Physical Observables In Canonical Quantum Gravity

Axel Maas

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 - Path integral formulation and symmetries
 - Canonical quantum gravity

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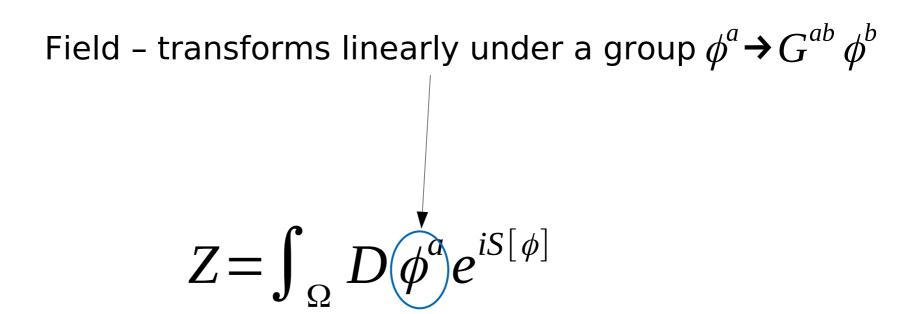
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 - Emergence of flat-space QFT

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- Fröhlich-Morchio-Strocchi mechanism
 - Emergence of flat-space QFT
- Connecting to other approaches

[Review: Maas'17]

 $Z = \int_{\Omega} D \phi^a e^{iS[\phi]}$

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- no anomalies

Action is invariant $S[\phi] = S[G\phi]$

iS[ϕ

 $D\phi$

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Integration range - contains all orbits $G \phi$

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$\langle \phi^b(x) \rangle = \int_{\Omega} D \phi^a e^{iS[\phi]} \phi^b(x)$

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- There is no preferred point on the group orbit
 - There is no absolute orientation/frame in the internal space
 - Does not change when averaging over position
 - There is no absolute charge

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$$\langle \phi^b(x) \rangle = \int_{\Omega} D \phi^a e^{iS[\phi]} \phi^b(x) = 0$$

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$$\langle \phi^b(x)\phi^c(y)\rangle = \int_{\Omega} D\phi^a e^{iS[\phi]}\phi^b(x)\phi^c(y)$$

• Relative charge measurement averaged over all possible starting points

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- Relative charge measurement averaged over all possible starting point
 - Vanishes because no preferred absolute starting point

[Review: Maas'17]

$\left\langle \delta_{bc} \phi^{b}(x) \phi^{c}(y) \right\rangle$ $= \int_{\Omega} D \phi^{a} e^{iS[\phi]} \delta_{bc} \phi^{b}(x) \phi^{c}(y)$

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 - Measures relative orientation
 - Created from an invariant tensor $\,\delta_{ab}\,$

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Field – transforms locally under a group $\phi^a(x) \rightarrow G^{ab}(x) \phi^b(x)$

 $Z = \int_{\Omega} D(\phi^{a}) e^{iS[\phi]}$

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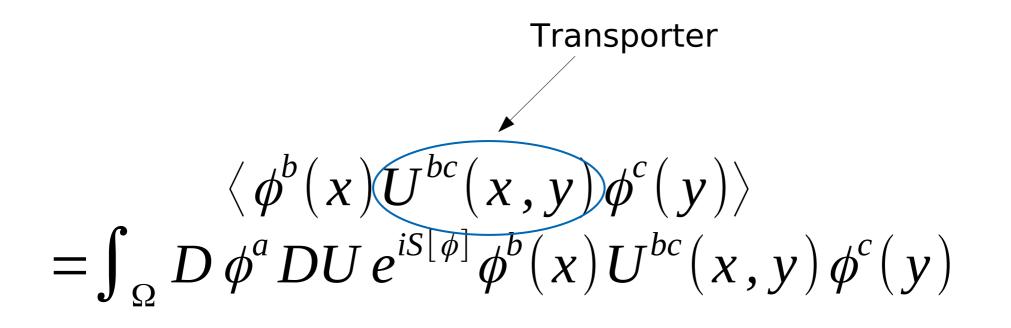
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- No longer invariant under gauge transformations
 - Vanishes just as any other non-invariant quantity

[Review: Maas'17]



•Transporter compensates gauge transformations

Implemented by gauge fields

[Review: Maas'17]

$\langle \phi^{b}(x) U^{bc}(x, y) \phi^{c}(y) \rangle$ = $\int_{\Omega} D \phi^{a} D U e^{iS[\phi]} \phi^{b}(x) U^{bc}(x, y) \phi^{c}(y) \neq 0$

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Reduced integration range

$= \int_{\Omega_c} \Phi^a DU W(U, \phi) e^{iS[\phi]} \phi^b(x) \phi^c(y) \neq 0$

- Reduction of integration region by gauge fixing
 - Arbitrary choice of coordinates
 - Weight factor to keep gauge-invariant quantities the same

• QFT setting – no strings or other non-QFT settings

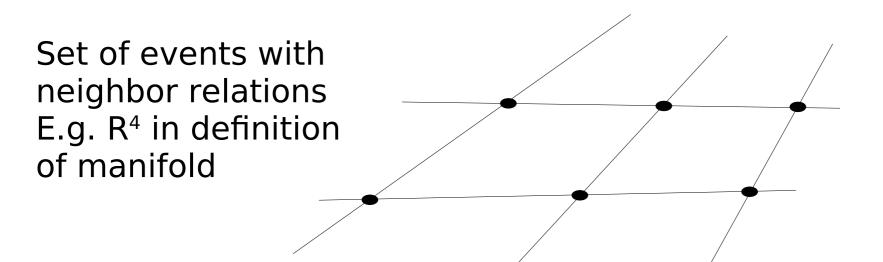
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- Diffeomorphism is like a non-standard gauge symmetry
 - Arbitrary local choices of coordinates do not affect observables – pure passive formulation
 - Physical observables must be manifestly invariant

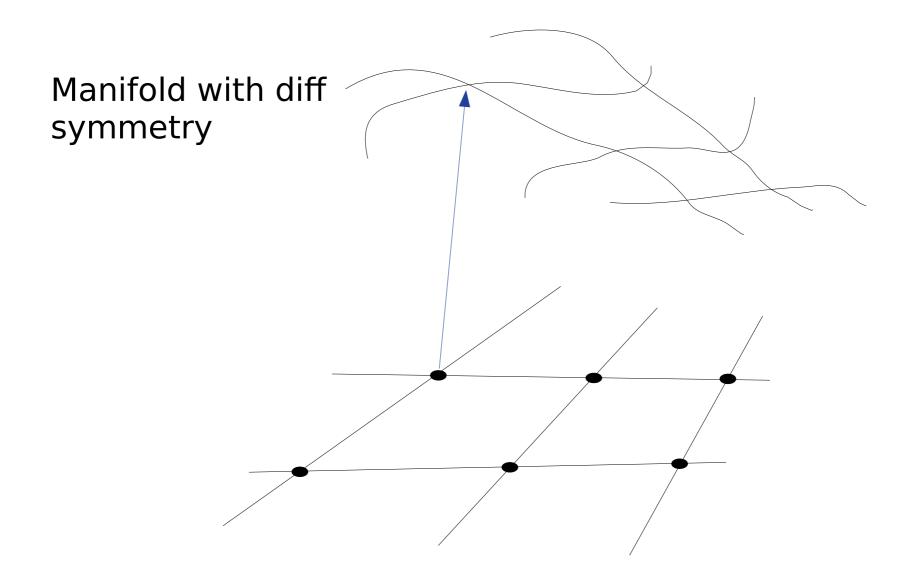
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 - Spin degeneracies and selection rules due to spin conservation
 - Global structure

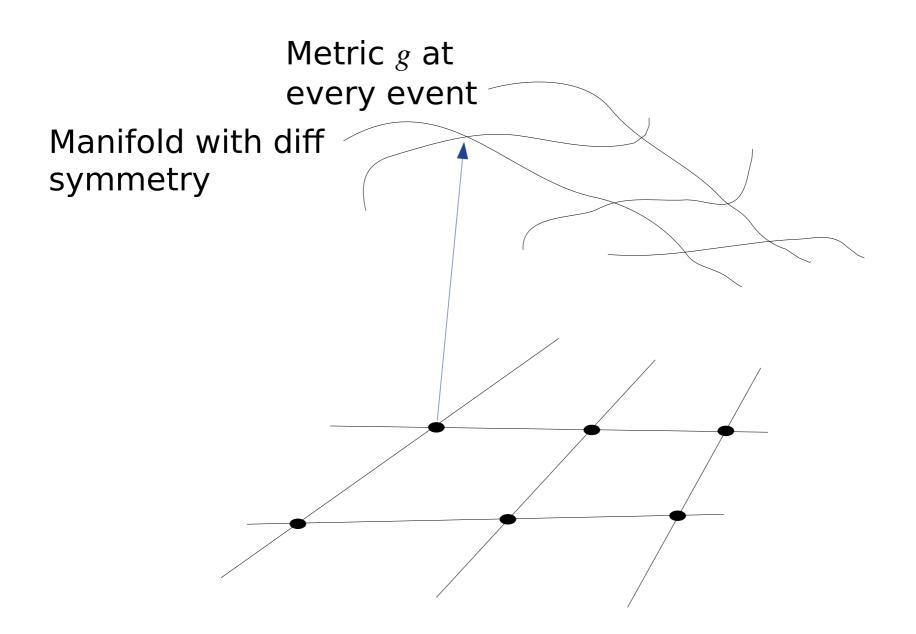
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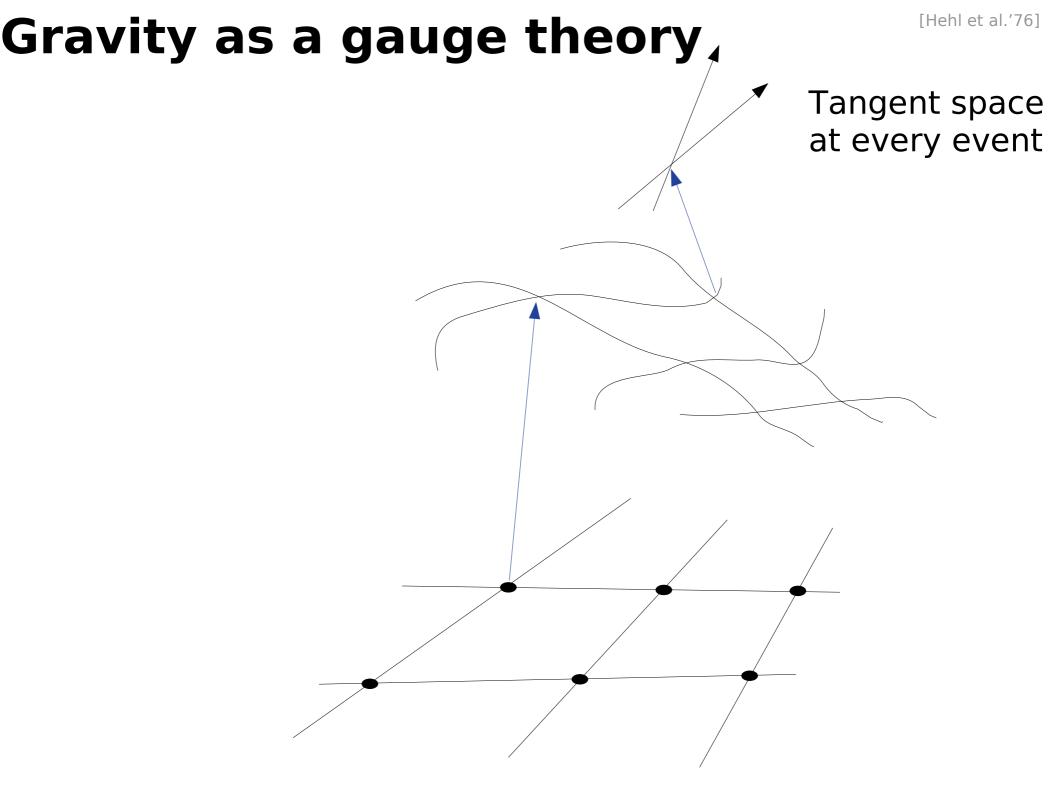
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- Particle physics gauge symmetries and global symmetries should remain the same

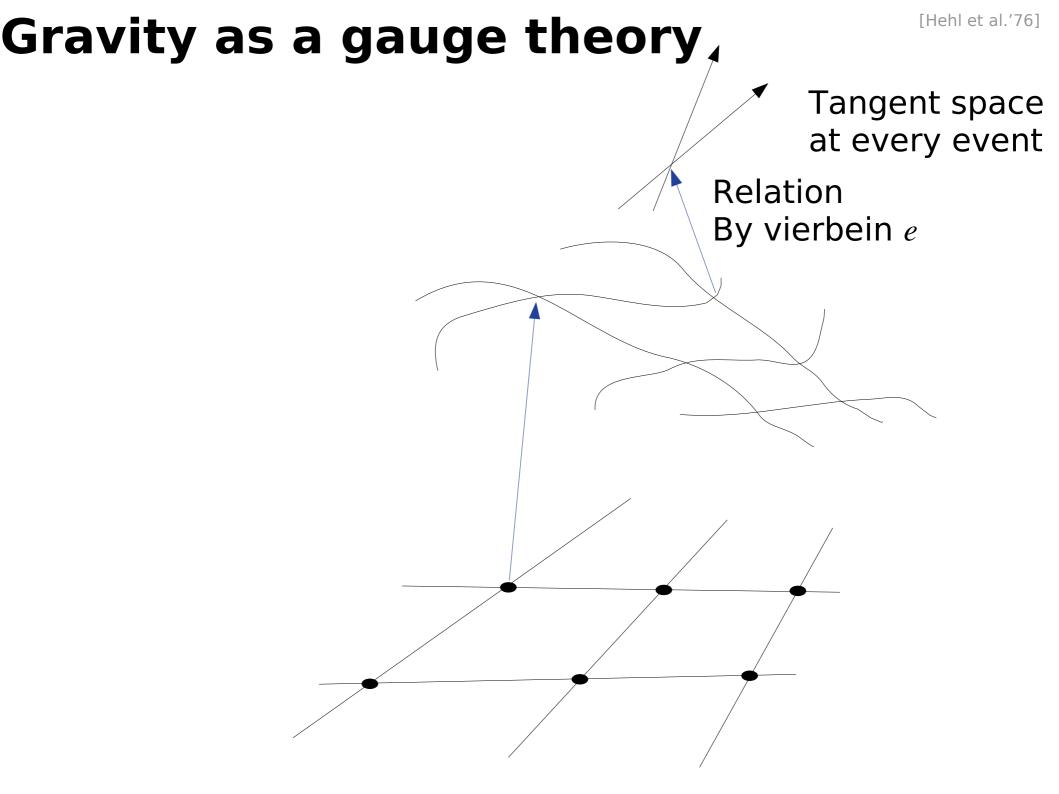
Set of events with neighbor relations

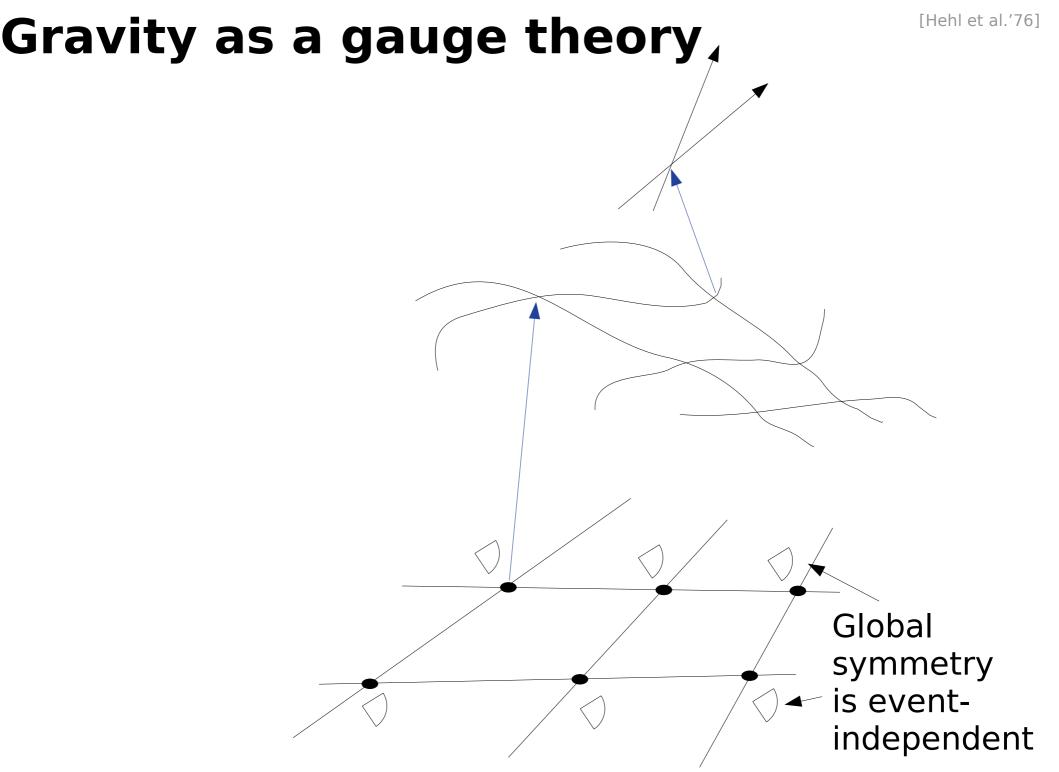


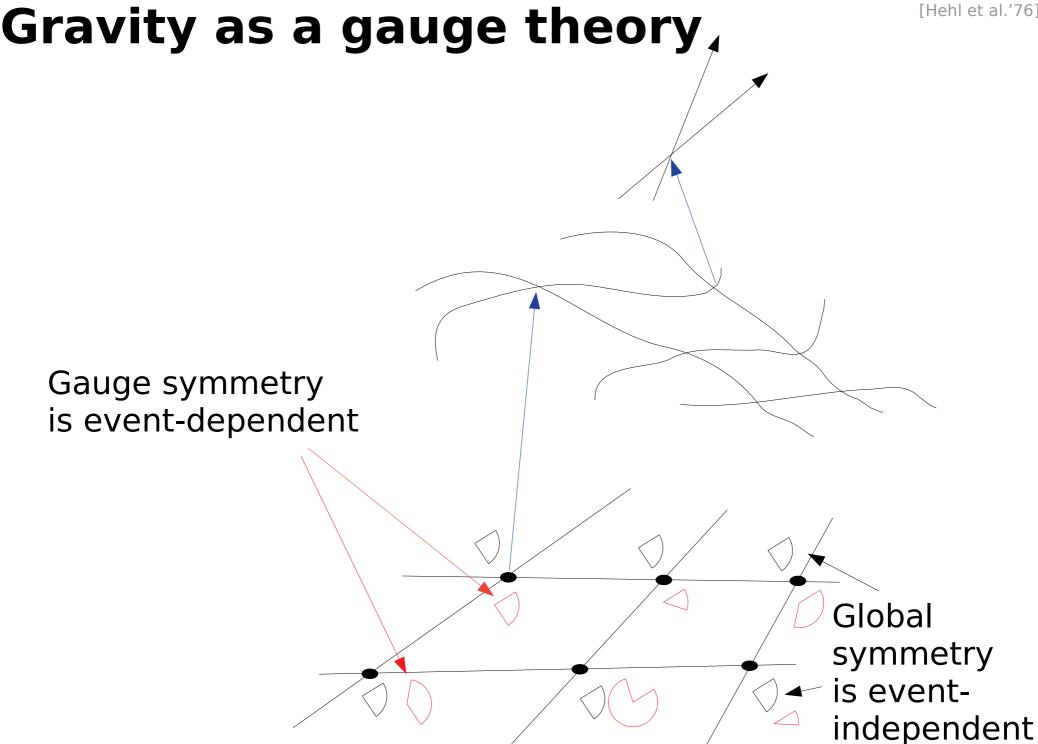










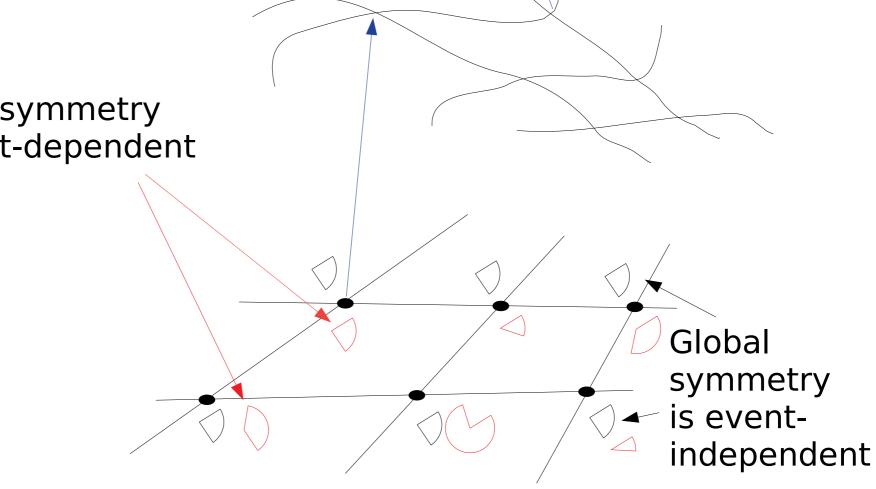




Gravity as a gauge theory

Internal symmetries act in internal spaces Global: One internal space Local: One space at every event

Gauge symmetry is event-dependent

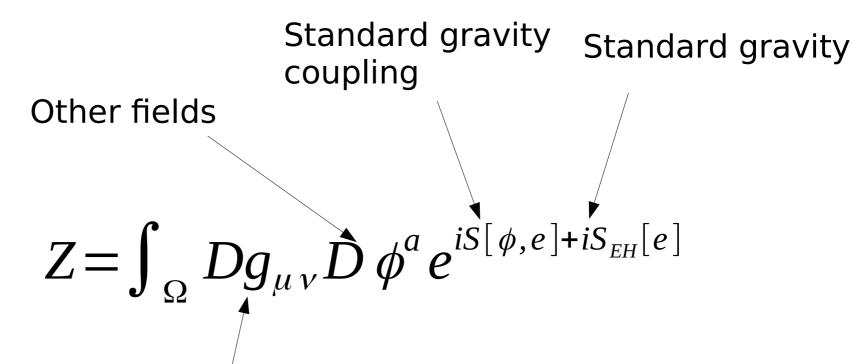


$Z = \int_{\Omega} Dg_{\mu\nu} D\phi^a e^{iS[\phi,e] + iS_{EH}[e]}$

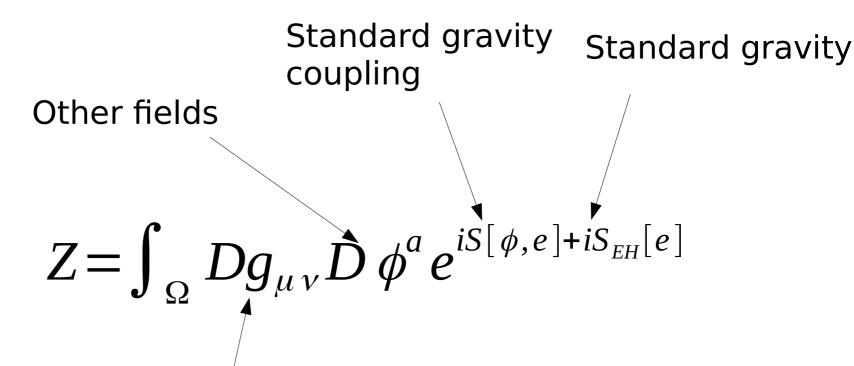
Standard gravity

$$Z = \int_{\Omega} Dg_{\mu\nu} D \phi^a e^{iS[\phi,e] + iS_{EH}[e]}$$

- Integration variable currently arbitrary choice
 - Here: Metric not relevant at leading order
 - Other choices (e.g. vierbein) possible



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- Otherwise standard
 - E.g. Asymptotic safety for ultraviolet stability

[Maas'19]

Dynamical formulation

$\langle O \rangle = \int_{\Omega} Dg_{\mu\nu} D \phi^a O e^{iS[\phi,e] + iS_{EH}[e]}$

[Maas'19]

Dynamical formulation

$0 \neq \langle O \rangle = \int_{\Omega} Dg_{\mu\nu} D \phi^{a} O e^{iS[\phi,e] + iS_{EH}[e]}$

$\bigoplus_{\mathbf{A}} \neq \langle O \rangle = \int_{\Omega} Dg_{\mu\nu} D \phi^{a} O e^{iS[\phi,e] + iS_{EH}[e]}$

Needs to be invariant

- Locally under Diffeomorphism
- Locally under Lorentz transformation
- Locally under gauge transformation
- Globally under custodial,... transformation to be non-zero

Space-time structure

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$$\frac{\langle \int d^d x \sqrt{\det g} R(x) \rangle}{\langle \int d^d x \sqrt{\det g} \rangle} = const$$

- No preferred events
 - Space-time on average homogenous and isotropic
 - Average space-time is flat or (anti-)de Sitter for canonical gravity
 - Invariants identify the particular type

- Consider a scalar particle
 - E.g. described by a scalar field
 - Completely invariant

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Completely scalar: Invariant under all symmetries

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Argument is the event, not the coordinate

Result depends on events

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 - Events not a useful argument

Some distance function

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- Distance is a quantum object: Expectation value
 - Needs a diff-invariant formulation

$$\langle O(x)O(y)\rangle = D(r(x,y))$$
$$r(x,y) = \langle \min_{z} \int_{x}^{y} d\lambda g_{\mu\nu} \frac{dz^{\mu}}{d\lambda} \frac{dz^{\nu}}{d\lambda} \rangle$$

 $\langle \mathbf{o}(\cdot) \rangle \mathbf{o}(\cdot) \rangle = \mathbf{p}(\cdot(\cdot))$

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$$\langle O(x)O(y)\rangle = D(r(x,y))$$
 Separate calculation
$$r(x,y) = \langle \min_{z} \int_{x}^{y} d\lambda g_{\mu\nu} \frac{dz^{\mu}}{d\lambda} \frac{dz^{\nu}}{d\lambda} \rangle$$

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 - Needs to be determined separately

Reduces the full dependence: Definition Dependence on events will only vanish if all events on the average are equal – probably true

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- Generalization of flat-space arguments

[Maas et al.'22]

What about cosmology?

• Big bang a preferred event - not possible!

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 $\langle O(x)P(x)...Q(y_1)...R(y_n)\rangle$

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- Distances between y_i space-like
- Evolution of a matter/curvature concentration

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- A universe is a scattering process

- [Fröhlich et al.'80,'81 Review: Maas'17]
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Review: Maas'171

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- Works very well in particle physics

[Fröhlich et al.'80,'81 Review: Maas'17]

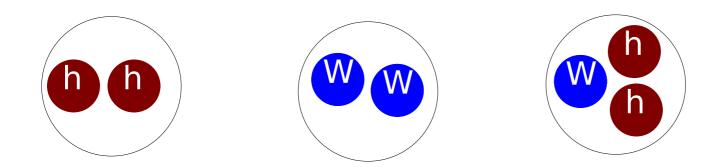
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 - Cannot be observable
- Gauge-invariant states are composite
 - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



[Fröhlich et al.'80,'81 Maas'12,'17]

1) Formulate gauge-invariant operator

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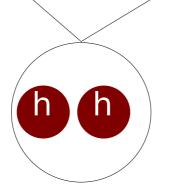
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Higgs field

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2 x Higgs mass: Scattering state

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Standard Perturbation Theory

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- [Fröhlich et al.'80,'81 Maas'12,'17 Maas & Sondenheimer'20 Dudal et al.'20]
- 1) Formulate gauge-invariant operator

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Deviations
$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+ (x)\eta(y) \rangle$$

 $+ v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$

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[Fröhlich et al.'80, Egger, Maas, Sondenheimer'17]

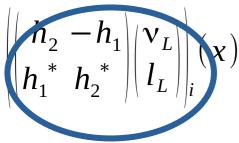
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$$\begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} _i (\mathbf{x})$$

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$$\begin{pmatrix} h_2 - h_1 \\ h_1^* & h_2^* \end{pmatrix} \begin{pmatrix} \mathbf{v}_L \\ l_L \end{pmatrix} (x)$$

• Gauge-invariant state, but custodial doublet

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
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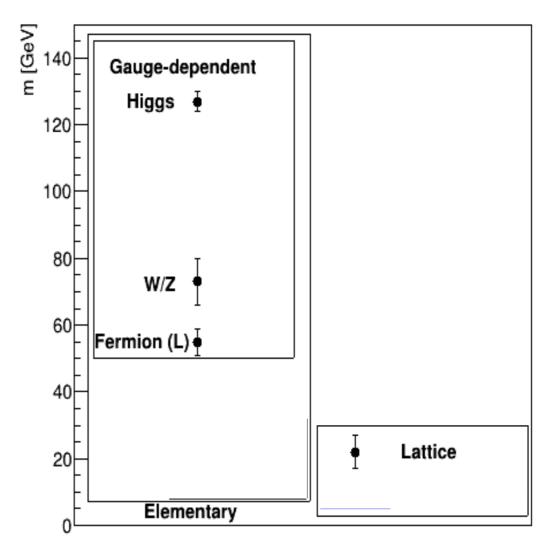
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- Can this be true? Lattice test

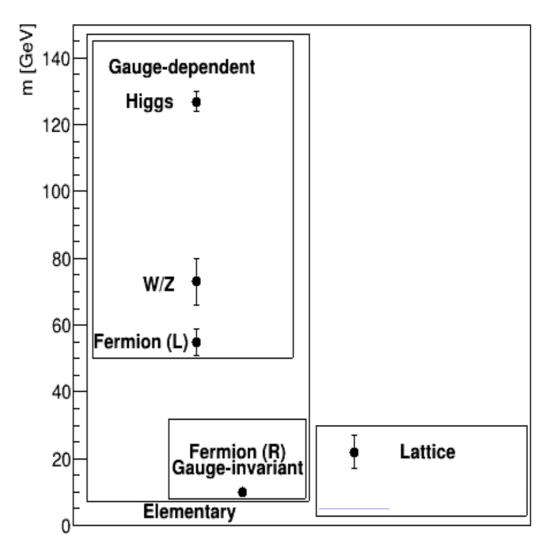
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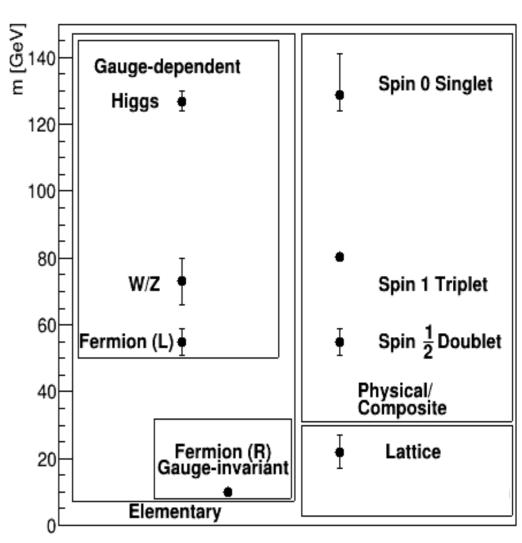
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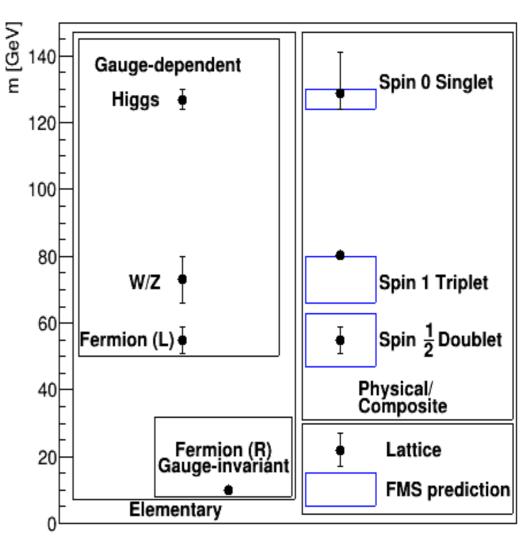
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 - Will avoid Bloch-Nordsieck violations
- Valence Higgs detectable at LHC?
 - Strong couplings to Higgs: tops, weak gauge bosons

Further consequences

- In SM physics: Quantitative changes
 - Anomalous couplings/form factors
 - (Small) differences in various kinematic regimes
 - More: See 1701.00182, 1811.03395, 2002.01688, 2008.07813, 2009.06671, 2204.02756, 2212.08470

Further consequences

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 - (Small) differences in various kinematic regimes
 - More: See 1701.00182, 1811.03395, 2002.01688, 2008.07813, 2009.06671, 2204.02756, 2212.08470
- In BSM physics: Sometimes qualitative changes
 - Even different spectrum
 - Affects viability of BSM Scenarios
 - More: See 1709.07477, 1804.04453, 1912.086680, 2002.08221, 2211.05812, 2211.16937

Applying FMS to gravity

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- FMS split after (convenient) gauge fixing
 - $g_{\mu\nu} = g^c_{\mu\nu} + \gamma_{\mu\nu}$
 - Classical part g^c is a metric, chosen to give exact (observed) curvature
 - Quantum part is needed (assumed) small

Details (and challenges)

- Classical metric needs to be useful
 - Should not have special events
 - Only flat and (anti-)de Sitter possible
 - Should satisfy gauge choice

[Maas et al.'22]

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 - Inverse fluctuation satisfies Dyson equation

$$\gamma^{\mu\nu} = -(g^c)^{\mu\sigma} \gamma_{\sigma\rho}((g^c)^{\rho\nu} + \gamma^{\rho\nu})$$

• Infinite series at tree-level

$$r(x,y) = \langle \min_{z} \int_{x}^{y} d\lambda g_{\mu\nu} \frac{dz^{\mu}}{d\lambda} \frac{dz^{\nu}}{d\lambda} \rangle$$

Application to distance between two events

$$r(x, y) = \langle \min_{z} \int_{x}^{y} d\lambda g_{\mu\nu} \frac{dz^{\mu}}{d\lambda} \frac{dz^{\nu}}{d\lambda} \rangle$$
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Classical geodesic distance

- Application to distance between two events
 - Yields to leading order classical distance

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Classical geodesic distance

Quantum corrections

- Application to distance between two events
 - Yields to leading order classical distance
 - Yields at leading-order classical space-time
 - Quantum corrections depends on events

 $\langle O(x)O(y) \rangle$

[Maas'19]

$\langle O(x)O(y)\rangle = D_c(r^c) + \sum (\delta r)^n \partial_r^n D_c(r) + \langle O(x)O(y)\rangle_{\gamma}$

Double expansion

$$\langle O(x)O(y)\rangle = D_c(r^c) + \sum (\delta r)^n \partial_r^n D_c(r) + \langle O(x)O(y)\rangle_{\gamma}$$

Leading term is $D_c = \langle O(x) O(y) \rangle_{g^c}$ flat space propagator

• Double expansion

Corrections from quantum distance effects

$$\langle O(x)O(y)\rangle = D_c(r^c) + \sum (\delta r)^n \partial_r^n D_c(r) + \langle O(x)O(y)\rangle_{\gamma}$$

$$D_c = \langle O(x) O(y) \rangle_{g^c}$$

- Double expansion
 - Quantum fluctuations in the argument

Corrections from metric fluctuations

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 - Consistent with EDT results [Dai'22]

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- Reduces to QFT at vanishing gravity
 - Higgs and W/Z mass in quantum gravity calculated

 Pure gravity excitation: Curvaturecurvature correlator

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$$\langle R(x)R(y)\rangle = D^{\mu\nu\rho\sigma}\langle \gamma_{\mu\nu}(x)\gamma_{\rho\sigma}(y)\rangle (d(x,y)) + O(\gamma^3)$$

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• In Minkowski space-time: No propagating mode at lowest order

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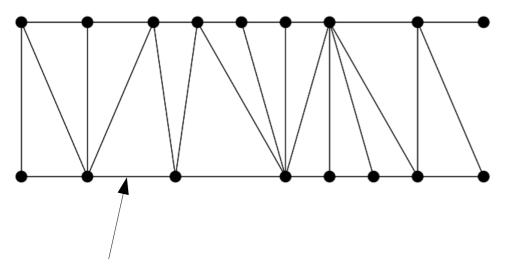
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Graviton propagator

- In Minkowski space-time: No propagating mode at lowest order
- Flat space: Better divergence properties

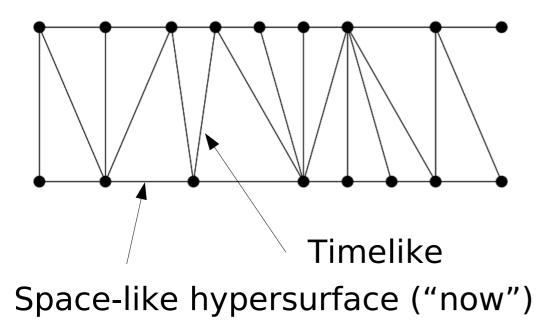
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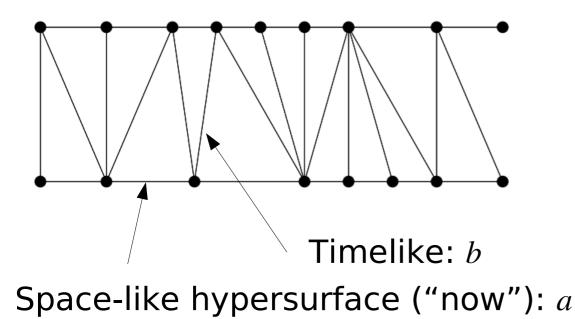


Space-like hypersurface ("now")

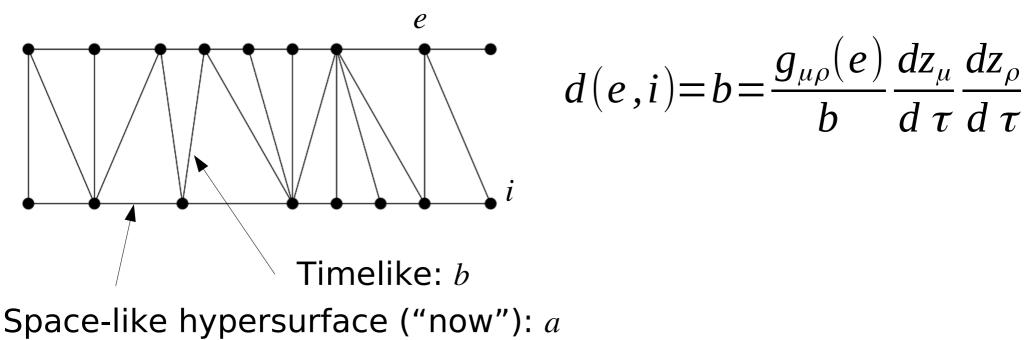
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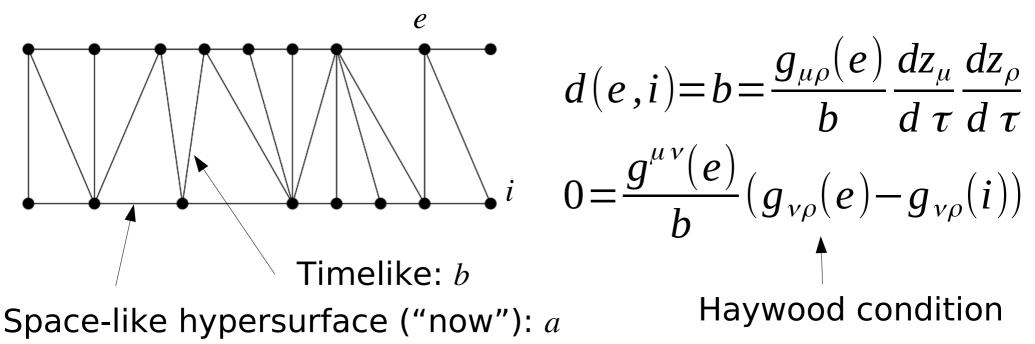
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Predictions for CDT

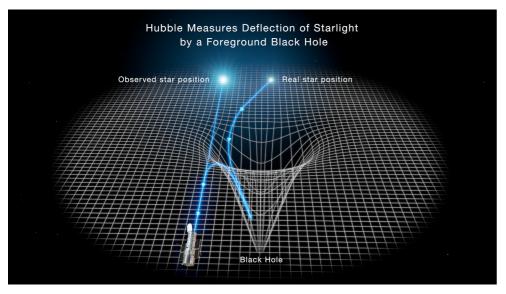
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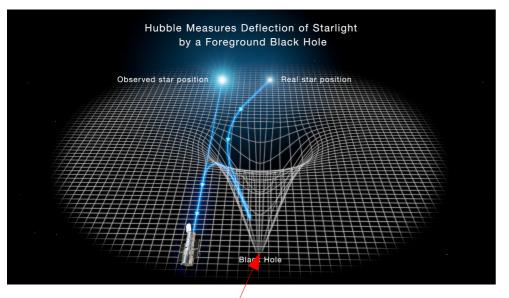
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- Geon propagator should behave as contracted metric propagator
 - As a function of the geodesic distance

- Macroscopic gravitational objects need to be build in the same way
 - Just like neutron stars from QCD

Classical picture of a black hole



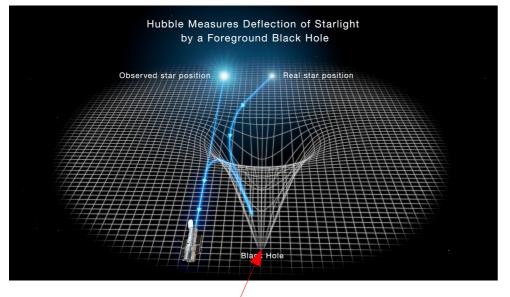
Classical picture of a black hole



But this is a special worldline, determining the full metric!

Not possible in a quantum expectation value.

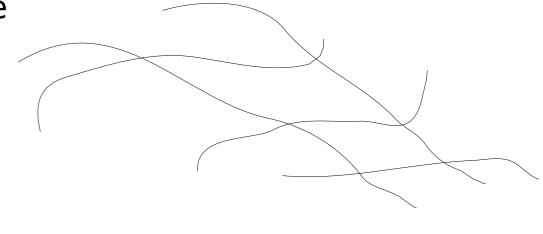
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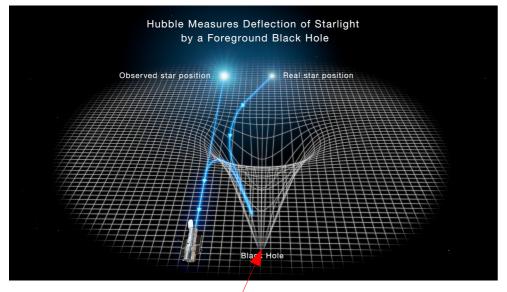
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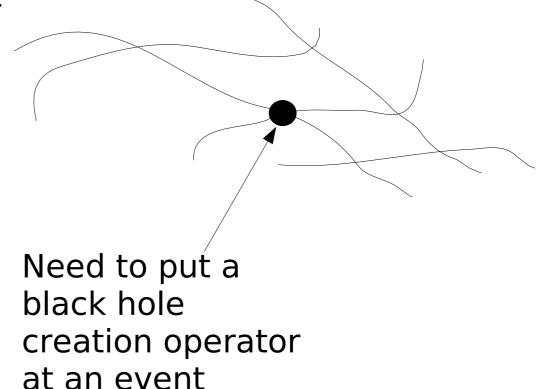
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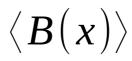


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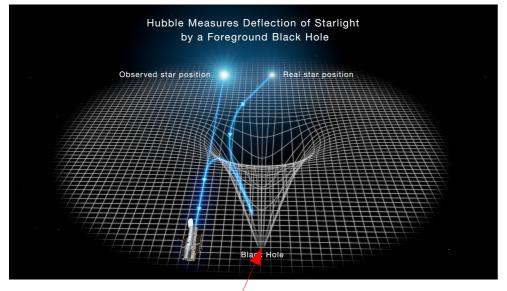
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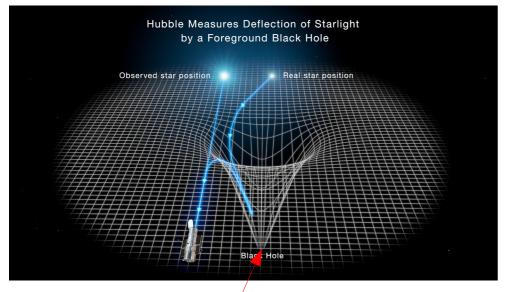
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Need to put a black hole creation operator at an event but it would still be a constant

 $\langle B(x) \rangle = d$

Classical picture of a black hole



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Need to put a Need to black hole correlate with creation operator e.g. curvature at an event but it would still be a constant

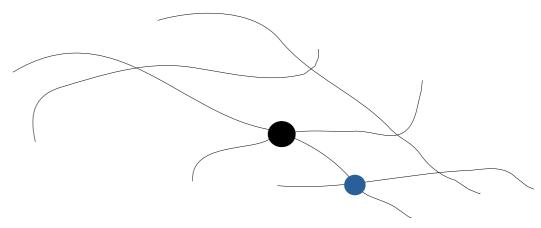
 $\langle B(x)R(y)\rangle = d(r(x,y))$

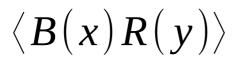
[Maas et al'22]

Views of black holes

Averaging over this!

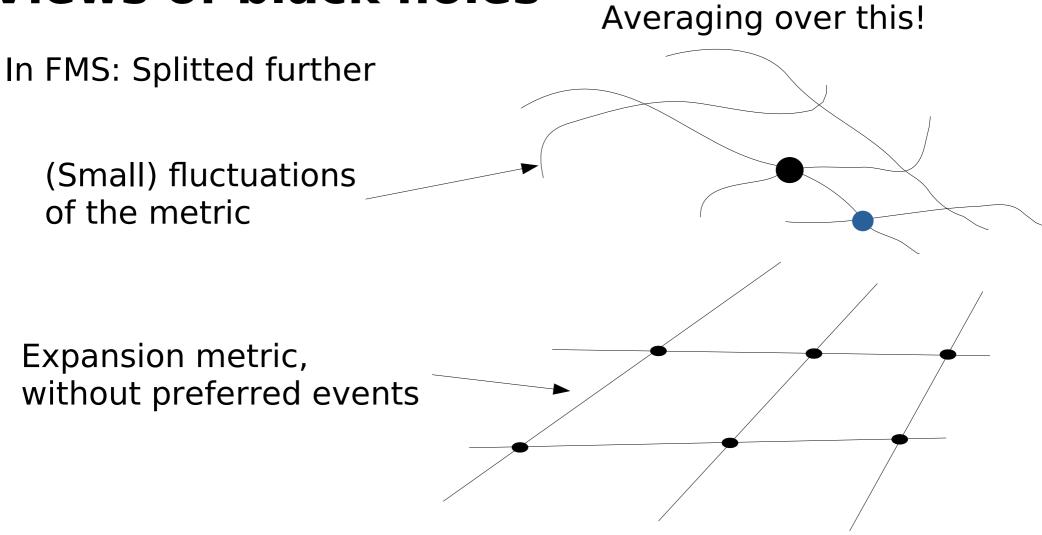
In FMS: Splitted further



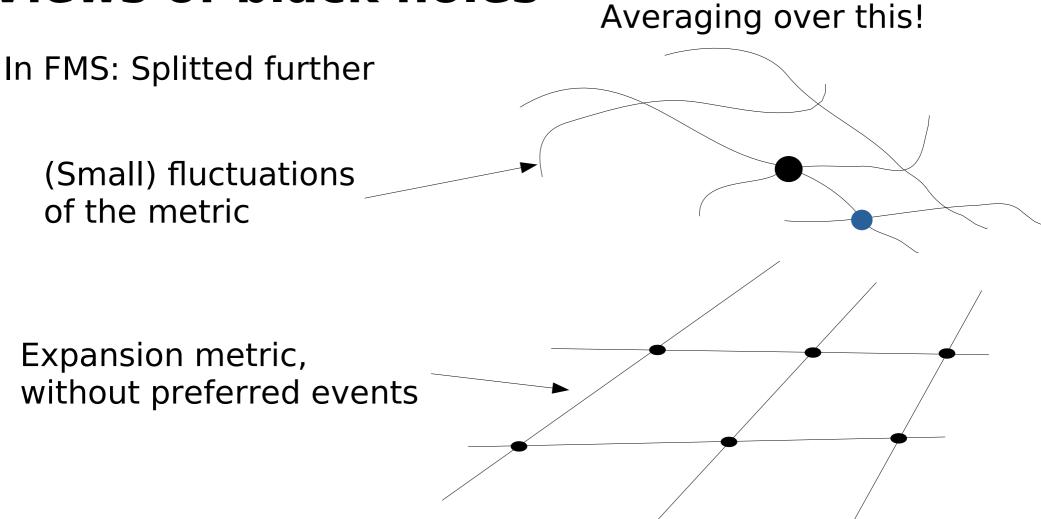


[Maas et al'22]

Views of black holes



 $\langle B(x)R(y)\rangle$



Calculate expansion:

 $\langle B(x)R(y)\rangle = d^{c}(r^{c}(x,y)) + quantum$ Classical field result

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 - Pure: Geon star, similar to neutron star

Summary

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 FMS mechanism allows estimates of quantum effects in a systematic expansion



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 FMS mechanism allows estimates of quantum effects in a systematic expansion

 Gives a new perspective on strong and quantum gravity



People involved



Christina Perner (Bachelor)



Markus Markl (Master)



Michael Müller (Master)



Bernd Riederer (PhD)



Simon Plätzer (Staff)

