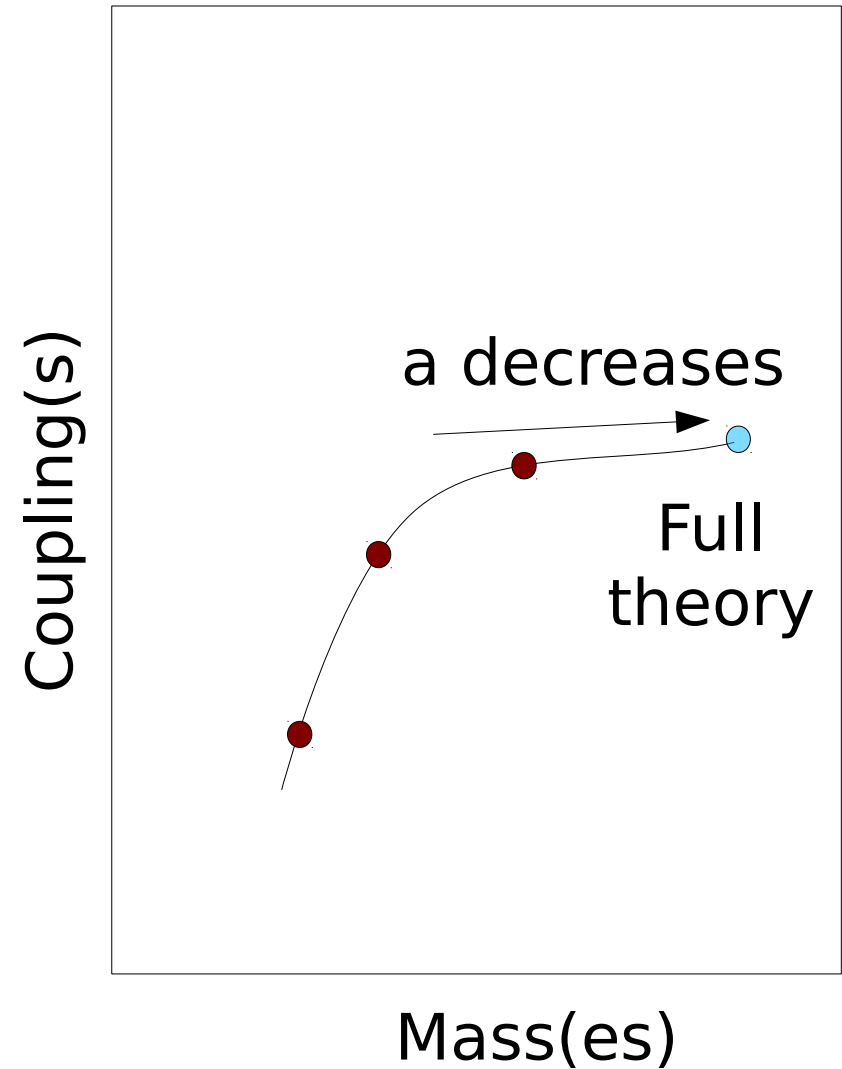
Lines of constant physics

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
Calculate for several a with all independent observables fixed
- “Lines of constant physics”



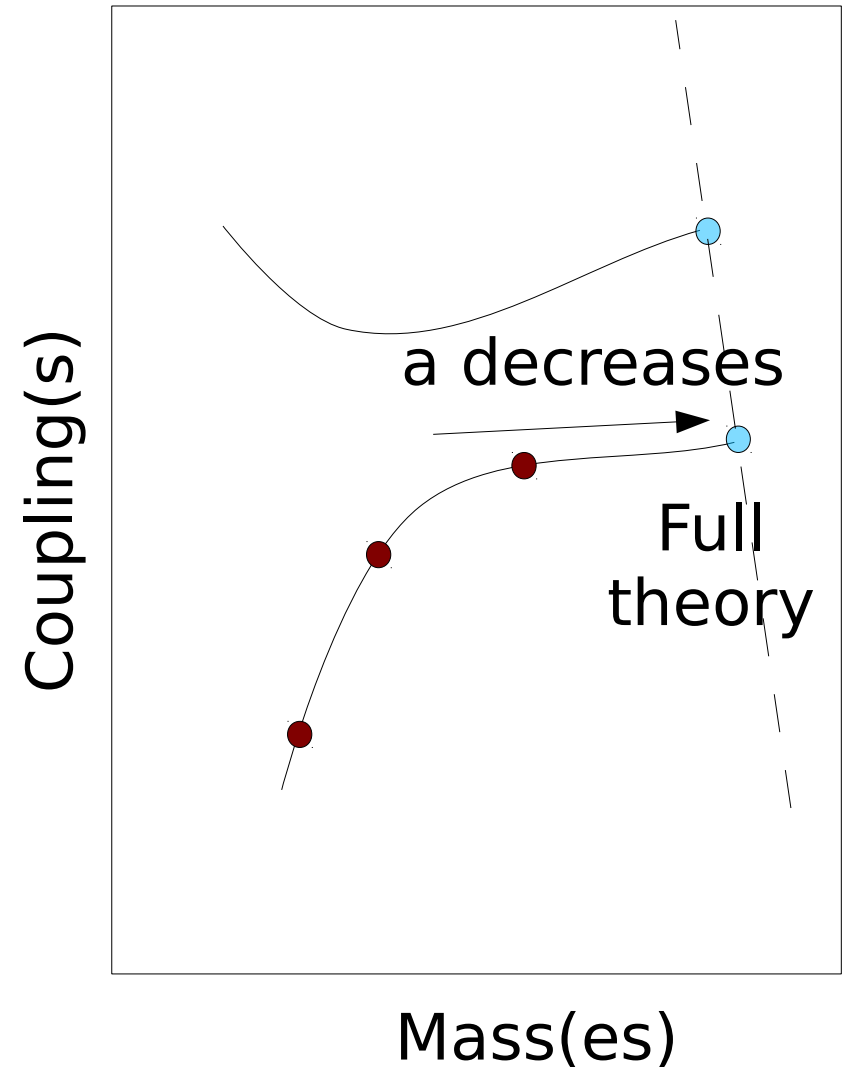
Lines of constant physics

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
Calculate for several a with all independent observables fixed
- “Lines of constant physics”



Lines of constant physics

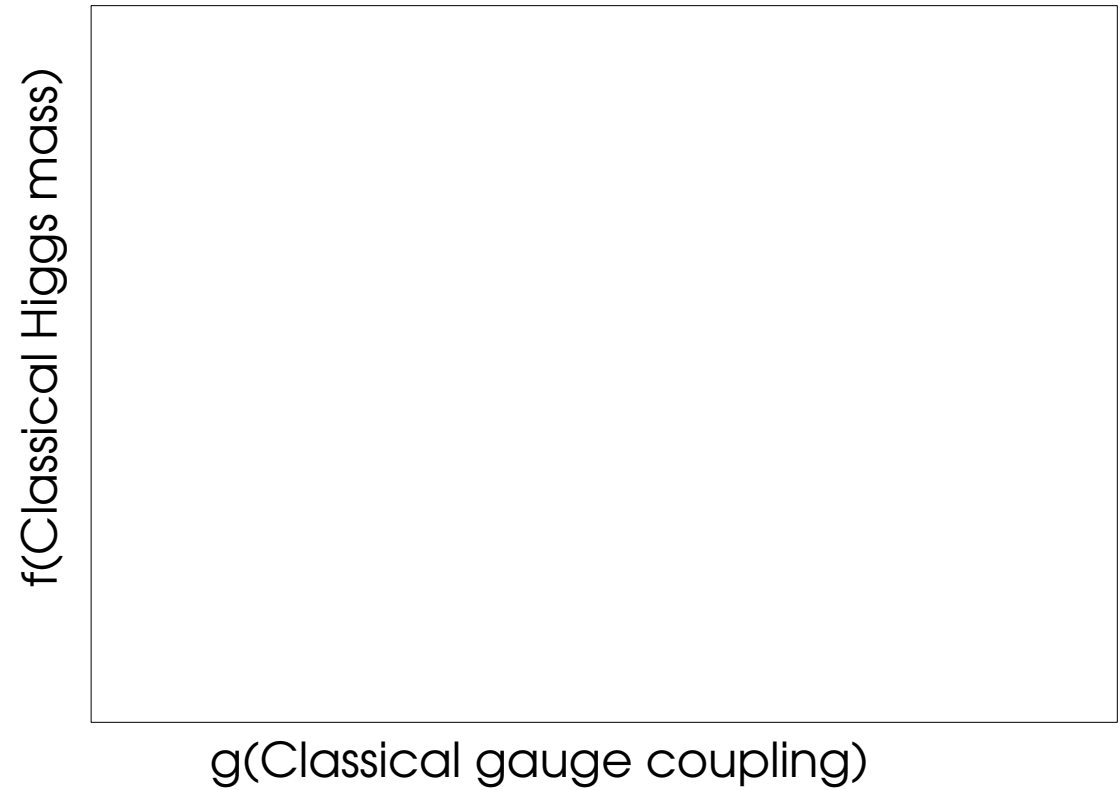
- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
Calculate for several a with all independent observables fixed
- “Lines of constant physics”
 - Different starting points yield different physics



Phase diagram

[Osterwalder et al.'78, Fradkin et al.'79
Caudy et al.'07]

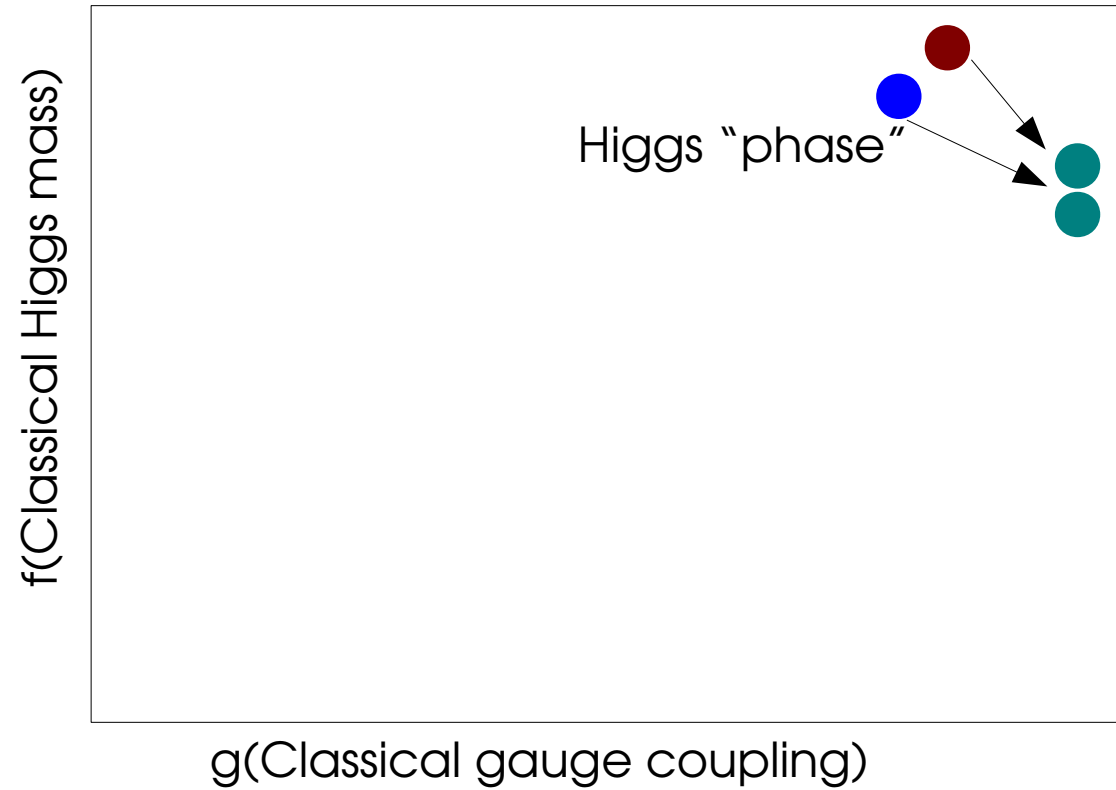
- (Lattice-regularized)
phase diagram



Phase diagram

[Osterwalder et al.'78, Fradkin et al.'79
Caudy et al.'07]

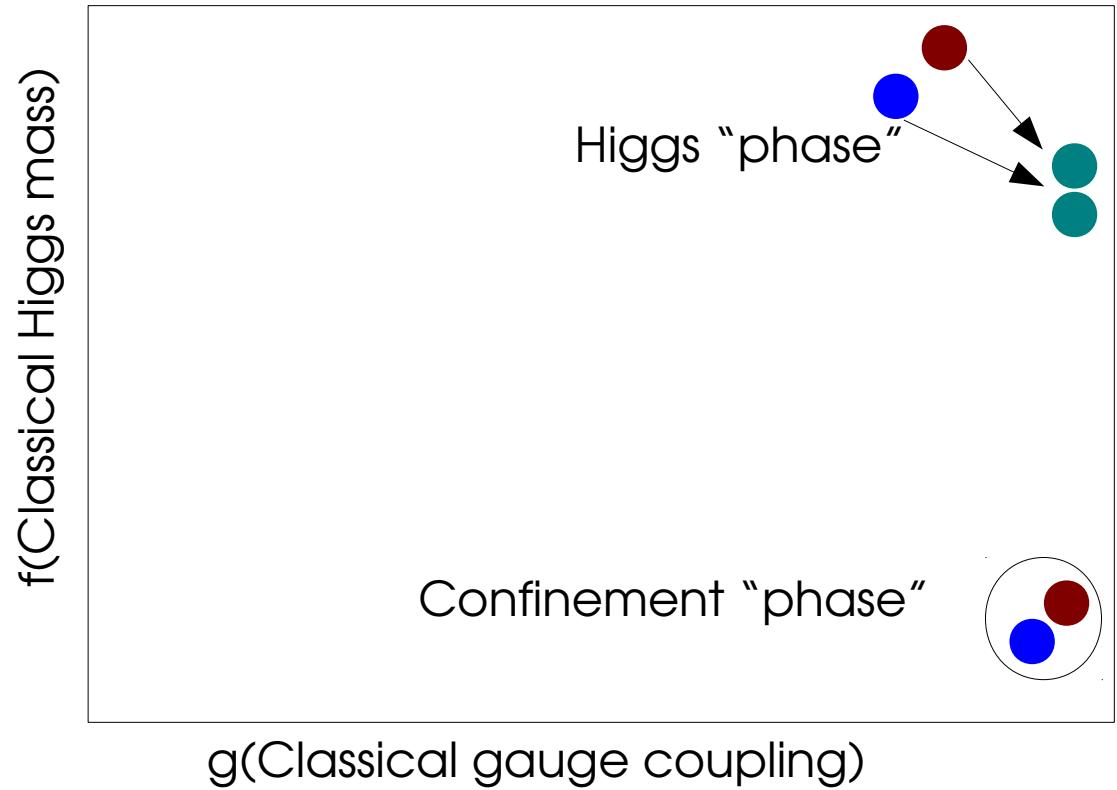
- (Lattice-regularized) phase diagram



Phase diagram

[Osterwalder et al.'78, Fradkin et al.'79
Caudy et al.'07]

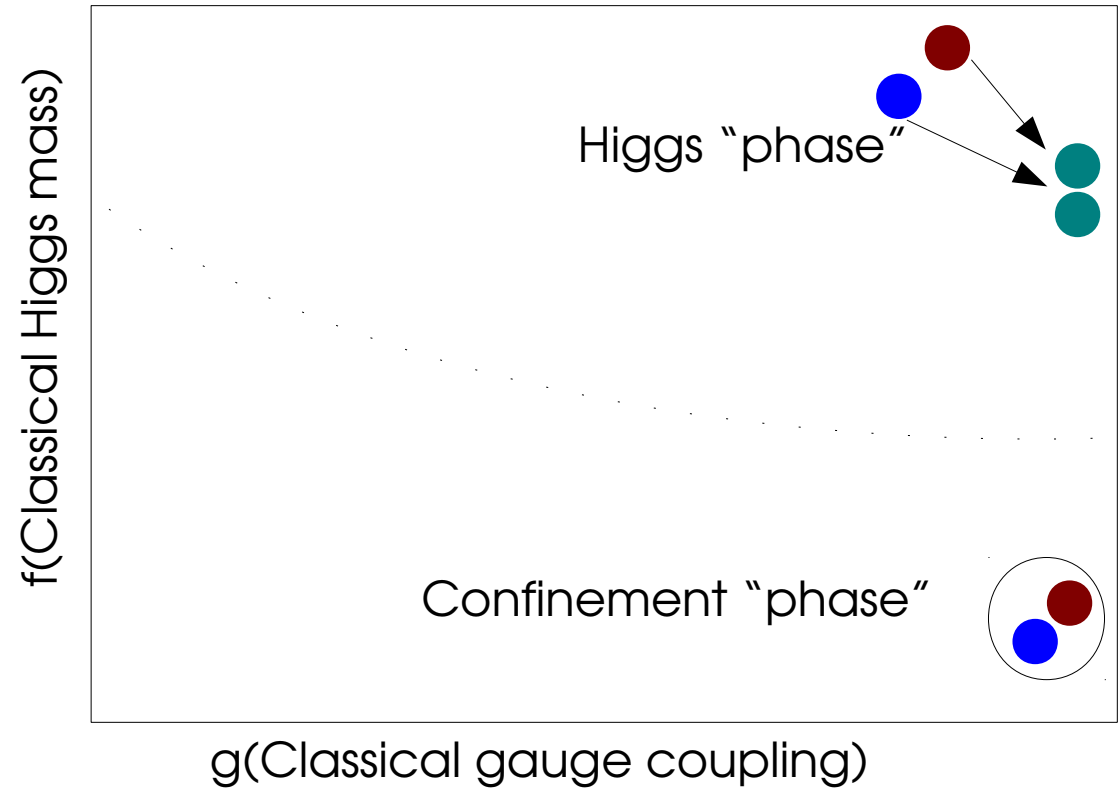
- (Lattice-regularized) phase diagram



Phase diagram

[Osterwalder et al.'78, Fradkin et al.'79
Caudy et al.'07]

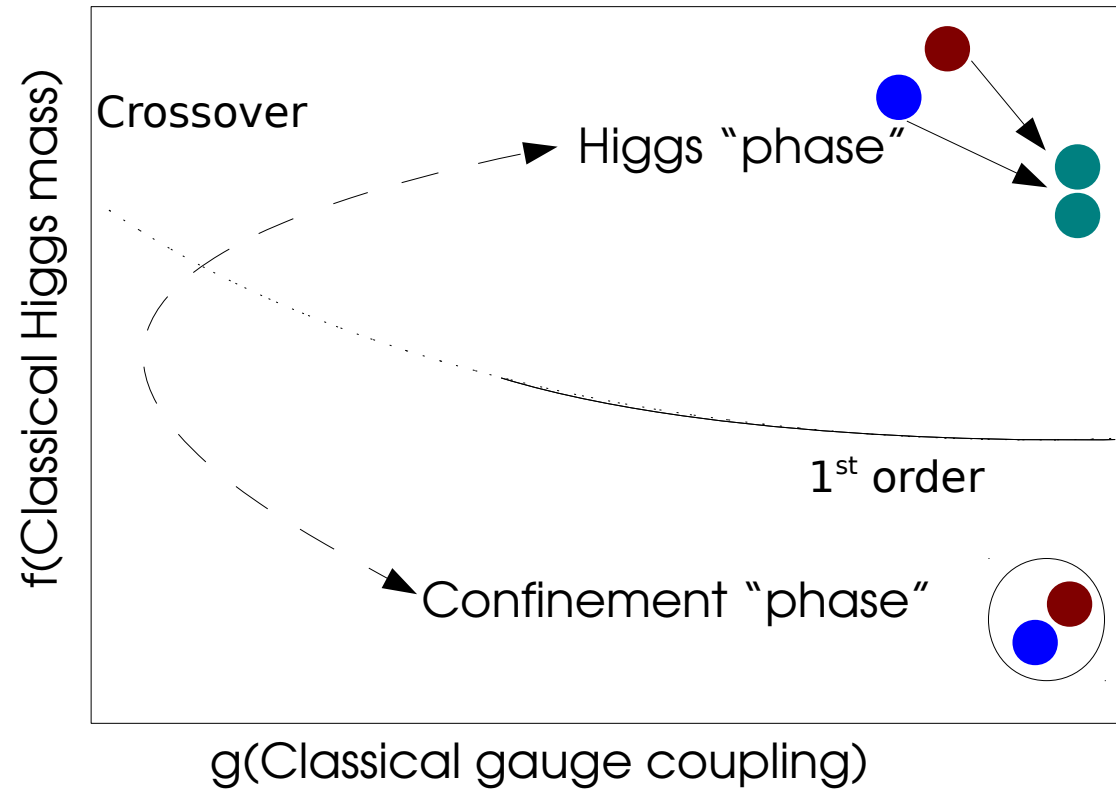
- (Lattice-regularized) phase diagram



Phase diagram

[Osterwalder et al.'78, Fradkin et al.'79
Caudy et al.'07]

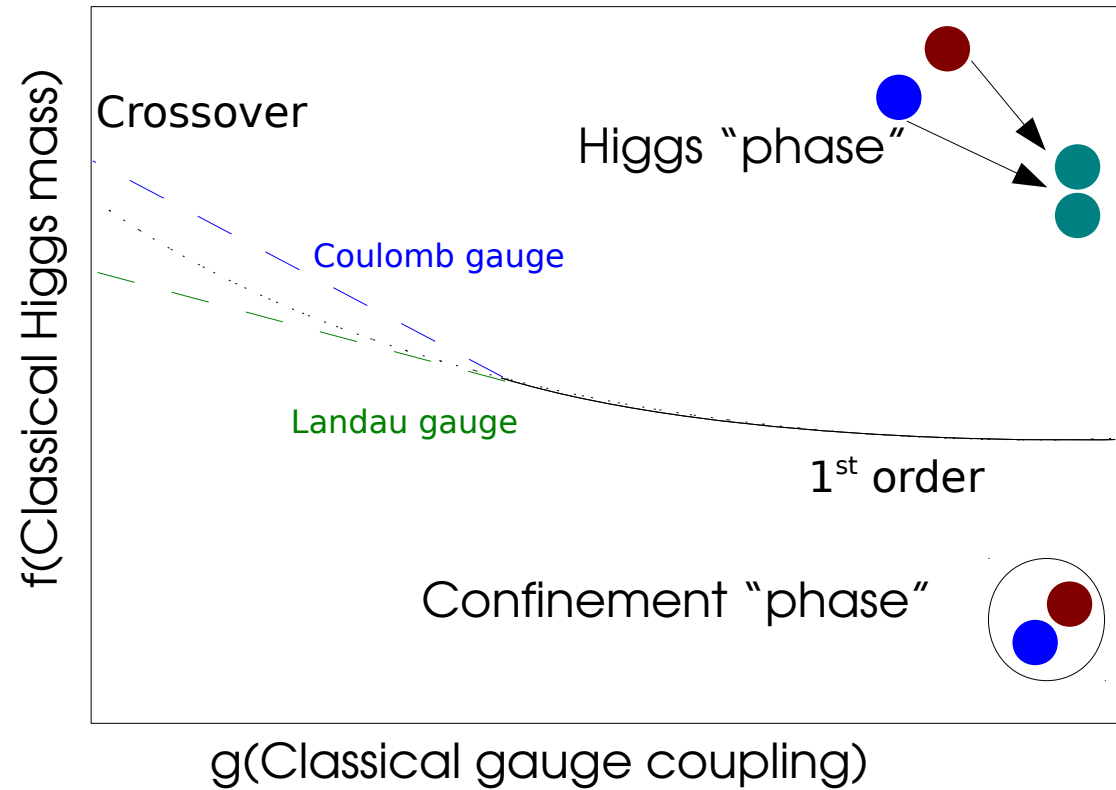
- (Lattice-regularized) phase diagram continuous



Phase diagram

[Osterwalder et al.'78, Fradkin et al.'79
Caudy et al.'07]

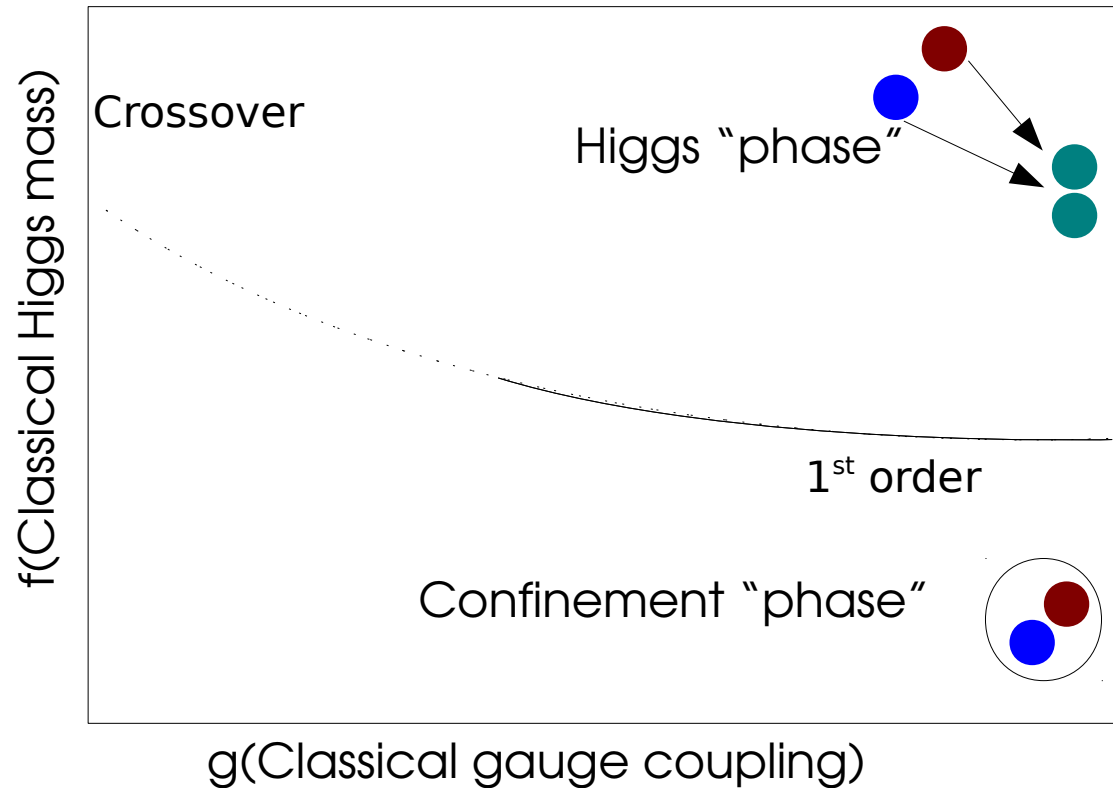
- (Lattice-regularized) phase diagram continuous
- Separation only in fixed gauges



Phase diagram

[Osterwalder et al.'78, Fradkin et al.'79
Caudy et al.'07]

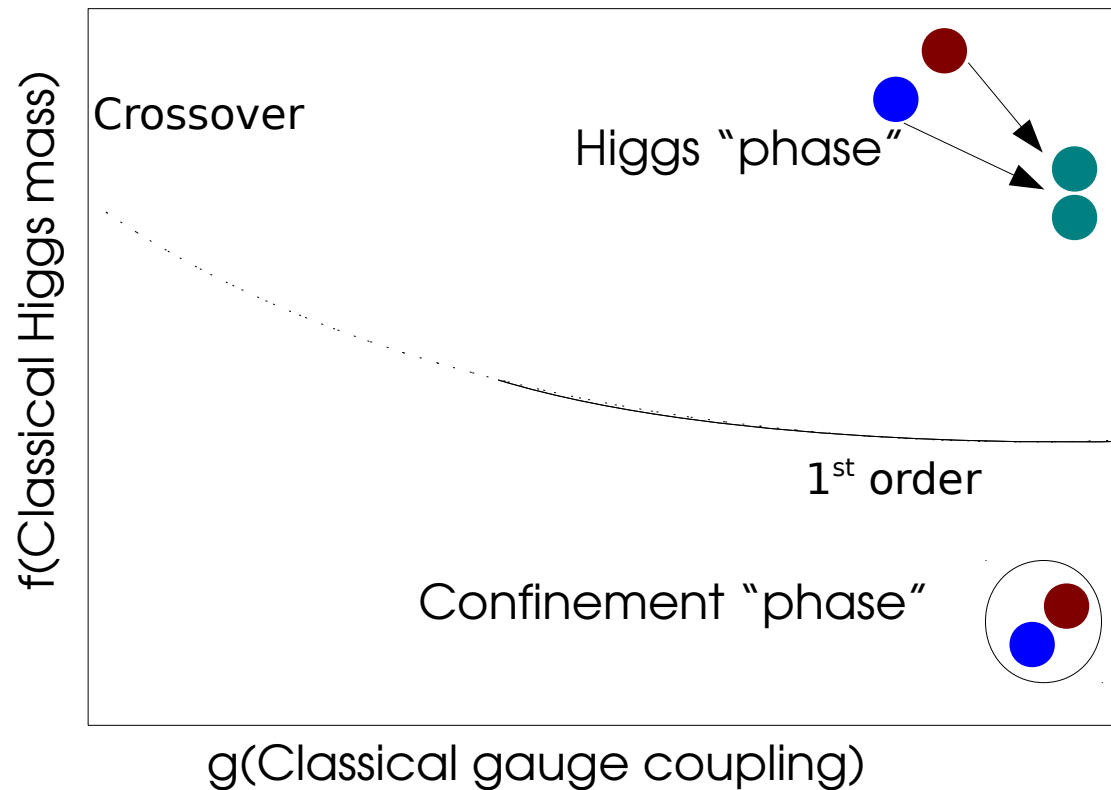
- (Lattice-regularized) phase diagram continuous
 - Separation only in fixed gauges
- Same asymptotic states in confinement and Higgs pseudo-phases
- Same asymptotic states irrespective of coupling strengths



Phase diagram

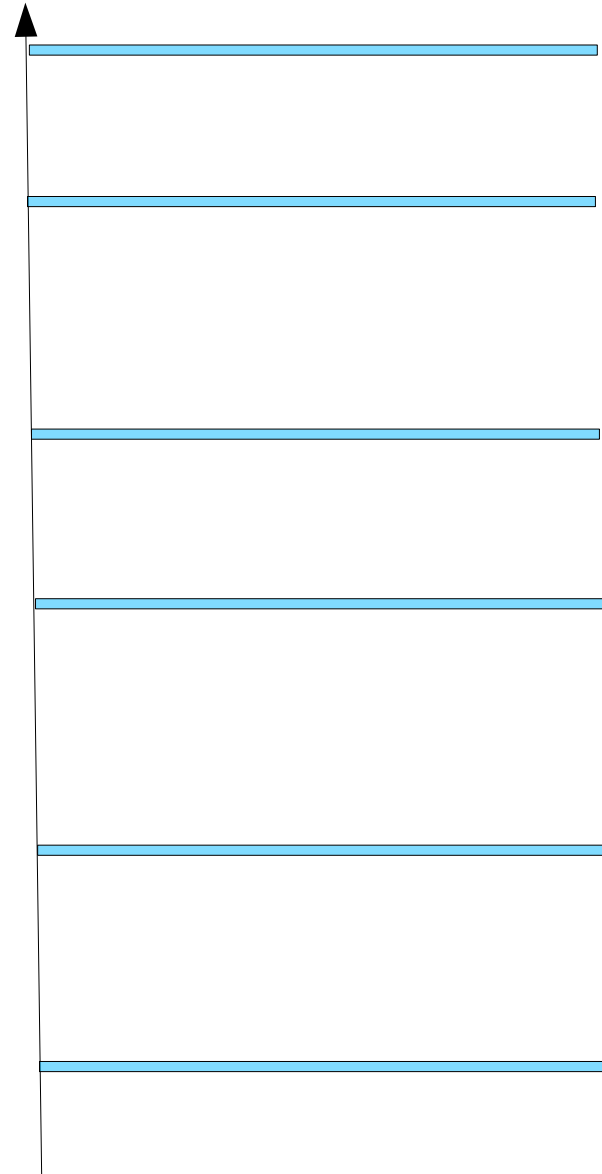
[Osterwalder et al.'78, Fradkin et al.'79
Caudy et al.'07]

- (Lattice-regularized) phase diagram continuous
 - Separation only in fixed gauges
- Same asymptotic states in confinement and Higgs pseudo-phases
- Same asymptotic states irrespective of coupling strengths
- Other states than 'Higgs' and 'W'?



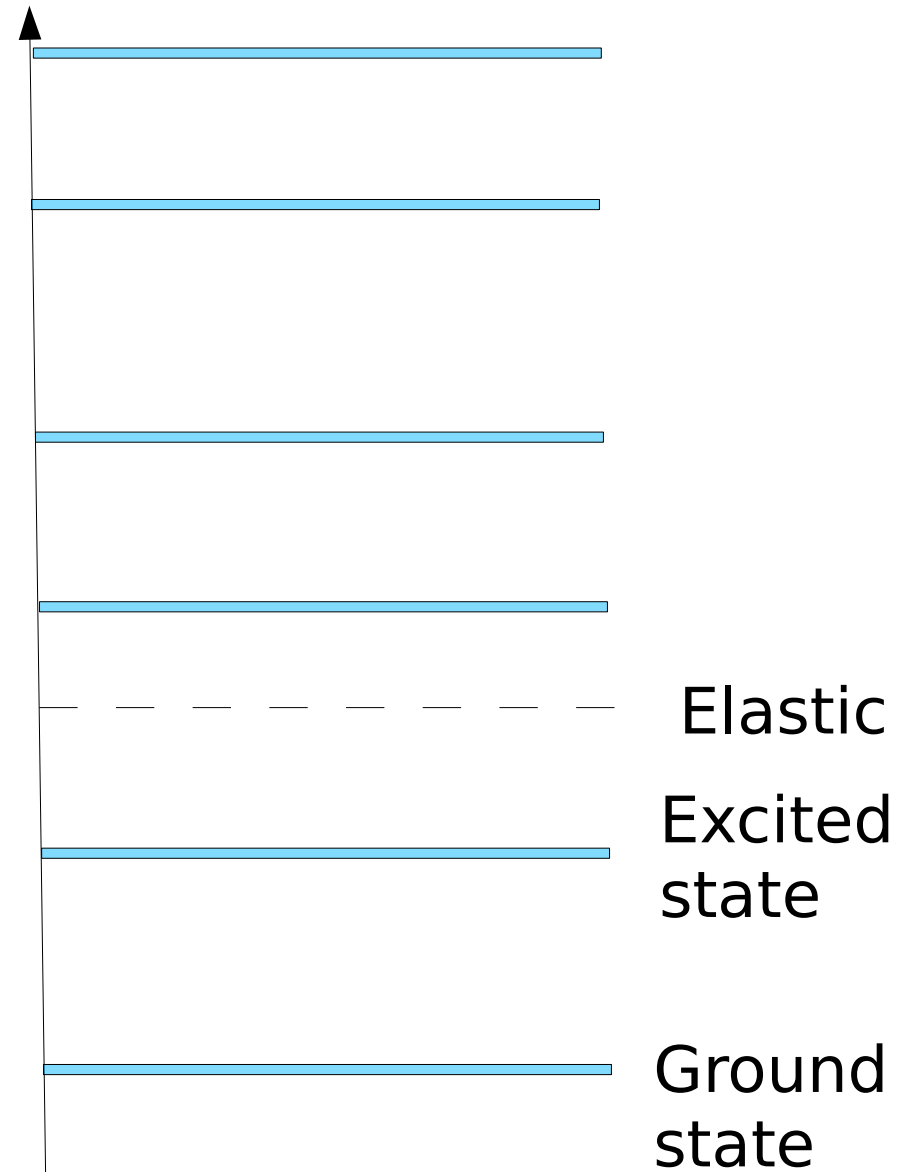
Excited states on the lattice

- Each quantum number channel has a spectrum
 - Discreet in a finite volume



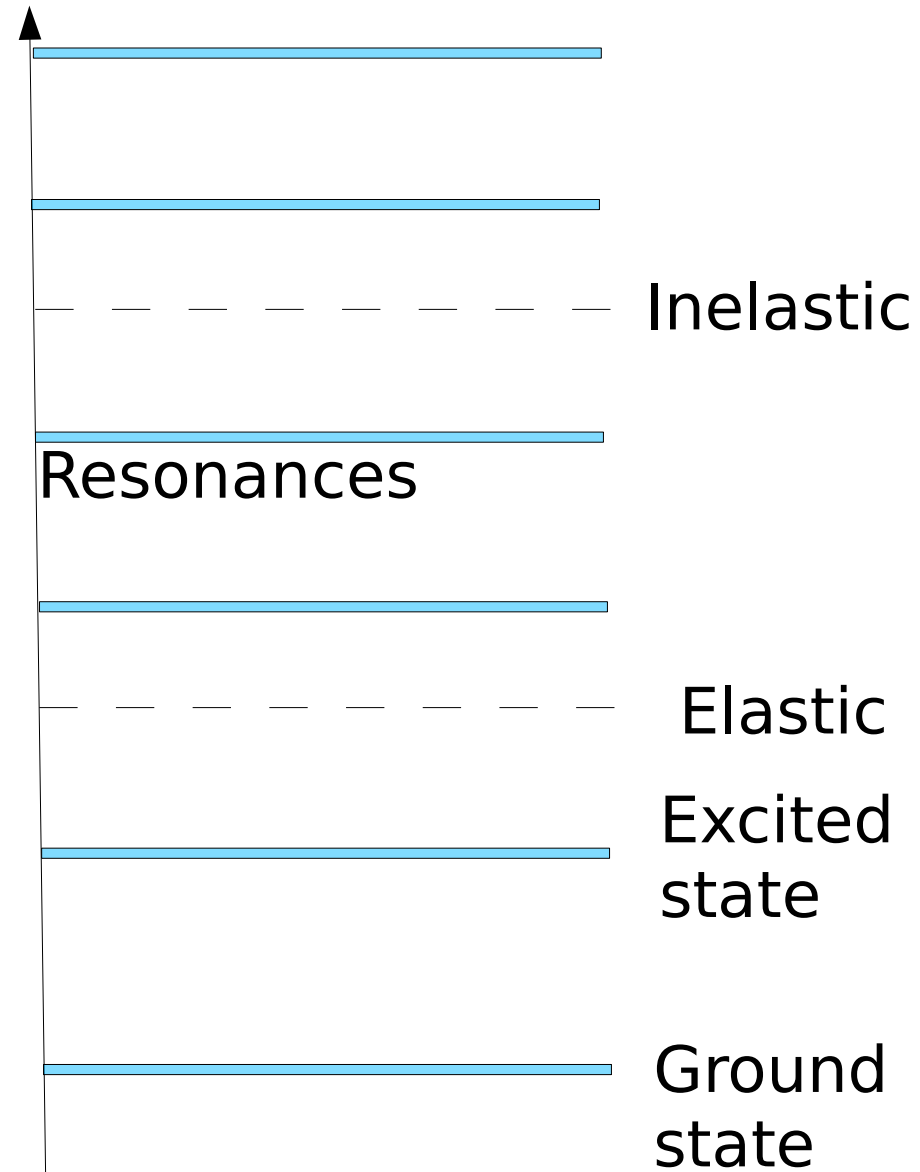
Excited states on the lattice

- Each quantum number channel has a spectrum
 - Discreet in a finite volume
- States can be either stable, excited states,



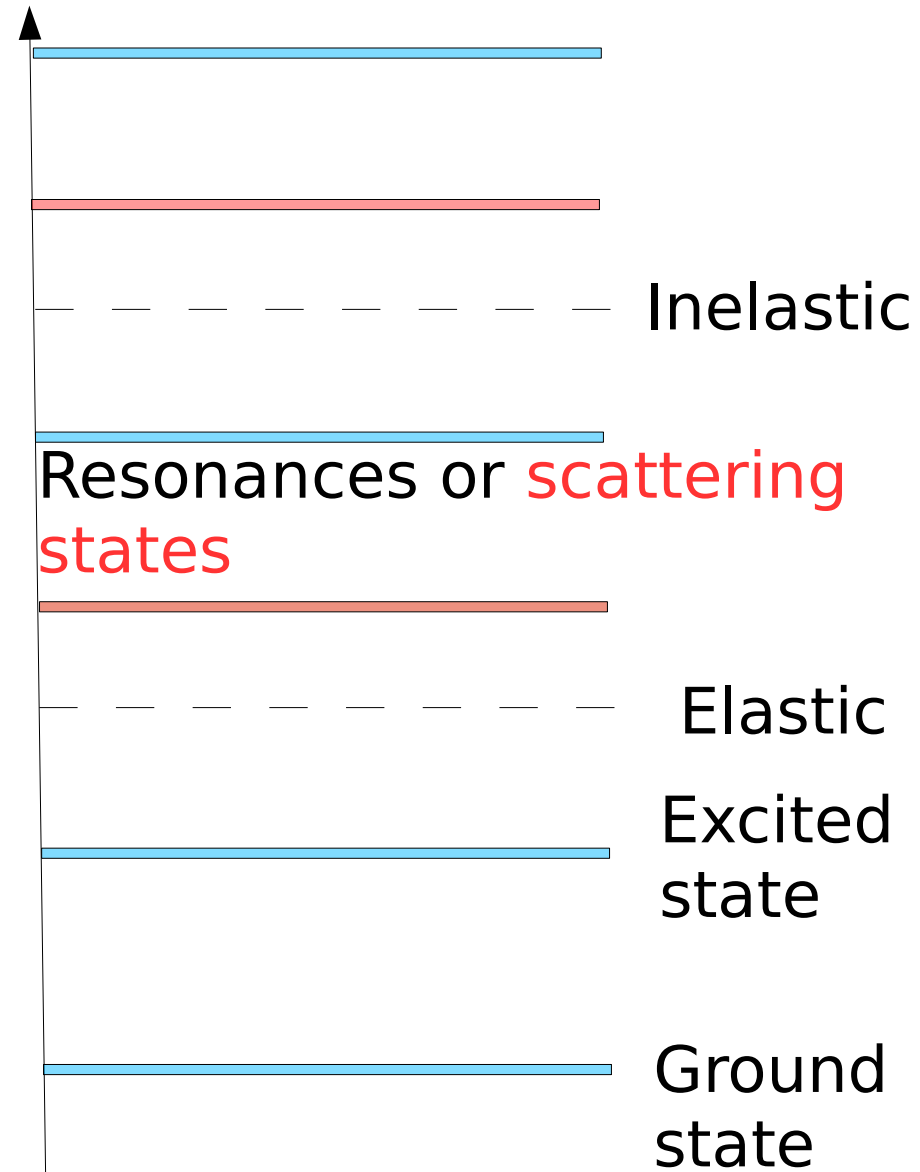
Excited states on the lattice

- Each quantum number channel has a spectrum
 - Discreet in a finite volume
- States can be either stable, excited states, resonances



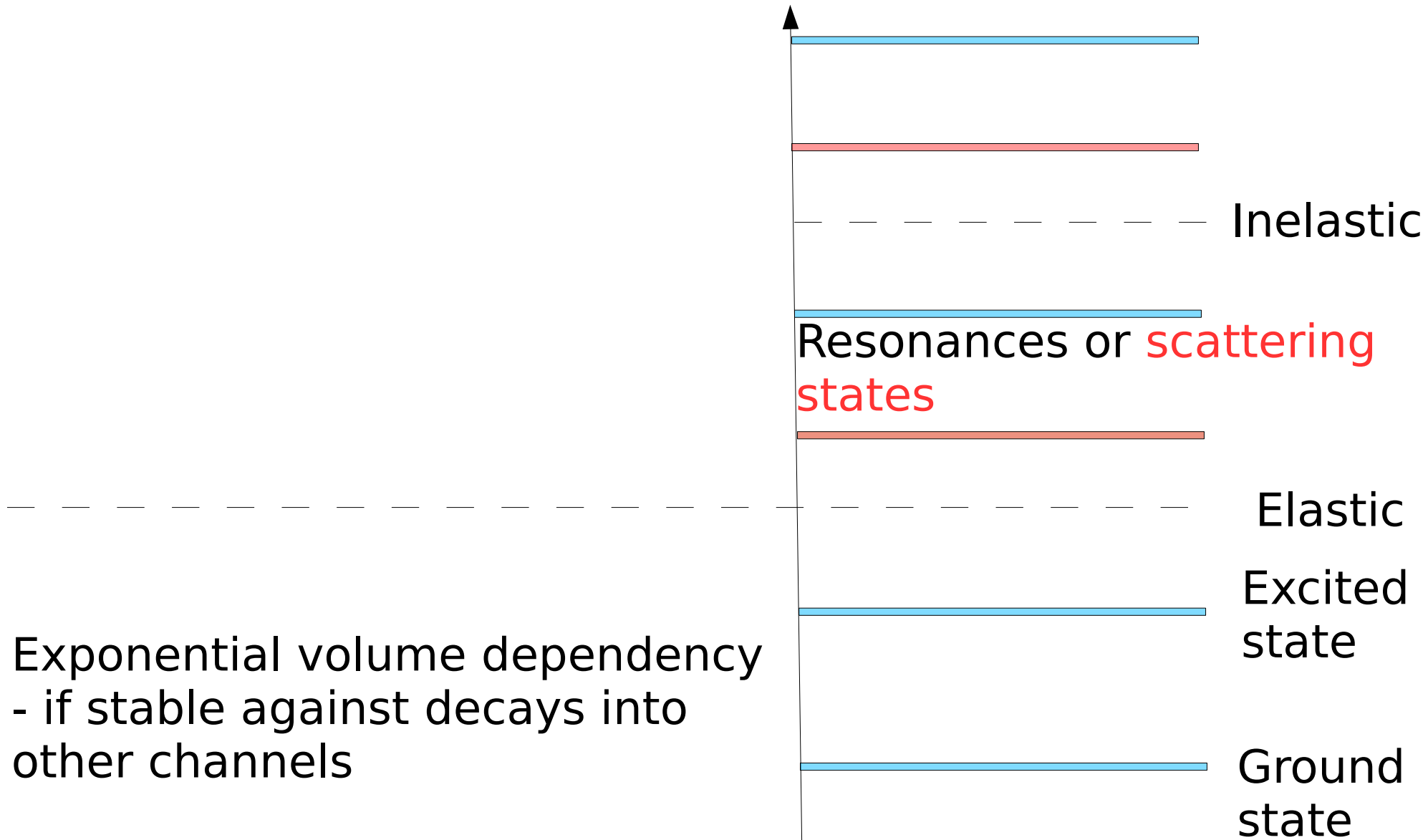
Excited states on the lattice

- Each quantum number channel has a spectrum
 - Discreet in a finite volume
- States can be either stable, excited states, resonances or scattering states



Excited states on the lattice

[Luescher'85,'86,'90,'91]

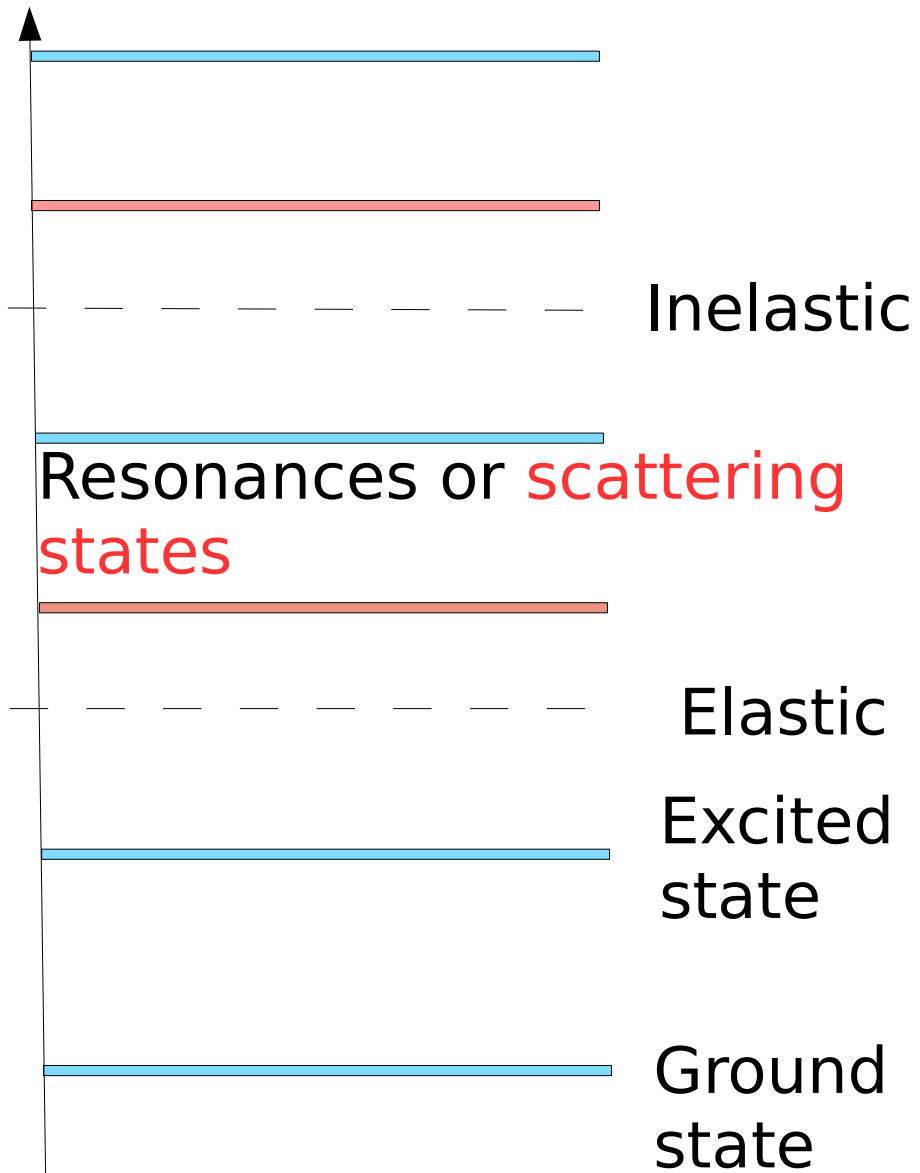


Excited states on the lattice

[Luescher'85,'86,'90,'91]

- Polynominal (inverse) volume dependence
- Width and nature from phase shifts below the inelastic threshold

Exponential volume dependency
- if stable against decays into other channels



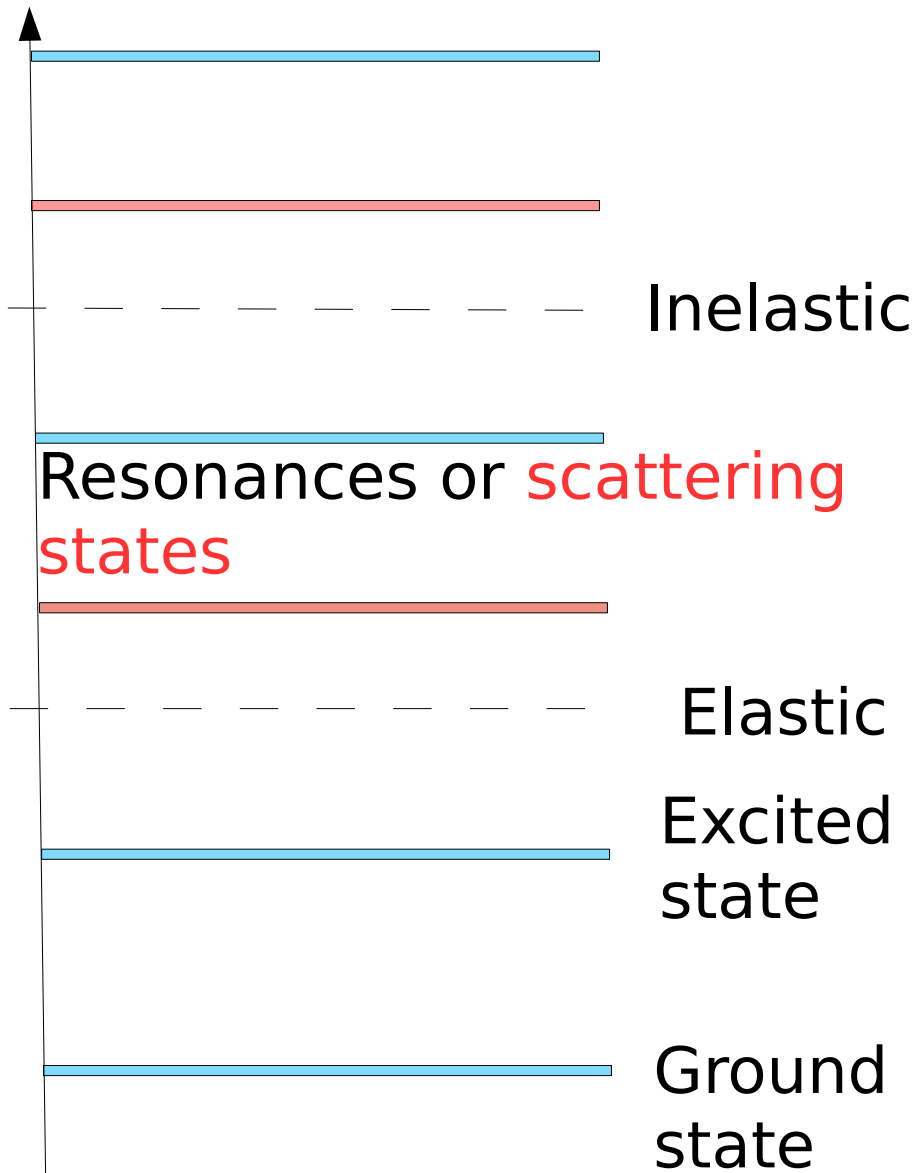
Excited states on the lattice

[Luescher'85,'86,'90,'91]

Above inelastic threshold still complicated

- Polynominal (inverse) volume dependence
- Width and nature from phase shifts below the inelastic threshold

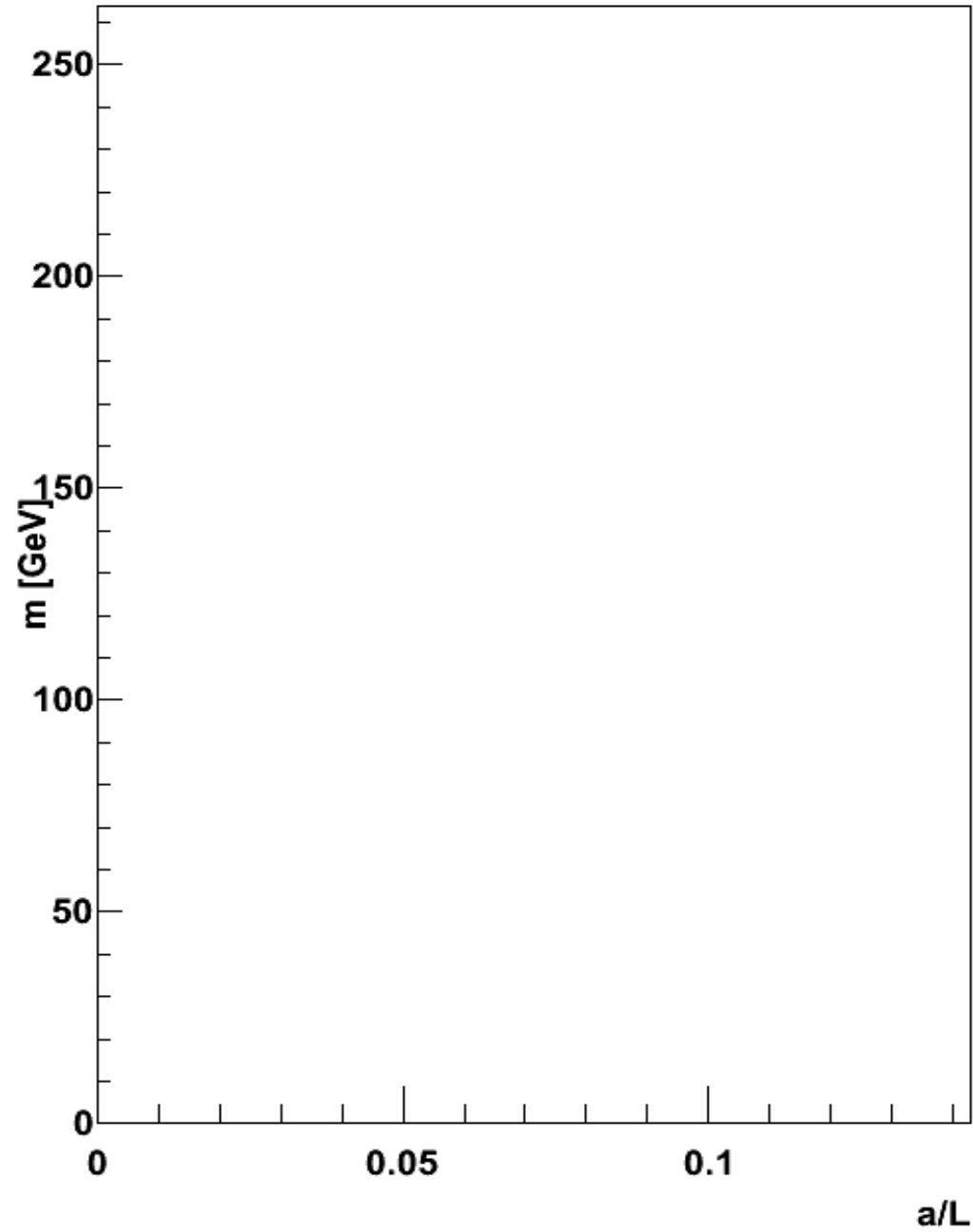
Exponential volume dependency
- if stable against decays into other channels



Excited states on the lattice

[Luescher'85,'86,'90,'91]

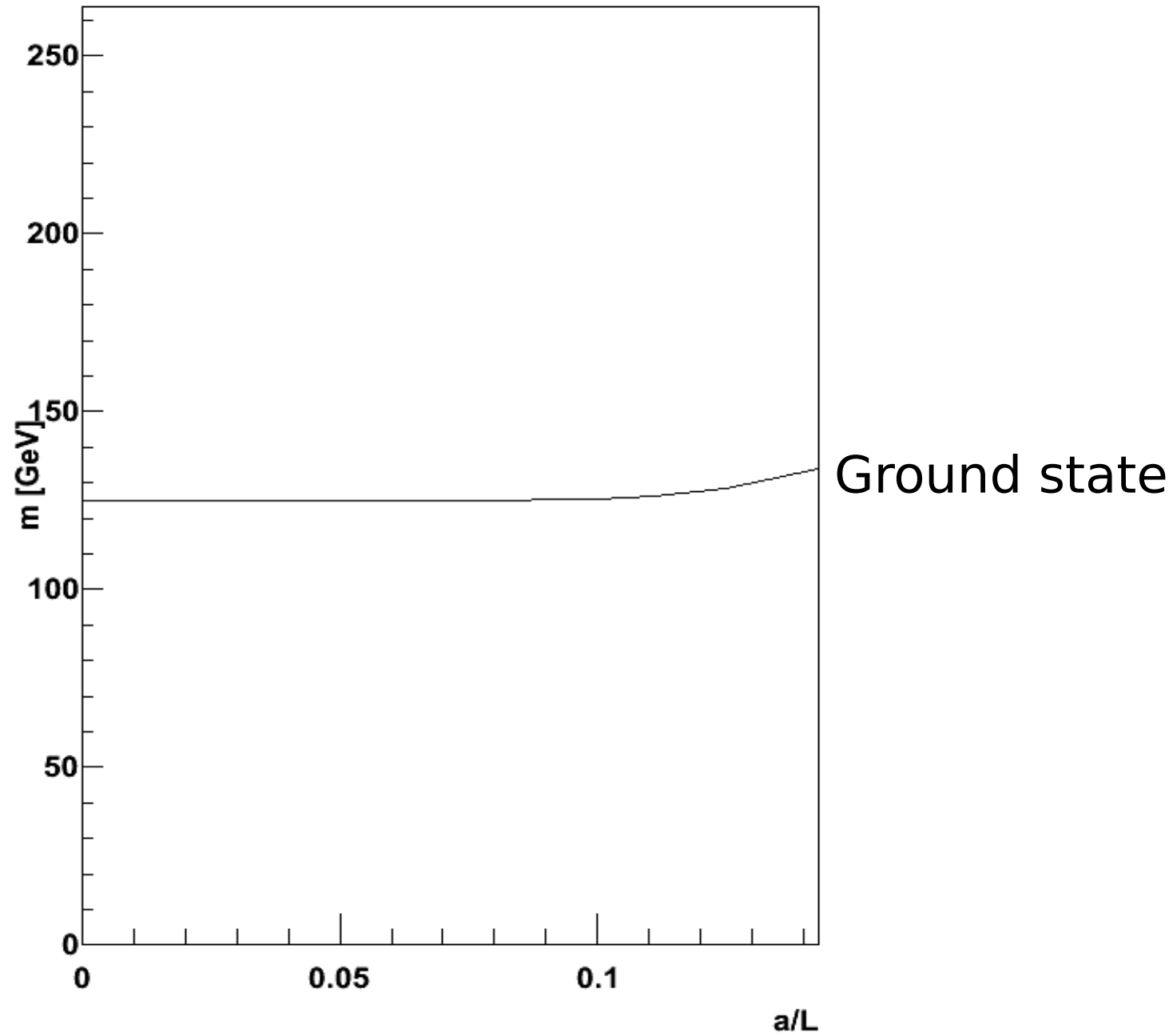
Spectrum



Excited states on the lattice

[Luescher'85,'86,'90,'91]

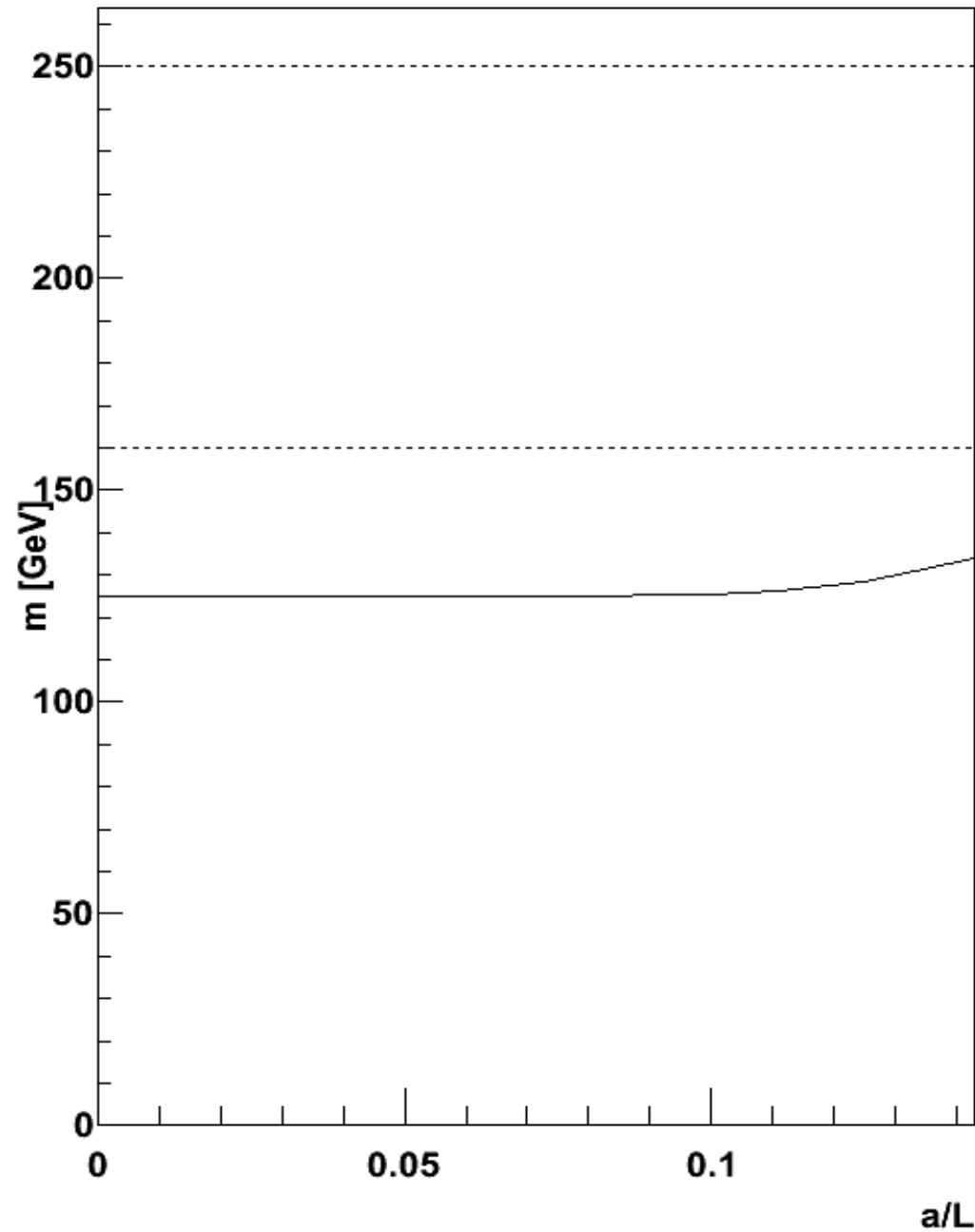
Spectrum



Excited states on the lattice

[Luescher'85,'86,'90,'91]

Spectrum



Inelastic threshold: $H \rightarrow 2H$

Elastic threshold: $H \rightarrow 2W$

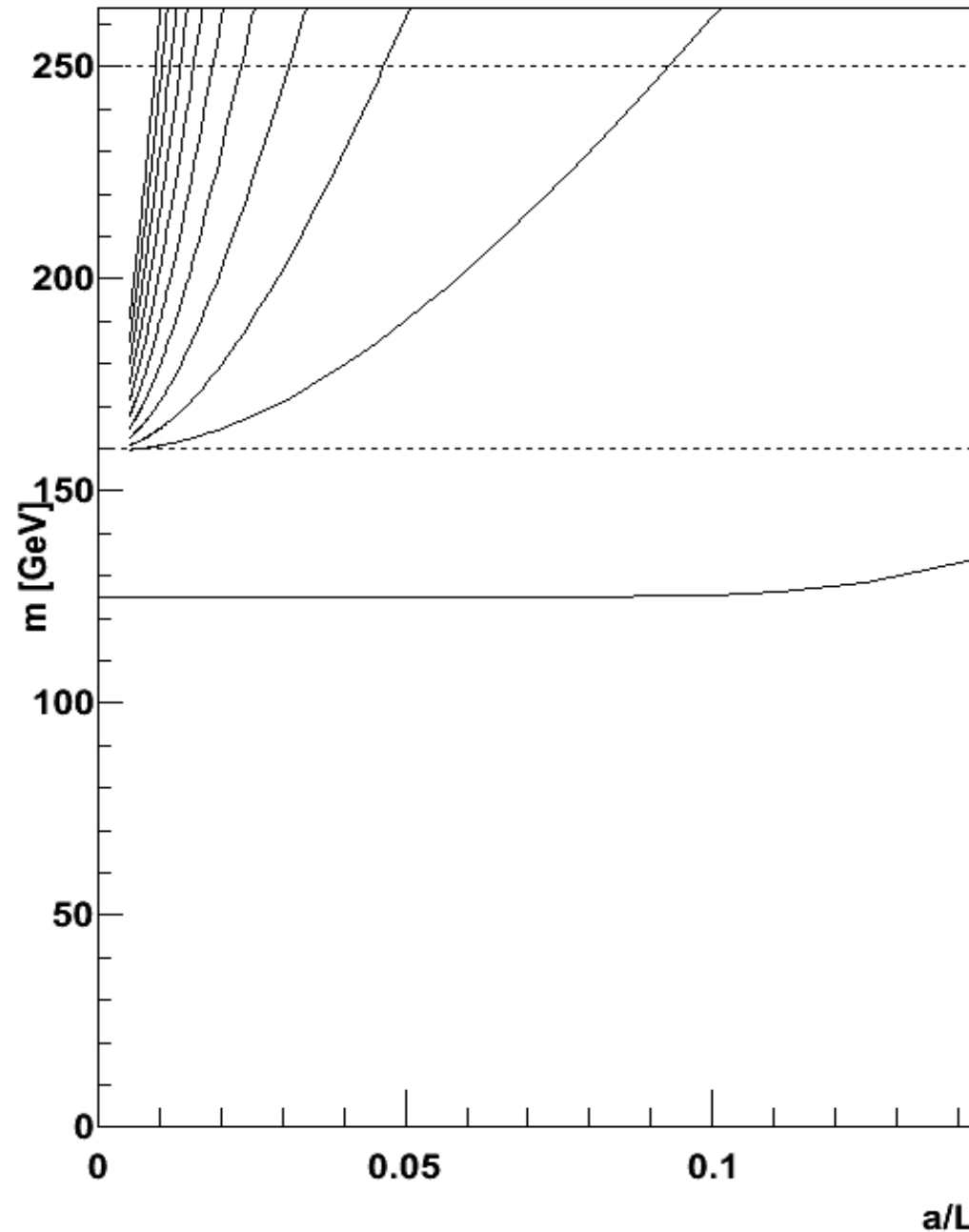
Ground state

Excited states on the lattice

[Luescher'85,'86,'90,'91]

Spectrum

Scattering states



Inelastic threshold: $H \rightarrow 2H$

Elastic threshold: $H \rightarrow 2W$

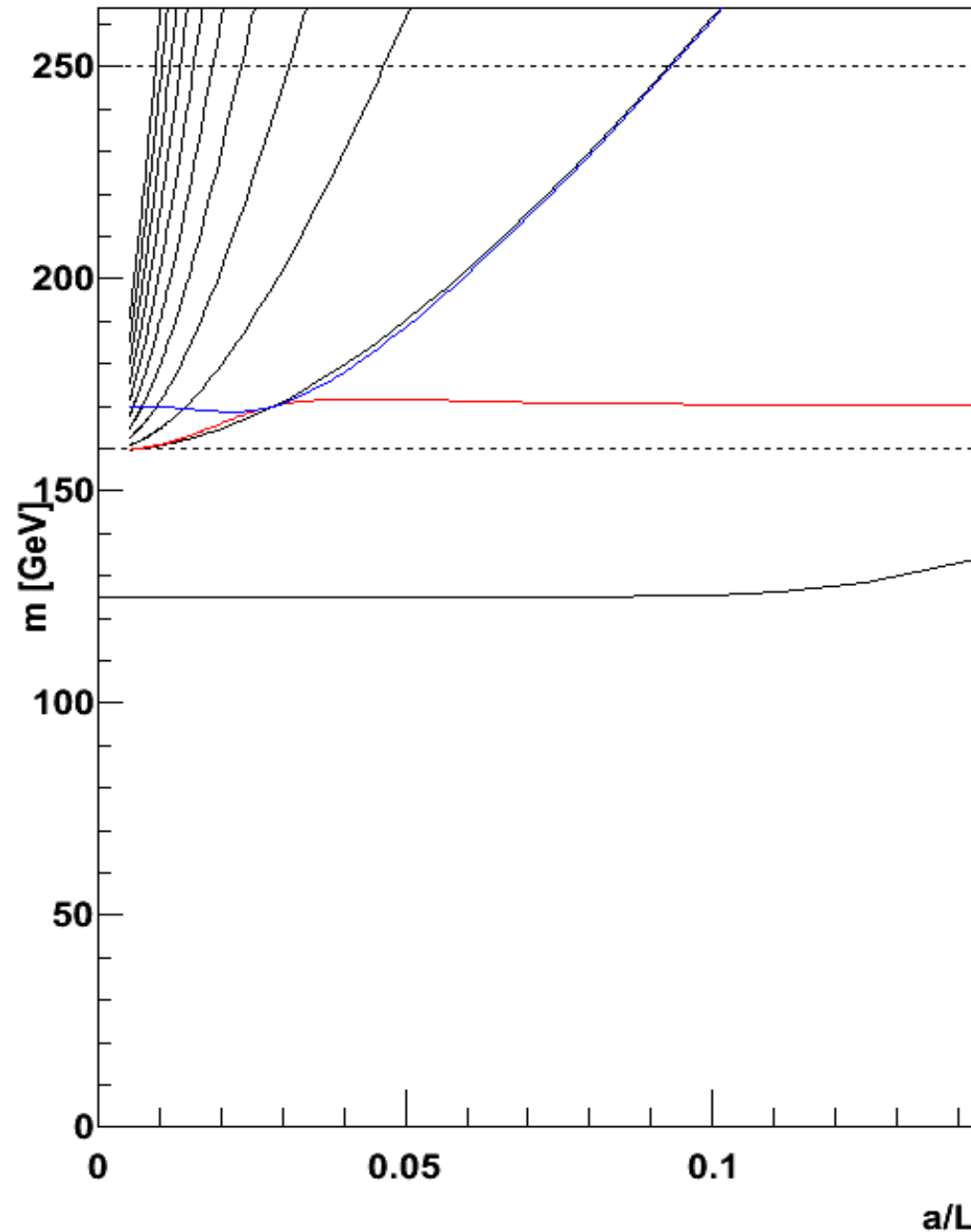
Ground state

Excited states on the lattice

[Luescher'85,'86,'90,'91]

Spectrum

Scattering states



Inelastic threshold: $H \rightarrow 2H$

Avoided level crossing
Identification and widths from
phase shifts

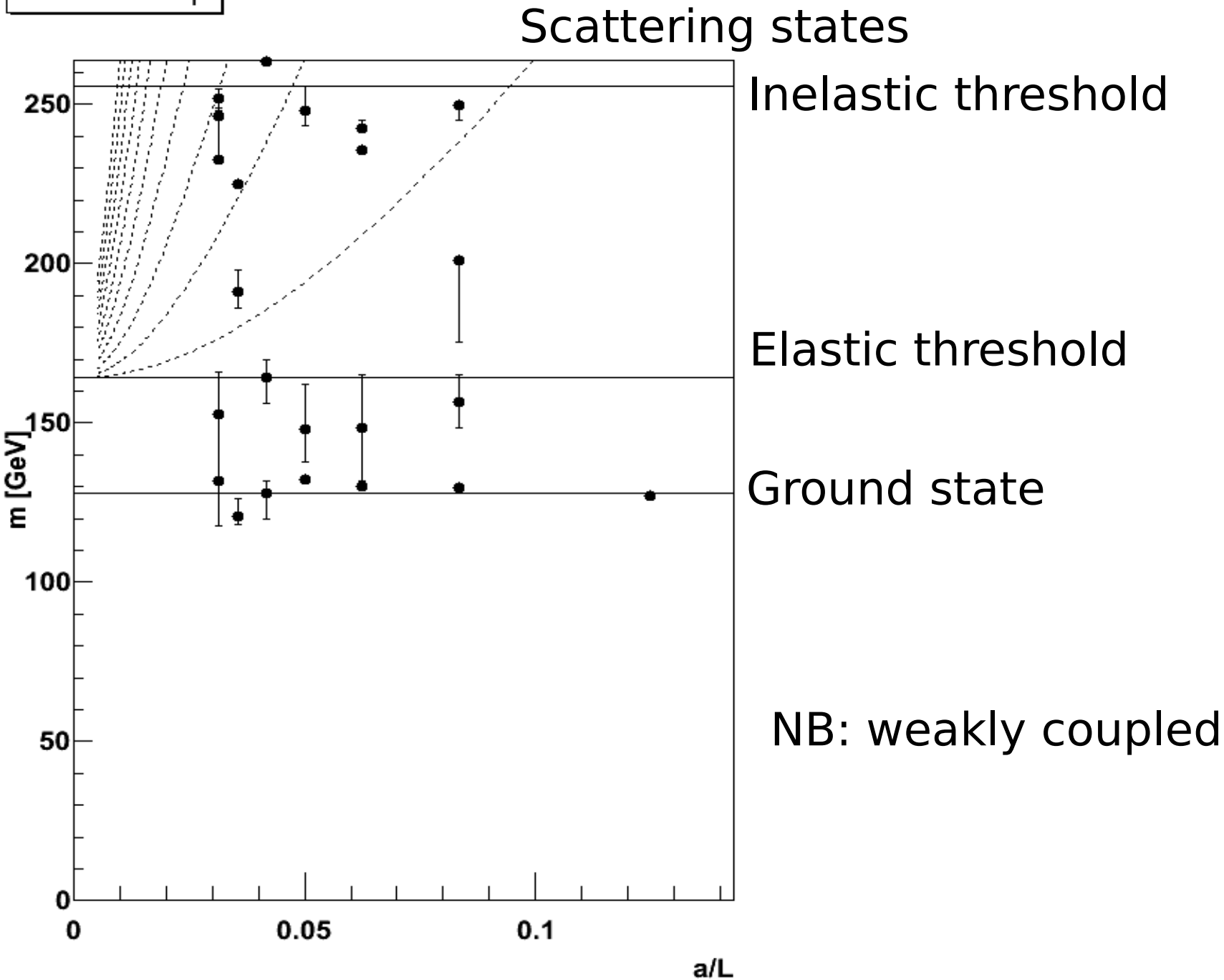
Elastic threshold: $H \rightarrow 2W$

Ground state

Search: Excited Higgs

[Maas et al.'14]

Spectrum 0_1^+



Search: Z'

[Maas et al.'14]

Spectrum 1_3^-

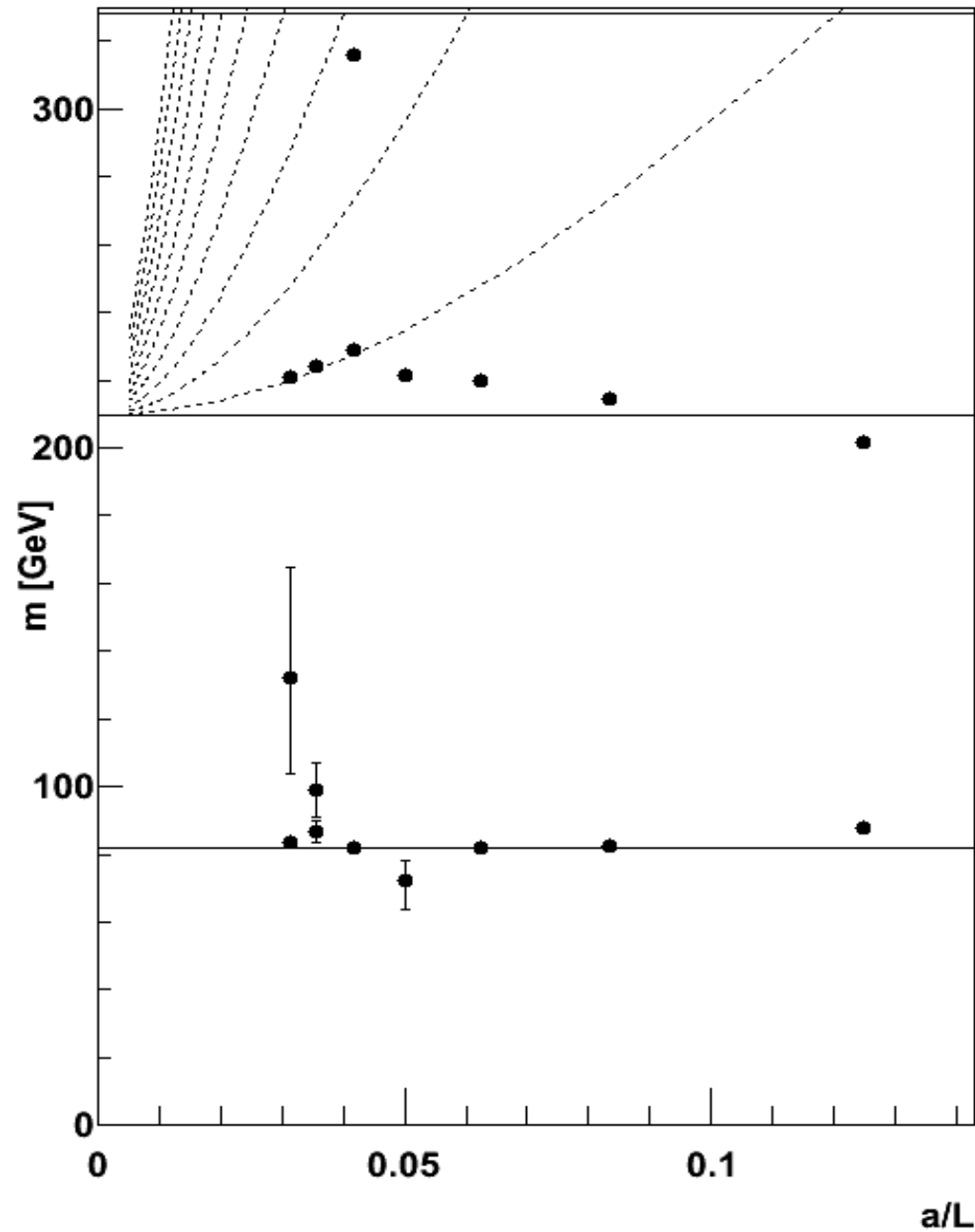
Scattering states

Inelastic threshold

Elastic threshold

Ground state

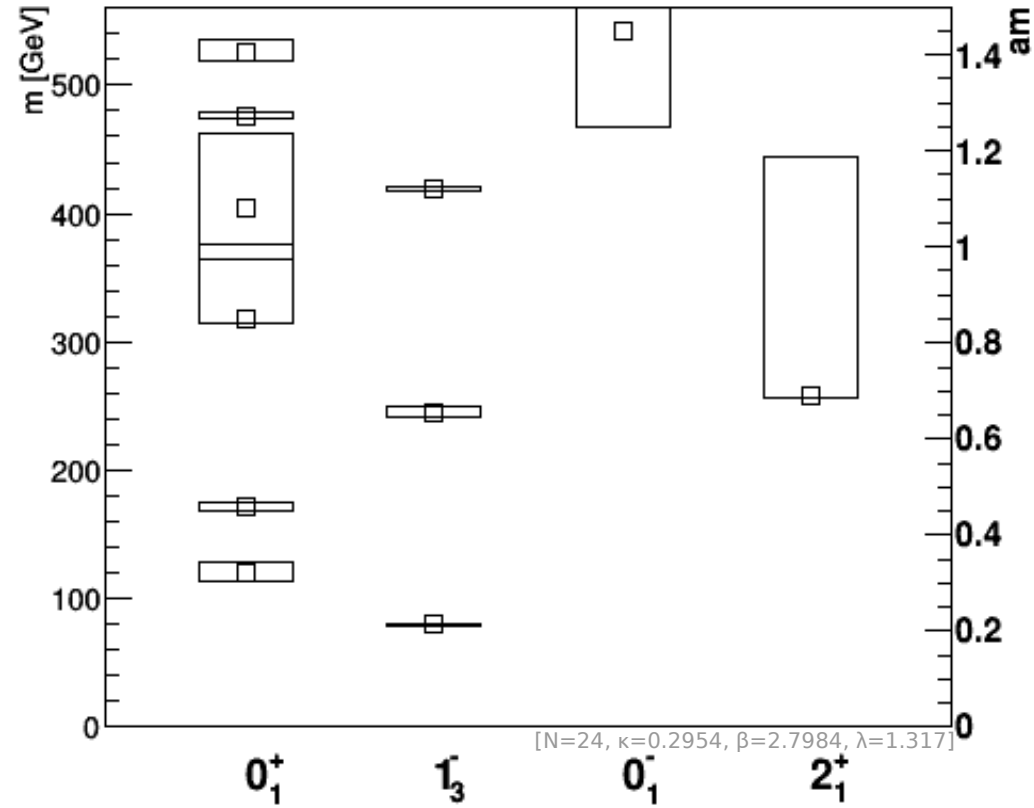
NB: weakly coupled



Typical spectra

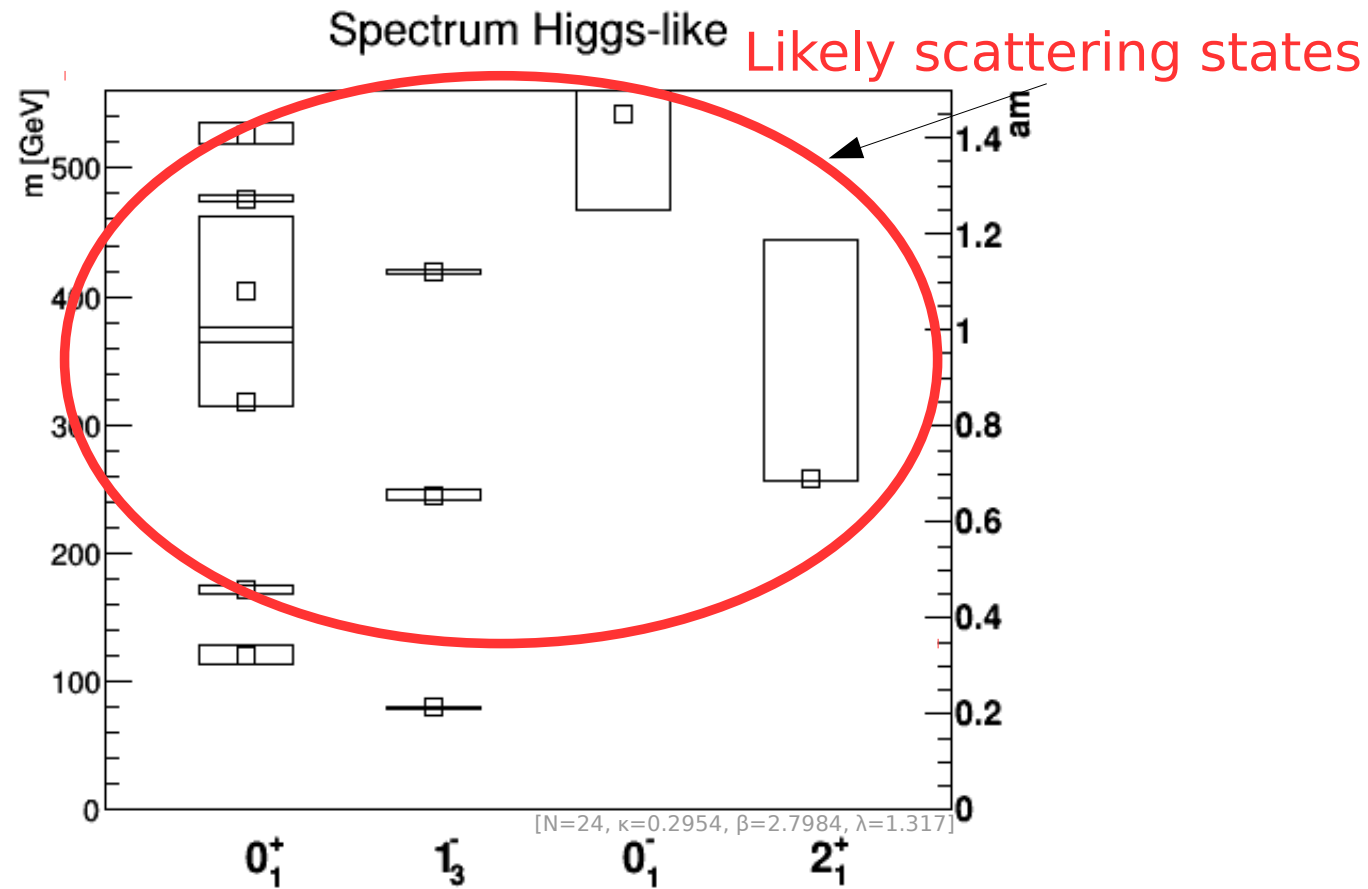
[Maas, Mufti '13,'14,
Evertz et al.'86, Langguth et al.'85,'86]

Spectrum Higgs-like



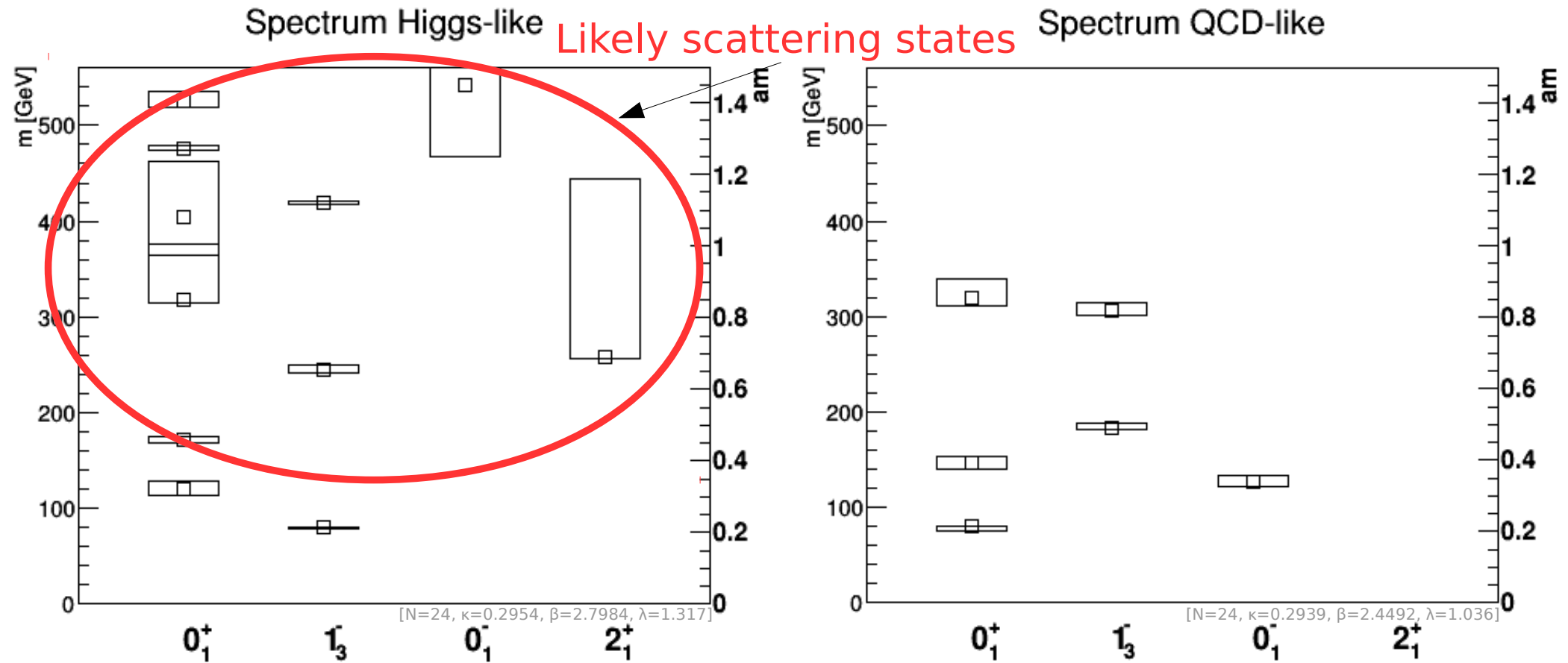
Typical spectra

[Maas, Mufti '13,'14,
Evertz et al.'86, Langguth et al.'85,'86]



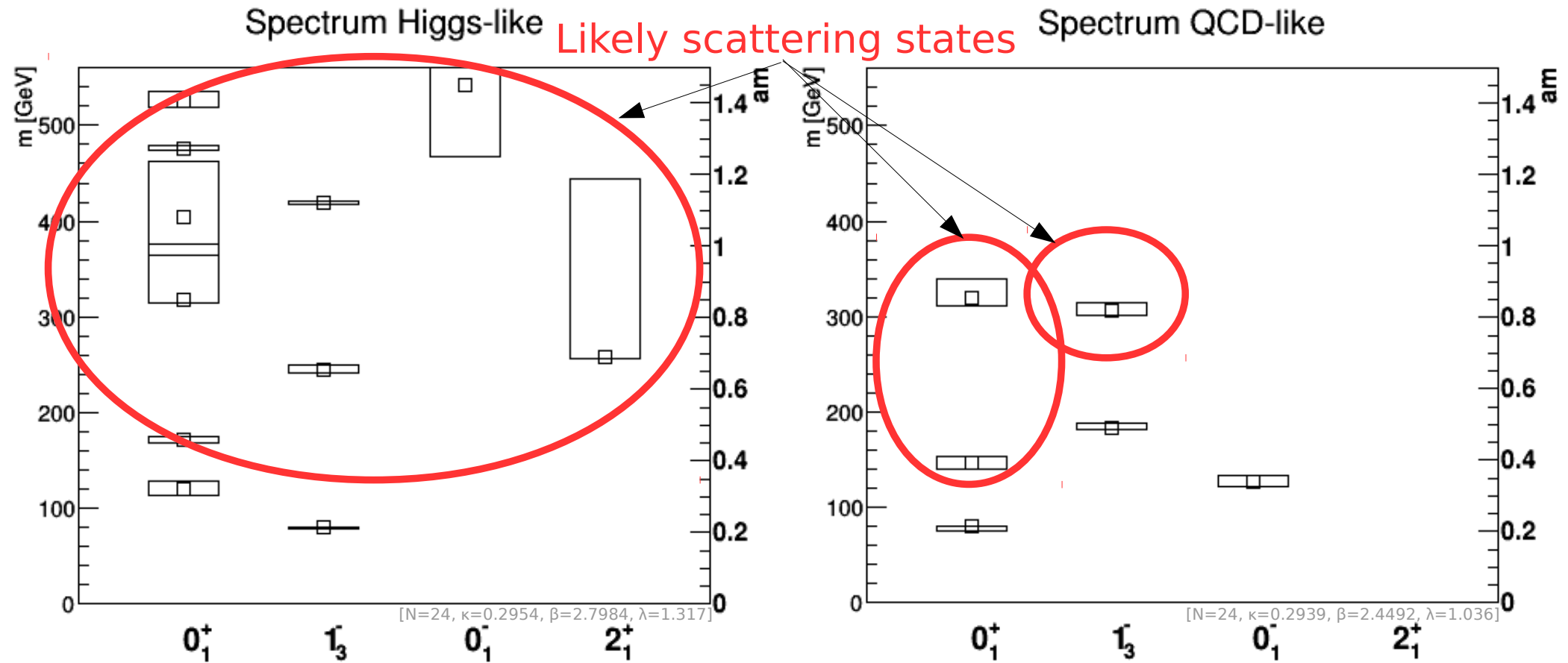
Typical spectra

[Maas, Mufti '13,'14,
Evertz et al.'86, Langguth et al.'85,'86]



Typical spectra

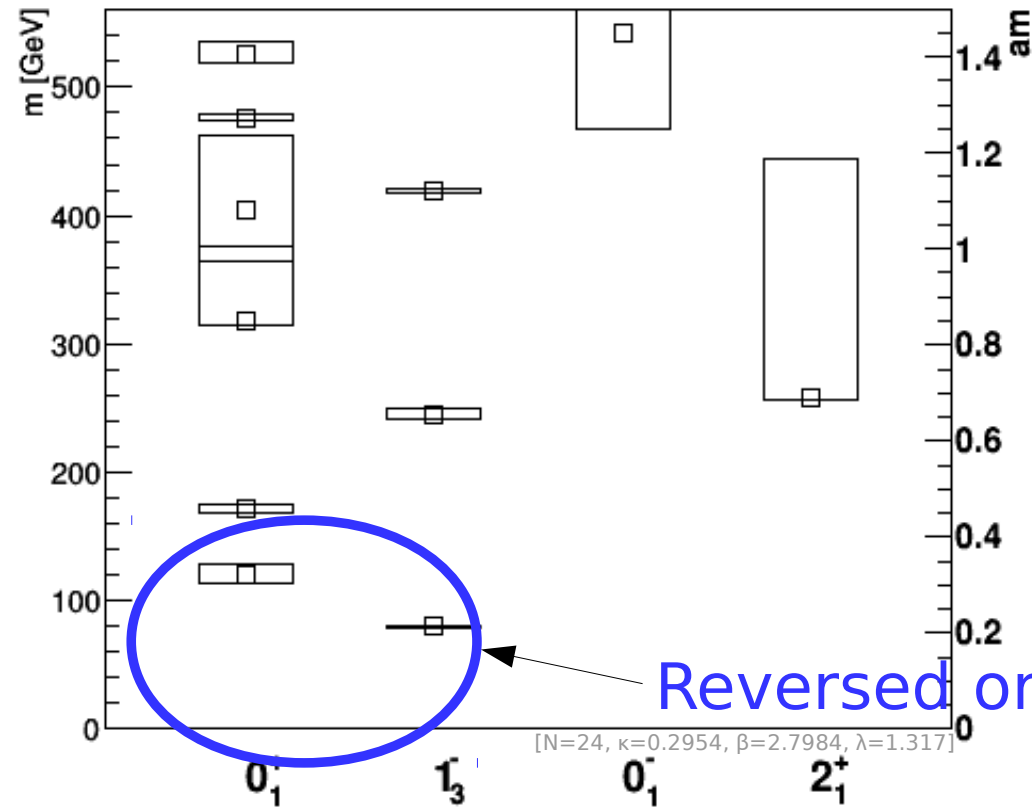
[Maas, Mufti '13,'14,
Evertz et al.'86, Langguth et al.'85,'86]



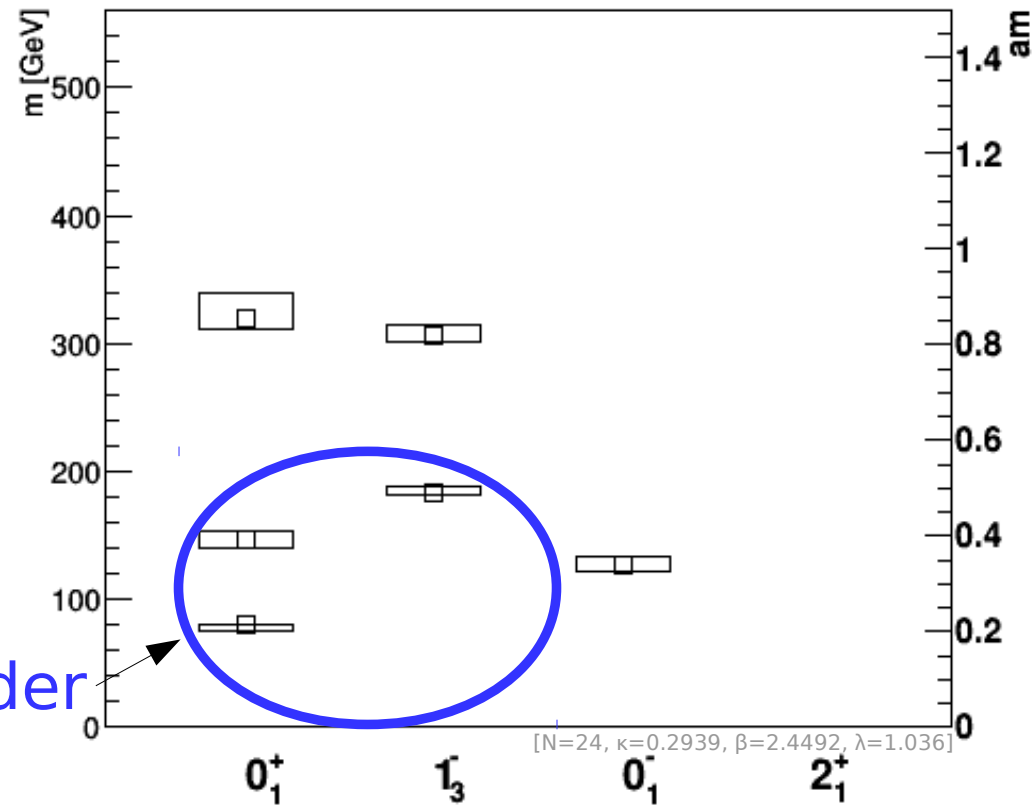
Typical spectra

[Maas, Mufti '13,'14,
Evertz et al.'86, Langguth et al.'85,'86]

Spectrum Higgs-like



Spectrum QCD-like

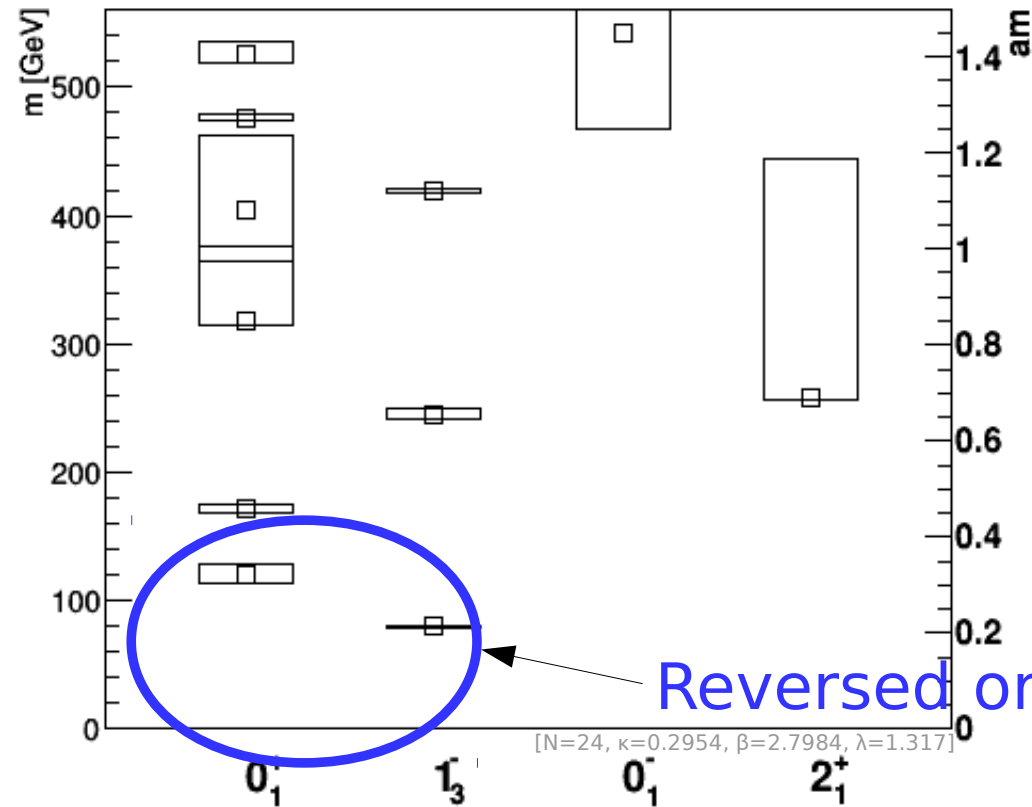


- Generically different low-lying spectra
 - 0^+ lighter in QCD-like region
 - 1^- lighter in Higgs-like region

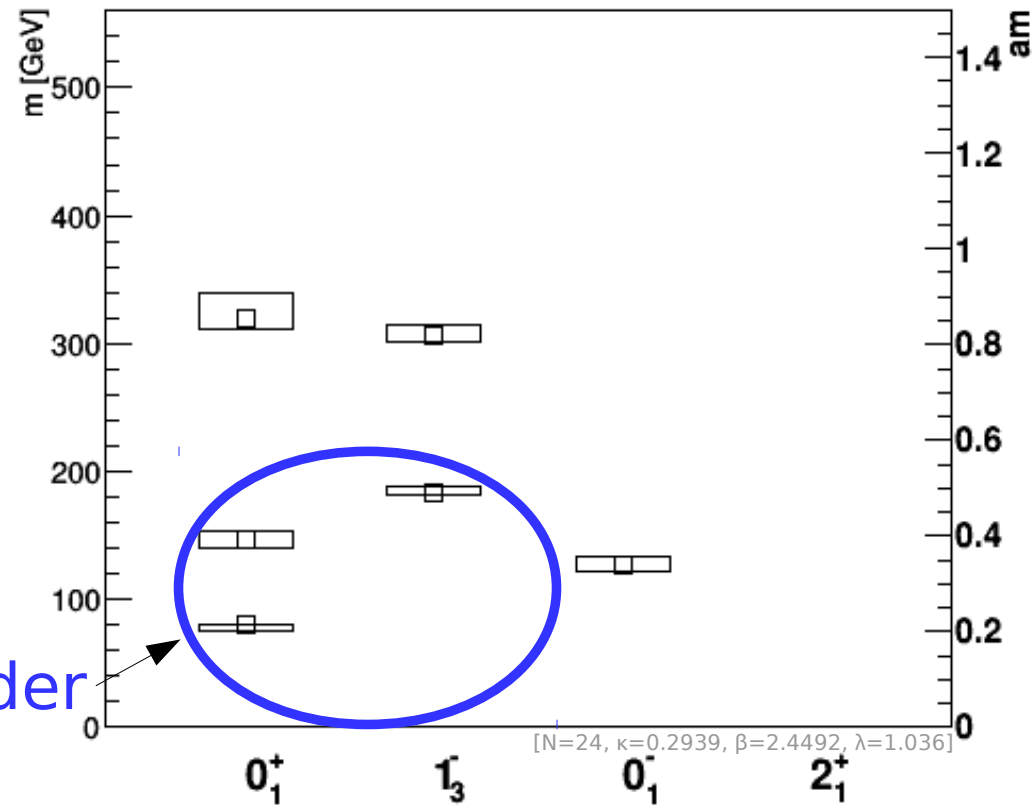
Typical spectra

[Maas, Mufti '13,'14,
Evertz et al.'86, Langguth et al.'85,'86]

Spectrum Higgs-like



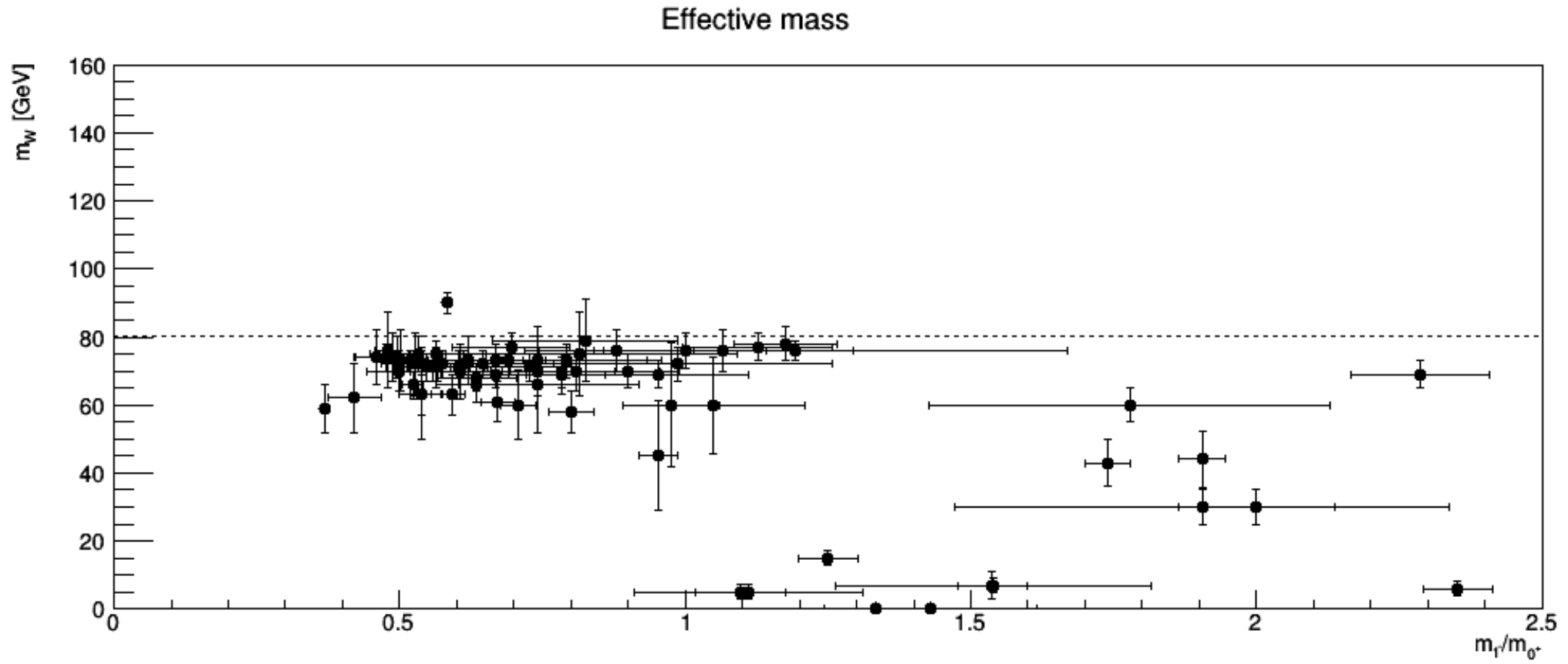
Spectrum QCD-like



- Generically different low-lying spectra
 - 0^+ lighter in QCD-like region
 - 1^- lighter in Higgs-like region
- Coincides with gauge-dependent definitions

Limits

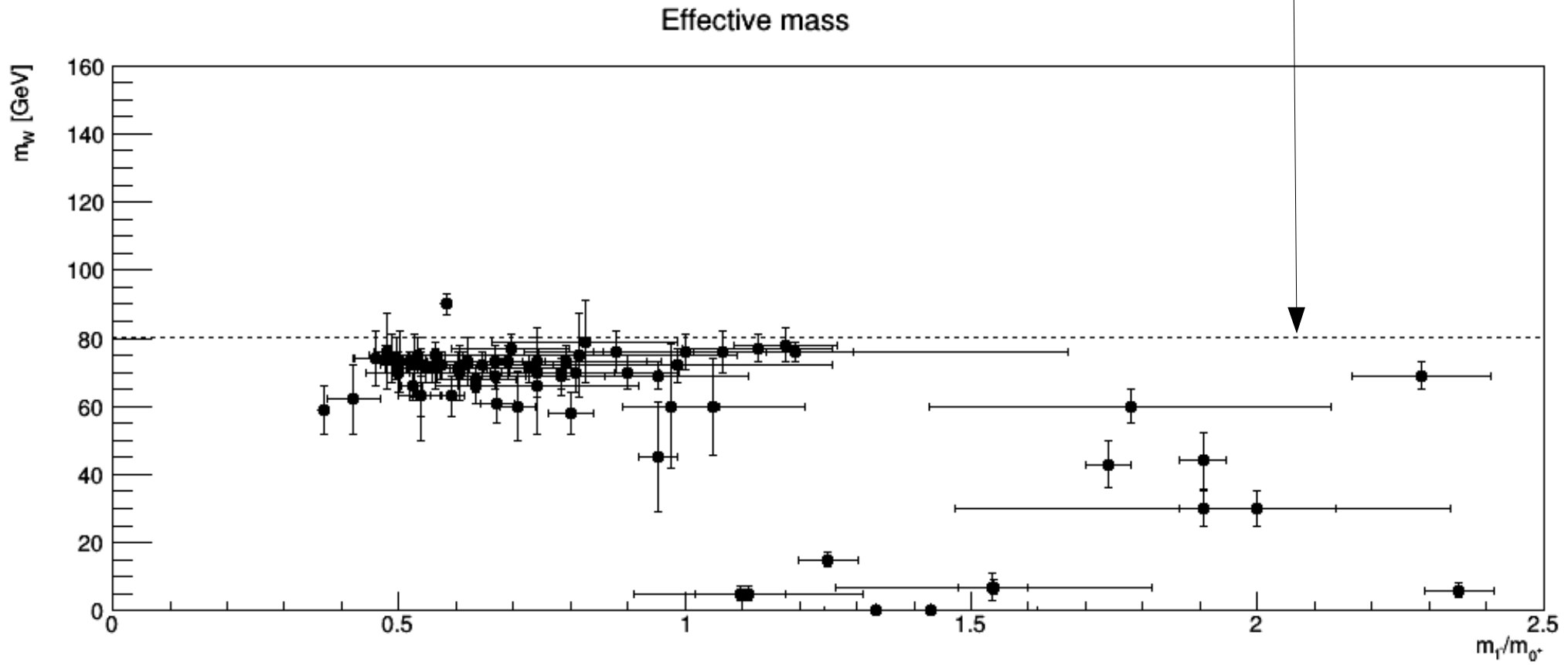
[Maas & Mufti'14]



Limits

[Maas & Mufti'14]

FMS prediction

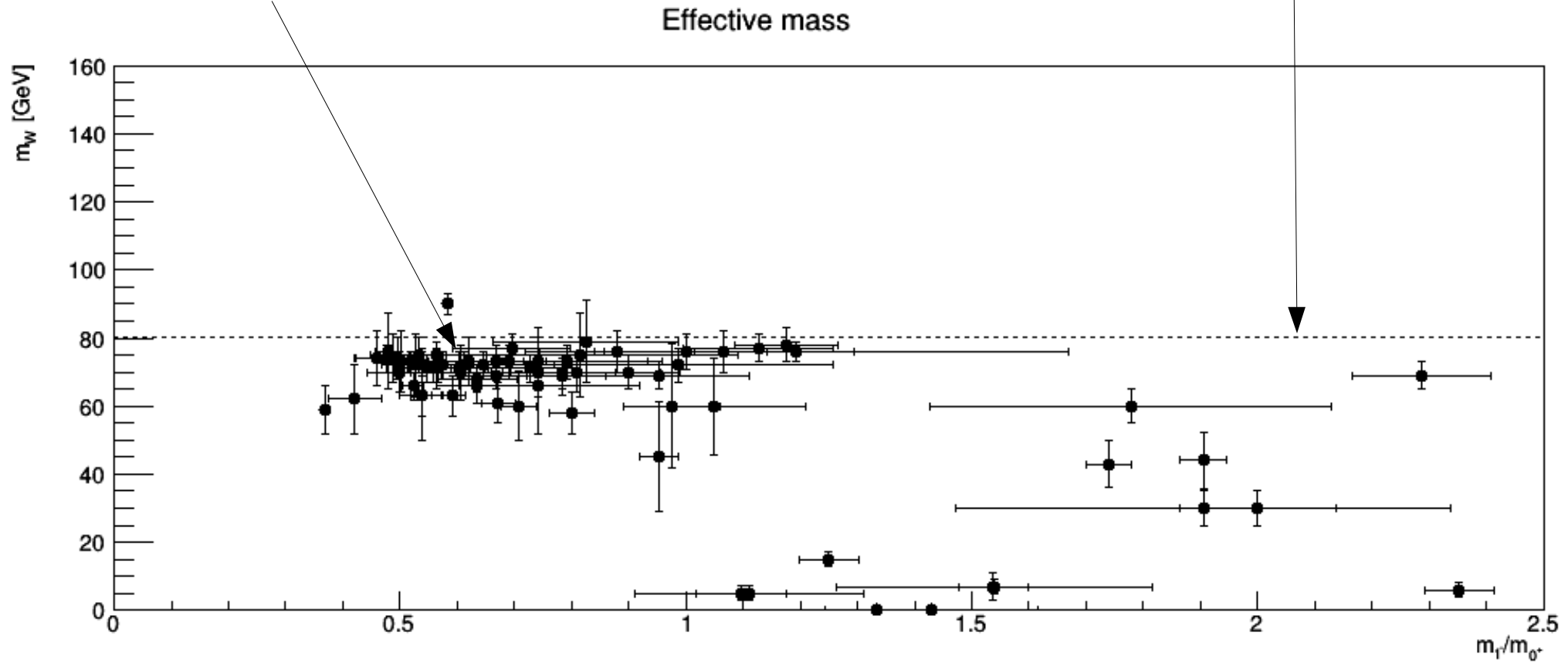


Limits

[Maas & Mufti'14]

Too low: Finite volume effect

FMS prediction

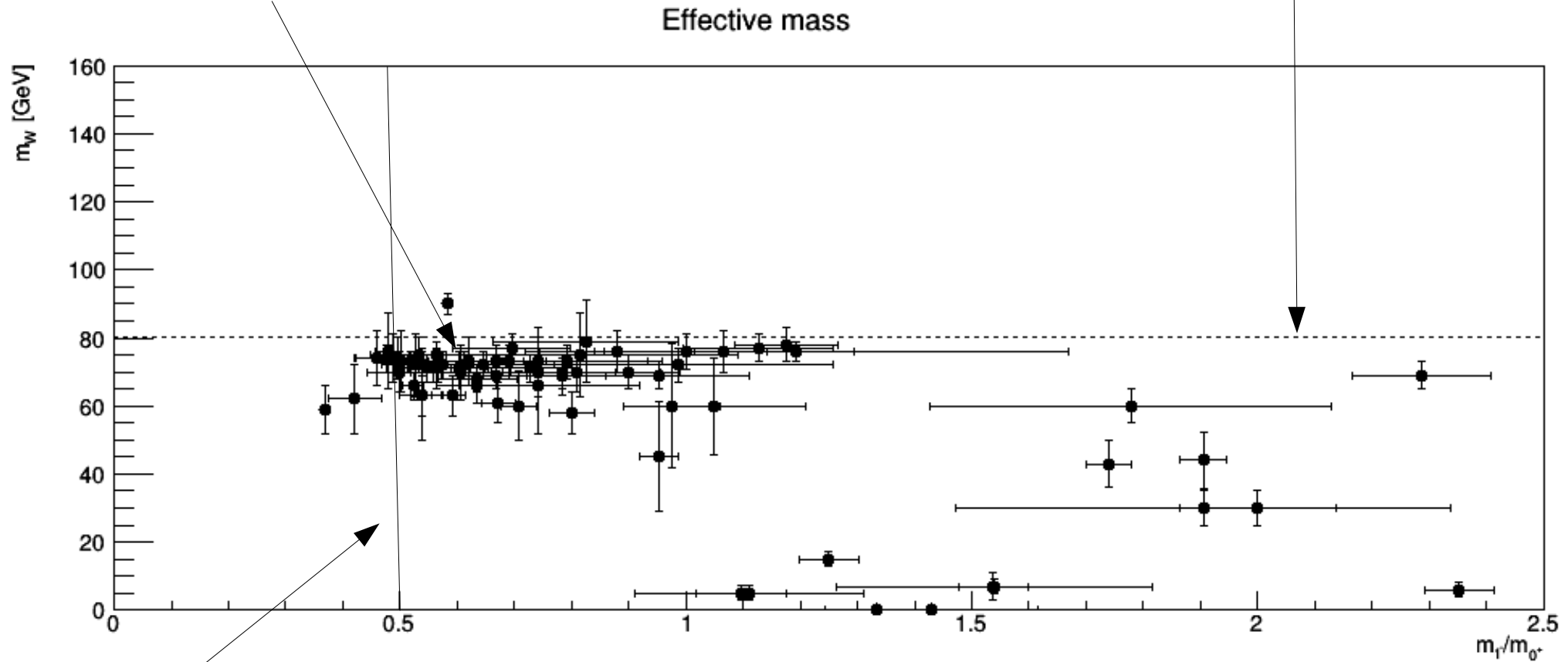


Limits

[Maas & Mufti'14]

Too low: Finite volume effect

FMS prediction



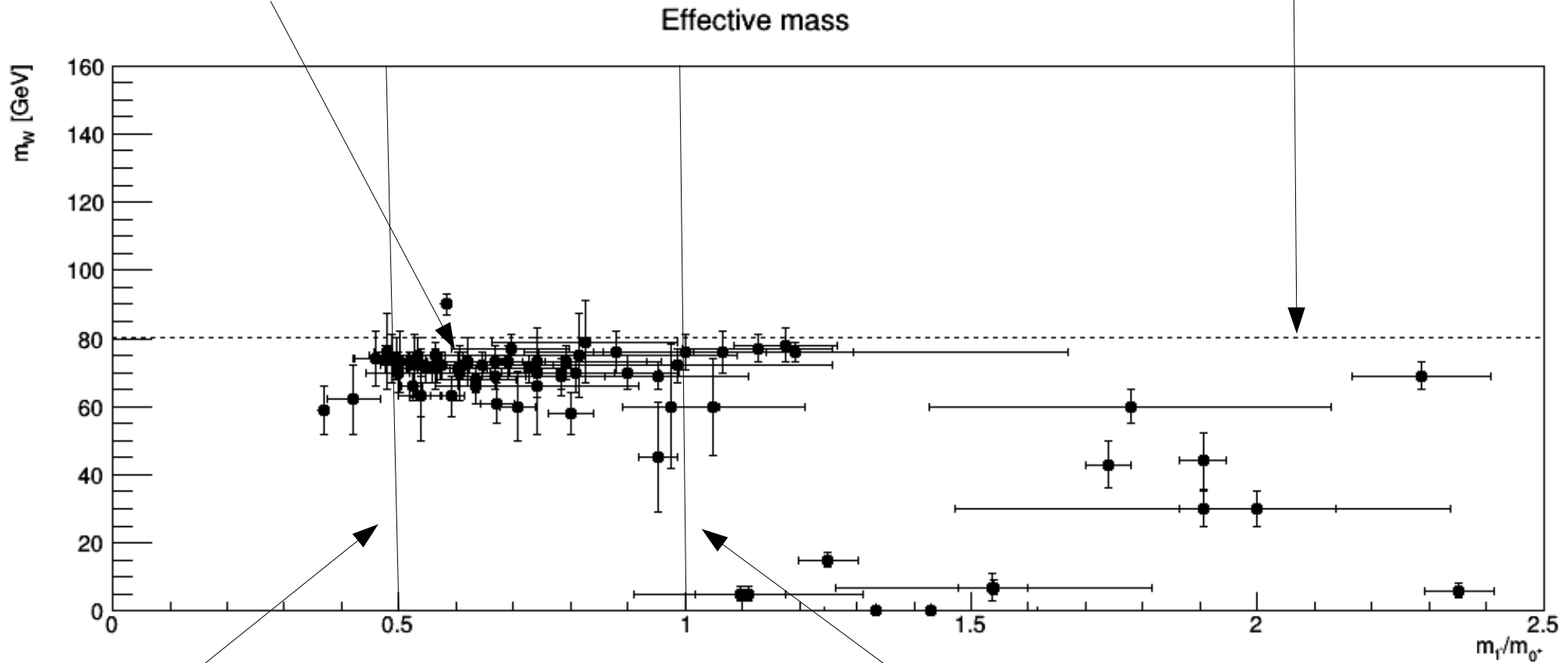
Elastic decay threshold
Higgs as resonance
Expensive, signal very bad

Limits

[Maas & Mufti'14]

Too low: Finite volume effect

FMS prediction

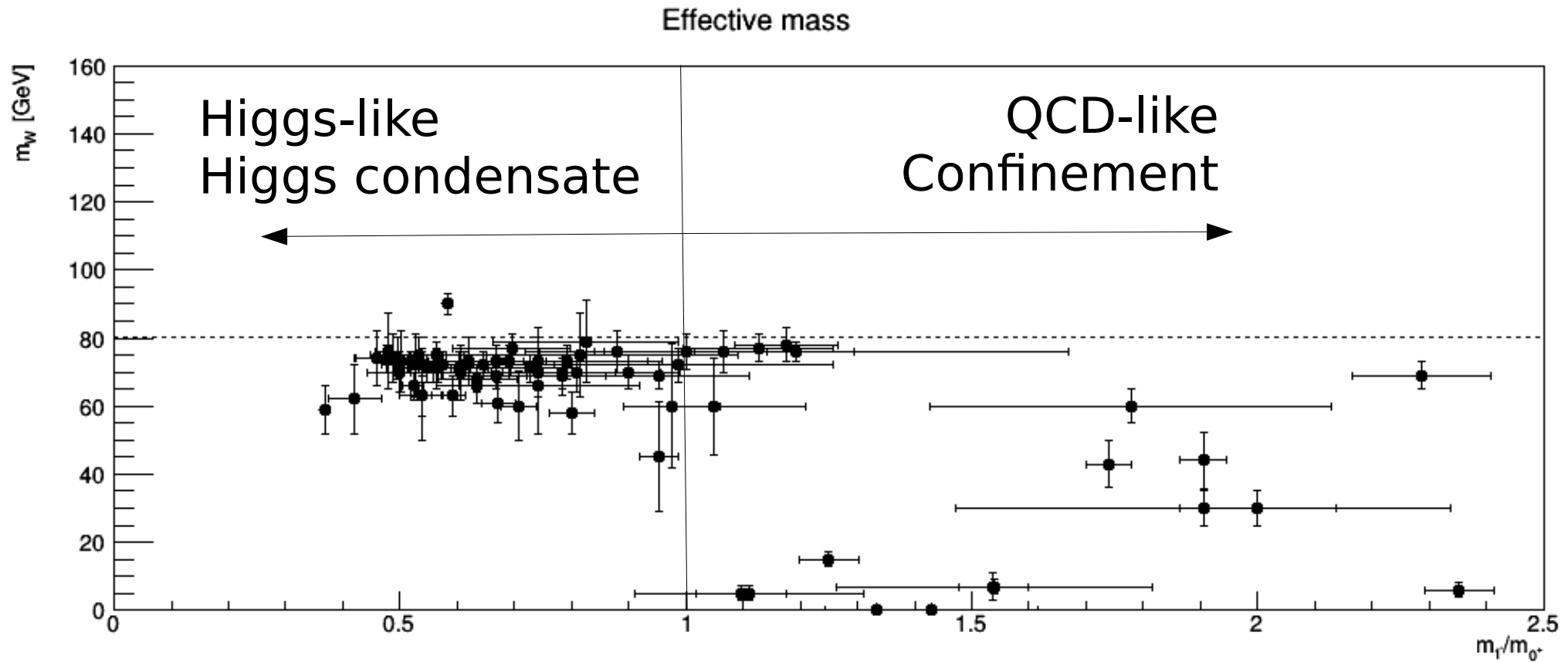


Elastic decay threshold
Higgs as resonance
Expensive, signal very bad

Higgs and W mass agrees
FMS stops working
So does Brout-Englert-Higgs!

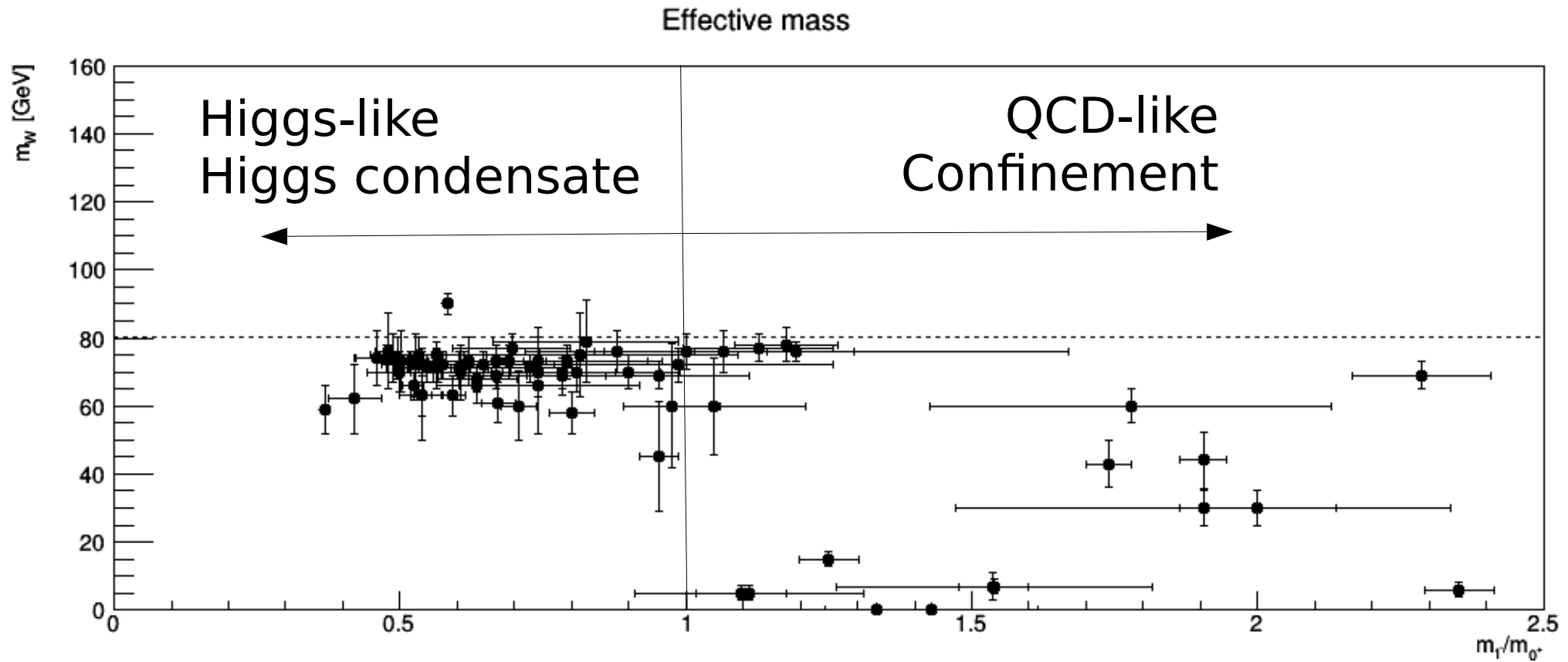
Limits

[Maas & Mufti'14]



Limits

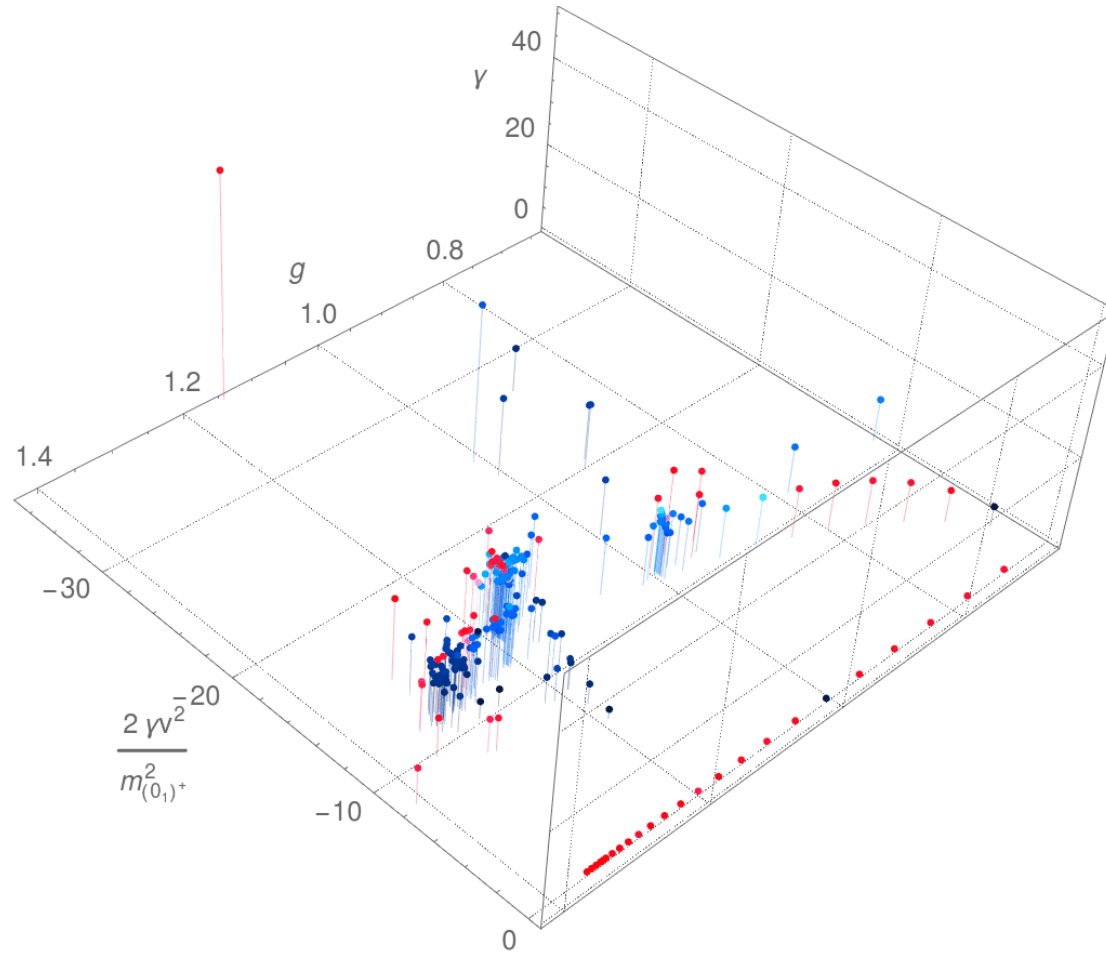
[Maas & Mufti'14]



Does not coincide with weak/strong coupling transitions!

Phase diagram

[Maas & Mufti'15]

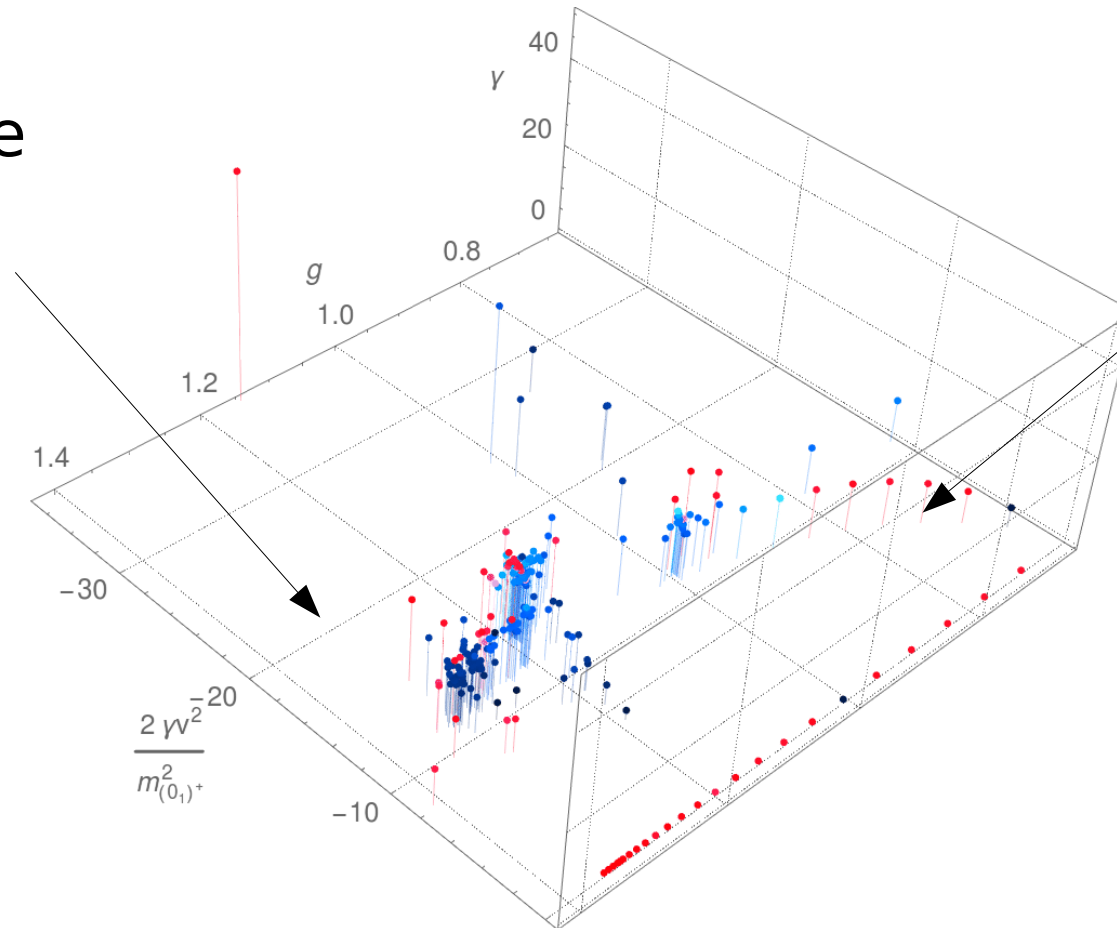


Phase diagram

[Maas & Mufti'15]

Higgs like

QCD like



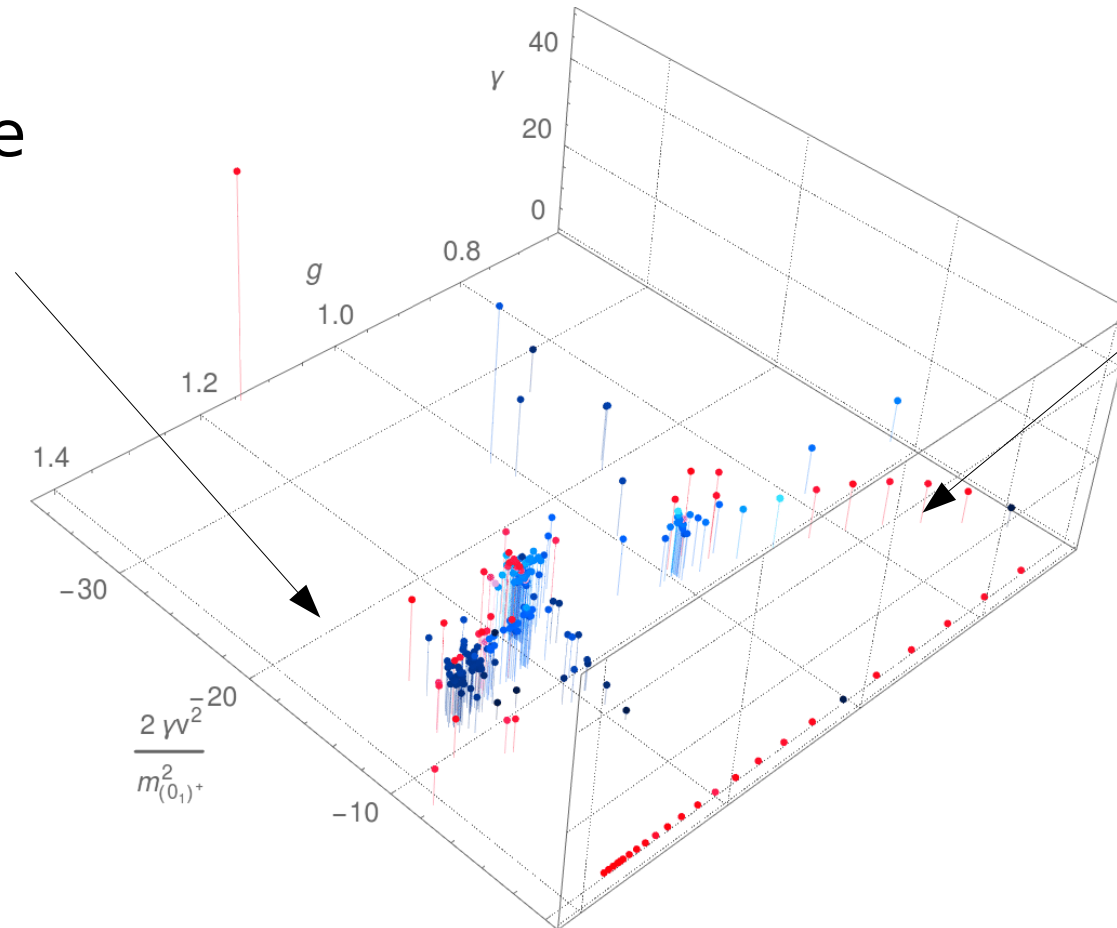
- Quantum effects remove BEH effect
 - Opposite does not happen

Phase diagram

[Maas & Mufti'15]

Higgs like

QCD like



- Quantum effects remove BEH effect
 - Opposite does not happen
- Interacting continuum limit? [Gies & Zambelli'15]

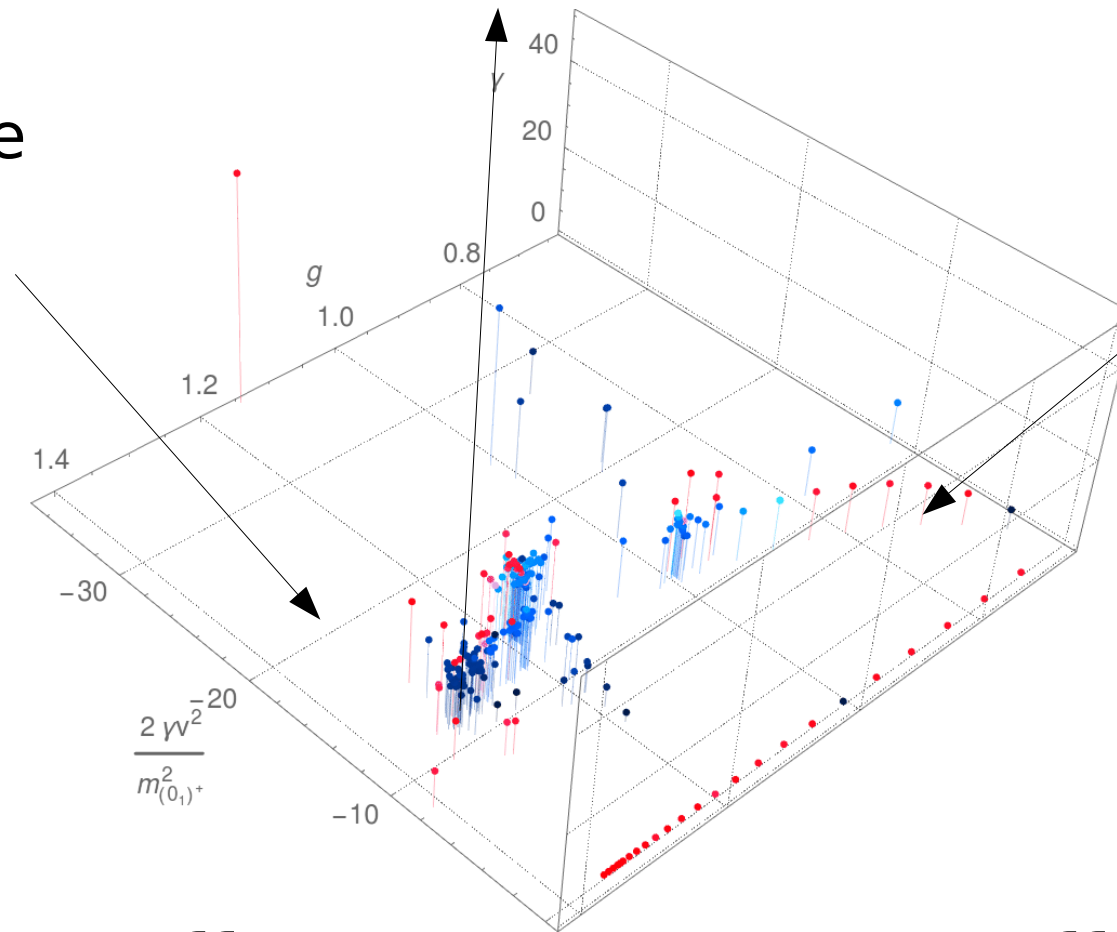
Phase diagram

[Maas & Mufti'15]

Critical end-line? [Bonati et al.'10]

Higgs like

QCD like



- Quantum effects remove BEH effect
 - Opposite does not happen
- Interacting continuum limit? [Gies & Zambelli'15]

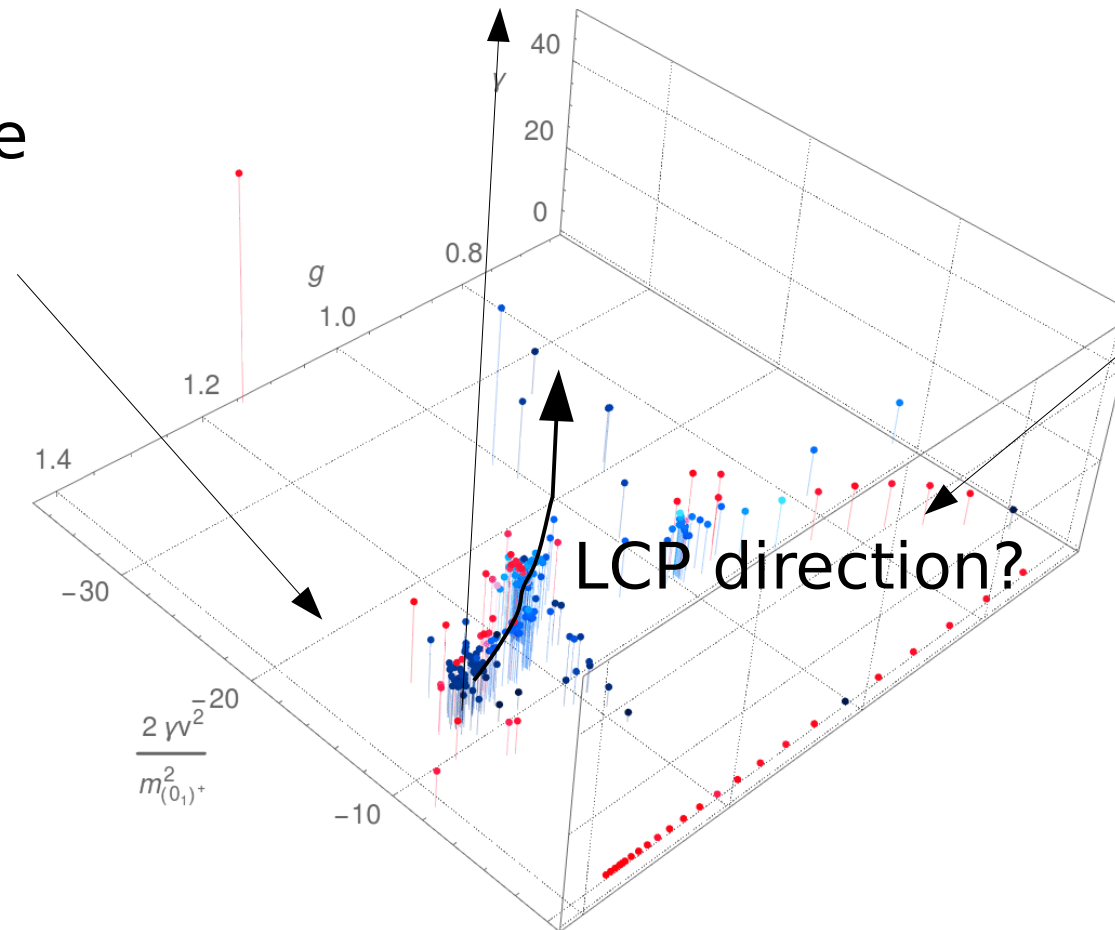
Phase diagram

[Maas & Mufti'15]

Critical end-line? [Bonati et al.'10]

Higgs like

QCD like



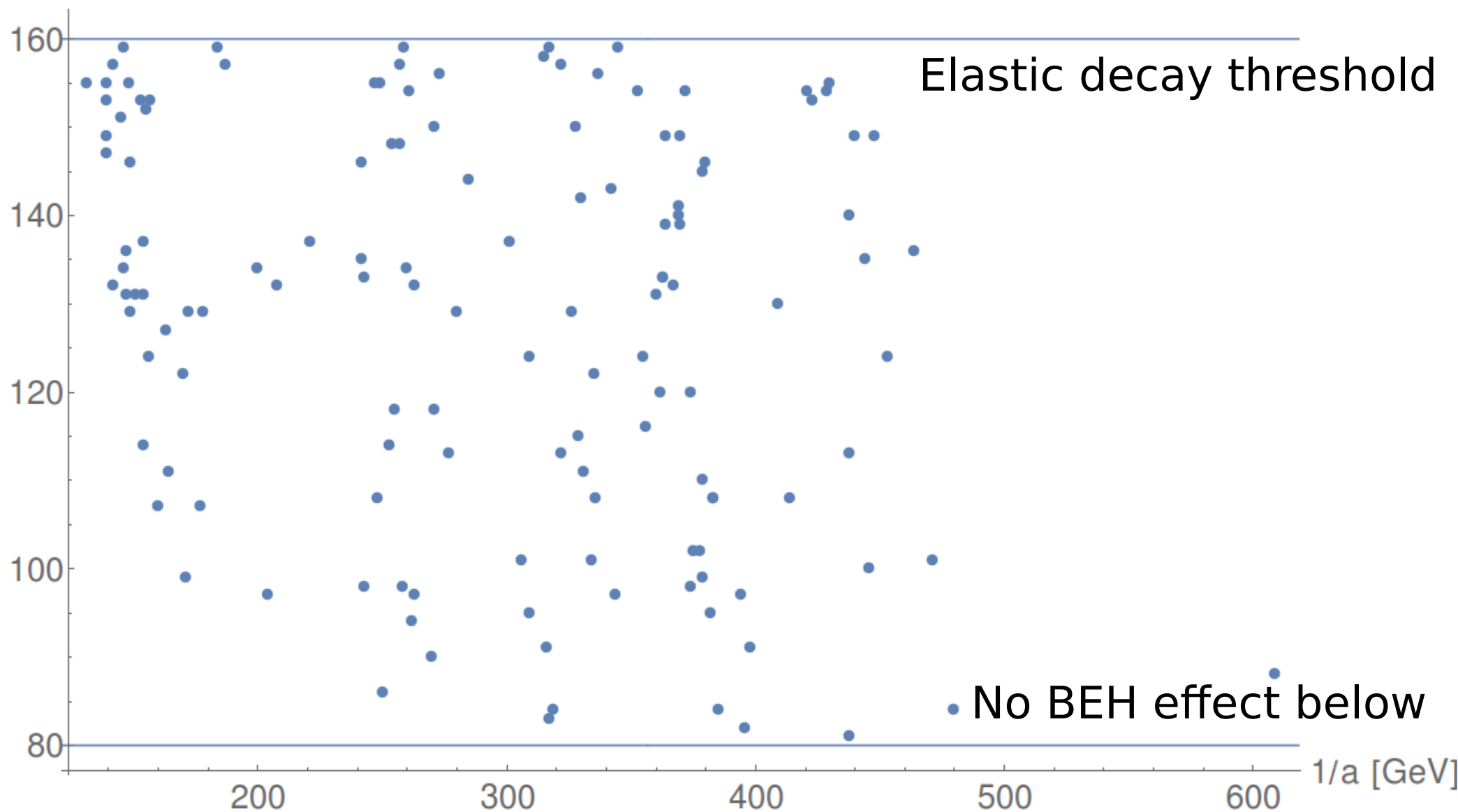
- Quantum effects remove BEH effect
 - Opposite does not happen
- Interacting continuum limit? [Gies & Zambelli'15]
 - LCP: 0^+ , 1^- masses, $\alpha(200 \text{ GeV})$ (miniMOM scheme)

Higgs mass

[Maas & Mufti'15]

$m_{0,+}$ [GeV]

Standard mass-cutoff plot



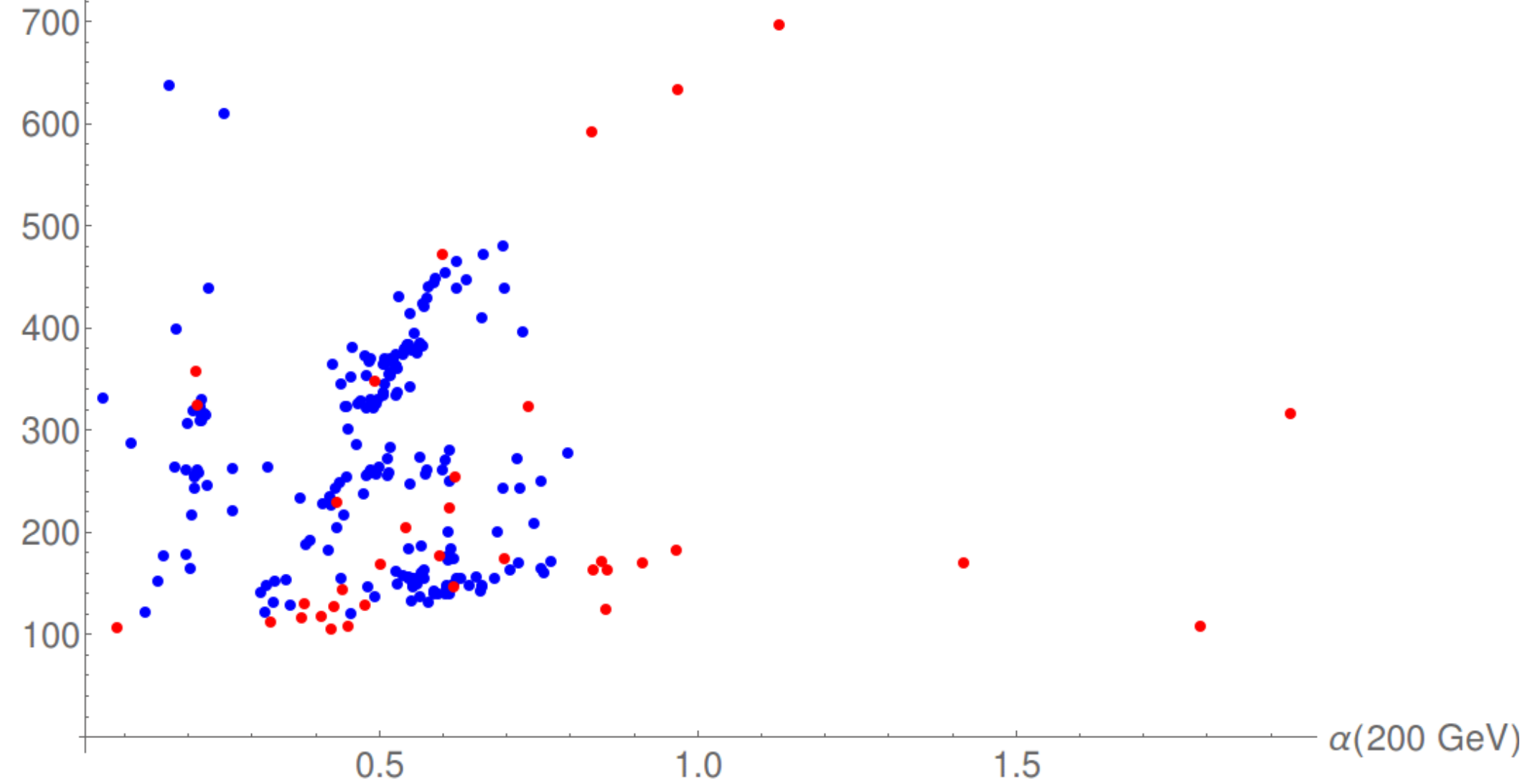
Elastic decay threshold

• No BEH effect below

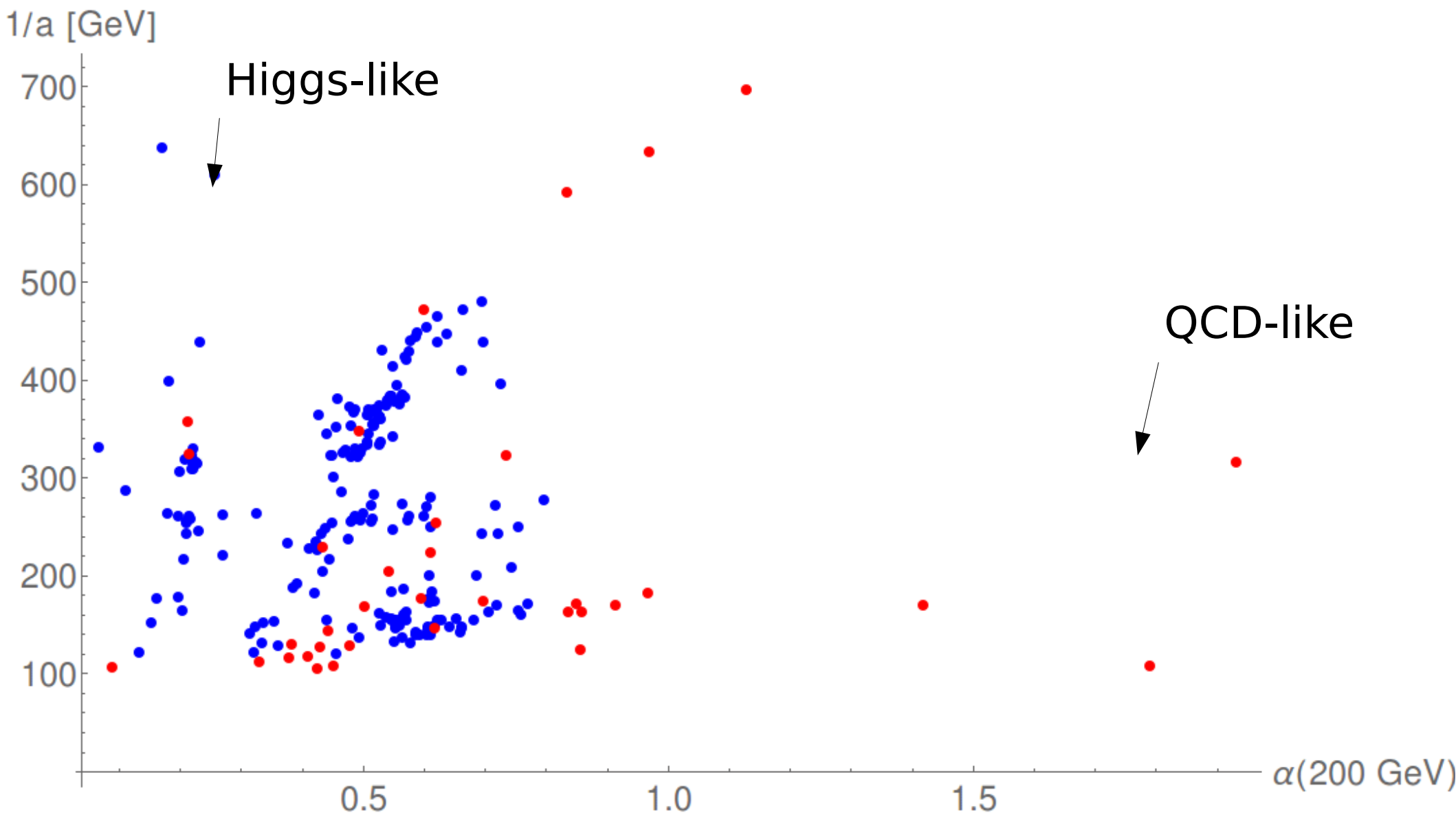
No strong dependence of mass range on cutoff - expected

Perturbative predictivity: Coupling

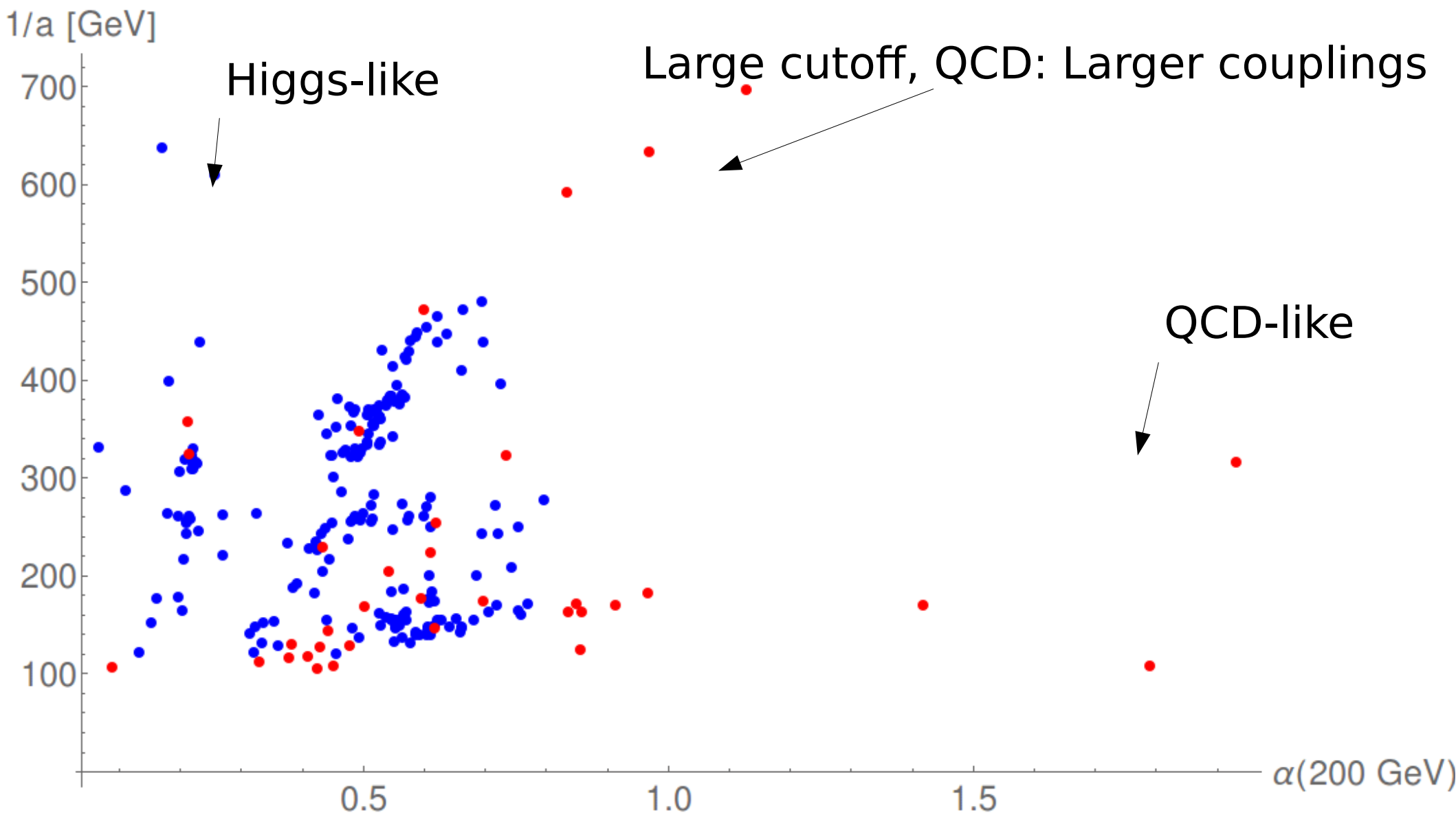
$1/a$ [GeV]



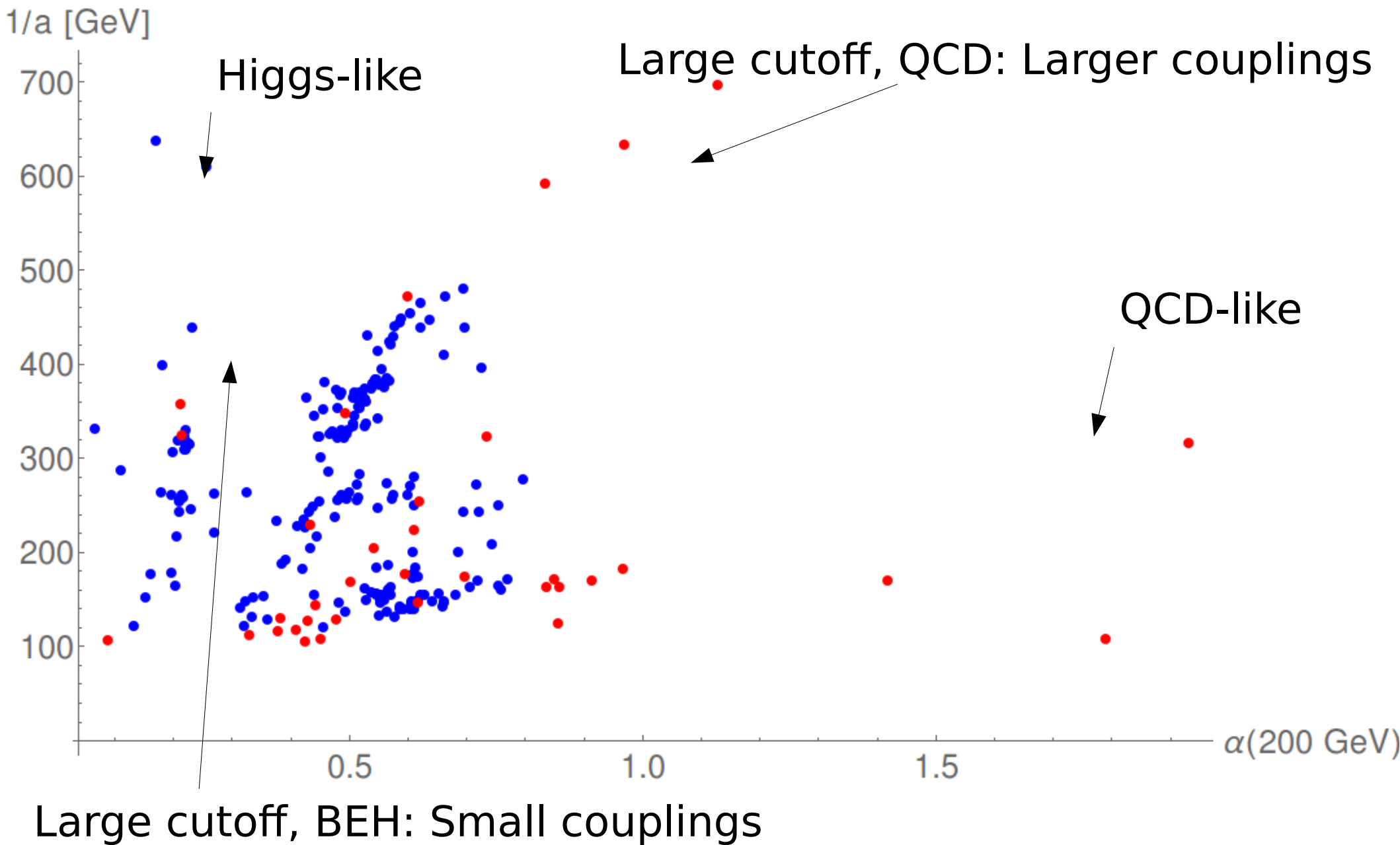
Perturbative predictivity: Coupling



Perturbative predictivity: Coupling

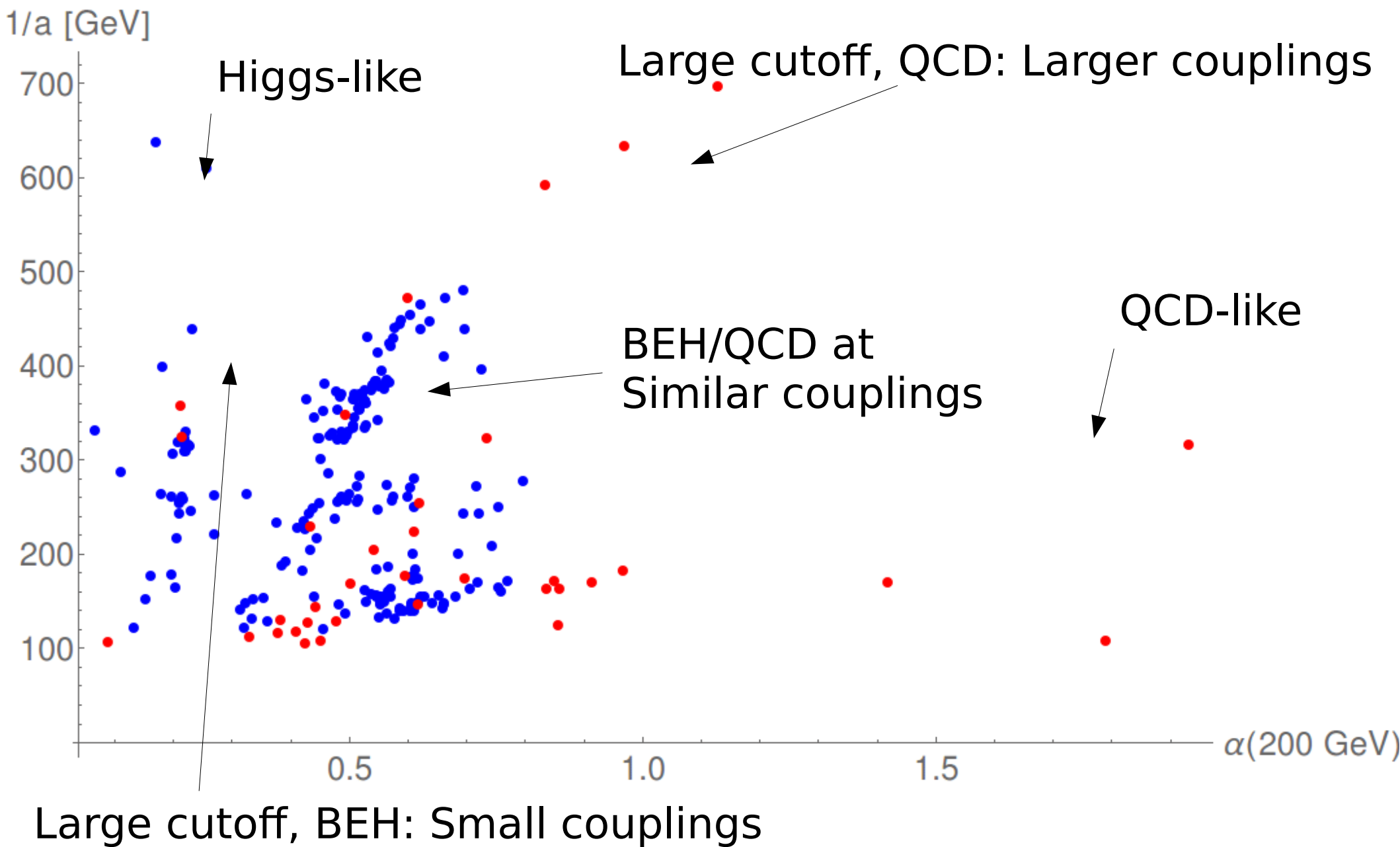


Perturbative predictivity: Coupling



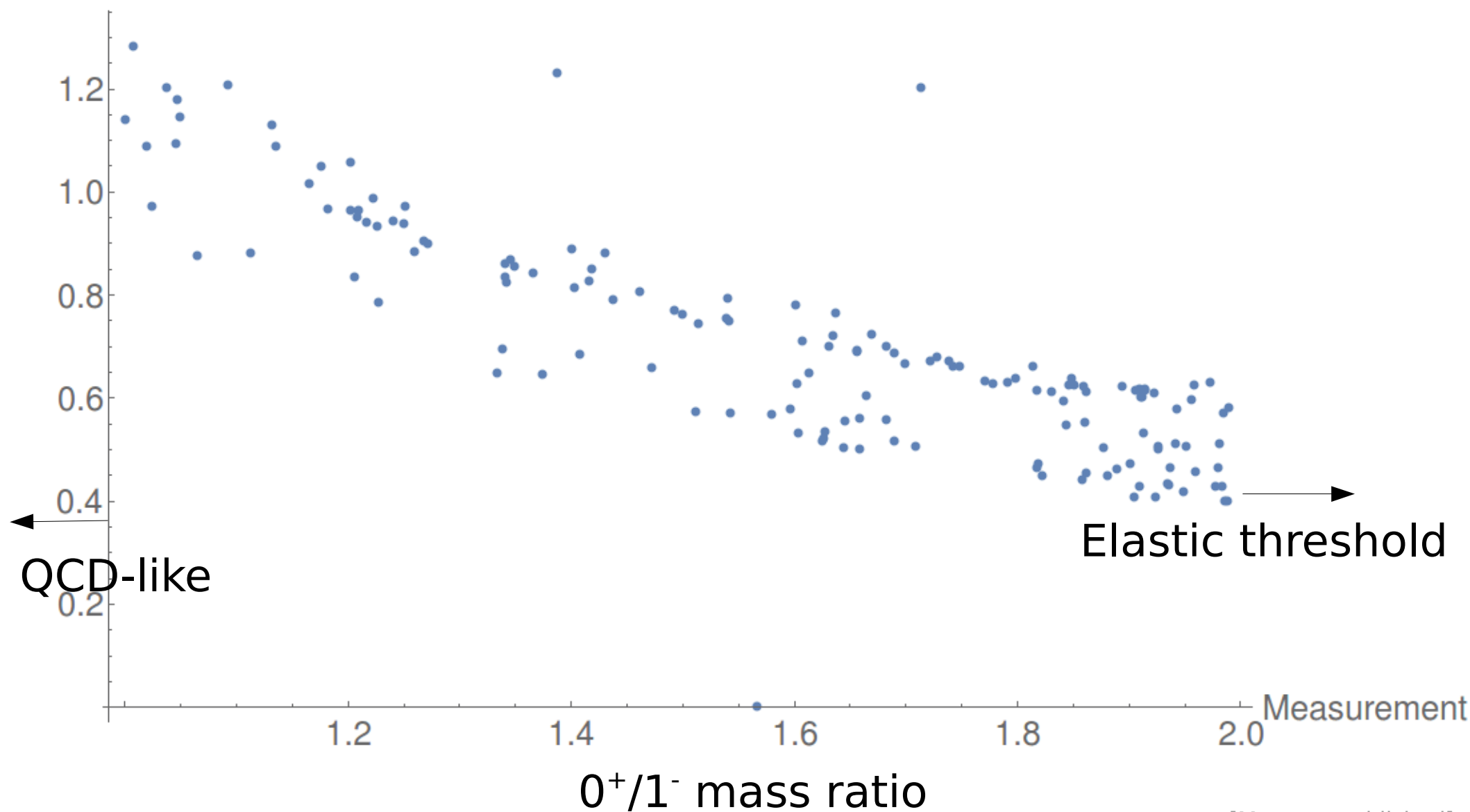
Large cutoff, BEH: Small couplings

Perturbative predictivity: Coupling



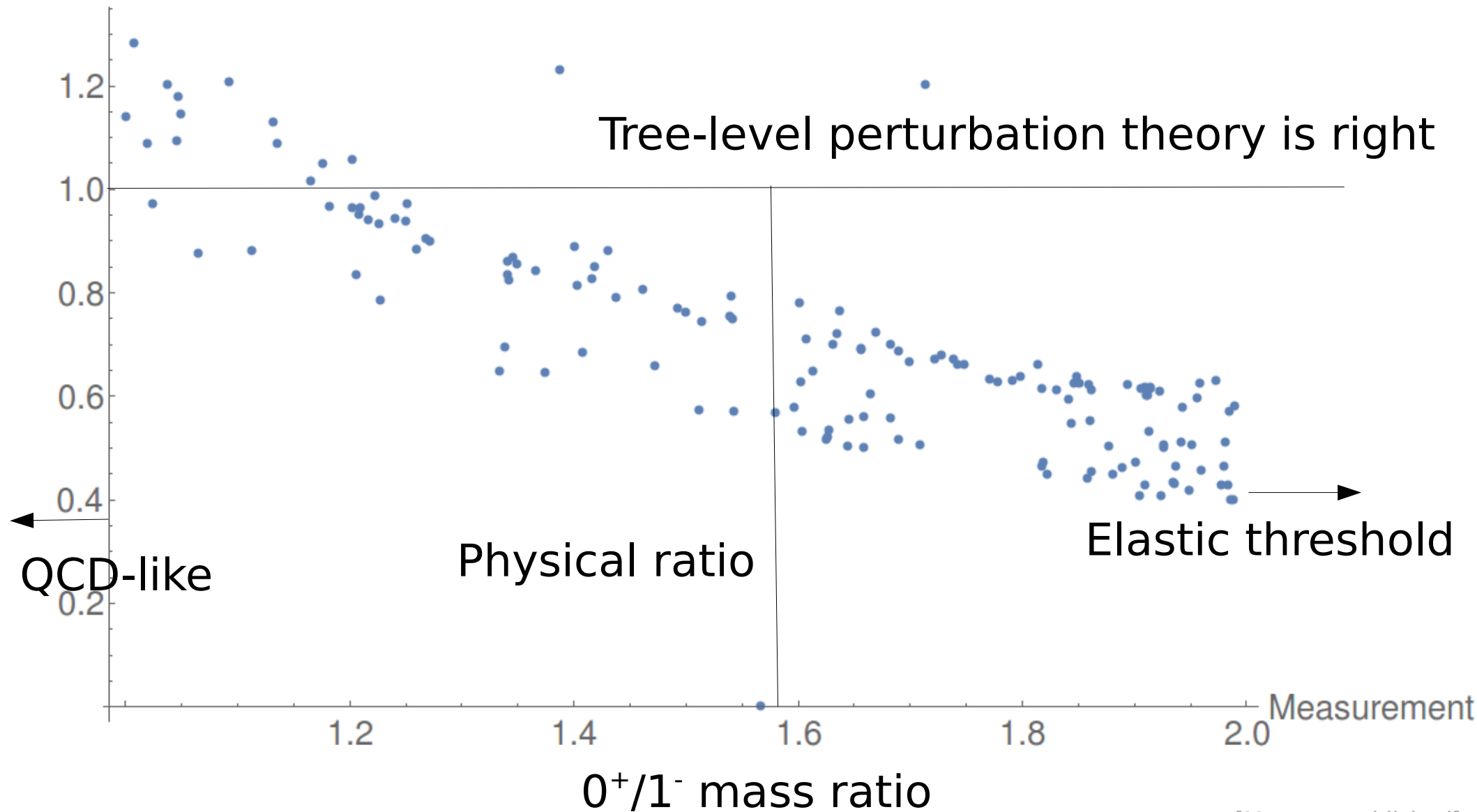
Perturbative predictivity: Mass ratios

Tree-level
Measurement



Perturbative predictivity: Mass ratios

Tree-level
Measurement



Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]
 - Fluctuations can invalidate it
 - Seen on the lattice – but SM is fine

Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]
 - Fluctuations can invalidate it
 - Seen on the lattice – but SM is fine
 - Local and global multiplet structure must fit

Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]
 - Fluctuations can invalidate it
 - Seen on the lattice – but SM is fine
 - Local and global multiplet structure must fit
- Has to be checked for BSM theories

Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]
 - Fluctuations can invalidate it
 - Seen on the lattice – but SM is fine
 - Local and global multiplet structure must fit
- Has to be checked for BSM theories
 - Without Higgs: More subtle [Maas'15]

Example 1: 2HDM

Like the standard model
Gauge-invariant and ordinary perturbation theory
coincide

Implications for 2HDM

[Maas'15,
Maas & Pedro'16]

- Additional Higgs doublet
- Enlarged custodial group

Implications for 2HDM

[Maas'15,
Maas & Pedro'16]

- Additional Higgs doublet
- Enlarged custodial group
- BEH Effect - FMS mechanism applicable
 - In a suitable basis, all condensates contained in a single doublet

Implications for 2HDM

[Maas'15,
Maas & Pedro'16]

- FMS states for maximal custodial group:
 - Scalar sector Singlet

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle \approx \text{const.} + \langle \eta_h^+(x) \eta_h(y) \rangle + O(\eta_h^3)$$

Implications for 2HDM

[Maas'15,
Maas & Pedro'16]

- FMS states for maximal custodial group:

- Scalar sector Singlet

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle \approx \text{const.} + \langle \eta_h^+(x) \eta_h(y) \rangle + O(\eta_h^3)$$

- Scalar Sector Quadruplet

$$\langle (a^+ \Gamma a)(x)(a^+ \Gamma a)(y) \rangle \approx \text{const.} + \langle \eta_a^+(x) \Gamma \eta_a(y) \rangle + O(\eta_a^3)$$

- Splitted into 1+3 states for broken group

Implications for 2HDM

[Maas'15,
Maas & Pedro'16]

- FMS states for maximal custodial group:

- Scalar sector Singlet

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle \approx \text{const.} + \langle \eta_h^+(x) \eta_h(y) \rangle + O(\eta_h^3)$$

- Scalar Sector Quadruplet

$$\langle (a^+ \Gamma a)(x)(a^+ \Gamma a)(y) \rangle \approx \text{const.} + \langle \eta_a^+(x) \Gamma \eta_a(y) \rangle + O(\eta_a^3)$$

- Splitted into 1+3 states for broken group

- Vector triplet

$$\langle (h^+ D_\mu h)(x)(h^+ D_\mu h)(y) \rangle \approx \text{const.} + \langle W_\mu(x) W_\mu(y) \rangle + O(\eta_h^3)$$

- All other states expand to scattering states

Implications for 2HDM

[Maas'15,
Maas & Pedro'16]

- FMS states for maximal custodial group:

- Scalar sector Singlet

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle \approx \text{const.} + \langle \eta_h^+(x) \eta_h(y) \rangle + O(\eta_h^3)$$

- Scalar Sector Quadruplet

$$\langle (a^+ \Gamma a)(x)(a^+ \Gamma a)(y) \rangle \approx \text{const.} + \langle \eta_a^+(x) \Gamma \eta_a(y) \rangle + O(\eta_a^3)$$

- Splitted into 1+3 states for broken group

- Vector triplet

$$\langle (h^+ D_\mu h)(x)(h^+ D_\mu h)(y) \rangle \approx \text{const.} + \langle W_\mu(x) W_\mu(y) \rangle + O(\eta_h^3)$$

- All other states expand to scattering states

- Validity: Requires non-perturbative check

- Discrete factor groups could yield doubling

Implications for 2HDM

[Maas'15,
Maas & Pedro'16]

- Additional Higgs doublet
- Enlarged custodial group
- BEH Effect - FMS mechanism applicable
 - In a suitable basis, all condensates contained in a single doublet
 - Yields again perturbative spectrum
 - Discrete factor groups may be a problem
- Key: Global multiplet structure diverse

Implications for 2HDM

[Maas'15,
Maas & Pedro'16]

- Additional Higgs doublet
- Enlarged custodial group
- BEH Effect - FMS mechanism applicable
 - In a suitable basis, all condensates contained in a single doublet
 - Yields again perturbative spectrum
 - Discrete factor groups may be a problem
- Key: Global multiplet structure diverse
- Size of fluctuations needs to be checked non-perturbatively!

Example 2: GUT-like structure

Gauge-invariant perturbation theory correct
and
different from ordinary perturbation theory

Implications for GUTs

[Maas'15
Törek & Maas '15, '16]

- GUTs: Large gauge group, small custodial group
 - Standard model structure: diagonal subgroup – not gauge-invariant

Implications for GUTs

[Maas'15
Törek & Maas '15, '16]

- GUTs: Large gauge group, small custodial group
 - Standard model structure: diagonal subgroup – not gauge-invariant
- Toy-GUT: $SU(3)$ broken to $SU(2)$

Implications for GUTs

[Maas'15
Törek & Maas '15, '16]

- GUTs: Large gauge group, small custodial group
 - Standard model structure: diagonal subgroup – not gauge-invariant
- Toy-GUT: $SU(3)$ broken to $SU(2)$
 - Perturbative spectrum
 - 1 massive Higgs, 3 massless and 5 (1 (heavier) + 4 (lighter)) massive vectors

Implications for GUTs

[Maas'15
Törek & Maas '15, '16]

- GUTs: Large gauge group, small custodial group
 - Standard model structure: diagonal subgroup – not gauge-invariant
- Toy-GUT: $SU(3)$ broken to $SU(2)$
 - Perturbative spectrum
 - 1 massive Higgs, 3 massless and 5 (1 (heavier) + 4 (lighter)) massive vectors
 - FMS spectrum
 - 1 massive scalar, 1 massive vector
 - Same masses as Higgs and heaviest gauge boson

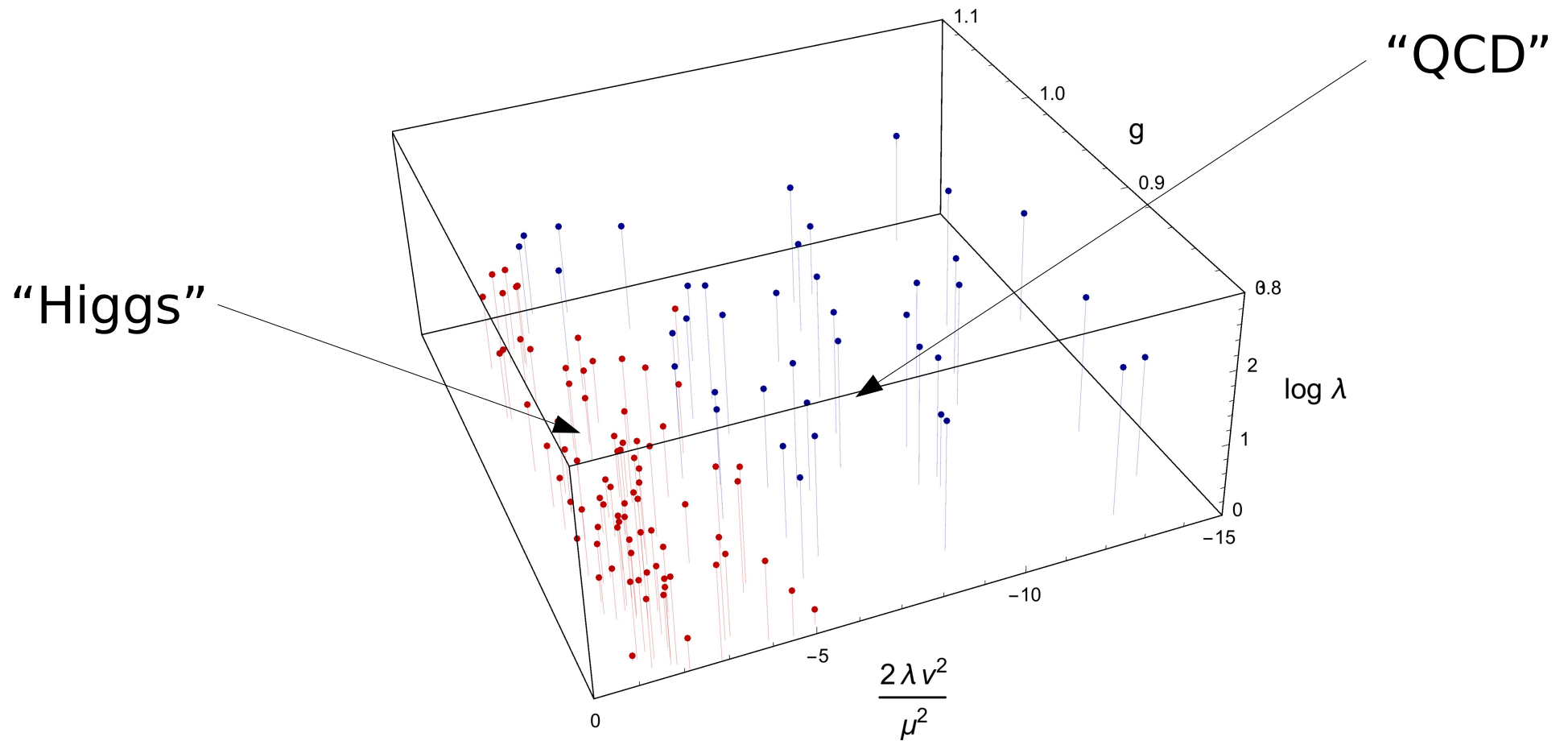
Implications for GUTs

[Maas'15
Törek & Maas '15, '16]

- GUTs: Large gauge group, small custodial group
 - Standard model structure: diagonal subgroup – not gauge-invariant
- Toy-GUT: $SU(3)$ broken to $SU(2)$
 - Perturbative spectrum
 - 1 massive Higgs, 3 massless and 5 (1 (heavier) + 4 (lighter)) massive vectors
 - FMS spectrum
 - 1 massive scalar, 1 massive vector
 - Same masses as Higgs and heaviest gauge boson
 - ...or something else?

Test for GUTs

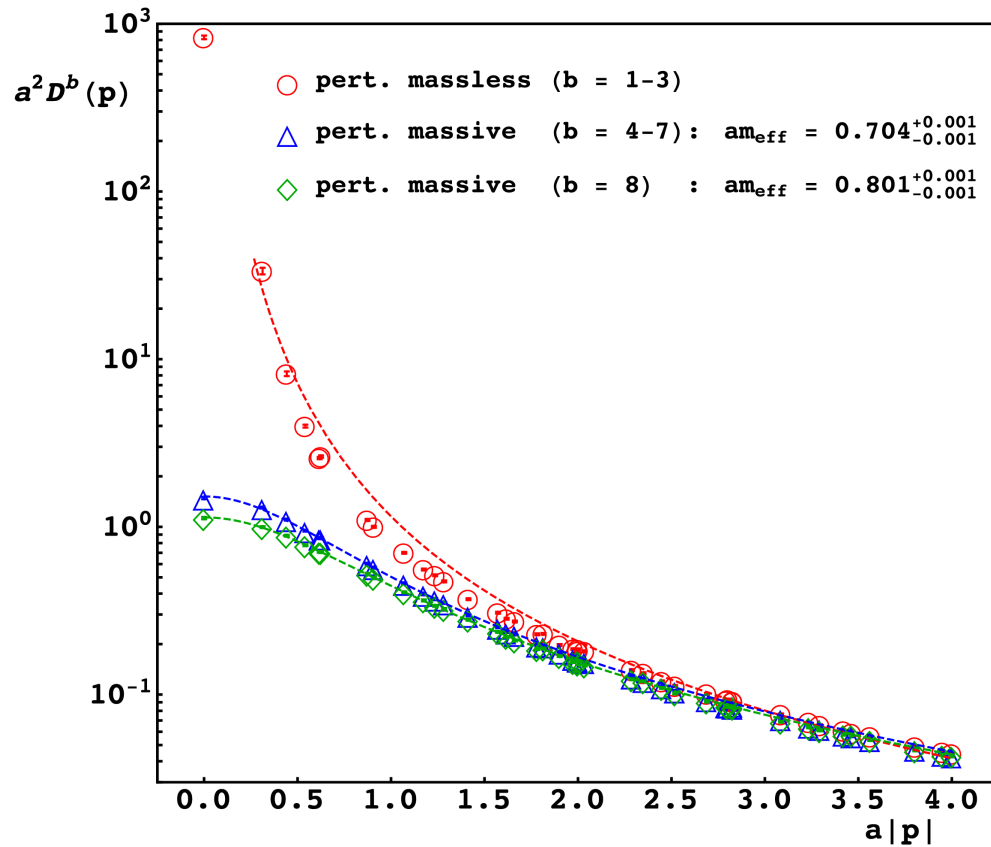
[Maas & Törek'16]



- Separation into Higgs-like and QCD-like

Test for GUTs

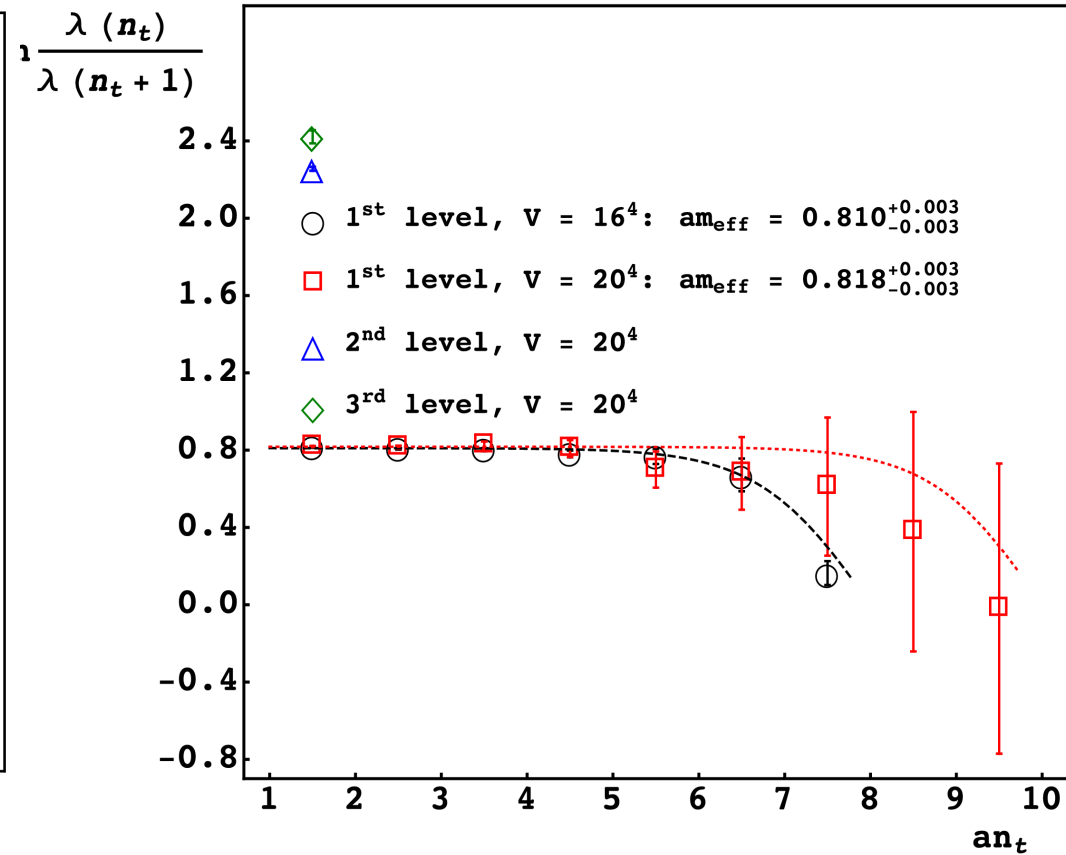
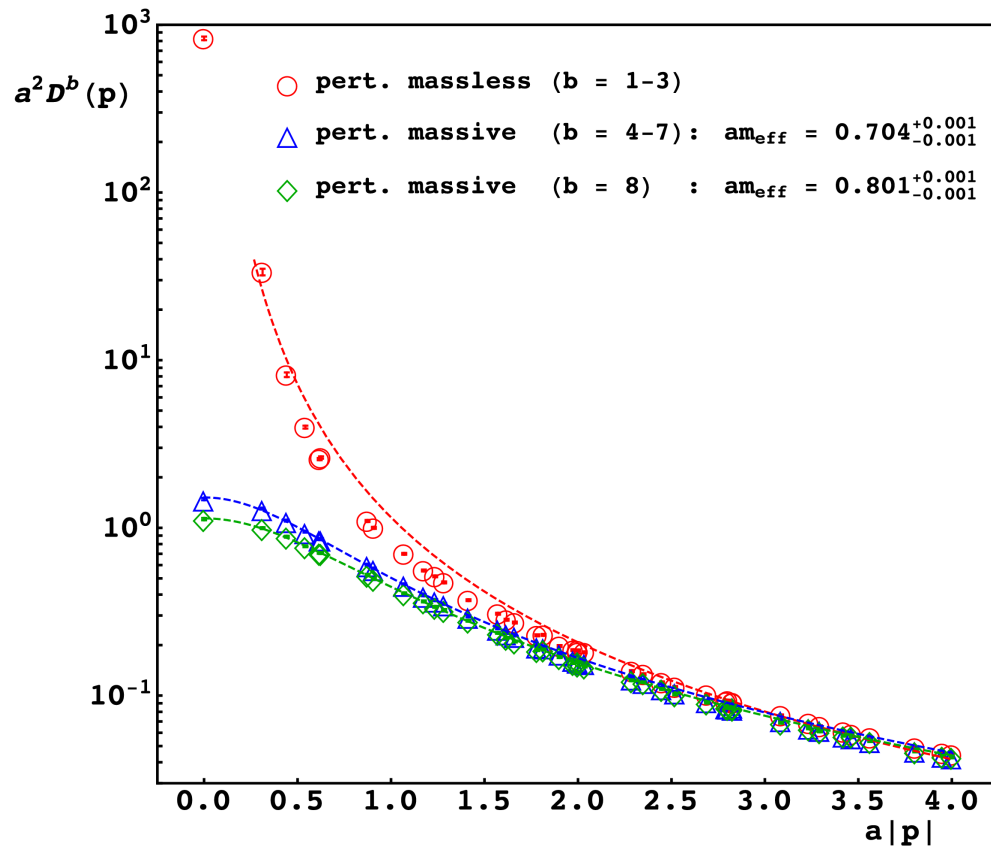
[Maas & Törek'16]



- Propagators almost tree-level
 - Expected splitting in gauge boson spectrum

Test for GUTs

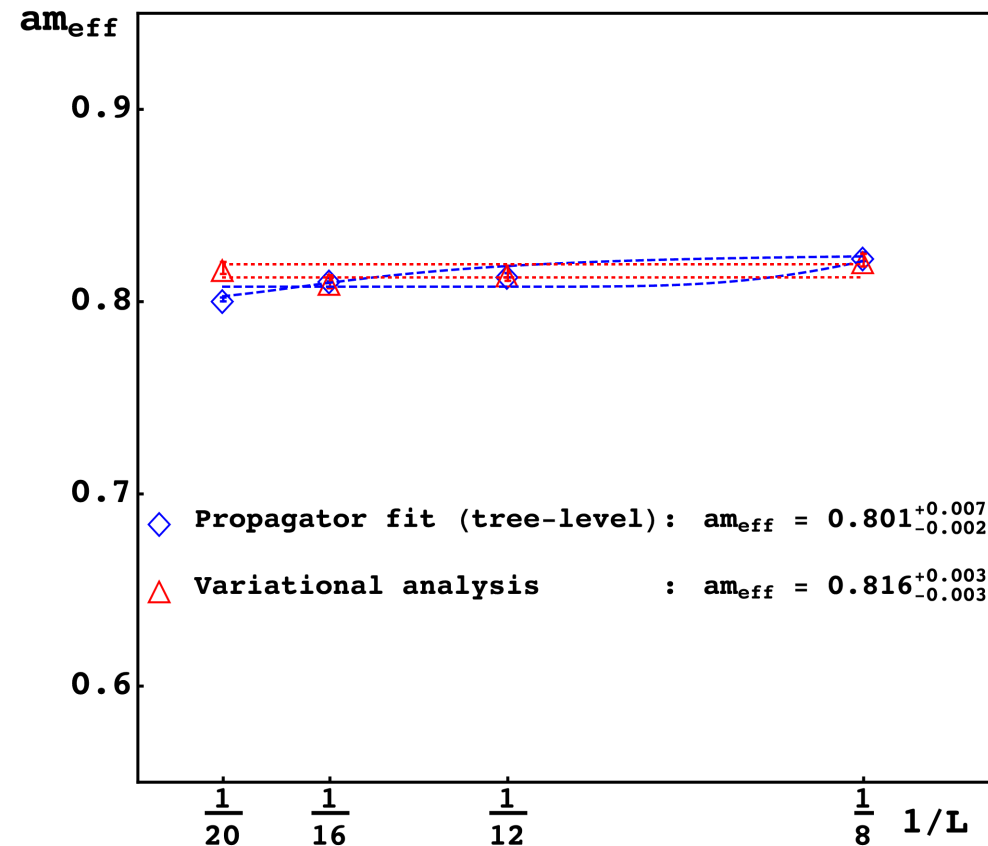
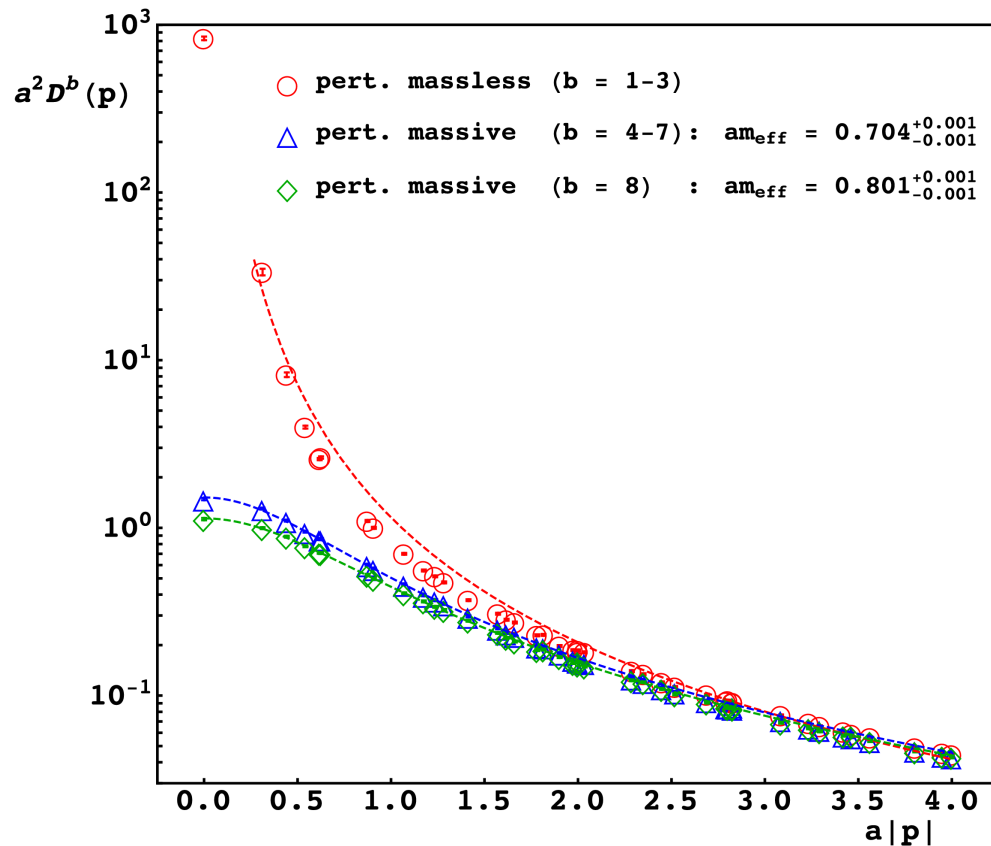
[Maas & Törek'16]



- Propagators almost tree-level
 - Expected splitting in gauge boson spectrum
- Physical vector: Massive, non-degenerate

Test for GUTs

[Maas & Törek'16]



- Propagators almost tree-level
 - Expected splitting in gauge boson spectrum
- Physical vector: Massive, non-degenerate
 - Agrees with FMS prediction

Example 3: Technicolor

No gauge-invariant perturbation theory
but
interesting implications

Implications for Technicolor

[Maas,'15]

- Higgs replaced by bound state of new fermions (techniquarks) and new gauge interaction (technicolor)

Implications for Technicolor

[Maas,'15]

- Higgs replaced by bound state of new fermions (techniquarks) and new gauge interaction (technicolor)
 - No BEH effect: FMS cannot work

Implications for Technicolor

[Maas,'15]

- Higgs replaced by bound state of new fermions (techniquarks) and new gauge interaction (technicolor)
 - No BEH effect: FMS cannot work
- Observable states must still be gauge-invariant
 - Needs to create Higgs and W/Z(!) signals by (new) bound states

Implications for Technicolor

[Maas,'15]

- Higgs replaced by bound state of new fermions (techniquarks) and new gauge interaction (technicolor)
 - No BEH effect: FMS cannot work
- Observable states must still be gauge-invariant
 - Needs to create Higgs and W/Z(!) signals by (new) bound states
 - Vectors must be lighter
 - Behavior not yet seen for strong interactions
 - Usually: Scalars and pseudoscalars

Summary

[Maas'12,'15
Törek & Maas'16]

- Observable spectrum must be gauge-invariant

Summary

[Maas'12,'15
Törek & Maas'16]

- Observable spectrum must be gauge-invariant
- In non-Abelian gauge theories: Bound states

Summary

[Maas'12,'15
Törek & Maas'16]

- Observable spectrum must be gauge-invariant
- In non-Abelian gauge theories: Bound states
- Gauge-invariant perturbation theory as a tool
 - Requires a Brout-Englert-Higgs effect
 - Yields the same results for the standard model
 - More robust
 - Mostly not much more complicated

Summary

[Maas'12,'15
Törek & Maas'16]

- Observable spectrum must be gauge-invariant
- In non-Abelian gauge theories: Bound states
- Gauge-invariant perturbation theory as a tool
 - Requires a Brout-Englert-Higgs effect
 - Yields the same results for the standard model
 - More robust
 - Mostly not much more complicated
- Applicable to beyond-the standard model
 - Structural requirement: Multiplets must match
 - Dynamical requirement: Small fluctuations
 - Verification requires non-perturbative methods

Advertisement



55th International Winter School on Theoretical Physics

Bound States and Resonances

13th-17th of February 2017

Lecturers: I. Belyaev, C. Fischer, C. Pica, S. Prelovsek, R. Roth, A. Szczepaniak

Admont, Styria, Austria

St. Goar 2017

Bound States in QCD and Beyond II

20th-23rd of February 2017

St. Goar, Germany

Registration open until 30th of November 2016