

Electroweak Sudakov Logarithms

Axel Maas

2nd of July 2024
Parton Showers and Resummation
Graz
Austria

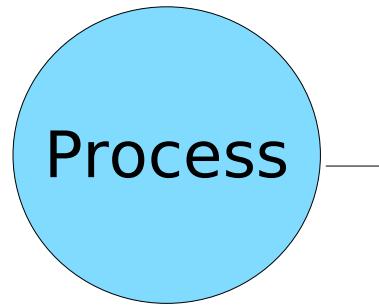


NAWI Graz
Natural Sciences

Österreichischer
Wissenschaftsfonds



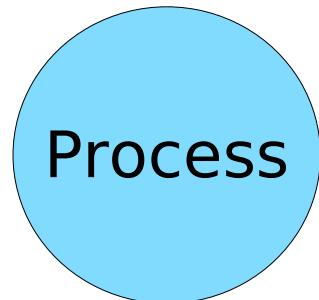
The issue in gauge theories



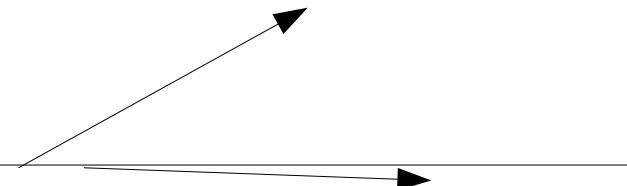
—

Incoming or
outgoing
gauged
particle

The issue in gauge theories



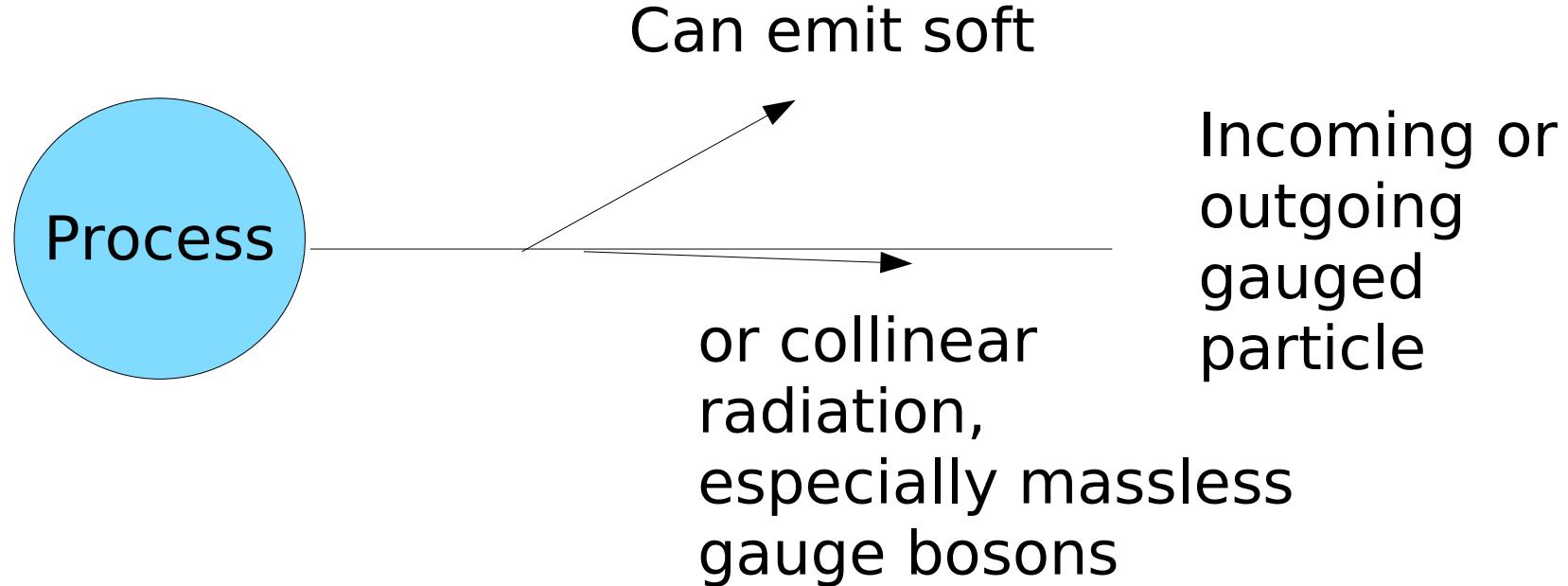
Can emit soft



or collinear
radiation,
especially massless
gauge bosons

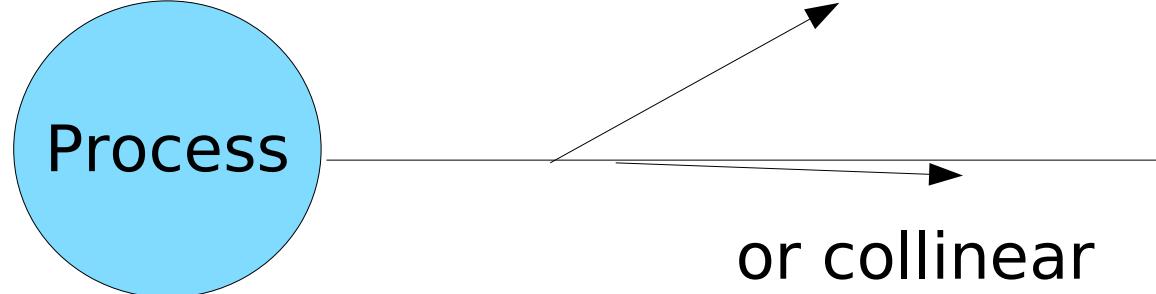
Incoming or
outgoing
gauged
particle

The issue in gauge theories



- Resummation yields (double) Sudakov logarithms: $\ln^2 \frac{S}{\Lambda}$
- Can dominate the cross sections
- Quick rise can obstruct convergence

Resolution in gauge theories



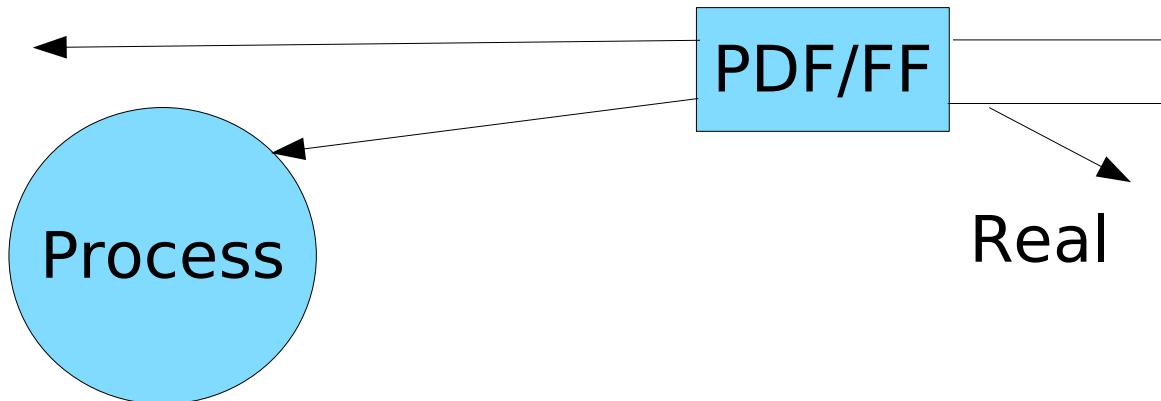
Can emit soft

or collinear
radiation,
especially massless
gauge bosons

Incoming or
outgoing
gauged
particle

- Involved particle is a gauge singlet

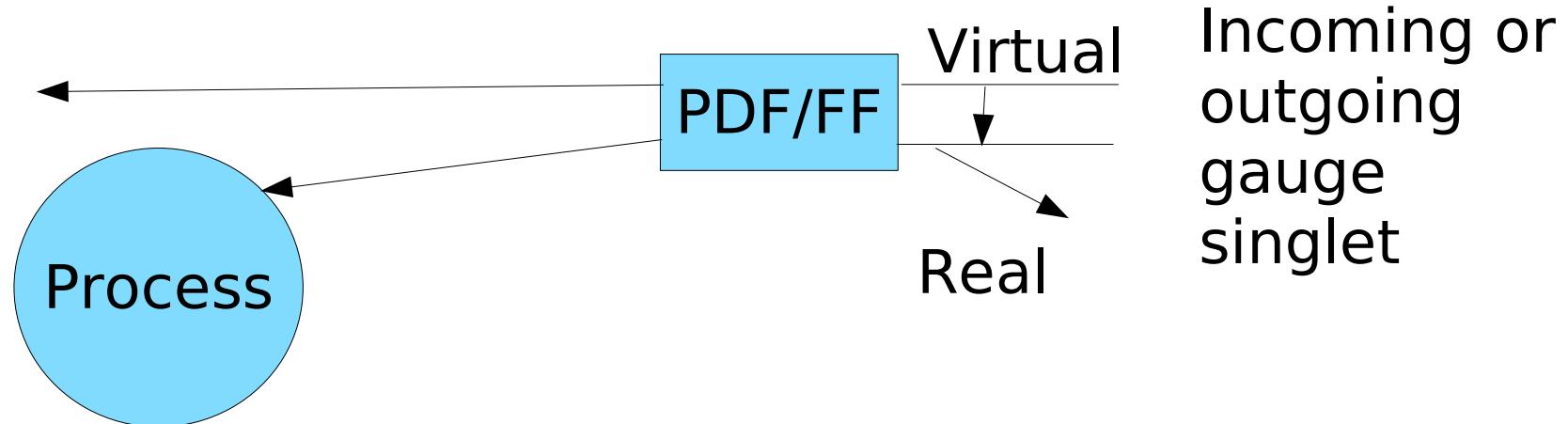
Resolution in gauge theories



Incoming or
outgoing
gauge
singlet

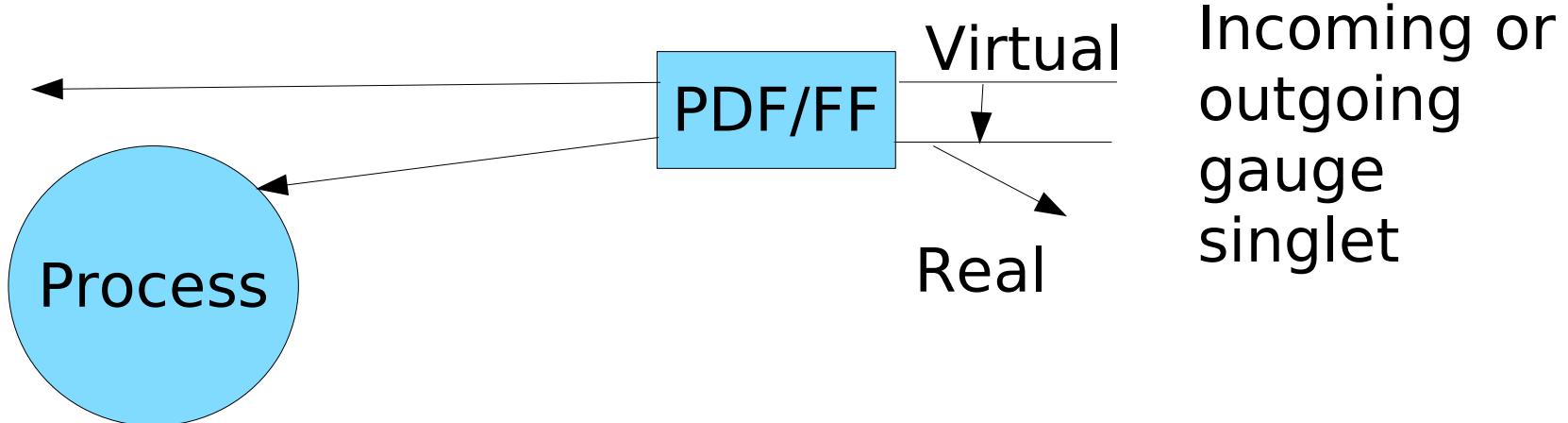
- Involved particle is a gauge singlet

Resolution in gauge theories



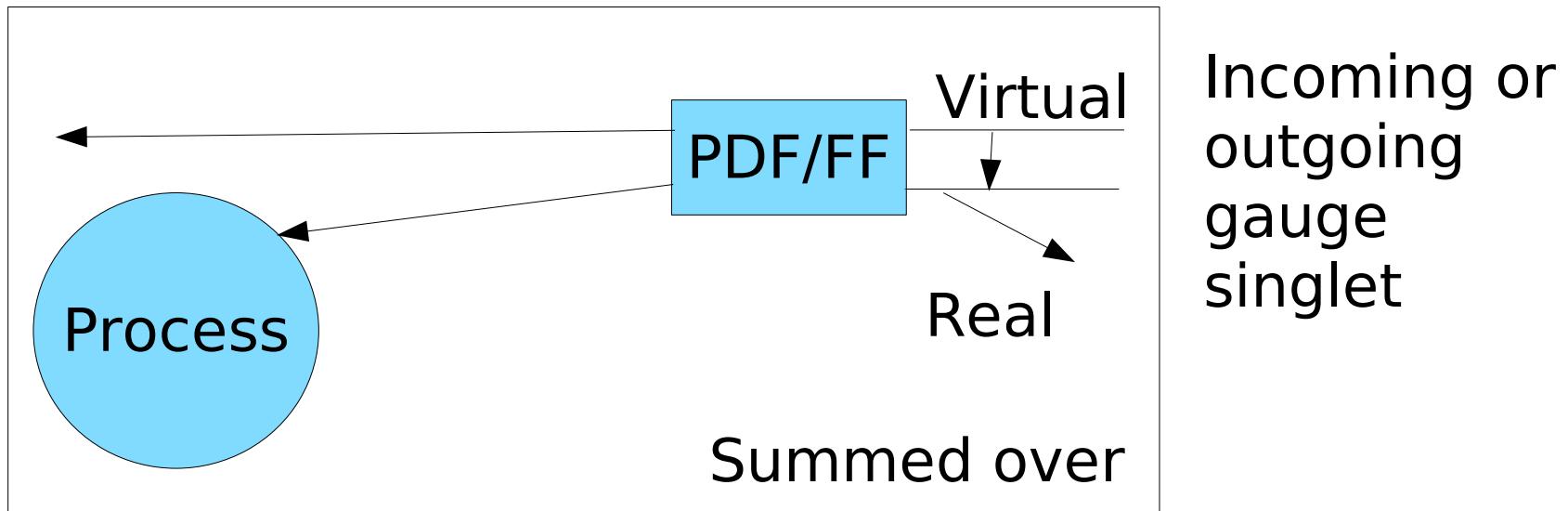
- Involved particle is a gauge singlet

Resolution in gauge theories



- Involved particle is a gauge singlet
 - Cancellations between real and virtual corrections
 - Cancel Sudakov logarithms

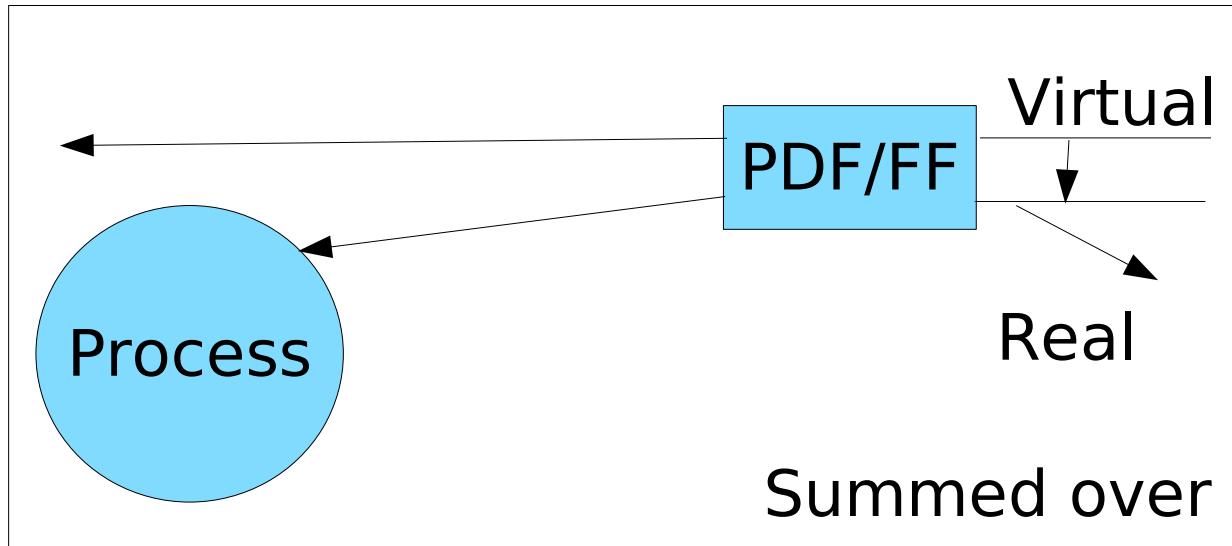
Resolution in gauge theories



- Involved particle is a gauge singlet
 - Cancellations between real and virtual corrections
 - Cancel Sudakov logarithms
- Initial state: Bloch-Nordsieck theorem/Initial state and final state: Kinoshita-Lee-Naunberg theorem
 - Required: Inclusive in the gauge charge

Invalidation in electroweak physics

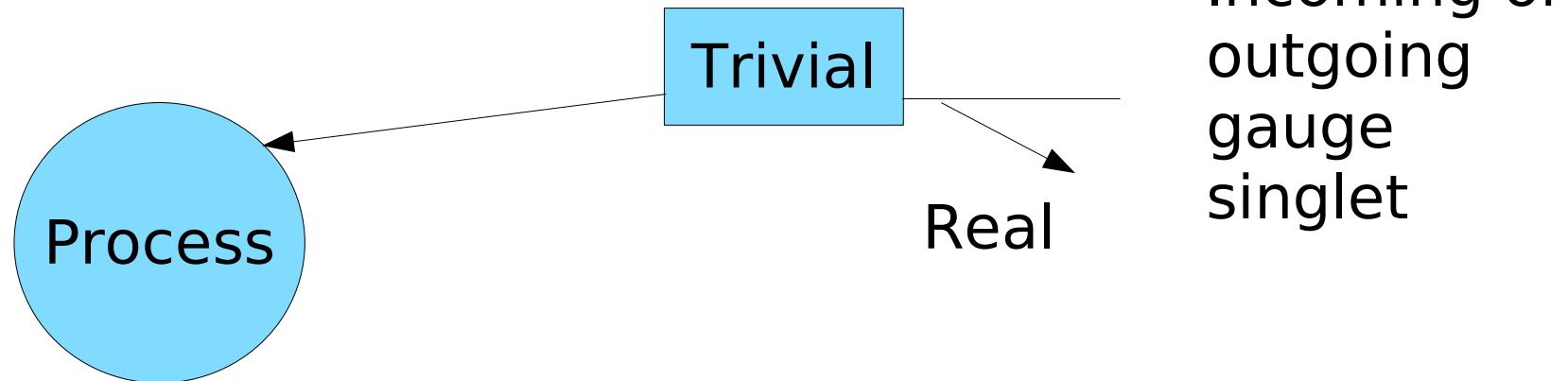
[Ciafaloni et al.'00,'22
Bauer et al.'18]



- Brout-Englert-Higgs effect “breaks” gauge symmetry

Invalidation in electroweak physics

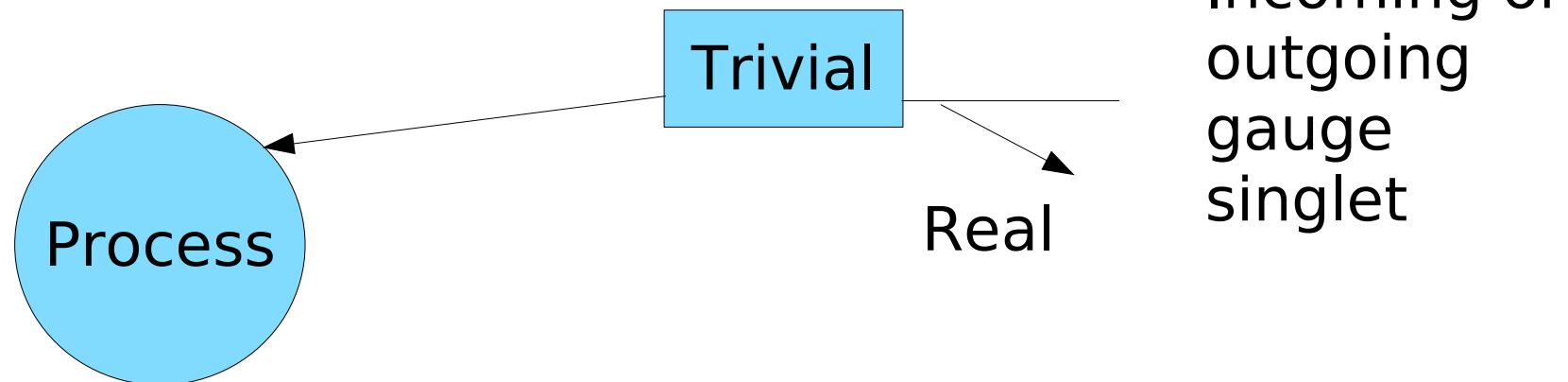
[Ciafaloni et al.'00,'22
Bauer et al.'18]



- Brout-Englert-Higgs effect “breaks” gauge symmetry
 - States like leptons become asymptotic states

Invalidation in electroweak physics

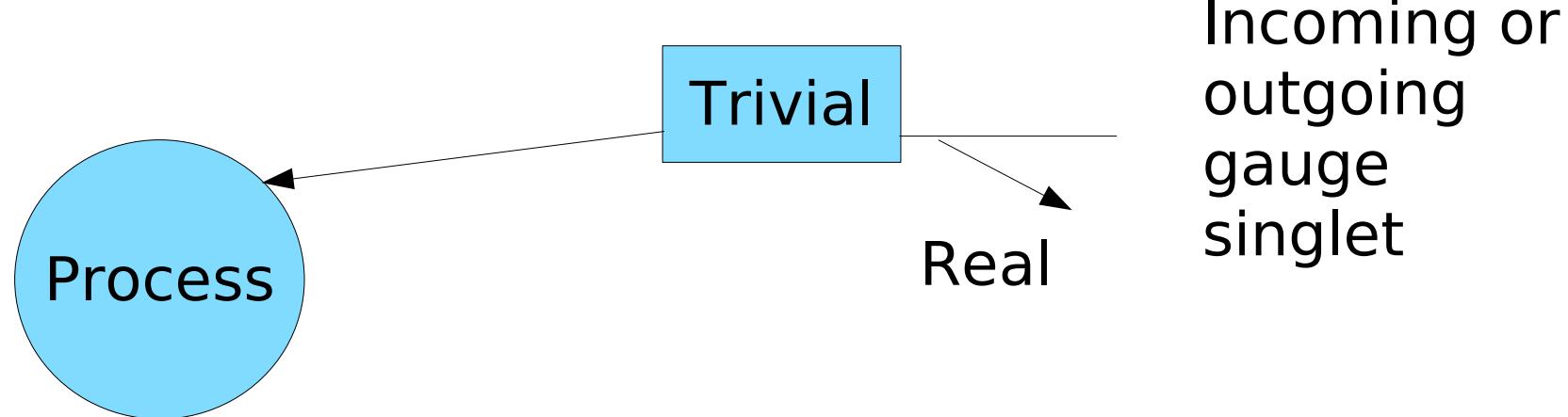
[Ciafaloni et al.'00,'22
Bauer et al.'18]



- Brout-Englert-Higgs effect “breaks” gauge symmetry
 - States like leptons become asymptotic states
 - BN and KLN theorems violated

Invalidation in electroweak physics

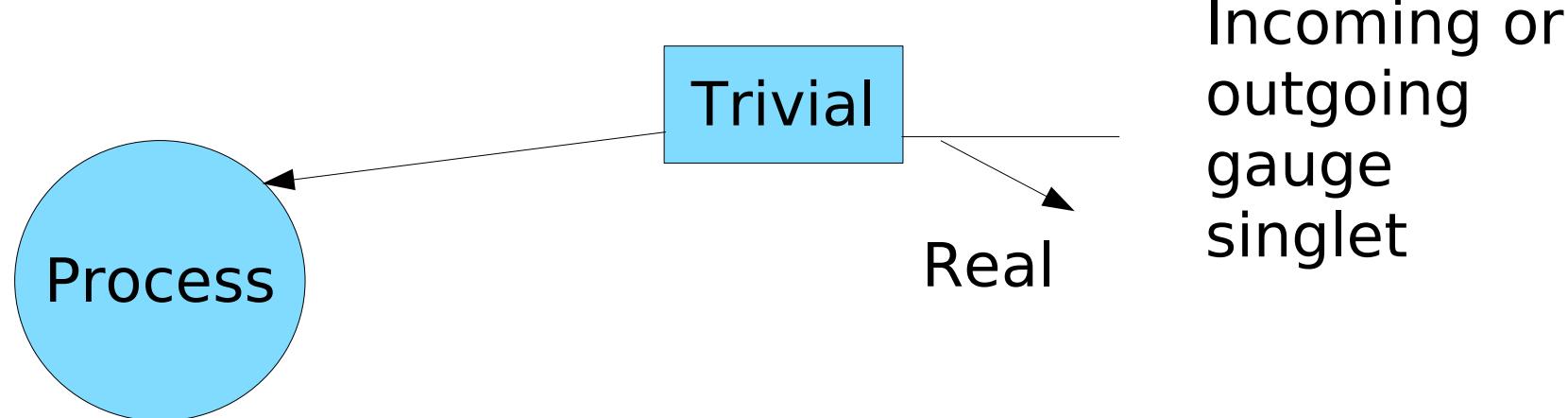
[Ciafaloni et al.'00,'22
Bauer et al.'18]



- Brout-Englert-Higgs effect “breaks” gauge symmetry
 - States like leptons become asymptotic states
 - BN and KLN theorems violated
- Double Sudakov logarithms suppressed: $\ln^2 \frac{S}{m_W^2}$
- Negligible at small energies

Invalidation in electroweak physics

[Ciafaloni et al.'00,'22
Bauer et al.'18]



- Brout-Englert-Higgs effect “breaks” gauge symmetry
 - States like leptons become asymptotic states
 - BN and KLN theorems violated
- Double Sudakov logarithms suppressed: $\ln^2 \frac{S}{m_W^2}$
- Negligible at small energies
- At LC@TeV: Same order as strong interactions
 - Swamped by jets of vector bosons and Higgs

Is this inevitable?

Is this inevitable?

- No!

Is this inevitable?

- No! - field theory to the rescue

Is this inevitable?

- No! - field theory to the rescue
- Problem: Asymptotic particles not a gauge singlet
 - Because of breaking the gauge symmetry

Is this inevitable?

[Fröhlich et al.'80,
Banks et al.'79]

- No! - field theory to the rescue
- Problem: Asymptotic particles not a gauge singlet
 - Because of breaking the gauge symmetry
 - But there is no physical gauge-symmetry breaking [Elitzur'75, Osterwalder & Seiler'77, Fradkin & Shenker'78]
 - Forbidden by Elitzur's theorem
 - Just a figure of speech
 - Actually just ordinary gauge-fixing

Is this inevitable?

[Fröhlich et al.'80,
Banks et al.'79]

- No! - field theory to the rescue
- Problem: Asymptotic particles not a gauge singlet
 - Because of breaking the gauge symmetry
 - But there is no physical gauge-symmetry breaking [Elitzur'75, Osterwalder & Seiler'77, Fradkin & Shenker'78]
 - Forbidden by Elitzur's theorem
 - Just a figure of speech
 - Actually just ordinary gauge-fixing
 - But it is well established?

Is this inevitable?

[Fröhlich et al.'80,
Banks et al.'79]

- No! - field theory to the rescue
- Problem: Asymptotic particles not a gauge singlet
 - Because of breaking the gauge symmetry
 - But there is no physical gauge-symmetry breaking [Elitzur'75, Osterwalder & Seiler'77, Fradkin & Shenker'78]
 - Forbidden by Elitzur's theorem
 - Just a figure of speech
 - Actually just ordinary gauge-fixing
- But it is well established?
 - Actually, a coincidence of the standard model
 - Subleading effects save the day

A toy model

A toy model: Higgs sector of the SM

A toy model: Higgs sector of the SM

- Consider an SU(2) with a fundamental scalar

A toy model: Higgs sector of the SM

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu}$$

$$W_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + g f_{bc}^a W_\mu^b W_\nu^c$$

- Ws 
- Coupling g and some numbers f^{abc}

A toy model: Higgs sector of the SM

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^\dagger D_{ik}^\mu h_k$$

$$W_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + g f_{bc}^a W_\mu^b W_\nu^c$$

$$D_\mu^{ij} = \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij}$$

- Ws W_μ^a 
- Higgs h_i 
- Coupling g and some numbers f^{abc} and t_a^{ij}

A toy model: Higgs sector of the SM

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^\dagger D_{ik}^\mu h_k + \lambda (h^a h_a^\dagger - v^2)^2$$

$$\begin{aligned} W_{\mu\nu}^a &= \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + g f_{bc}^a W_\mu^b W_\nu^c \\ D_\mu^{ij} &= \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij} \end{aligned}$$

- Ws W_μ^a 
- Higgs h_i 
- Couplings g, v, λ and some numbers f^{abc} and t_a^{ij}

A toy model: Higgs sector of the SM

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^\dagger D_{ik}^\mu h_k + \lambda (h^a h_a^\dagger - v^2)^2$$

$$\begin{aligned} W_{\mu\nu}^a &= \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + g f_{bc}^a W_\mu^b W_\nu^c \\ D_\mu^{ij} &= \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij} \end{aligned}$$

- Ws W_μ^a
- Higgs h_i
- Couplings g, v, λ and some numbers f^{abc} and t_a^{ij}
- Parameters selected for a BEH effect

A toy model: Symmetries

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^+ D_{ik}^\mu h_k + \lambda (h^a h_a^+ - v^2)^2$$

$$\begin{aligned} W_{\mu\nu}^a &= \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + g f_{bc}^a W_\mu^b W_\nu^c \\ D_\mu^{ij} &= \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij} \end{aligned}$$

A toy model: Symmetries

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^+ D_{ik}^\mu h_k + \lambda (h^a h_a^+ - v^2)^2$$

$$\begin{aligned} W_{\mu\nu}^a &= \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + g f_{bc}^a W_\mu^b W_\nu^c \\ D_\mu^{ij} &= \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij} \end{aligned}$$

- Local SU(2) gauge symmetry

$$W_\mu^a \rightarrow W_\mu^a + (\delta_b^a \partial_\mu - g f_{bc}^a W_\mu^c) \Phi^b \quad h_i \rightarrow h_i + g t_a^{ij} \Phi^a h_j$$

A toy model: Symmetries

- Consider an SU(2) with a fundamental scalar
- Essentially the standard model Higgs

$$L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^+ D_{ik}^\mu h_k + \lambda (h^a h_a^+ - v^2)^2$$

$$\begin{aligned} W_{\mu\nu}^a &= \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + g f_{bc}^a W_\mu^b W_\nu^c \\ D_\mu^{ij} &= \delta^{ij} \partial_\mu - ig W_\mu^a t_a^{ij} \end{aligned}$$

- Local SU(2) gauge symmetry

$$W_\mu^a \rightarrow W_\mu^a + (\delta_b^a \partial_\mu - g f_{bc}^a W_\mu^c) \phi^b \quad h_i \rightarrow h_i + g t_a^{ij} \phi^a h_j$$

- Global SU(2) custodial (flavor) symmetry

- Acts as (right-)transformation on the scalar field only

$$W_\mu^a \rightarrow W_\mu^a \quad h \rightarrow h \Omega$$

Textbook approach

- Choose parameters to get a Brout-Englert-Higgs effect

Textbook approach

- Choose parameters to get a Brout-Englert-Higgs effect
- Minimize the classical action

Textbook approach

- Choose parameters to get a Brout-Englert-Higgs effect
- Minimize the classical action
- Choose a suitable gauge and obtain ‘spontaneous gauge symmetry breaking’: $SU(2) \rightarrow 1$

Textbook approach

- Choose parameters to get a Brout-Englert-Higgs effect
- Minimize the classical action
- Choose a suitable gauge and obtain ‘spontaneous gauge symmetry breaking’: $SU(2) \rightarrow 1$
- Get masses and degeneracies at tree-level

Textbook approach

- Choose parameters to get a Brout-Englert-Higgs effect
- Minimize the classical action
- Choose a suitable gauge and obtain ‘spontaneous gauge symmetry breaking’: $SU(2) \rightarrow 1$
- Get masses and degeneracies at tree-level
- Perform perturbation theory

Physical spectrum

Perturbation theory

Mass

0

Physical spectrum

Perturbation theory

Scalar

fixed charge

Mass

0

Custodial singlet



Physical spectrum

Perturbation theory

Scalar

Vector

fixed charge

gauge triplet

Mass

0

Both custodial singlets



Physical states

[Fröhlich et al.'80,
Banks et al.'79]

- No “real” breaking

Physical states

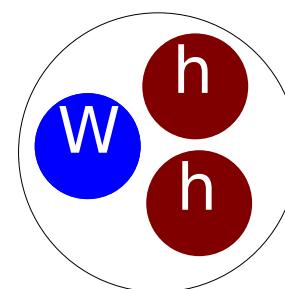
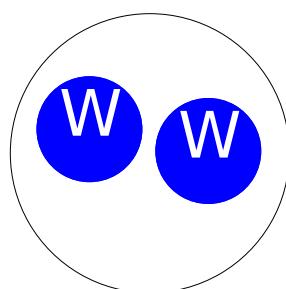
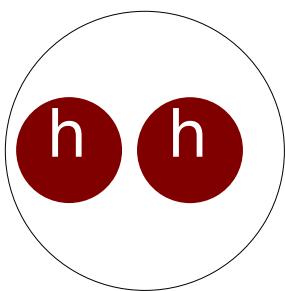
[Fröhlich et al.'80,
Banks et al.'79]

- No “real” breaking
- Physical particles are gauge-invariant particles
 - **Cannot** be the elementary particles
 - Non-Abelian nature is relevant

Physical states

[Fröhlich et al.'80,
Banks et al.'79]

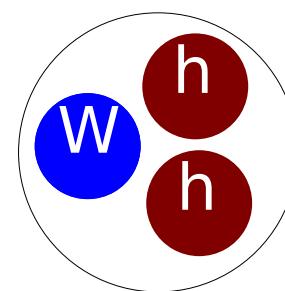
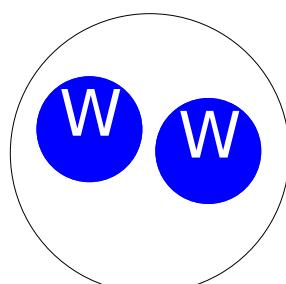
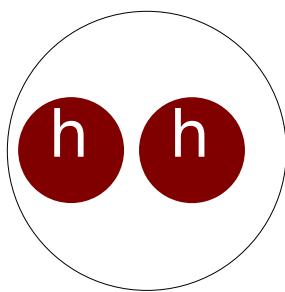
- No “real” breaking
- Physical particles are gauge-invariant particles
 - **Cannot** be the elementary particles
 - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
 - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



Physical states

[Fröhlich et al.'80,
Banks et al.'79]

- No “real” breaking
- Physical particles are gauge-invariant particles
 - **Cannot** be the elementary particles
 - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
 - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Has nothing to do with weak coupling
 - Think QED (hydrogen atom!)

Physical spectrum

Perturbation theory

Scalar

Vector

fixed charge

gauge triplet

Mass

0

Both custodial singlets



Remember: Experiment tells that somehow the left is correct!

Physical spectrum

Perturbation theory

Scalar

fixed charge

Vector

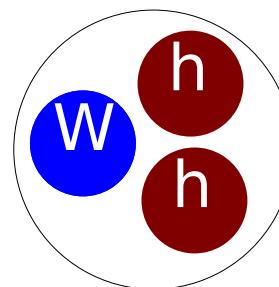
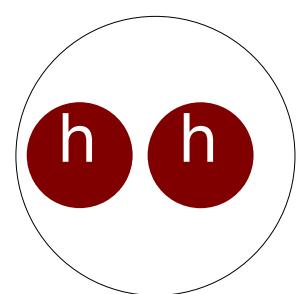
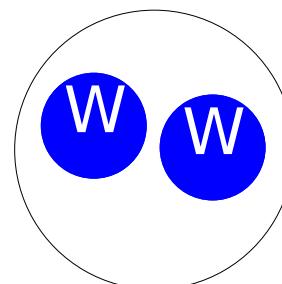
gauge triplet

Mass

0

Both custodial singlets

Composite (bound) states



Experiment tells that somehow the left is correct
Theory say the right is correct

Physical spectrum

Perturbation theory

Scalar

fixed charge

Vector

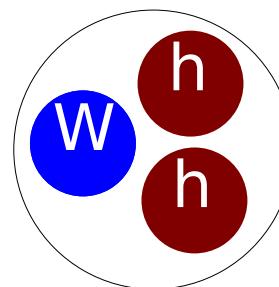
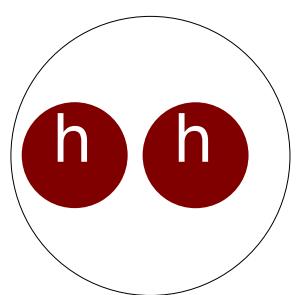
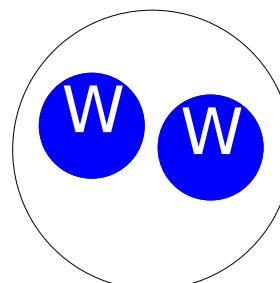
gauge triplet

Mass



Both custodial singlets

Composite (bound) states



Experiment tells that somehow the left is correct

Theory say the right is correct

There must exist a relation that both are correct

Physical particles

[Fröhlich et al.'80,'81,
Maas & Törek'16,'18,
Maas, Sondenheimer & Törek'17]

- J^{PC} and custodial charge only quantum numbers

Physical particles

[Fröhlich et al.'80,'81,
Maas & Törek'16,'18,
Maas, Sondenheimer & Törek'17]

- J^{PC} and custodial charge only quantum numbers
 - Different from perturbation theory
 - Operators limited to asymptotic, elementary, gauge-dependent states

Physical particles

[Fröhlich et al.'80,'81,
Maas & Törek'16,'18,
Maas, Sondenheimer & Törek'17]

- J^{PC} and custodial charge only quantum numbers
 - Different from perturbation theory
 - Operators limited to asymptotic, elementary, gauge-dependent states
 - Formulate gauge-invariant, composite operators

Physical particles

[Fröhlich et al.'80,'81,
Maas & Törek'16,'18,
Maas, Sondenheimer & Törek'17]

- J^{PC} and custodial charge only quantum numbers
 - Different from perturbation theory
 - Operators limited to asymptotic, elementary, gauge-dependent states
 - Formulate gauge-invariant, composite operators
 - Bound state structure

Physical particles

[Fröhlich et al.'80,'81,
Maas & Törek'16,'18,
Maas, Sondenheimer & Törek'17]

- J^{PC} and custodial charge only quantum numbers
 - Different from perturbation theory
 - Operators limited to asymptotic, elementary, gauge-dependent states
 - Formulate gauge-invariant, composite operators
 - Bound state structure – non-perturbative methods!

Physical particles

[Fröhlich et al.'80,'81,
Maas & Törek'16,'18,
Maas, Sondenheimer & Törek'17]

- J^{PC} and custodial charge only quantum numbers
 - Different from perturbation theory
 - Operators limited to asymptotic, elementary, gauge-dependent states
 - Formulate gauge-invariant, composite operators
 - Bound state structure - non-perturbative methods! - Lattice

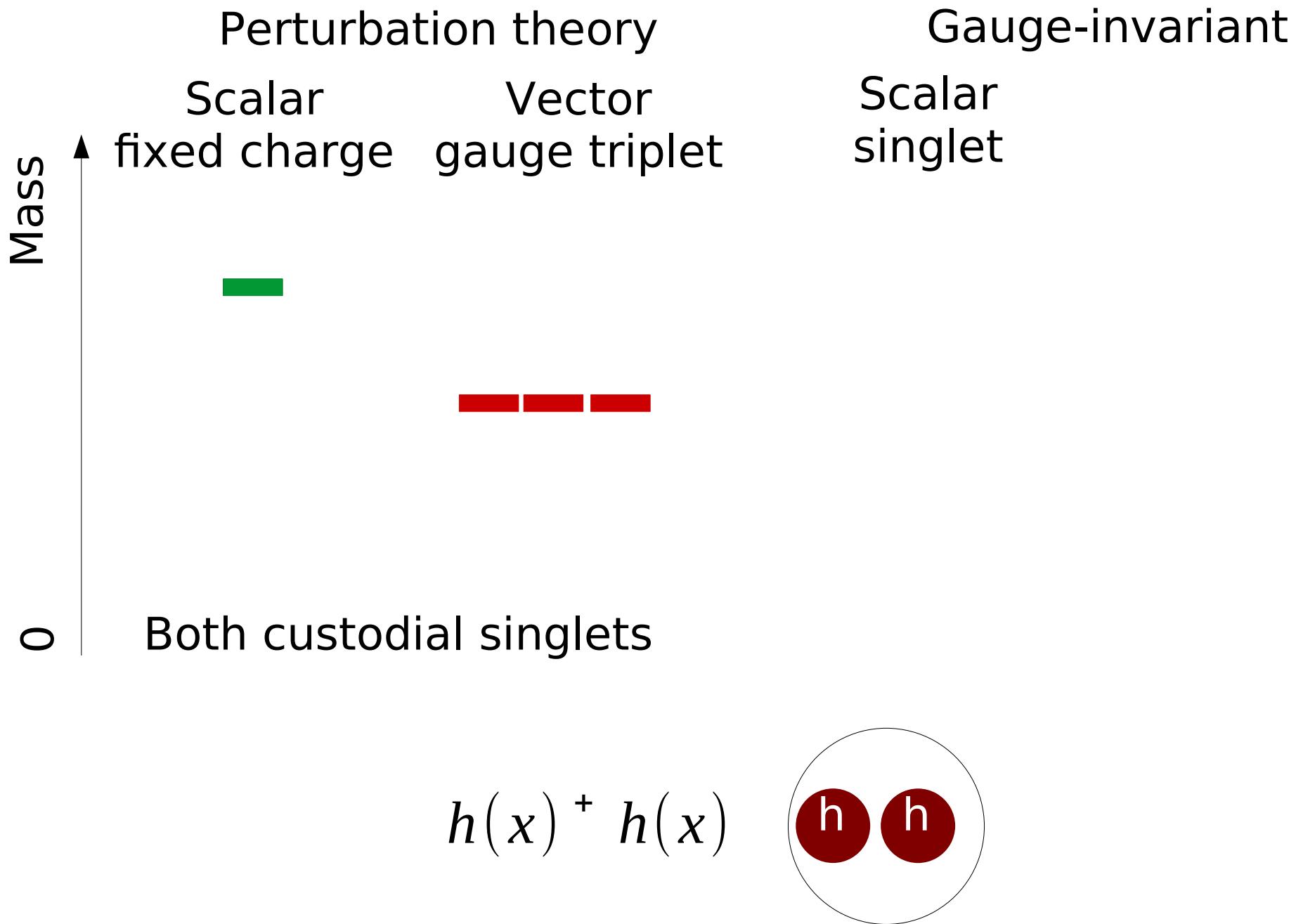
Physical particles

[Fröhlich et al.'80,'81,
Maas & Törek'16,'18,
Maas, Sondenheimer & Törek'17]

- J^{PC} and custodial charge only quantum numbers
 - Different from perturbation theory
 - Operators limited to asymptotic, elementary, gauge-dependent states
 - Formulate gauge-invariant, composite operators
 - Bound state structure - non-perturbative methods! - Lattice
 - Standard lattice spectroscopy problem
 - Standard methods
 - Smearing, variational analysis, systematic error analysis etc.
 - Very large statistics ($>10^5$ configurations)

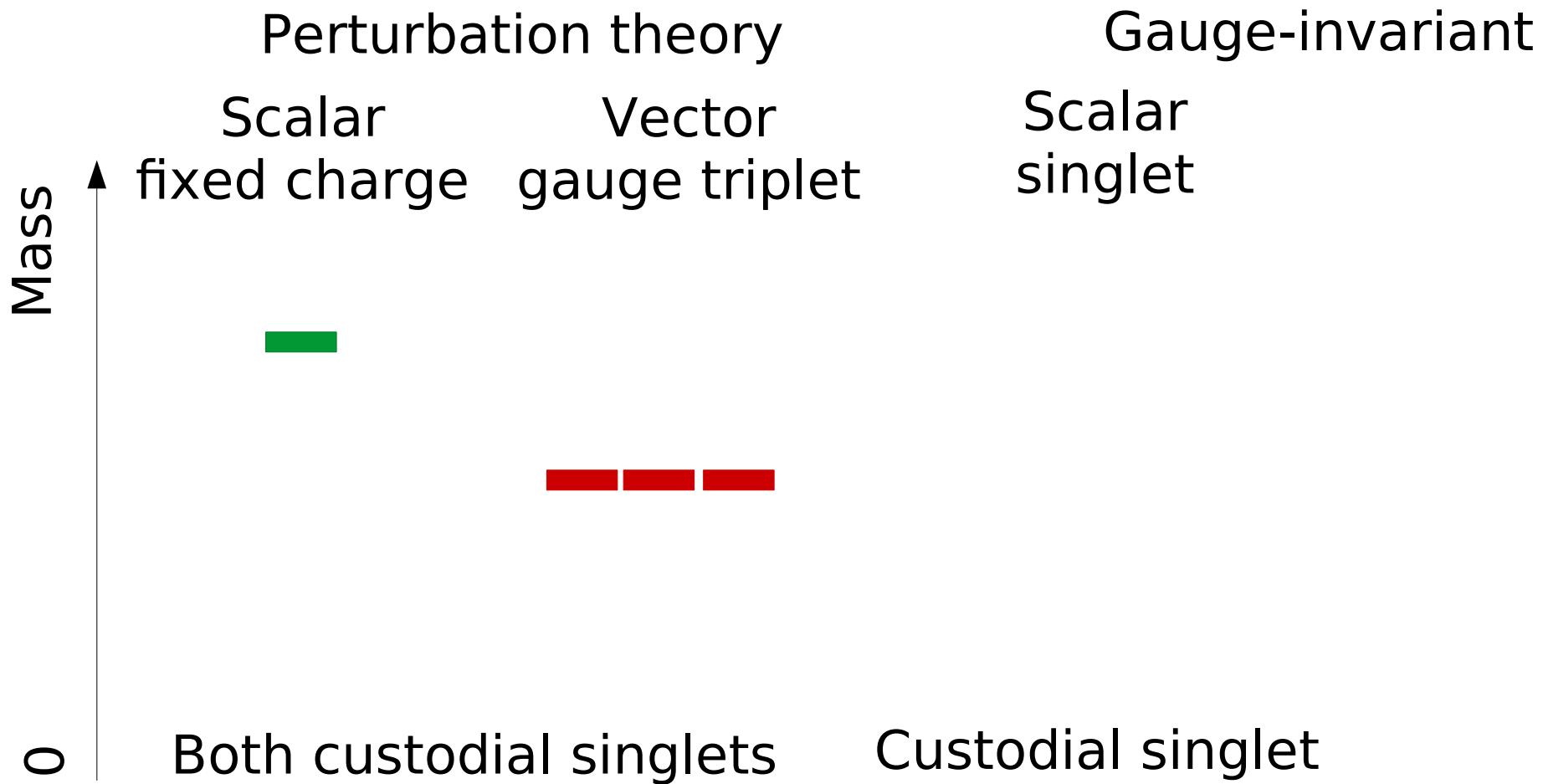
Physical spectrum

[Maas'12, Maas & Mufti'14]

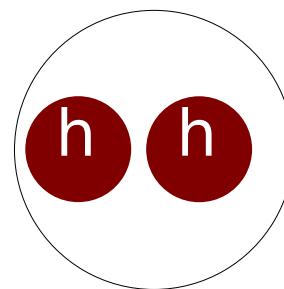


Physical spectrum

[Maas'12, Maas & Mufti'14]

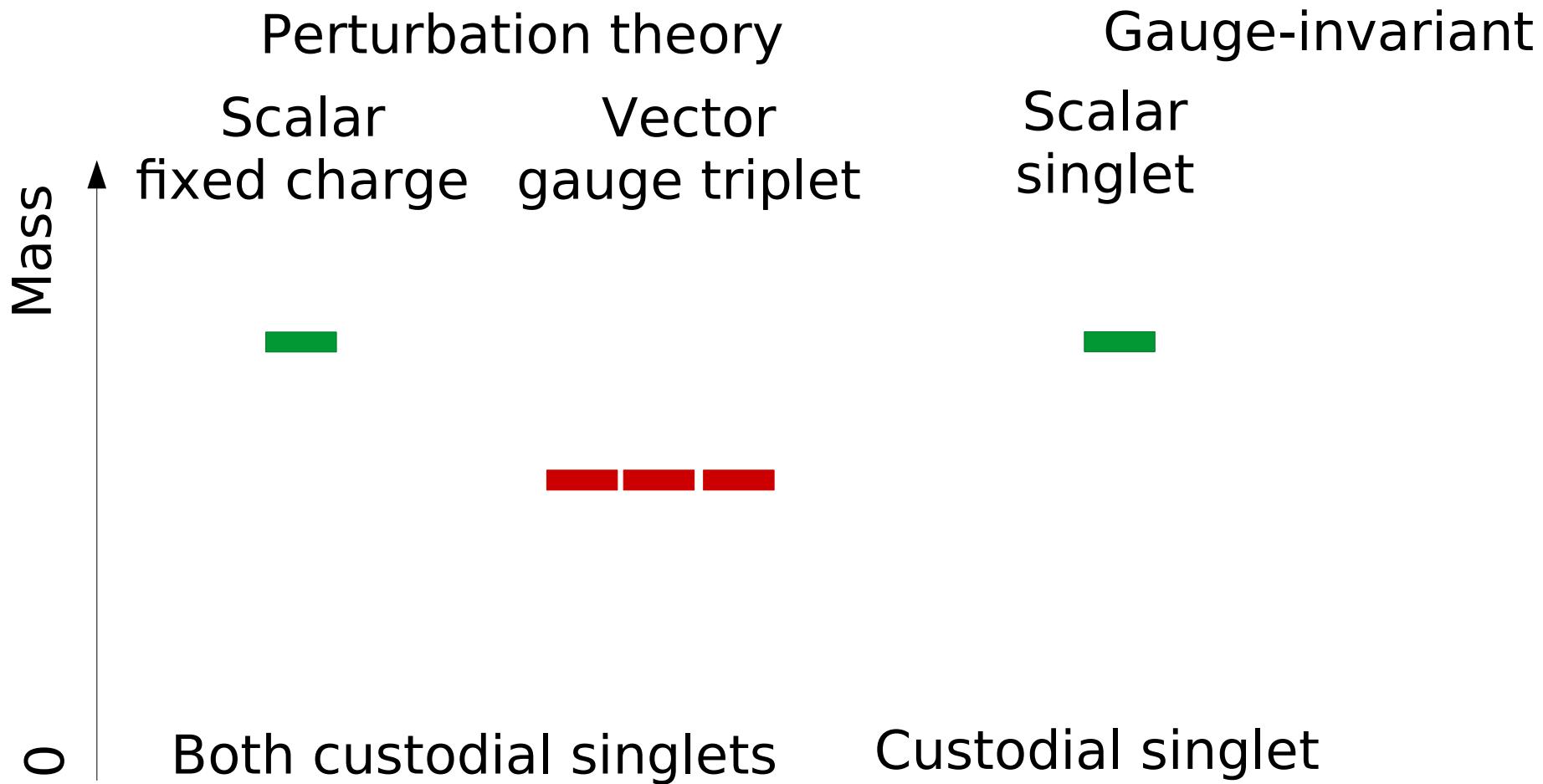


$$h(x)^+ h(x)$$



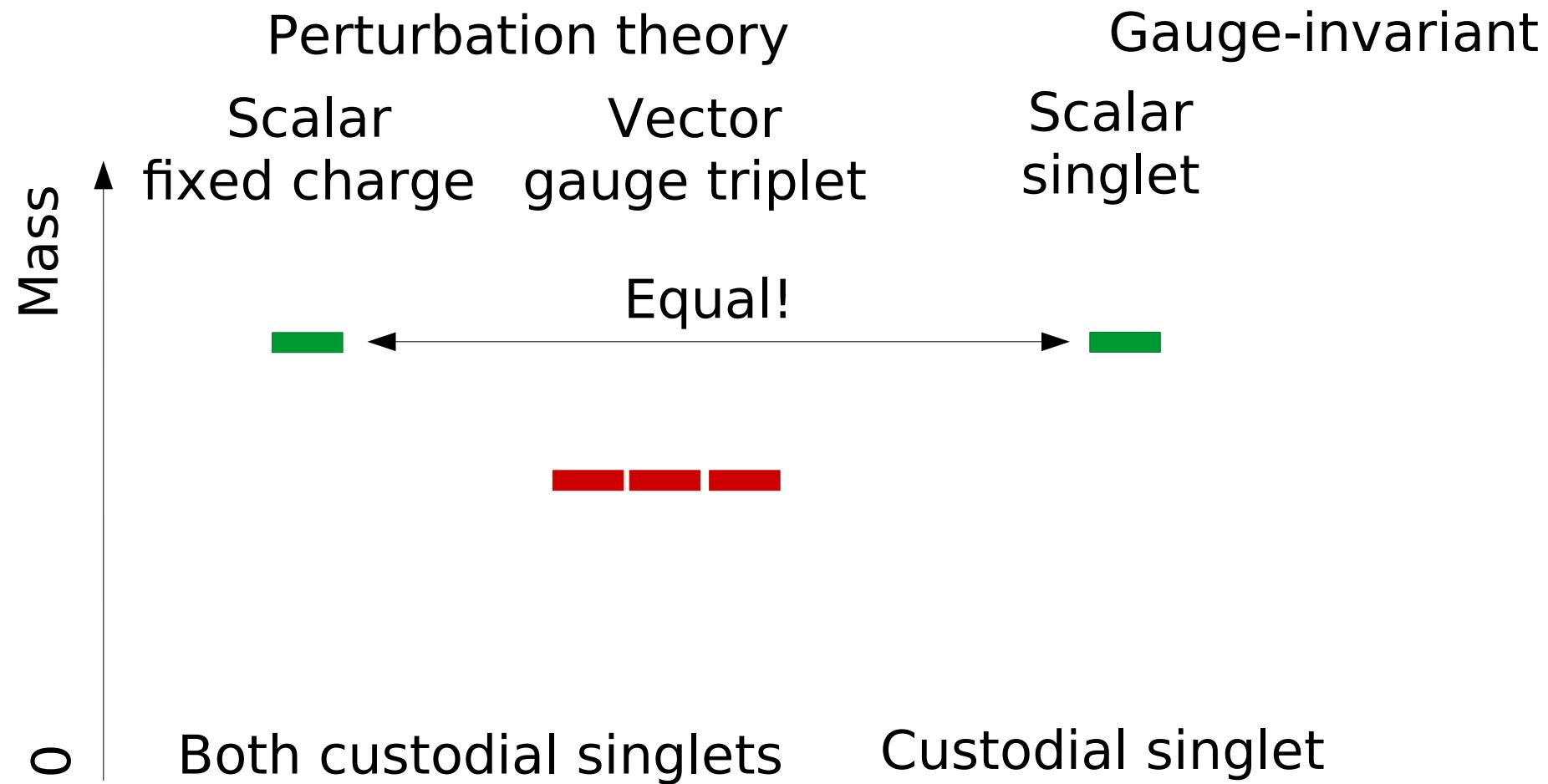
Physical spectrum

[Maas'12, Maas & Mufti'14]



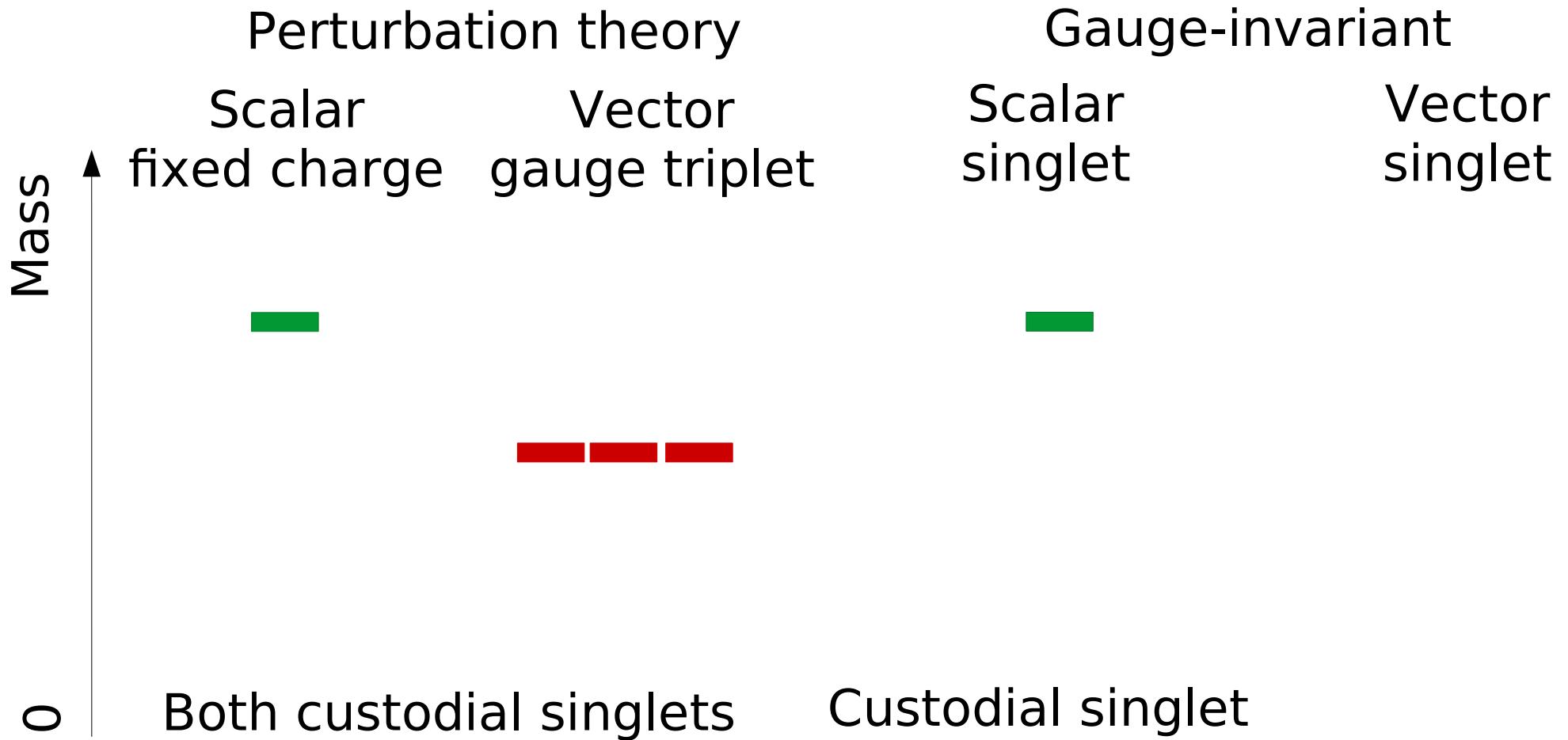
Physical spectrum

[Maas'12, Maas & Mufti'14]

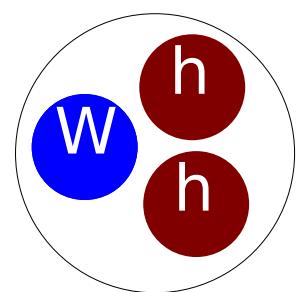


Physical spectrum

[Maas'12, Maas & Mufti'14]

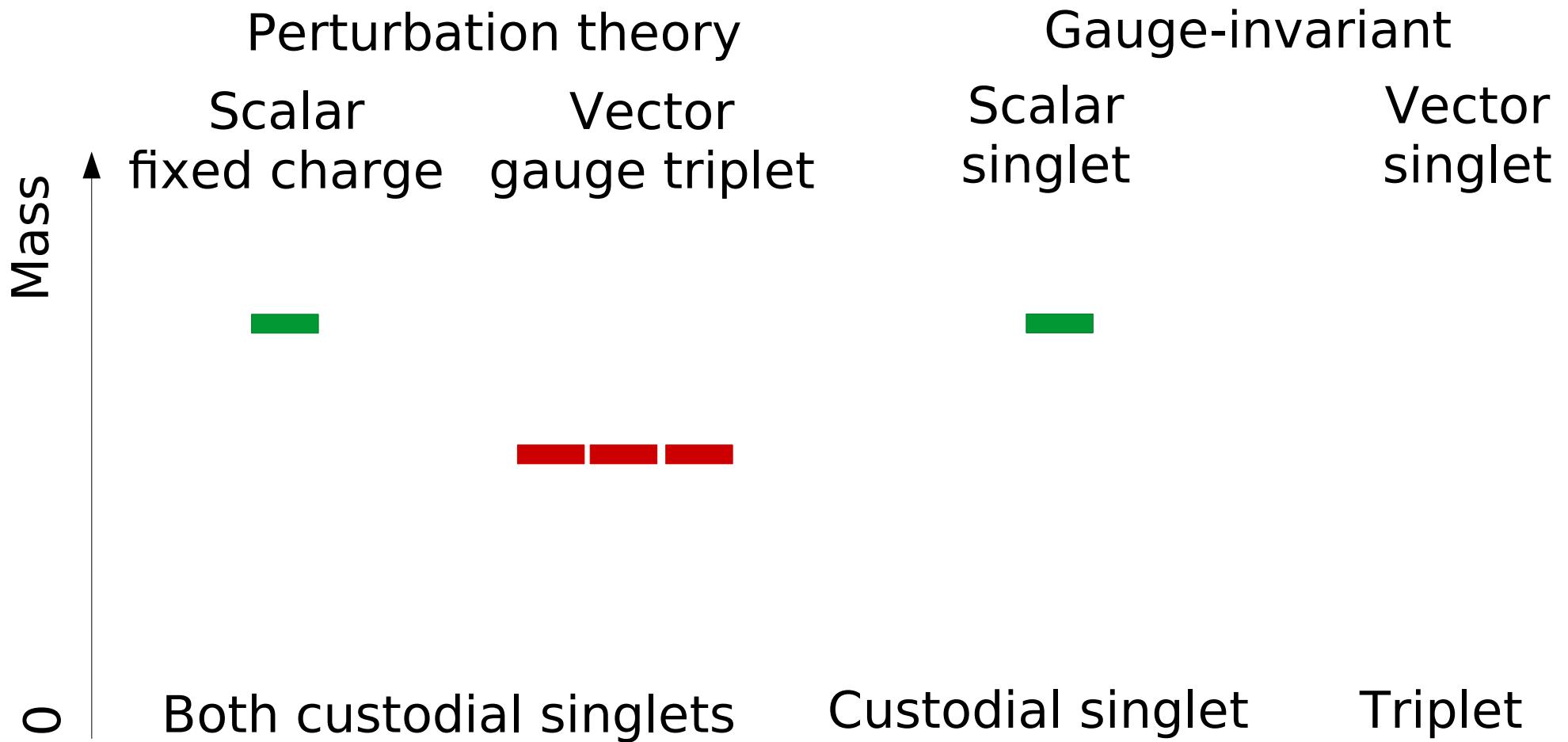


$$tr t^a \frac{h^+}{\sqrt{h^+ h}} D_\mu \frac{h}{\sqrt{h^+ h}}$$

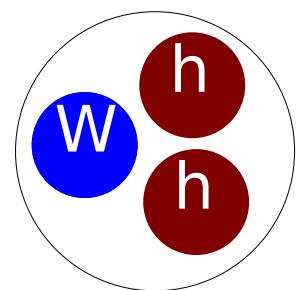


Physical spectrum

[Maas'12, Maas & Mufti'14]

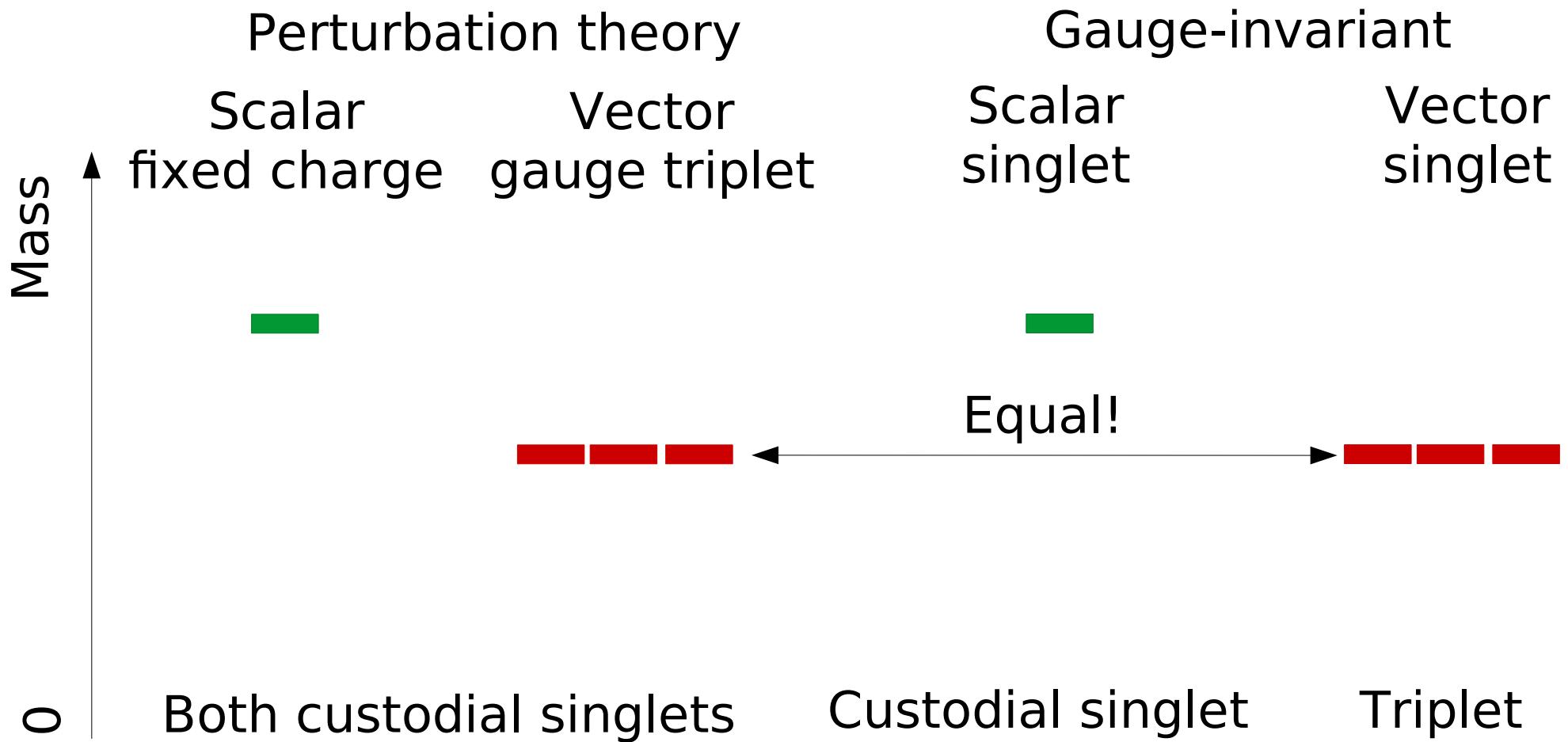


$$tr t^a \frac{h^+}{\sqrt{h^+ h}} D_\mu \frac{h}{\sqrt{h^+ h}}$$



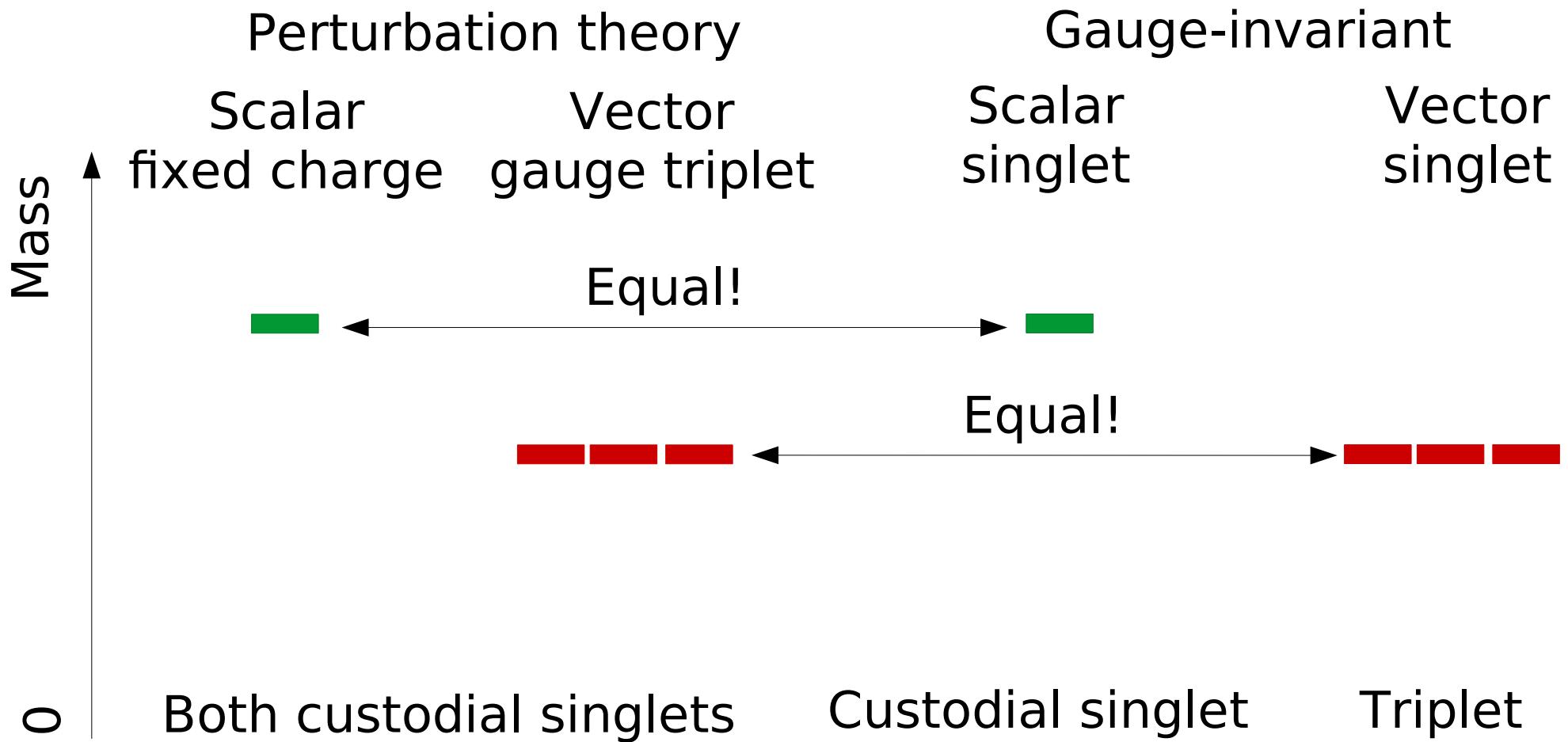
Physical spectrum

[Maas'12, Maas & Mufti'14]



Physical spectrum

[Maas'12, Maas & Mufti'14]



Why?

How to make predictions

[Fröhlich et al.'80,'81,
Maas & Törek'16,'18,
Maas, Sondenheimer & Törek'17
Maas & Sondenheimer '20]

- J^{PC} and custodial charge only quantum numbers
 - Different from perturbation theory
 - Operators limited to asymptotic, elementary, gauge-dependent states
 - Formulate gauge-invariant, composite operators
 - Bound state structure – non-perturbative methods?

How to make predictions

[Fröhlich et al.'80,'81,
Maas & Törek'16,'18,
Maas, Sondenheimer & Törek'17
Maas & Sondenheimer '20]

- J^{PC} and custodial charge only quantum numbers
 - Different from perturbation theory
 - Operators limited to asymptotic, elementary, gauge-dependent states
 - Formulate gauge-invariant, composite operators
 - Bound state structure – non-perturbative methods?
 - But coupling is still weak and there is a BEH

How to make predictions

[Fröhlich et al.'80,'81,
Maas & Törek'16,'18,
Maas, Sondenheimer & Törek'17
Maas & Sondenheimer '20]

- J^{PC} and custodial charge only quantum numbers
 - Different from perturbation theory
 - Operators limited to asymptotic, elementary, gauge-dependent states
 - Formulate gauge-invariant, composite operators
 - Bound state structure – non-perturbative methods?
 - But coupling is still weak and there is a BEH
 - Perform double expansion [Fröhlich et al.'80, Maas'12]
 - Vacuum expectation value (FMS mechanism)
 - Standard expansion in couplings
 - Together: Augmented perturbation theory

Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

0^+ singlet: $\langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle$



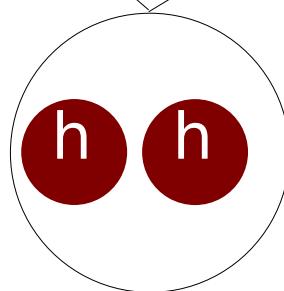
Higgs field

Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

0^+ singlet: $\langle (h^+ h)(x) (h^+ h)(y) \rangle$



Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

0^+ singlet: $\langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle$

2) Expand Higgs field around fluctuations $h = v + \eta$

Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

0^+ singlet: $\langle (h^+ h)(x)(h^+ h)(y) \rangle$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\begin{aligned}\langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x)\eta(y) \rangle \\ &+ v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle\end{aligned}$$

Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^+ h)(x)(h^+ h)(y) \rangle$$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x)\eta(y) \rangle \\ &+ v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle \end{aligned}$$

3) Standard perturbation theory

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x)\eta(y) \rangle \\ &+ \langle \eta^+(x)\eta(y) \rangle \langle \eta^+(x)\eta(y) \rangle + O(g, \lambda) \end{aligned}$$

Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^+ h)(x)(h^+ h)(y) \rangle$$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x)\eta(y) \rangle \\ &+ v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle \end{aligned}$$

3) Standard perturbation theory

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x)\eta(y) \rangle \\ &+ \langle \eta^+(x)\eta(y) \rangle \langle \eta^+(x)\eta(y) \rangle + O(g, \lambda) \end{aligned}$$

4) Compare poles on both sides

Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^+ h)(x)(h^+ h)(y) \rangle$$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x)\eta(y) \rangle \\ &\quad + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle \end{aligned}$$

3) Standard perturbation theory

Bound state mass \rightarrow

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x)\eta(y) \rangle \\ &\quad + \langle \eta^+(x)\eta(y) \rangle \langle \eta^+(x)\eta(y) \rangle + O(g, \lambda) \end{aligned}$$

4) Compare poles on both sides

Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^+ h)(x)(h^+ h)(y) \rangle$$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x)\eta(y) \rangle \\ &+ v \langle \eta^+\eta^2 + \eta^{+2}\eta \rangle + \langle \eta^{+2}\eta^2 \rangle \end{aligned}$$

3) Standard perturbation theory

Bound
state

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+(x)\eta(y) \rangle$$

mass

$$+ \langle \eta^+(x)\eta(y) \rangle \langle \eta^+(x)\eta(y) \rangle + O(g, \lambda)$$

Trivial two-particle state

4) Compare poles on both sides

Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^+ h)(x)(h^+ h)(y) \rangle$$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x) \eta(y) \rangle \\ &+ v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle \end{aligned}$$

3) Standard perturbation theory

Bound state mass

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x) \eta(y) \rangle \\ &+ \langle \eta^+(x) \eta(y) \rangle \langle \eta^+(x) \eta(y) \rangle + O(g, \lambda) \end{aligned}$$

Higgs mass

4) Compare poles on both sides

Augmented perturbation theory

[Fröhlich et al.'80, '81
Maas'12, '17]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^+ h)(x)(h^+ h)(y) \rangle$$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x)\eta(y) \rangle \\ &+ v \langle \eta^+\eta^2 + \eta^{+2}\eta \rangle + \langle \eta^{+2}\eta^2 \rangle \end{aligned}$$

Standard
Perturbation
Theory

3) Standard perturbation theory

Bound
state
mass

$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x)\eta(y) \rangle \\ &+ \langle \eta^+(x)\eta(y) \rangle \langle \eta^+(x)\eta(y) \rangle + O(g, \lambda) \end{aligned}$$

4) Compare poles on both sides

Augmented perturbation theory

[Fröhlich et al.'80,'81
Maas'12,'17
Maas & Sondenheimer'20]

1) Formulate gauge-invariant operator

$$0^+ \text{ singlet: } \langle (h^+ h)(x)(h^+ h)(y) \rangle$$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+(x)\eta(y) \rangle$$

What about this?

$$+v\langle \eta^+\eta^2 + \eta^{+2}\eta \rangle + \langle \eta^{+2}\eta^2 \rangle$$

3) Standard perturbation theory

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+(x)\eta(y) \rangle$$
$$+ \langle \eta^+(x)\eta(y) \rangle \langle \eta^+(x)\eta(y) \rangle + O(g, \lambda)$$

4) Compare poles on both sides

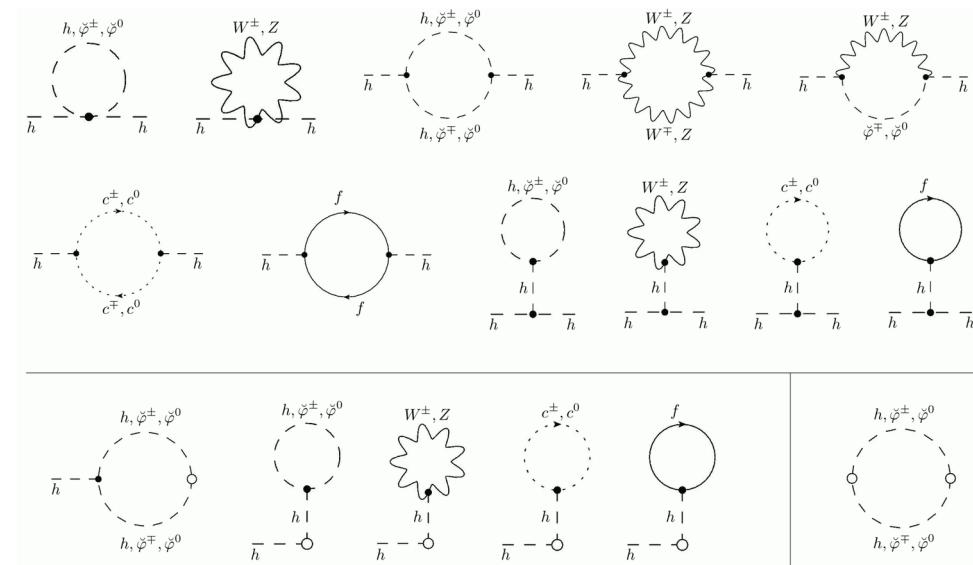
Consequences: The Higgs

[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]

$$\begin{aligned}\langle (h^+ h)(x) (h^+ h)(y) \rangle &= v^2 \langle \eta^+(x) \eta(y) \rangle \\ &+ v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle\end{aligned}$$

Consequences: The Higgs

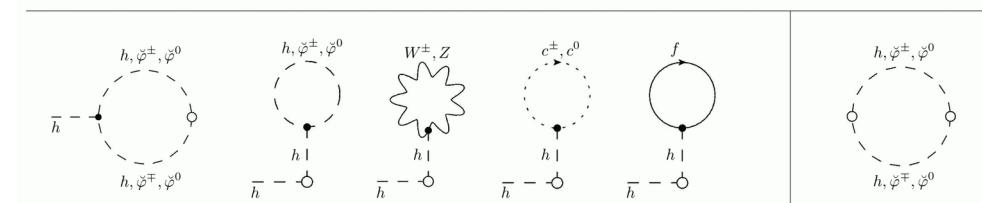
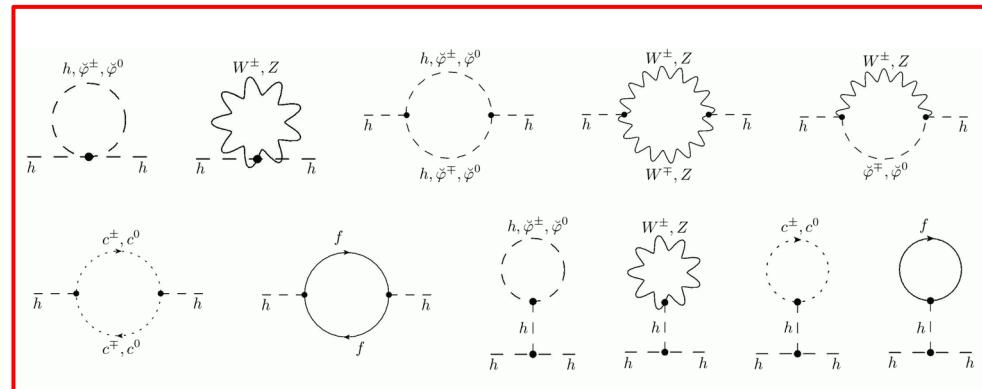
[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]



$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x) \eta(y) \rangle \\ &+ v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle \end{aligned}$$

Consequences: The Higgs

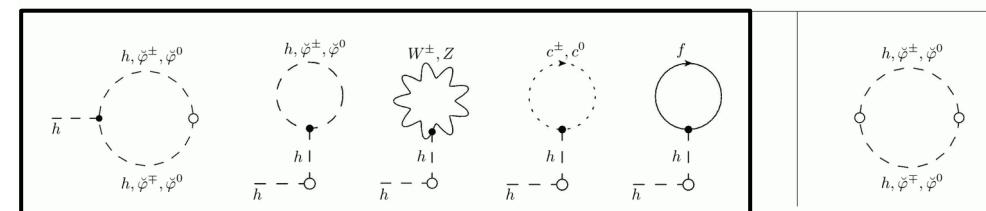
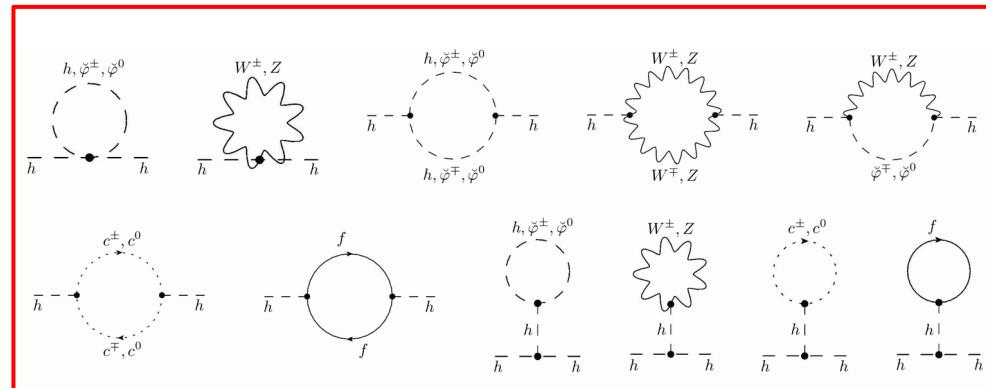
[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]



$$\begin{aligned} \langle (h^\dagger h)(x)(h^\dagger h)(y) \rangle &= v^2 \langle \eta^\dagger(x) \eta(y) \rangle \\ &+ v \langle \eta^\dagger \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle \end{aligned}$$

Consequences: The Higgs

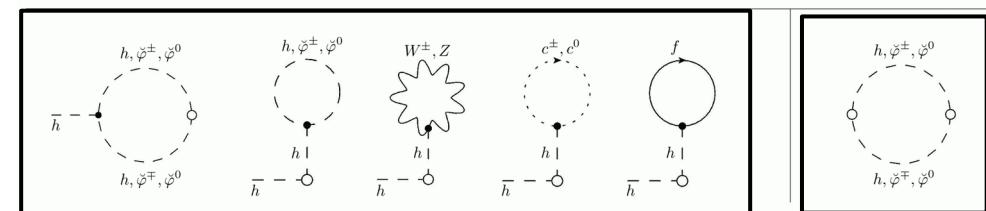
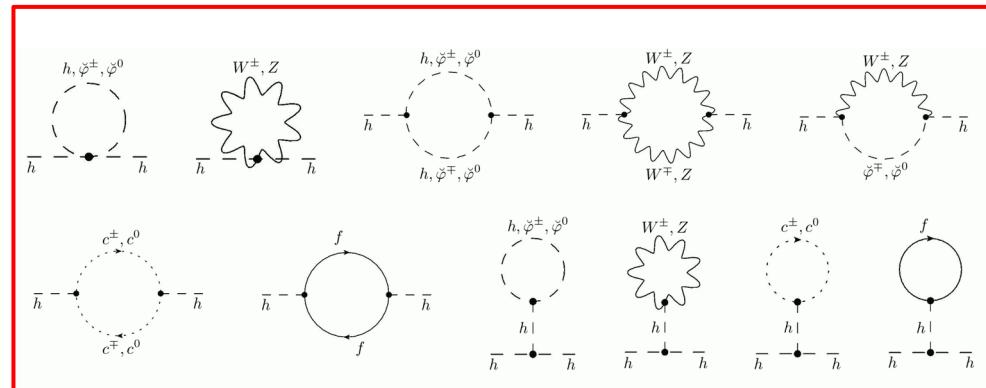
[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]



$$\begin{aligned} \langle (h^+ h)(x)(h^+ h)(y) \rangle &= v^2 \langle \eta^+(x) \eta(y) \rangle \\ &+ v \left[\langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle \right] \end{aligned}$$

Consequences: The Higgs

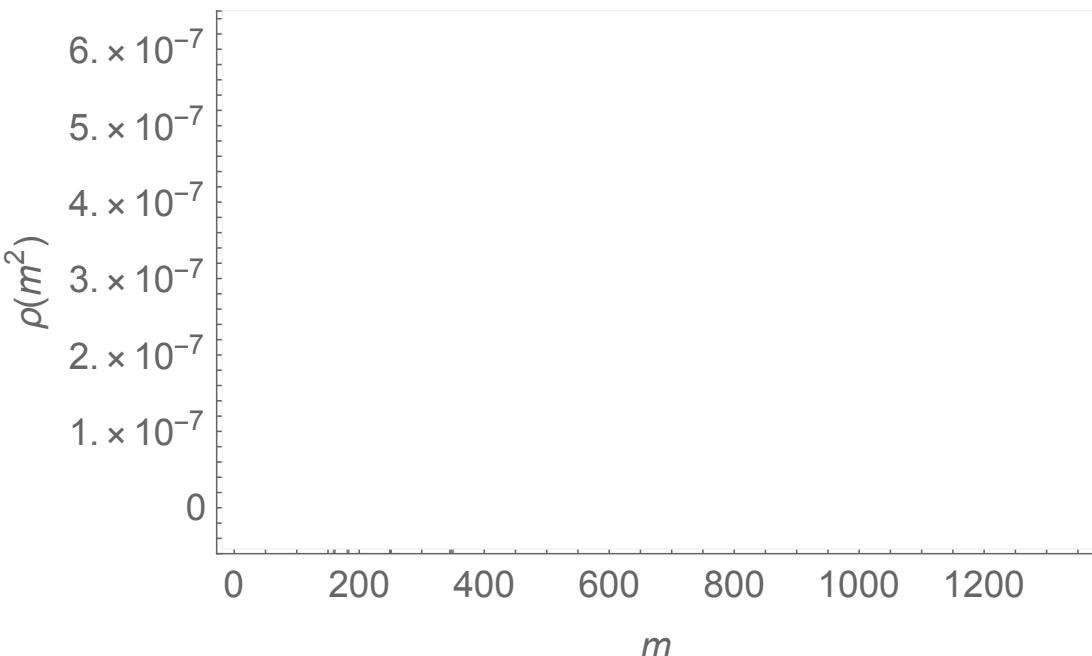
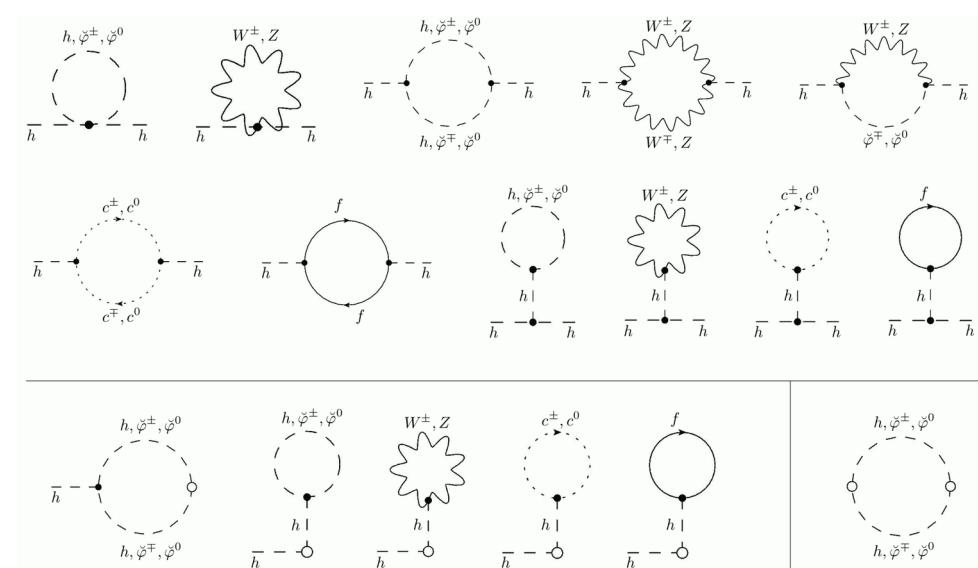
[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]



$$\begin{aligned} \langle (h^+ h)(x) (h^+ h)(y) \rangle &= v^2 \langle \eta^+(x) \eta(y) \rangle \\ &+ v \left[\langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle \right] \end{aligned}$$

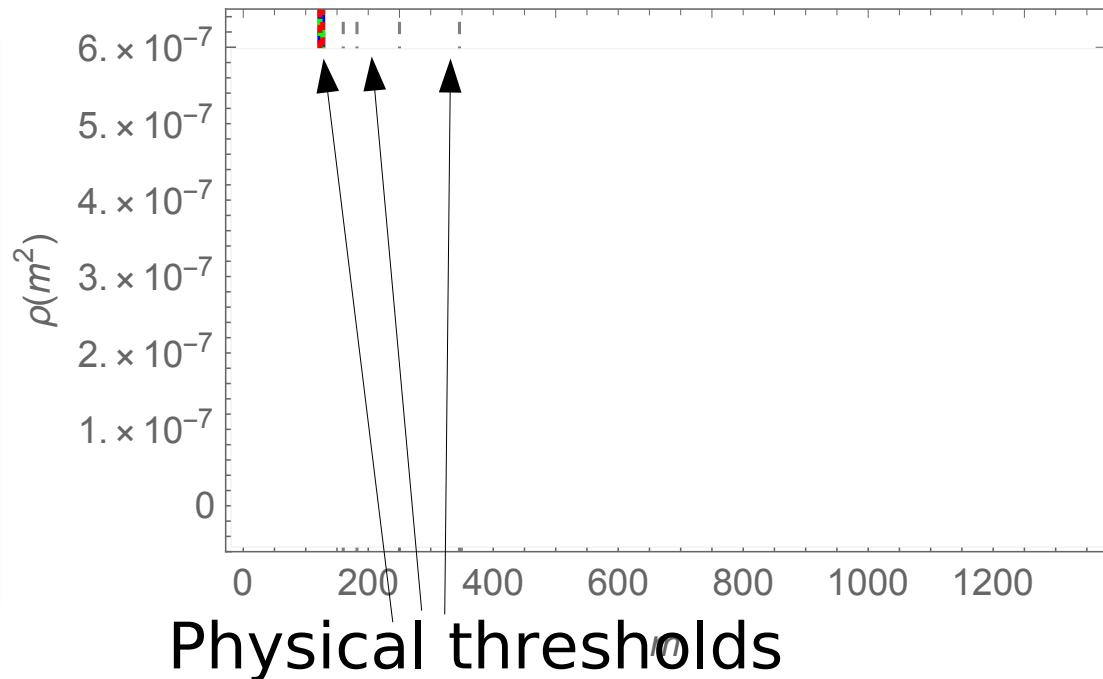
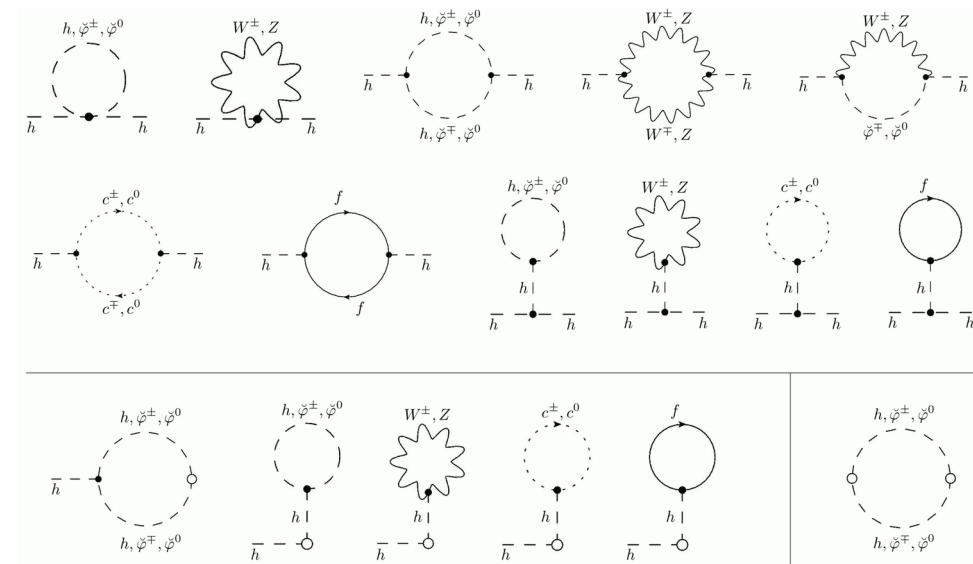
Consequences: The Higgs

[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]



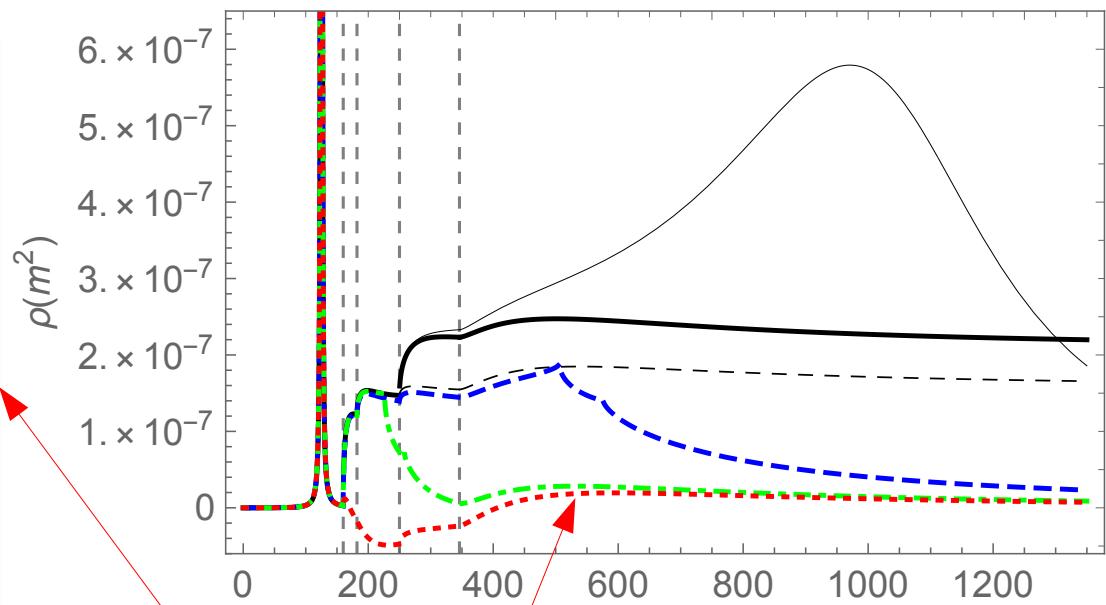
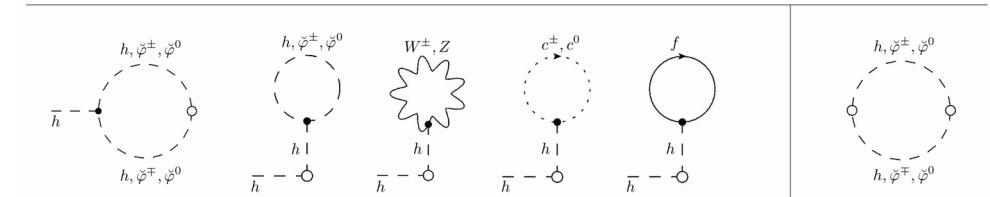
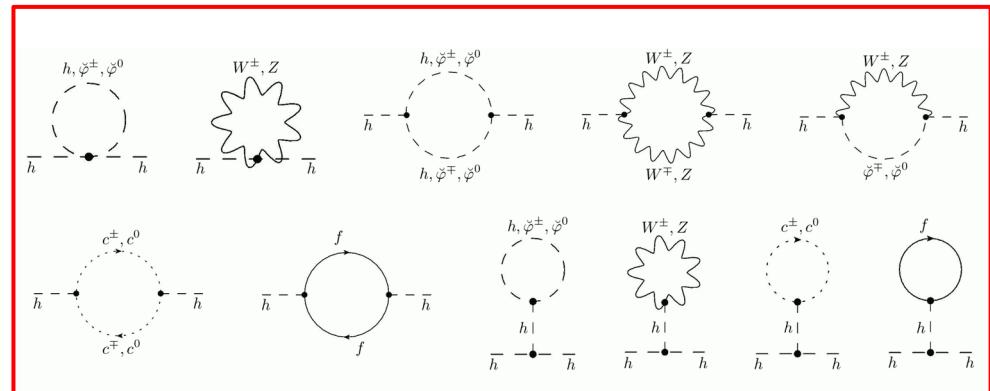
Consequences: The Higgs

[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]



Consequences: The Higgs

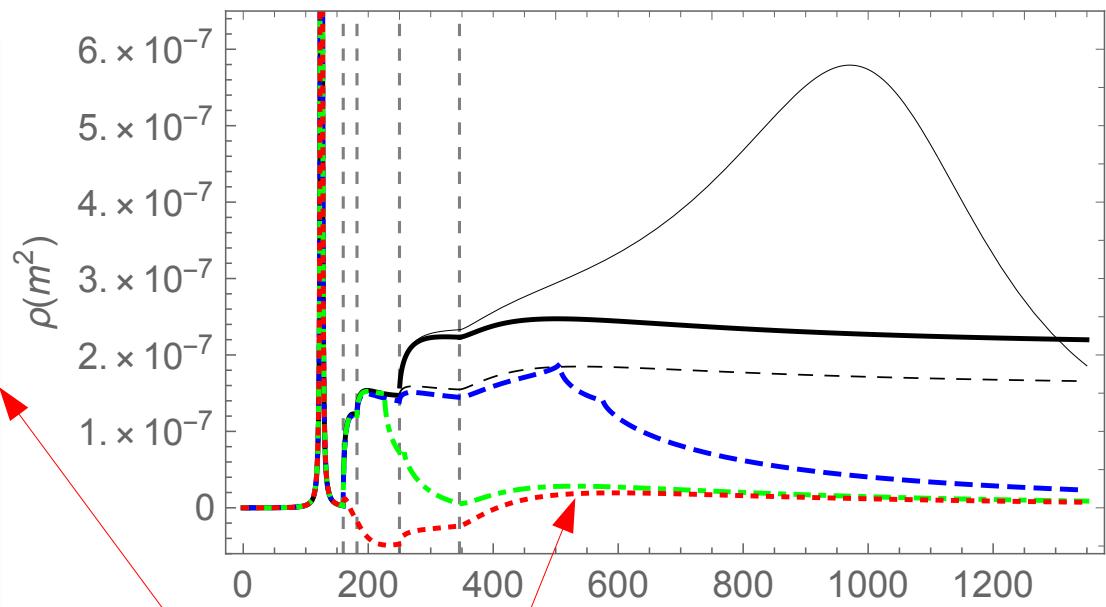
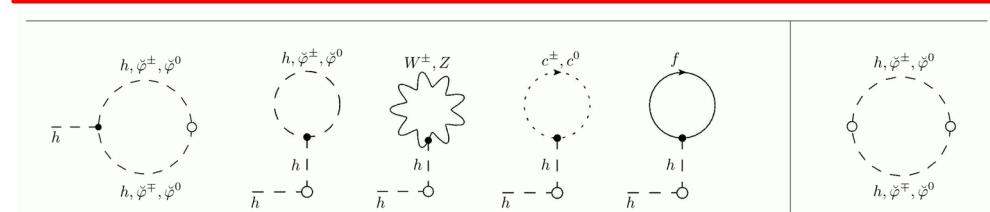
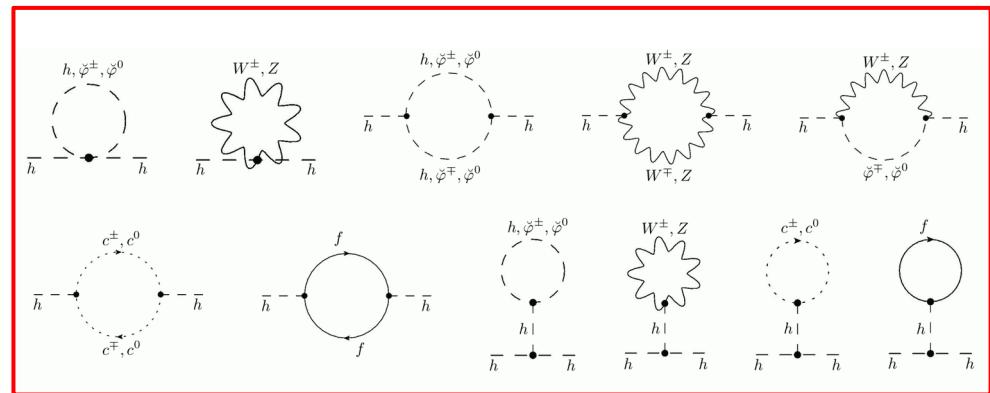
[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]



Gauge-dependent

Consequences: The Higgs

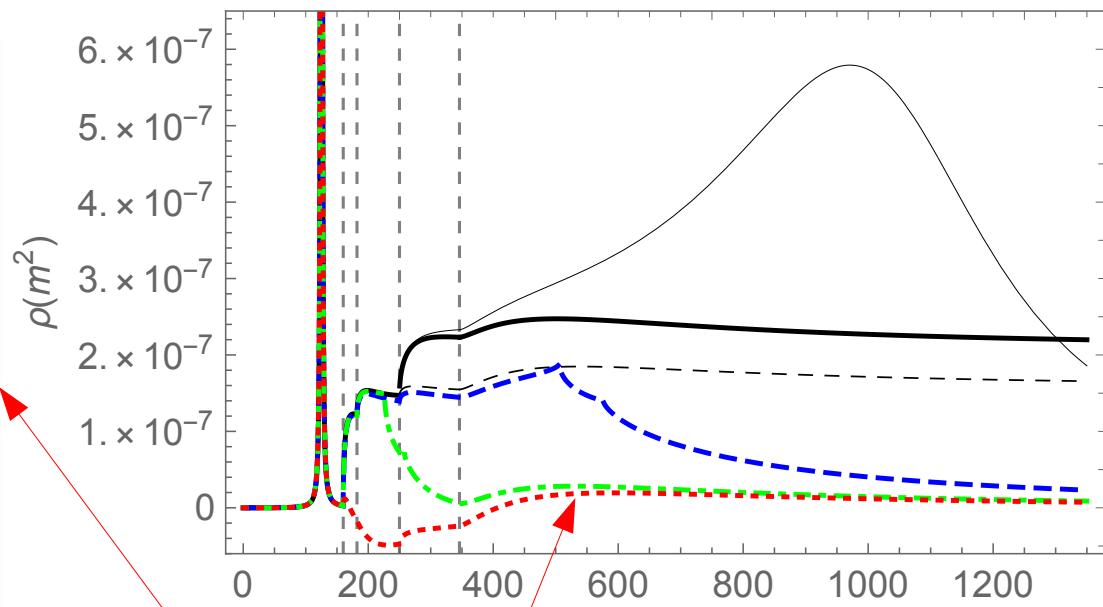
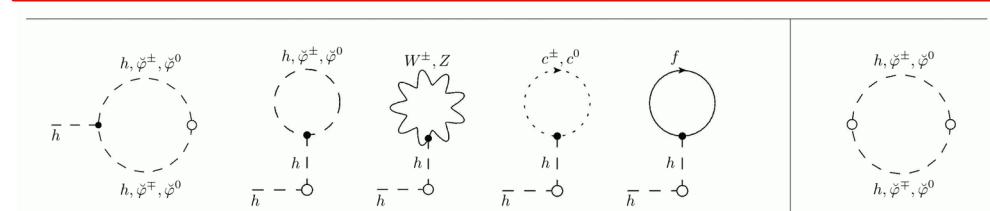
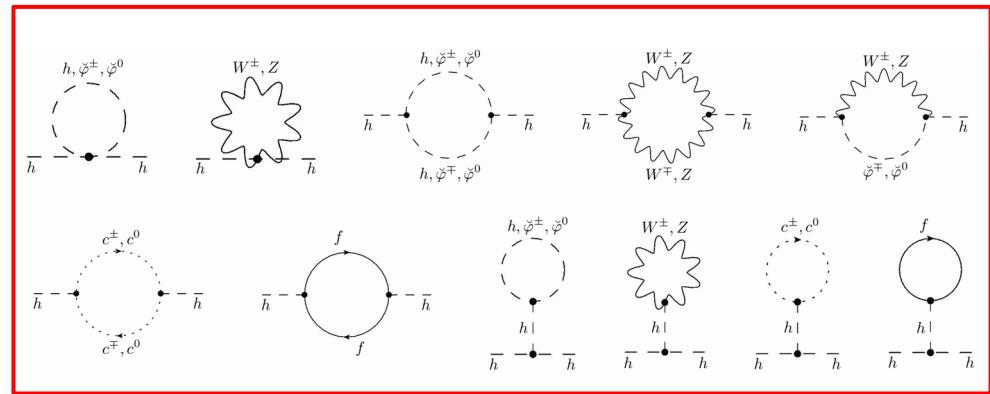
[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]



**Gauge-dependent
Unphysical features:
Positivity violation
Additional thresholds**

Consequences: The Higgs

[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]

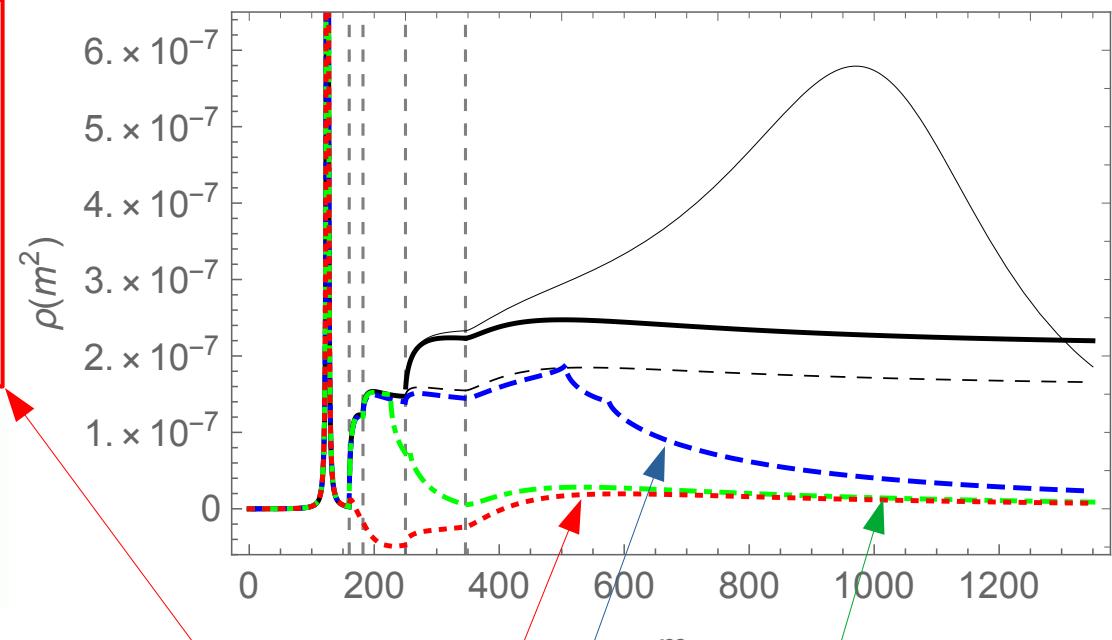
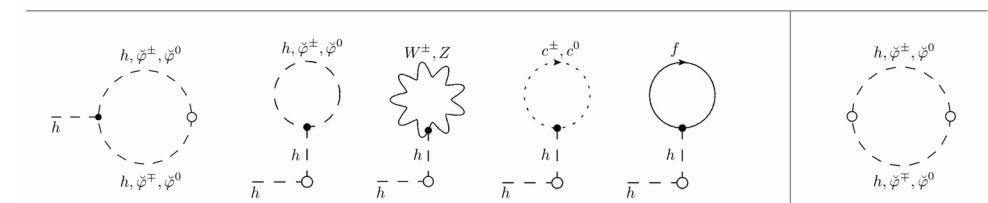
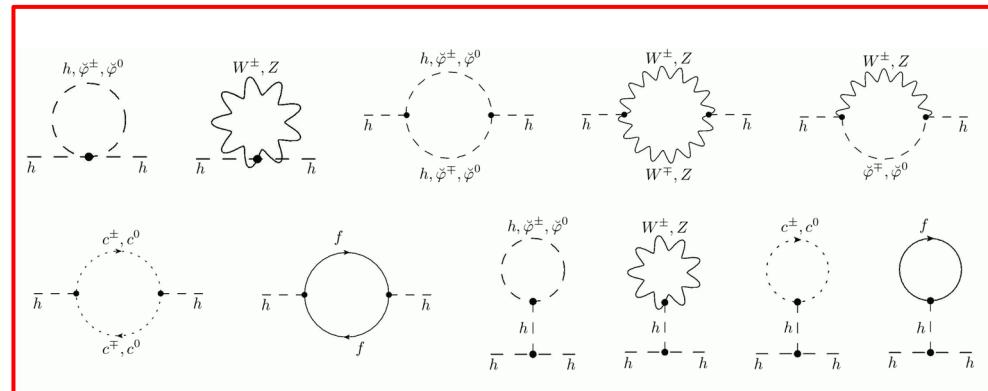


Gauge-dependent
Unphysical features:
Positivity violation
Additional thresholds

Not a consequence
of instability: Occurs even
for an asymptotically stable
Higgs in a toy theory

Consequences: The Higgs

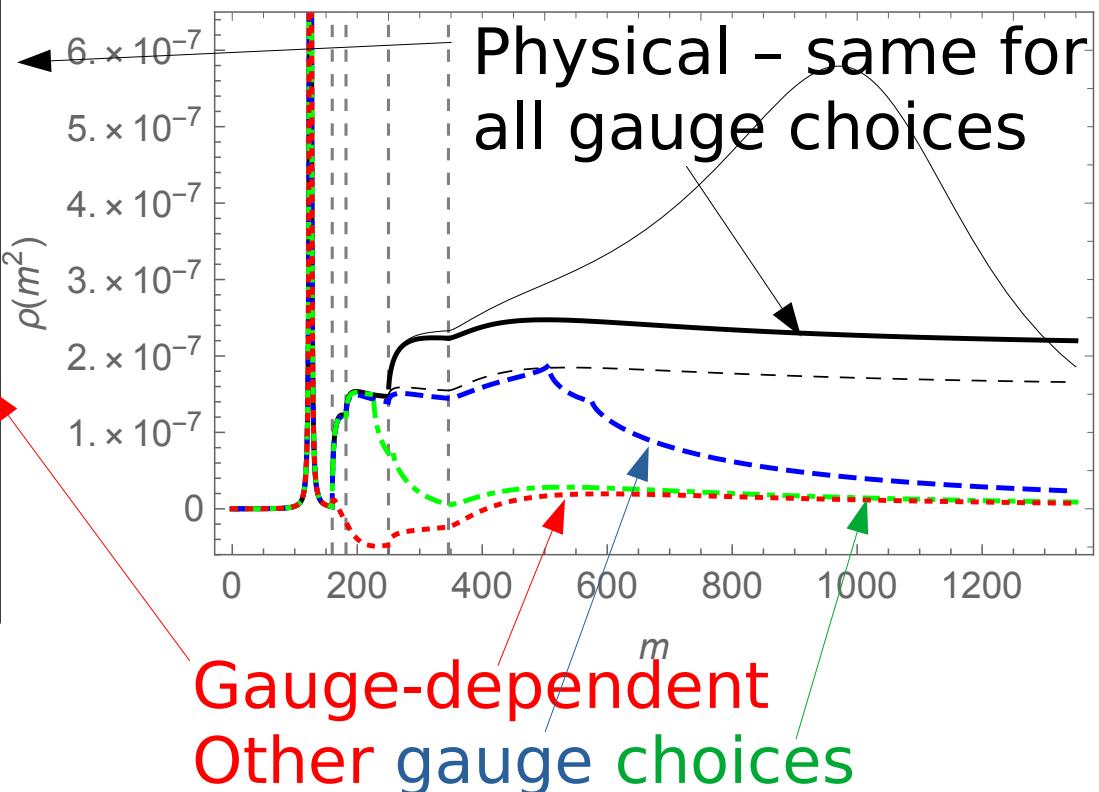
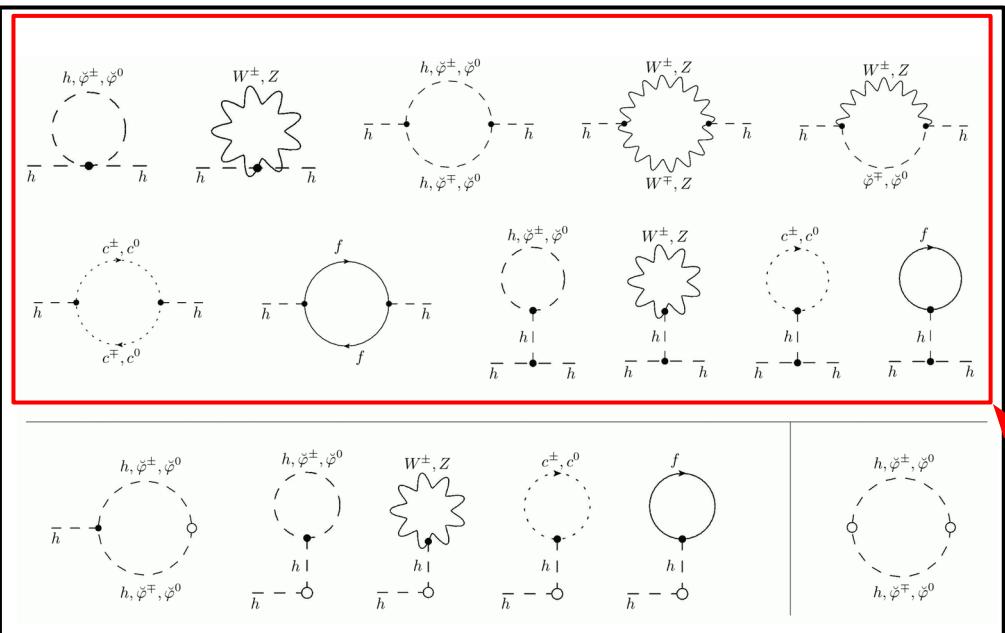
[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]



Gauge-dependent
Other gauge choices

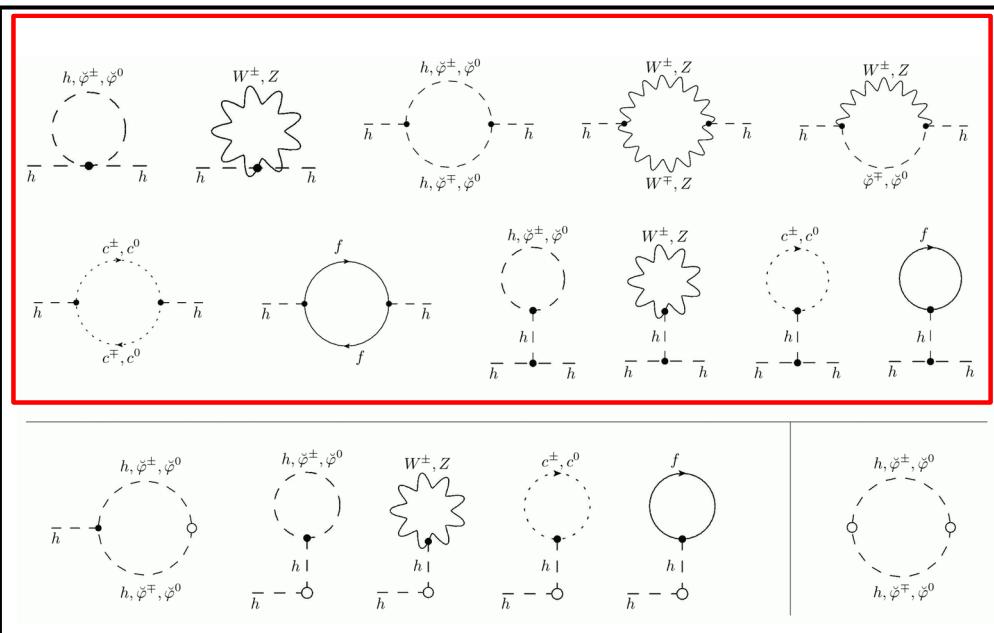
Consequences: The Higgs

[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20]



Consequences: The Higgs

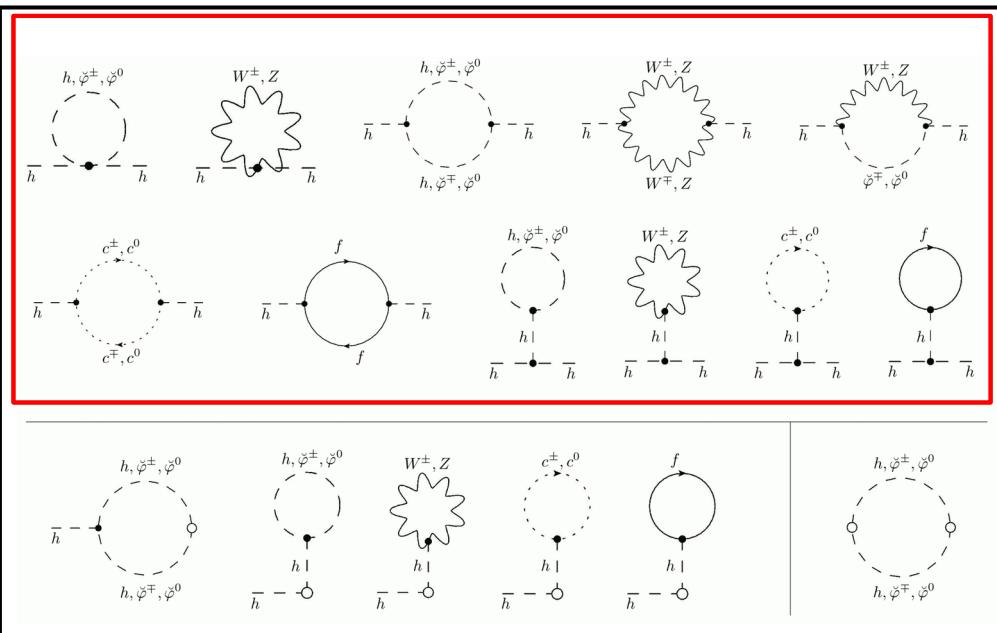
[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20
Maas et al. unpublished]



Same structure repeats itself in decays and scattering processes

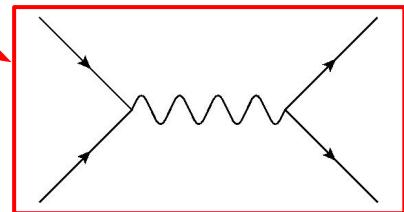
Consequences: The Higgs

[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20
Maas et al. unpublished]



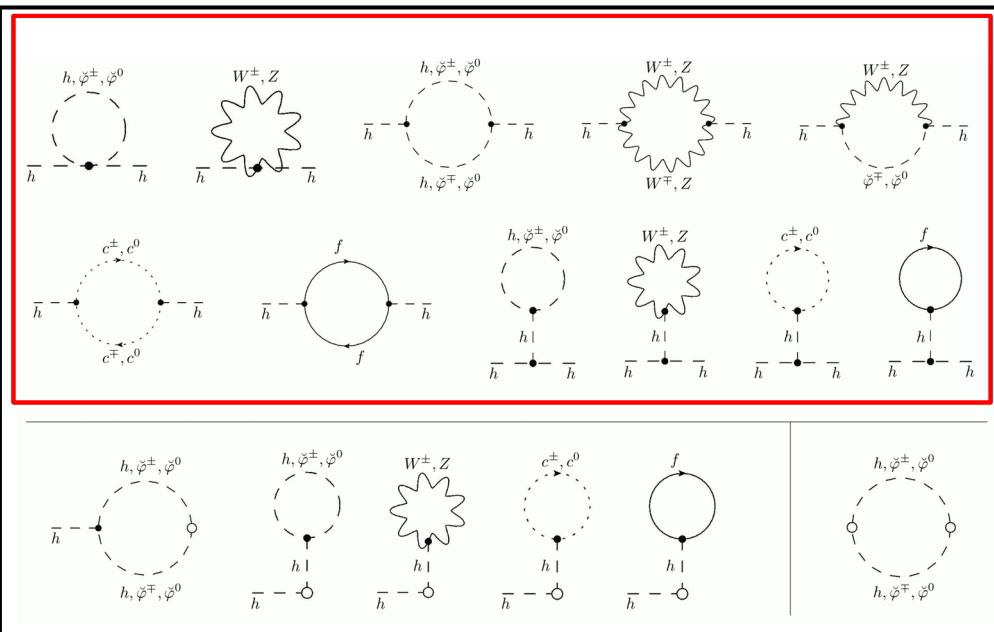
Same structure repeats itself in decays and scattering processes

LO: Standard perturbation theory



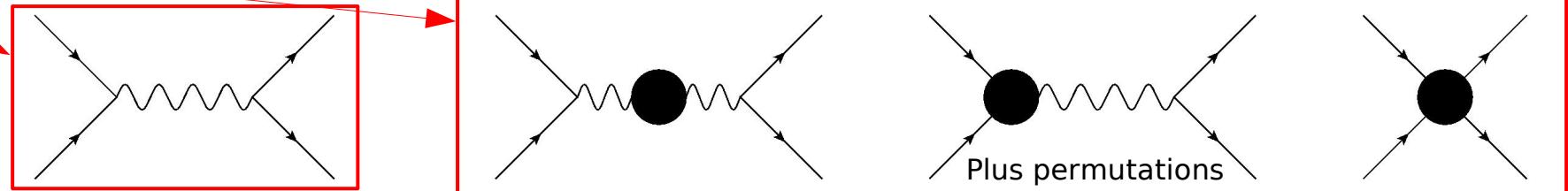
Consequences: The Higgs

[Maas'12,'17
Maas & Sondenheimer'20
Dudal et al.'20
Maas et al. unpublished]



Same structure repeats itself in decays and scattering processes

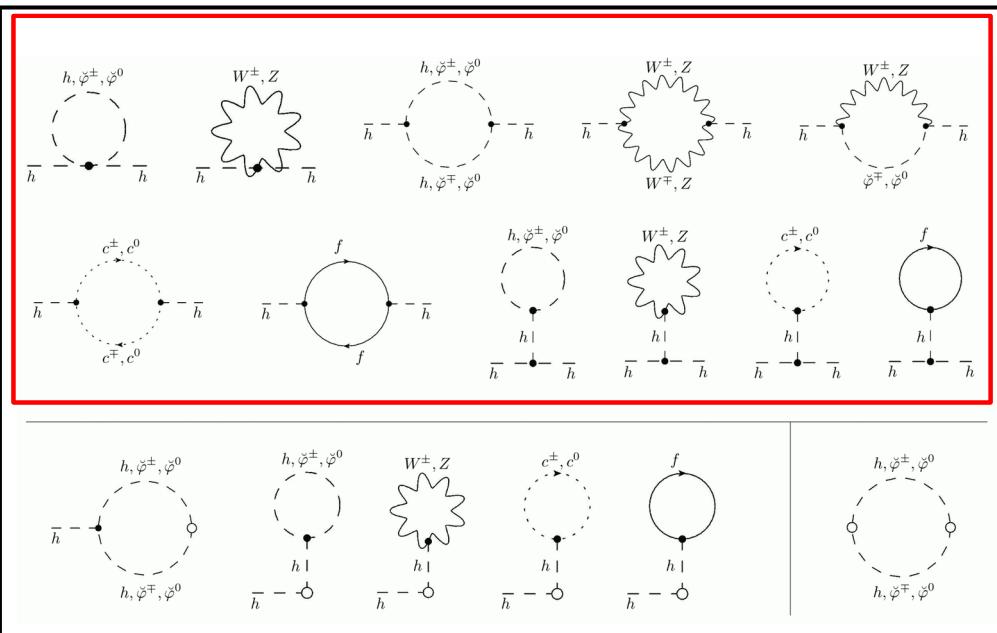
LO, NLO: Standard perturbation theory



Plus permutations

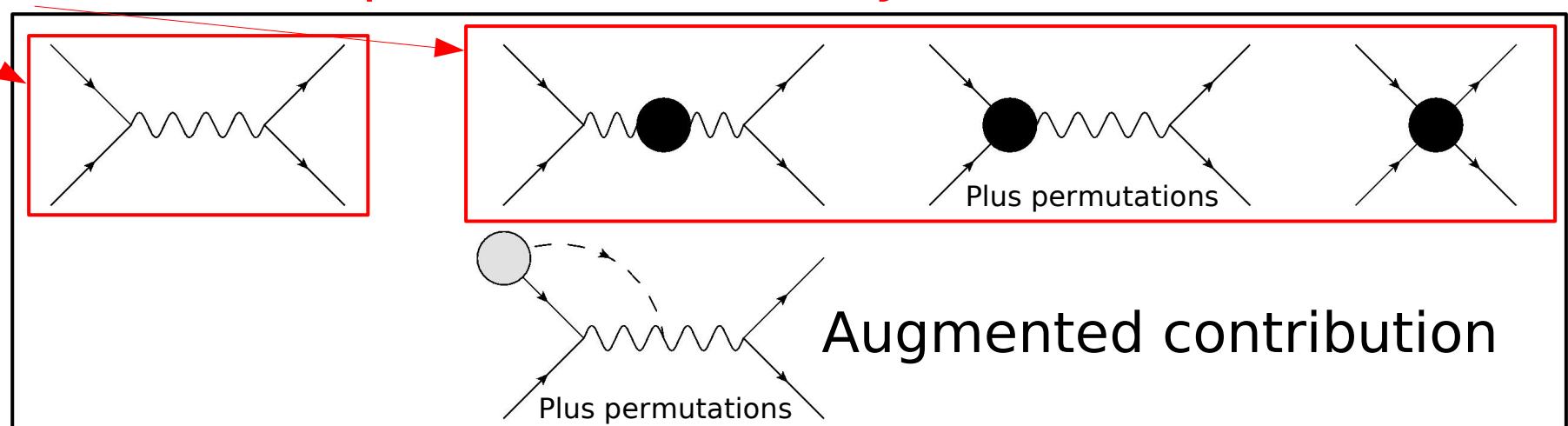
Consequences: The Higgs

[Maas'12,'17
 Maas & Sondenheimer'20
 Dudal et al.'20
 Maas et al. unpublished]



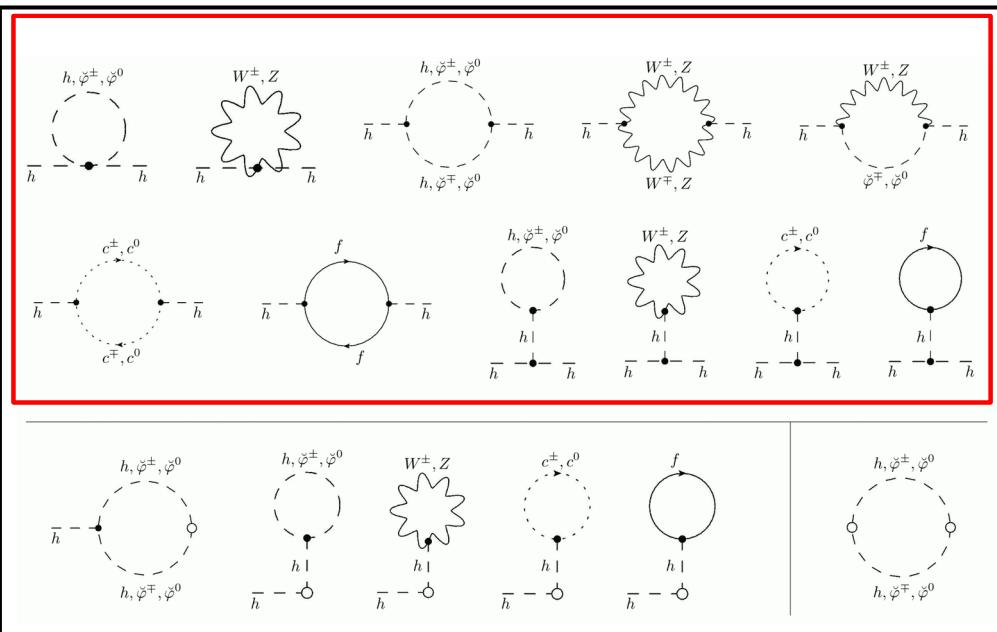
Same structure repeats itself in decays and scattering processes

LO, NLO: Standard perturbation theory



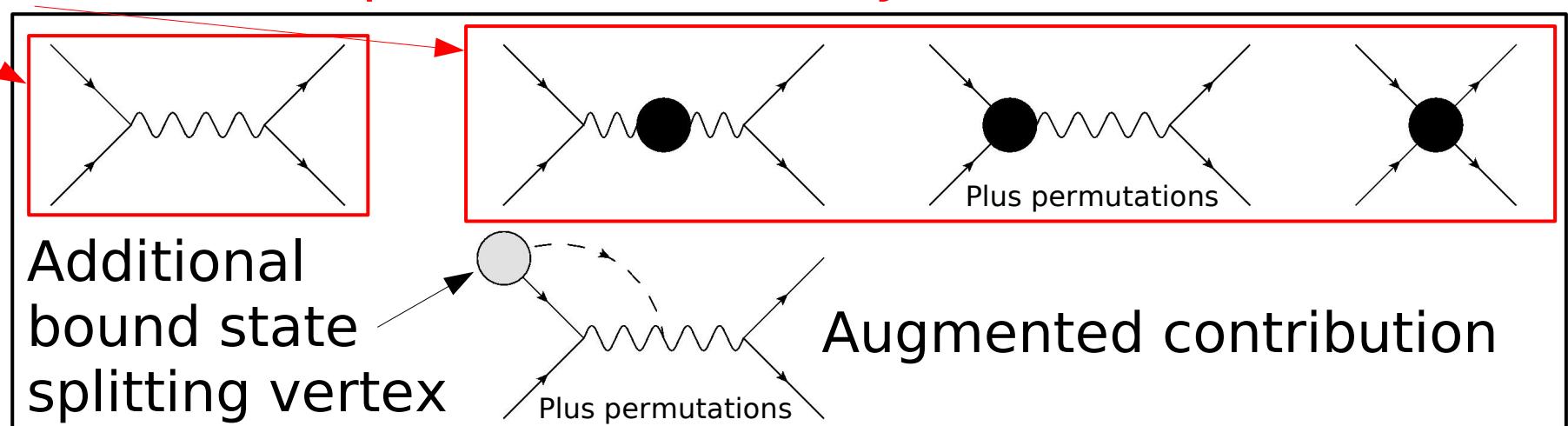
Consequences: The Higgs

[Maas'12,'17
 Maas & Sondenheimer'20
 Dudal et al.'20
 Maas et al. unpublished]



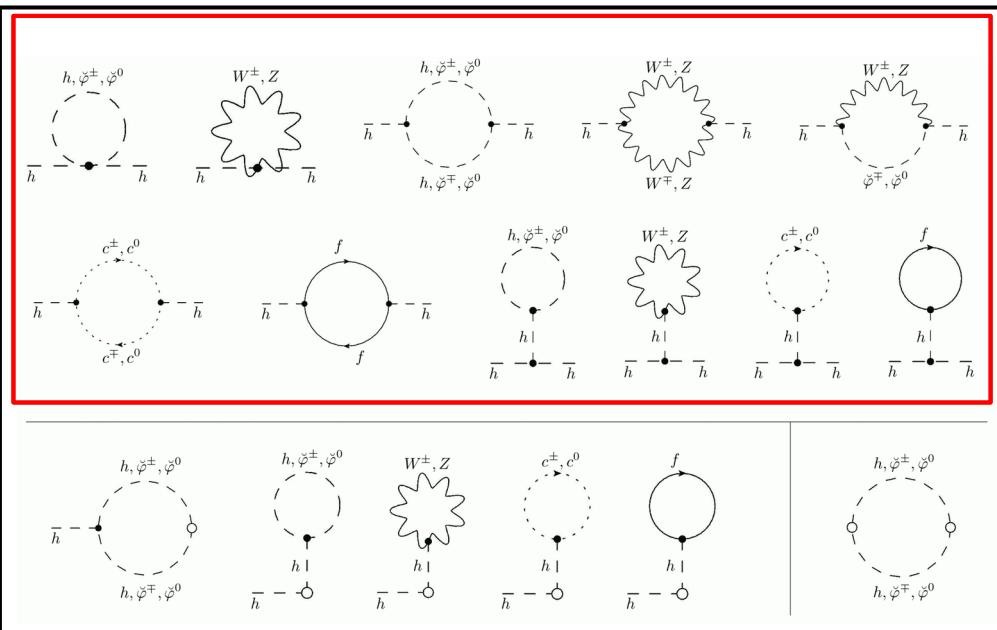
Same structure repeats itself in decays and scattering processes

LO, NLO: Standard perturbation theory



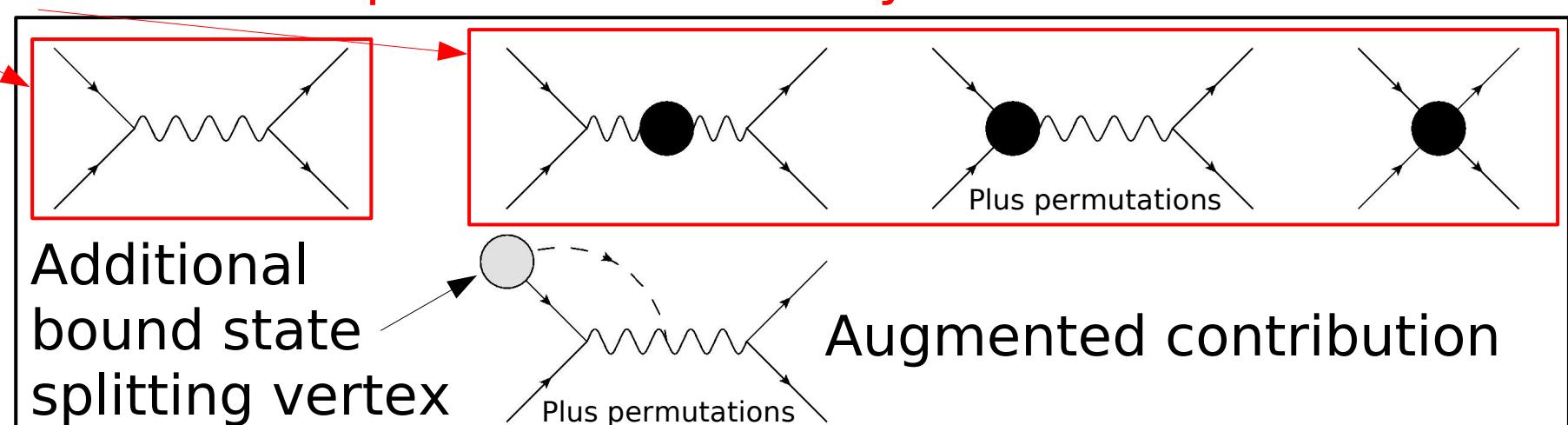
Consequences: The Higgs

[Maas'12,'17
 Maas & Sondenheimer'20
 Dudal et al.'20
 Maas et al. unpublished]



Same structure repeats itself in decays and scattering processes

LO, NLO: Standard perturbation theory



Augmented perturbation theory only augments Feynman rules

What about the vector?

[Fröhlich et al.'80,'81
Maas'12]

What about the vector?

[Fröhlich et al.'80,'81
Maas'12]

1) Formulate gauge-invariant operator

1⁻ triplet: $\langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle$

What about the vector?

[Fröhlich et al.'80,'81
Maas'12]

1) Formulate gauge-invariant operator

1⁻ triplet: $\langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle$

2) Expand Higgs field around fluctuations $h = v + \eta$

What about the vector?

[Fröhlich et al.'80,'81
Maas'12]

1) Formulate gauge-invariant operator

1⁻ triplet: $\langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle = v^2 c_{ij}^{ab} \langle W_\mu^a(x) W_\mu^b(y) \rangle + \dots$$

What about the vector?

[Fröhlich et al.'80,'81
Maas'12]

1) Formulate gauge-invariant operator

1⁻ triplet: $\langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle = v^2 c_{ij}^{ab} \langle W_\mu^a(x) W^\mu_b(y) \rangle + \dots$$

Matrix from
group structure

What about the vector?

[Fröhlich et al.'80, '81
Maas'12]

1) Formulate gauge-invariant operator

1⁻ triplet: $\langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\begin{aligned} \langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle &= v^2 c_{ij}^{ab} \langle W_\mu^a(x) W_\mu^b(y)^\mu \rangle + \dots \\ &= v^2 \langle W_\mu^i W_\mu^j \rangle + \dots \end{aligned}$$

Matrix from
group structure

What about the vector?

[Fröhlich et al.'80,'81
Maas'12]

1) Formulate gauge-invariant operator

1⁻ triplet: $\langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\begin{aligned} \langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle &= v^2 c_{ij}^{ab} \langle W_\mu^a(x) W_\mu^b(y)^\mu \rangle + \dots \\ &= v^2 \langle W_\mu^i W_\mu^j \rangle + \dots \end{aligned}$$

c projects custodial states to gauge states

Matrix from group structure

What about the vector?

[Fröhlich et al.'80, '81
Maas'12]

1) Formulate gauge-invariant operator

1⁻ triplet: $\langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle$

2) Expand Higgs field around fluctuations $h = v + \eta$

$$\begin{aligned} \langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle &= v^2 c_{ij}^{ab} \langle W_\mu^a(x) W^\mu_b(y) \rangle + \dots \\ &= v^2 \langle W_\mu^i W_\mu^j \rangle + \dots \end{aligned}$$

c projects custodial states to gauge states

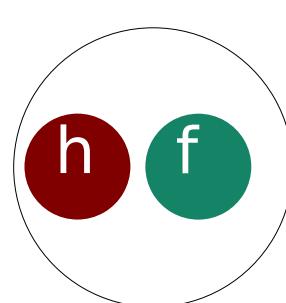
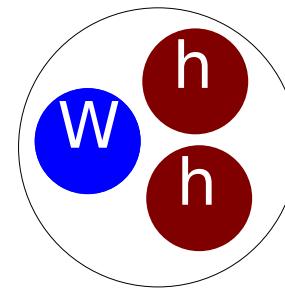
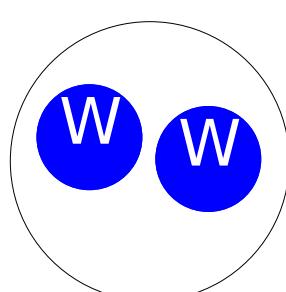
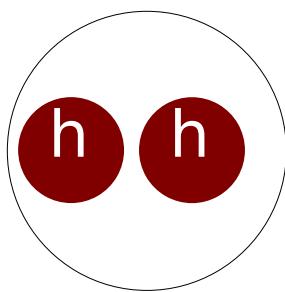
Exactly one gauge boson for every physical state

Matrix from group structure

Physical states

[Fröhlich et al.'80,
Banks et al.'79]

- No “real” breaking
- Physical particles are gauge-invariant particles
 - **Cannot** be the elementary particles
 - Non-Abelian nature is relevant
- Need more than one particle: Composite particles
 - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Has nothing to do with weak coupling
 - Think QED (hydrogen atom!)

Flavor

[Fröhlich et al.'80,
Egger, Maas, Sondenheimer'17]

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable

Flavor

[Fröhlich et al.'80,
Egger, Maas, Sondenheimer'17]

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state

$$\left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} \begin{pmatrix} v_L \\ l_L \end{pmatrix} \right\rangle_i (x)$$

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state

$$\left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} \begin{pmatrix} v_L \\ l_L \end{pmatrix} \right|_i (x)$$

- Gauge-invariant state

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state

$$\begin{pmatrix} h_2 - h_1 \\ h_1^* \ h_2^* \end{pmatrix} \begin{pmatrix} v_L \\ l_L \end{pmatrix}_i(x)$$

- Gauge-invariant state, but custodial doublet

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state

$$\left| \begin{pmatrix} h_2 - h_1 \\ h_1^* h_2^* \end{pmatrix} \begin{pmatrix} v_L \\ l_L \end{pmatrix} \right|_i (x) + \left| \begin{pmatrix} h_2 - h_1 \\ h_1^* h_2^* \end{pmatrix} \begin{pmatrix} v_L \\ l_L \end{pmatrix} \right|_j (y) \right|$$

- Gauge-invariant state, but custodial doublet

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state – FMS applicable

$$\left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_i (x) + \left| \begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_j (y) \approx v^2 \left| \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix}_i (x) + \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix}_j (y) \right| + O(\eta)$$

$h = v + \eta$

- Gauge-invariant state, but custodial doublet

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state – FMS applicable

$$\left| \begin{pmatrix} h_2 - h_1 \\ h_1^* h_2^* \end{pmatrix} \begin{pmatrix} v_L \\ l_L \end{pmatrix} \right|_i (x) + \left| \begin{pmatrix} h_2 - h_1 \\ h_1^* h_2^* \end{pmatrix} \begin{pmatrix} v_L \\ l_L \end{pmatrix} \right|_j (y) \right|^{h=v+\eta} \approx v^2 \left| \begin{pmatrix} v_L \\ l_L \end{pmatrix}_i (x) + \begin{pmatrix} v_L \\ l_L \end{pmatrix}_j (y) \right| + O(\eta)$$

- Gauge-invariant state, but custodial doublet
- Yukawa terms break custodial symmetry
 - Different masses for doublet members

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state – **FMS applicable**

$$\left| \begin{pmatrix} h_2 - h_1 \\ h_1^* h_2^* \end{pmatrix} \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_i (x) + \left| \begin{pmatrix} h_2 - h_1 \\ h_1^* h_2^* \end{pmatrix} \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_j (y) \right|^{h=v+\eta} \approx v^2 \left| \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_i (x) + \left| \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} \right|_j (y) \right| + O(\eta)$$

- Gauge-invariant state, but **custodial doublet**
- Yukawa terms break custodial symmetry
 - Different masses for doublet members
- Extends non-trivially to hadrons

Flavor on the lattice

[Afferrante,Maas,Sondenheimer,Törek'20]

- Only mock-up standard model
 - Compressed mass scales
 - One generation
 - Degenerate leptons and neutrinos
 - Dirac fermions: left/right-handed non-degenerate
 - Quenched

Flavor on the lattice

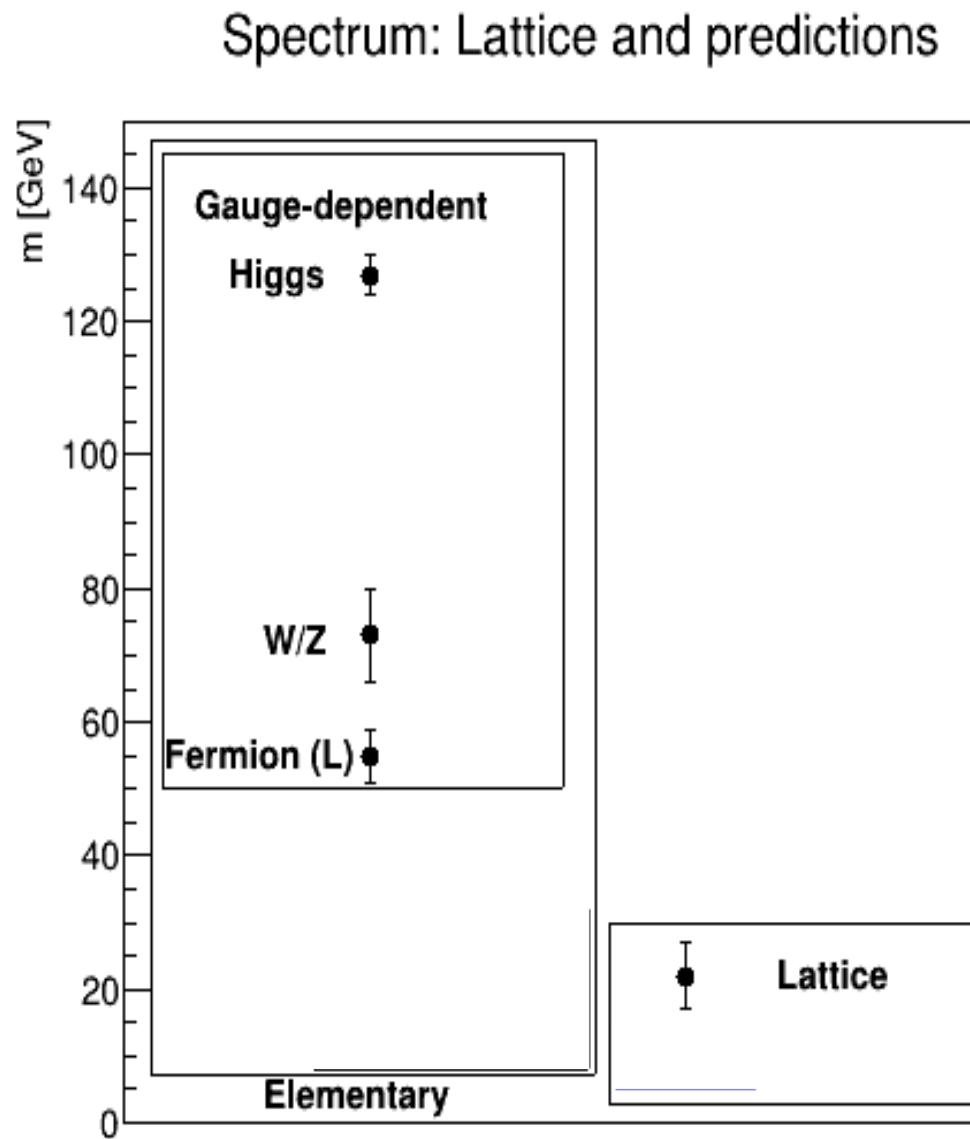
[Afferrante,Maas,Sondenheimer,Törek'20]

- Only mock-up standard model
 - Compressed mass scales
 - One generation
 - Degenerate leptons and neutrinos
 - Dirac fermions: left/right-handed non-degenerate
 - Quenched
- Same qualitative outcome
 - FMS construction
 - Mass defect
 - Flavor and custodial symmetry patterns

Flavor on the lattice

[Afferrante,Maas,Sondenheimer,Törek'20]

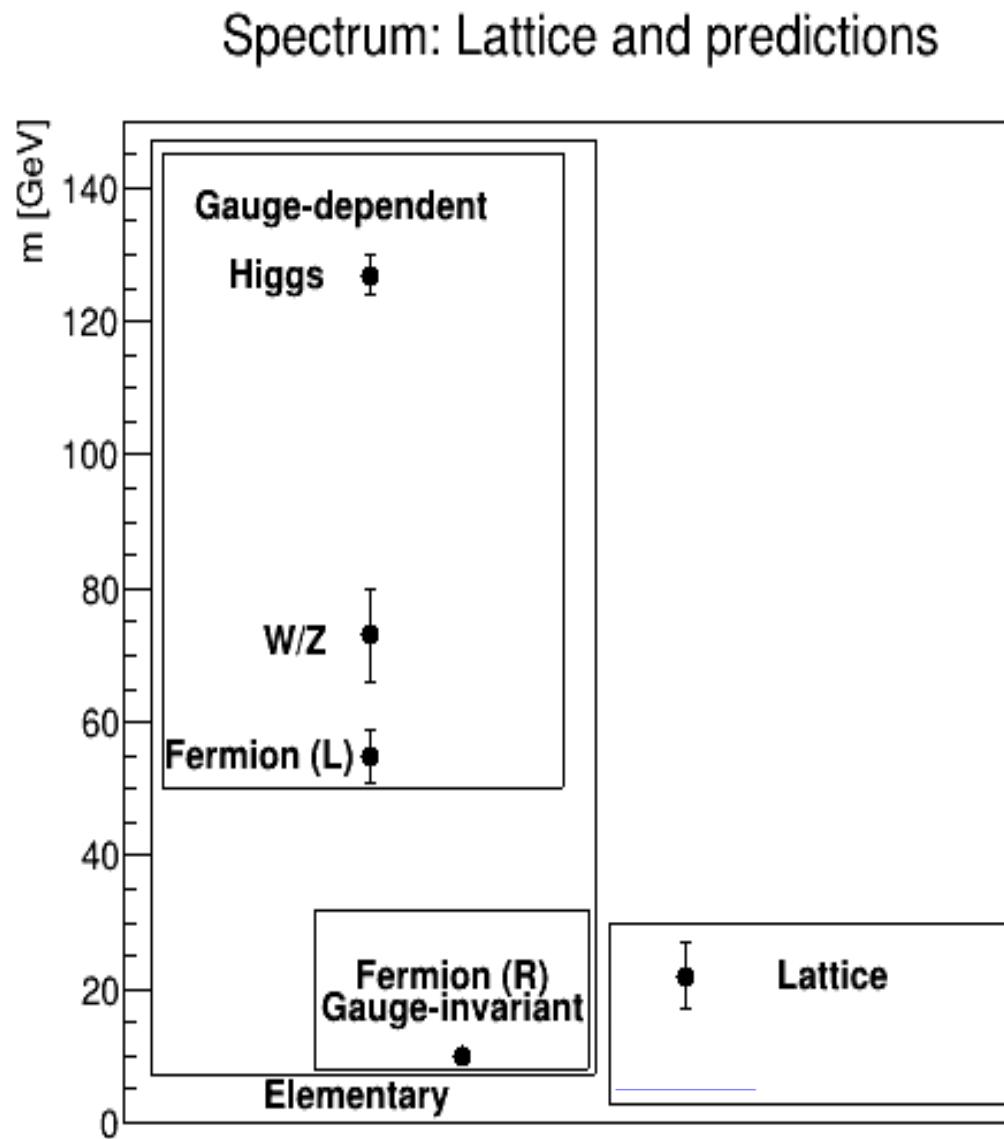
- Only mock-up standard model
 - Compressed mass scales
 - One generation
 - Degenerate leptons and neutrinos
 - Dirac fermions: left/right-handed non-degenerate
 - Quenched
- Same qualitative outcome
 - FMS construction
 - Mass defect
 - Flavor and custodial symmetry patterns



Flavor on the lattice

[Afferrante,Maas,Sondenheimer,Török'20]

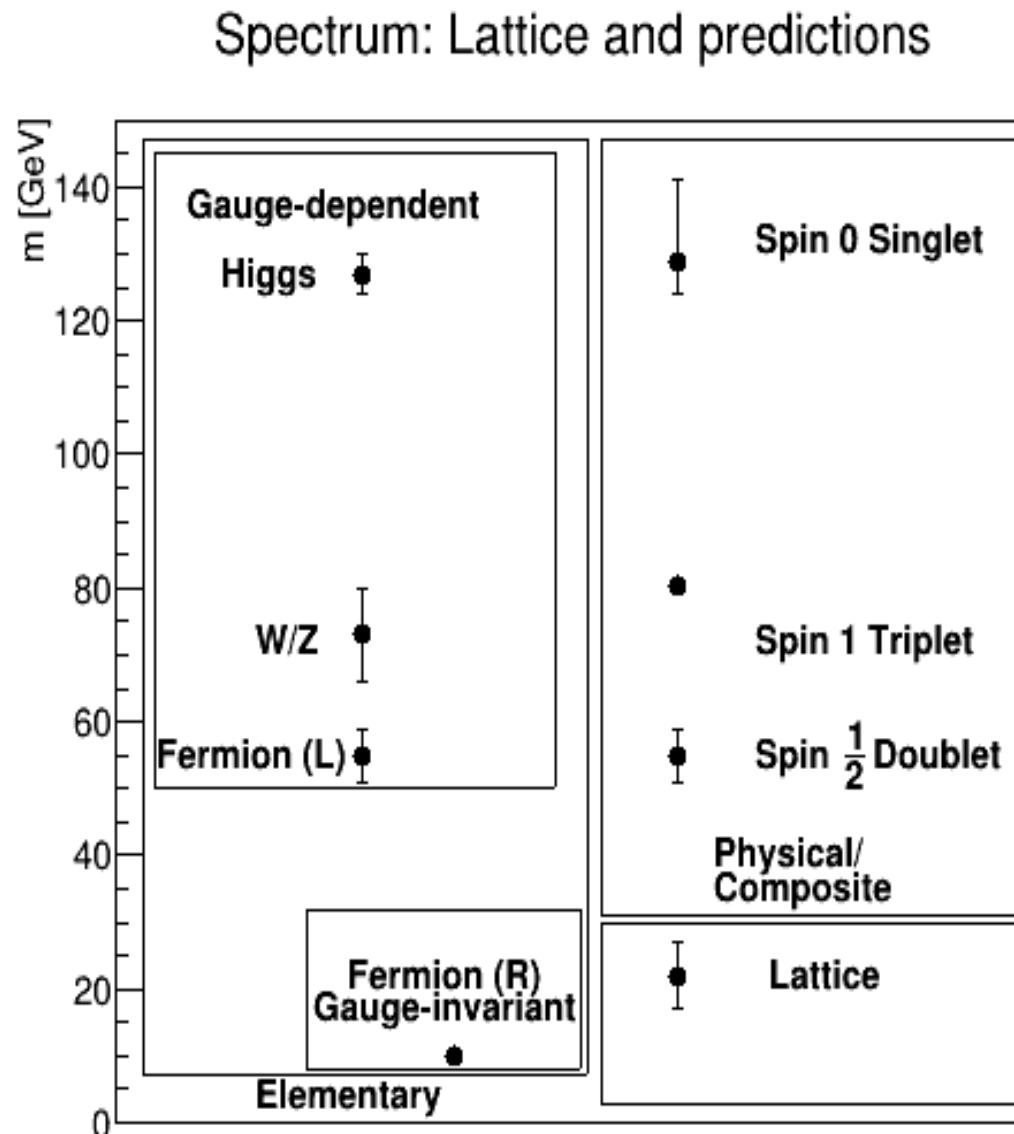
- Only mock-up standard model
 - Compressed mass scales
 - One generation
 - Degenerate leptons and neutrinos
 - Dirac fermions: left/right-handed non-degenerate
 - Quenched
- Same qualitative outcome
 - FMS construction
 - Mass defect
 - Flavor and custodial symmetry patterns



Flavor on the lattice

[Afferrante,Maas,Sondenheimer,Törek'20]

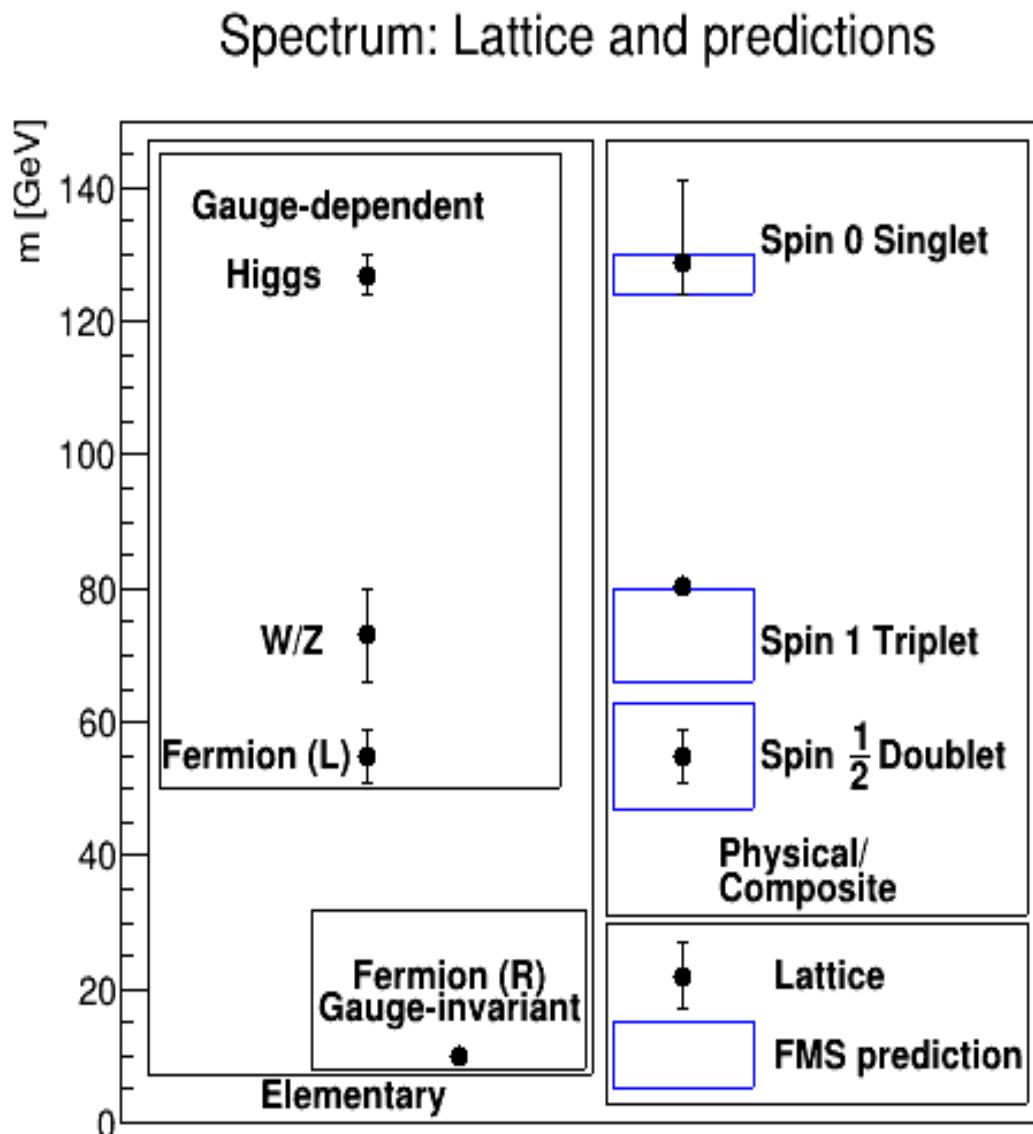
- Only mock-up standard model
 - Compressed mass scales
 - One generation
 - Degenerate leptons and neutrinos
 - Dirac fermions: left/right-handed non-degenerate
 - Quenched
- Same qualitative outcome
 - FMS construction
 - Mass defect
 - Flavor and custodial symmetry patterns



Flavor on the lattice

[Afferrante,Maas,Sondenheimer,Törek'20]

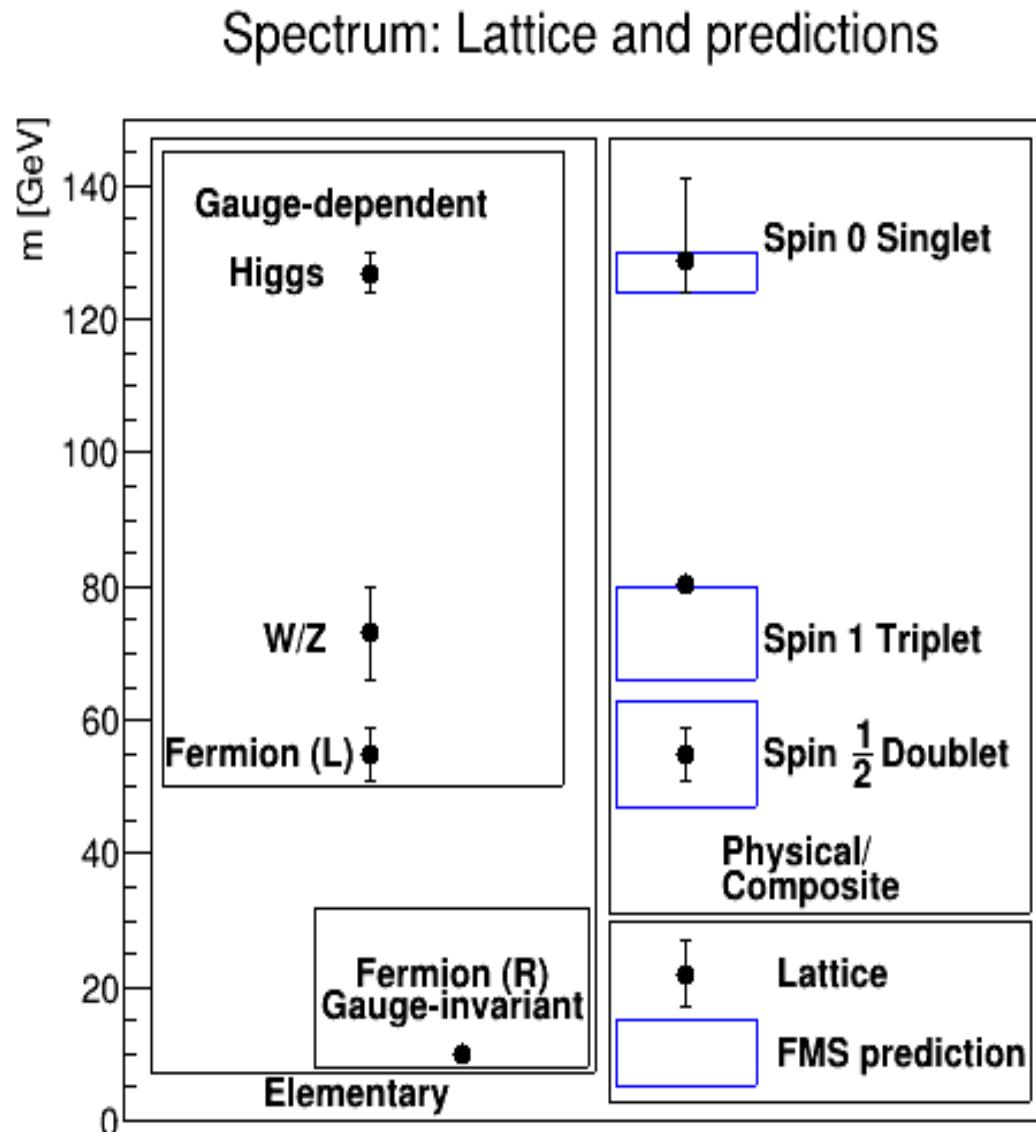
- Only mock-up standard model
 - Compressed mass scales
 - One generation
 - Degenerate leptons and neutrinos
 - Dirac fermions: left/right-handed non-degenerate
 - Quenched
- Same qualitative outcome
 - FMS construction
 - Mass defect
 - Flavor and custodial symmetry patterns
- Supports FMS prediction



Flavor on the lattice

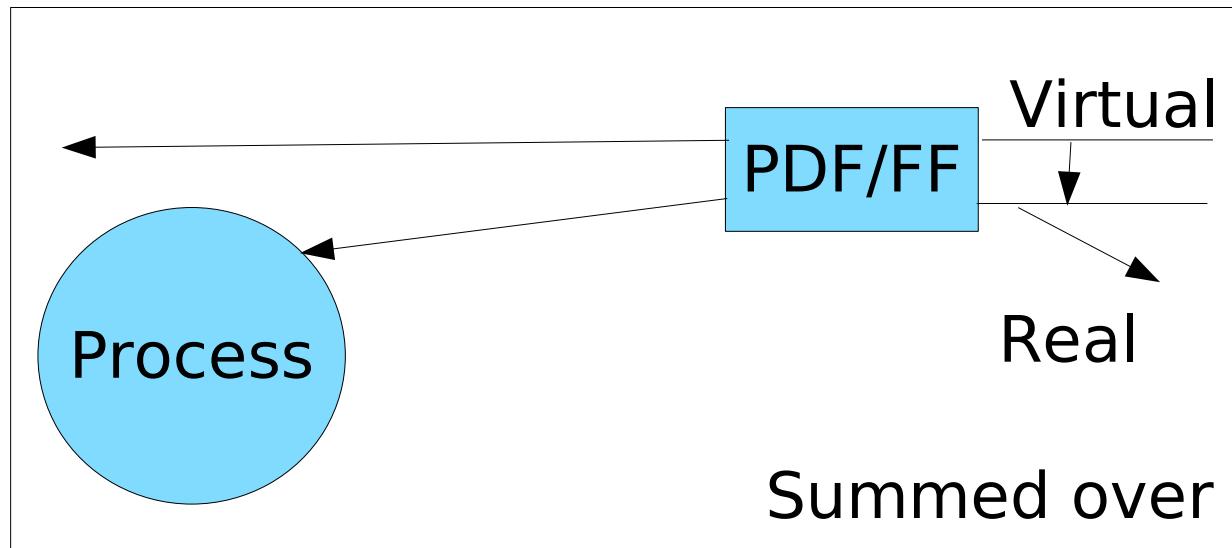
[Afferrante,Maas,Sondenheimer,Törek'20]

- Only mock-up standard model
 - Compressed mass scales
 - One generation
 - Degenerate leptons and neutrinos
 - Dirac fermions: left/right-handed non-degenerate
 - Quenched
- Same qualitative outcome
 - FMS construction
 - Mass defect
 - Flavor and custodial symmetry patterns
- Supports FMS prediction – grant for unquenching '24-'28



Invalidation in electroweak physics

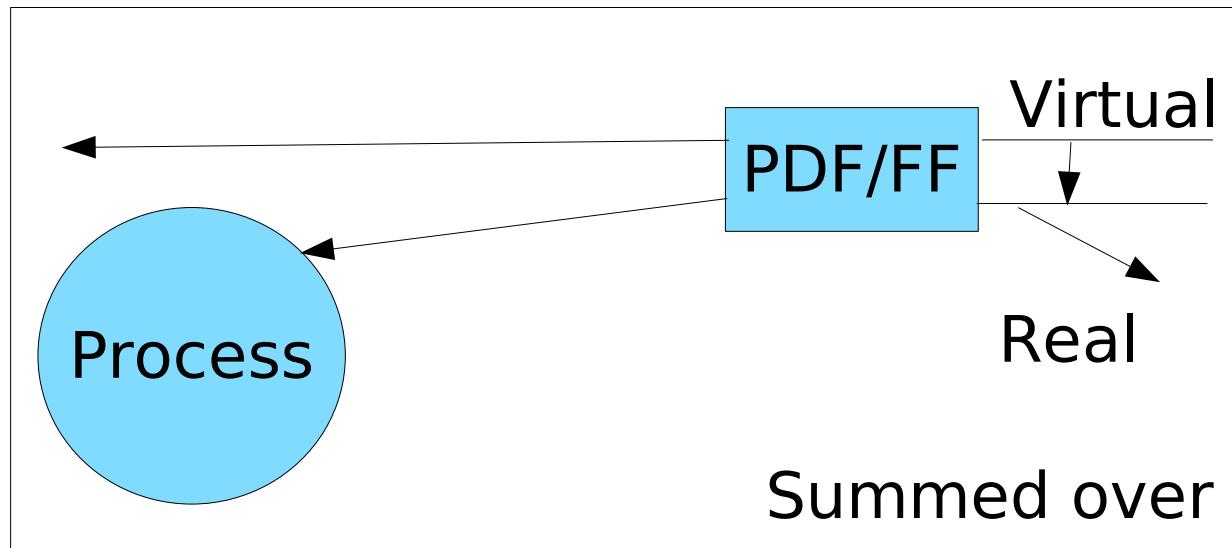
[Maas et al.'22]



- Particles are again (electroweak) gauge-singlets

Invalidation in electroweak physics

[Maas et al.'22]



Incoming or
outgoing
gauge
singlet

- Particles are again (electroweak) gauge-singlets
- Low energy: FMS expansion

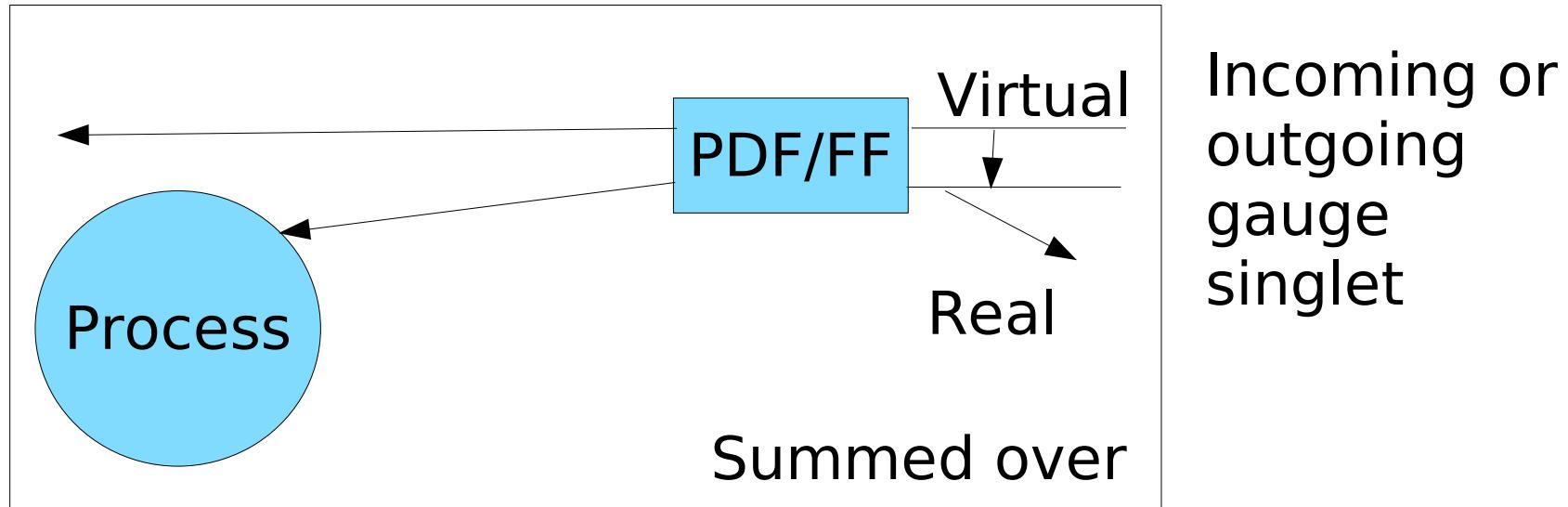
$$\langle h e h e | h \mu h \mu \rangle = \langle e e | \mu \mu \rangle + \langle \eta \eta \rangle \langle e e | \mu \mu \rangle + \langle e e \rangle \langle \eta \eta | \mu \mu \rangle + \dots$$

Standard
result

Irrelevant at low energies

Invalidation in electroweak physics

[Maas et al.'22]

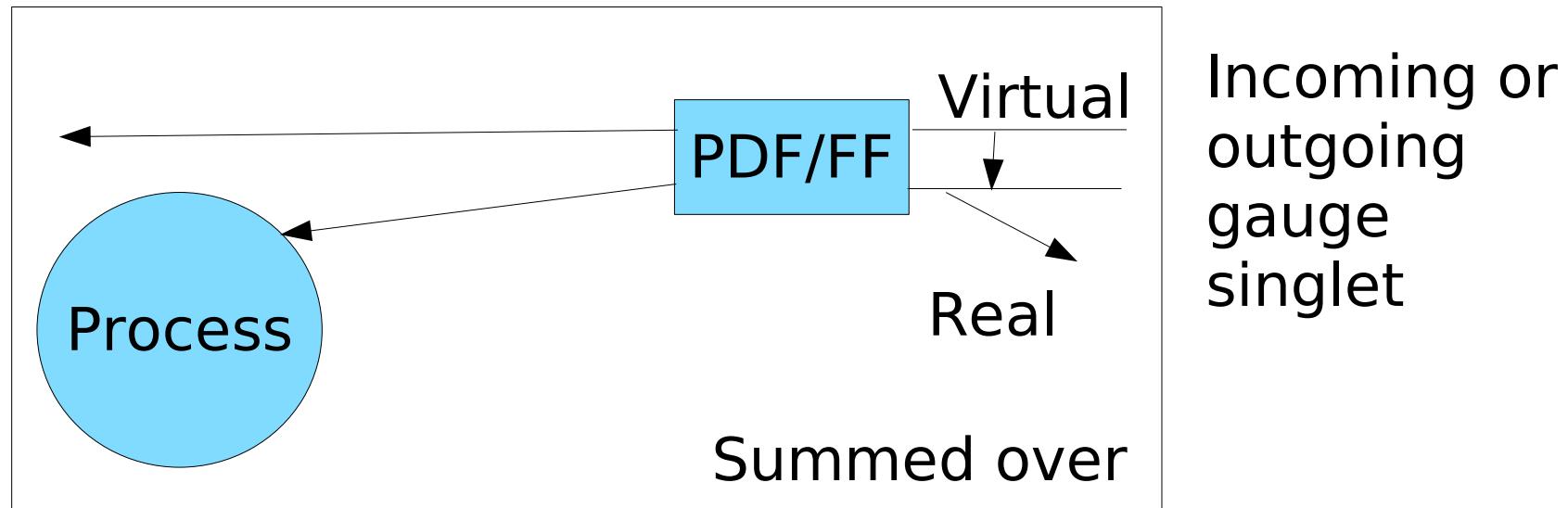


- Particles are again (electroweak) gauge-singlets
- Low energy: FMS expansion

$$\langle h e h e | X \rangle = \langle e e | X \rangle + \langle \eta \eta \rangle \langle e e | X \rangle + \langle e e \rangle \langle \eta \eta | X \rangle + \dots$$

Invalidation in electroweak physics

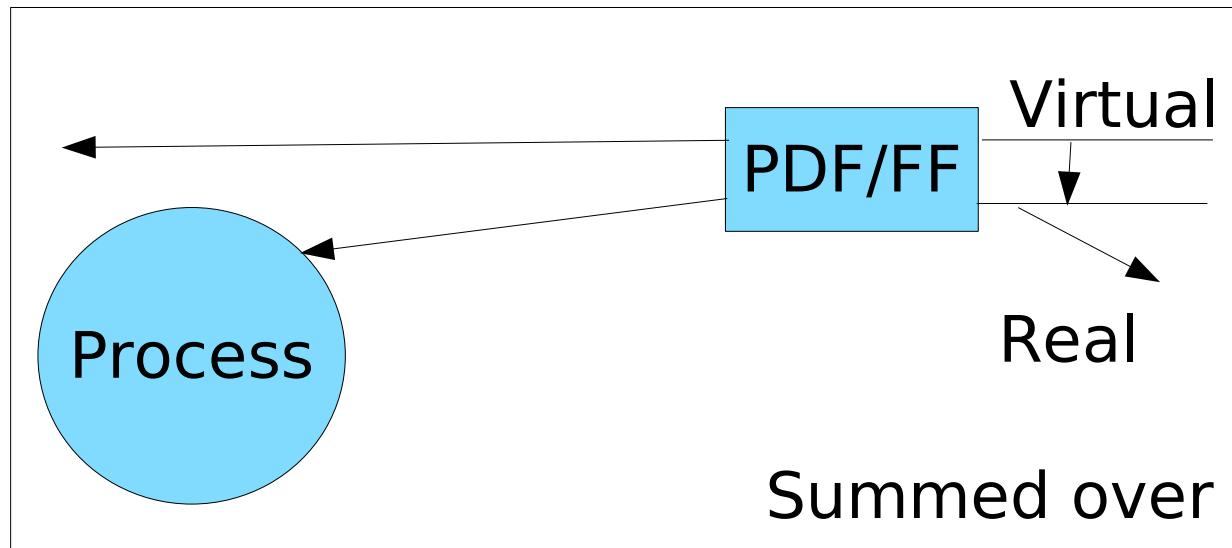
[Maas et al.'22]



- Particles are again (electroweak) gauge-singlets
- Low energy: FMS expansion
$$\langle h e h e | X \rangle = \langle e e | X \rangle + \langle \eta \eta \rangle \langle e e | X \rangle + \langle e e \rangle \langle \eta \eta | X \rangle + \dots$$
- High energies: Weakly inclusive because of Higgs

Invalidation in electroweak physics

[Maas et al.'22]



- Particles are again (electroweak) gauge-singlets

- Low energy: FMS expansion

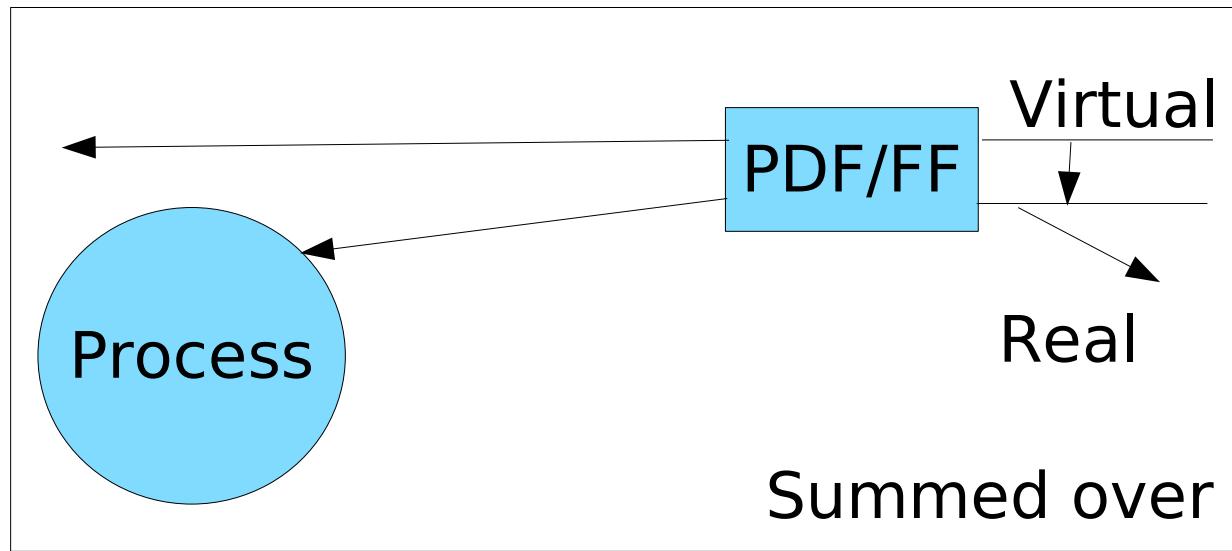
$$\langle h e h e | X \rangle = \langle e e | X \rangle + \langle \eta \eta \rangle \langle e e | X \rangle + \langle e e \rangle \langle \eta \eta | X \rangle + \dots$$

- High energies: Weakly inclusive because of Higgs

$$\sigma_{\Psi_L^2 \Psi_L^2 \rightarrow X}^{LO} = \sigma_{l_L l_L \rightarrow X}^{LO} + \sigma_{l_L \nu_L \rightarrow X}^{LO} + \sigma_{\bar{\nu}_L l_L \rightarrow X}^{LO} + \sigma_{\bar{\nu}_L \nu_L \rightarrow X}^{LO} + \text{higher order}$$

Invalidation in electroweak physics

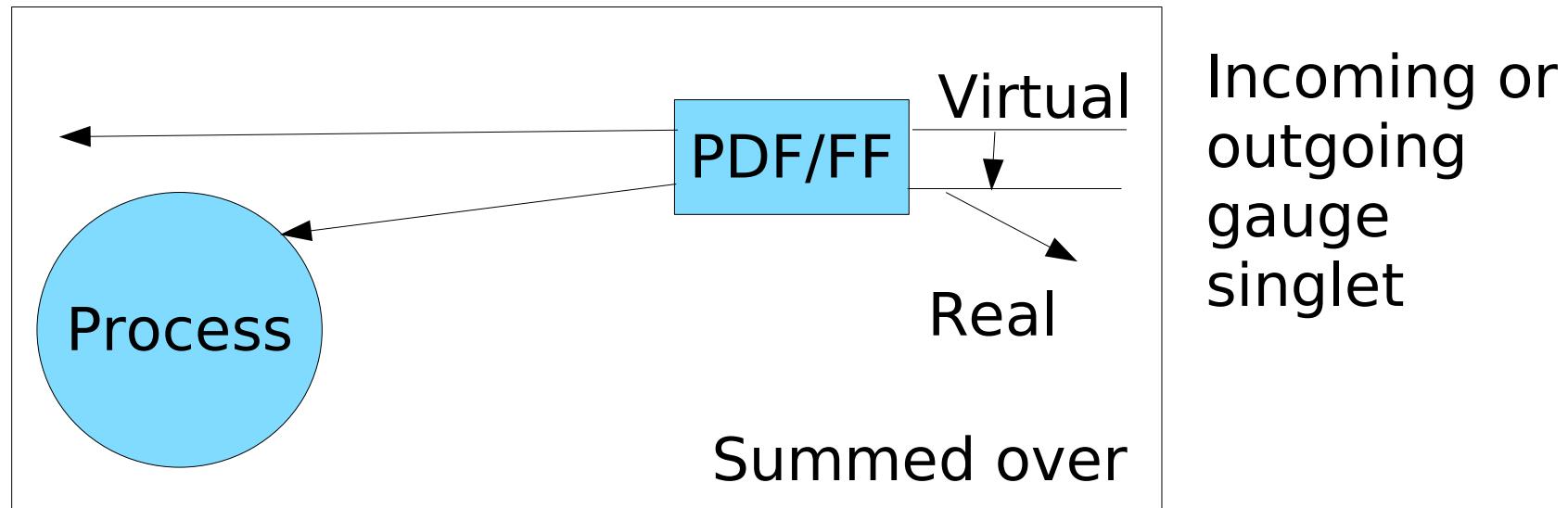
[Maas et al.'22]



- Particles are again (electroweak) gauge-singlets
- Low energy: FMS expansion
 $\langle h e h e | X \rangle = \langle e e | X \rangle + \langle \eta \eta \rangle \langle e e | X \rangle + \langle e e \rangle \langle \eta \eta | X \rangle + \dots$
- High energies: Weakly inclusive because of Higgs
 $\sigma_{\Psi_L^2 \Psi_L^2 \rightarrow X}^{LO} = \sigma_{l_L^- l_L \rightarrow X}^{LO} + \sigma_{l_L^- \nu_L \rightarrow X}^{LO} + \sigma_{\bar{\nu}_L l_L \rightarrow X}^{LO} + \sigma_{\bar{\nu}_L \nu_L \rightarrow X}^{LO} + \text{higher order}$
 - Restores BN theorem and KLN theorem: No Sudakov

Invalidation in electroweak physics

[Maas et al.'22]



- Particles are again (electroweak) gauge-singlets
- Low energy: FMS expansion
 $\langle h e h e | X \rangle = \langle e e | X \rangle + \langle \eta \eta \rangle \langle e e | X \rangle + \langle e e \rangle \langle \eta \eta | X \rangle + \dots$
- High energies: Weakly inclusive because of Higgs
 $\sigma_{\Psi_L^2 \Psi_L^2 \rightarrow X}^{LO} = \sigma_{l_L^- l_L \rightarrow X}^{LO} + \sigma_{l_L^- \nu_L \rightarrow X}^{LO} + \sigma_{\bar{\nu}_L l_L \rightarrow X}^{LO} + \sigma_{\bar{\nu}_L \nu_L \rightarrow X}^{LO} + \text{higher order}$
 - Restores BN theorem and KLN theorem: No Sudakov
 - Interesting consequences for PDFs/FFs

Summary

- Full gauge invariance also for weak interactions

Summary

- Full gauge invariance also for weak interactions
- Relevant for electroweak resummation at high energies (\rightarrow FCC/FLC)
 - Comparable to strong corrections

Summary

- Full gauge invariance also for weak interactions
- Relevant for electroweak resummation at high energies (\rightarrow FCC/FLC)
 - Comparable to strong corrections
- Effect suppressed at low energies because of standard model structure
 - Different in BSM physics