

Geons as gravity balls

Axel Maas

 **@axelmaas**

20th of August 2025
ÖPG/SPG Meeting
Vienna
Austria



NAWI Graz
Natural Sciences

FWF Österreichischer
Wissenschaftsfonds



Setup

- QFT setting – no strings or other non-QFT structures

Setup

$$Z = \int_{\Omega} Dg_{\mu\nu} D\phi^a e^{iS[\phi, e] + iS_{EH}[e]}$$

- QFT setting – no strings or other non-QFT structures

Setup

Standard gravity


$$Z = \int_{\Omega} Dg_{\mu\nu} D\phi^a e^{iS[\phi, e] + iS_{EH}[e]}$$

- QFT setting – no strings or other non-QFT structures

Setup

Standard gravity


$$Z = \int_{\Omega} Dg_{\mu\nu} D\phi^a e^{iS[\phi, e] + iS_{EH}[e]}$$

- QFT setting – no strings or other non-QFT structures
- Diffeomorphism is like a gauge symmetry [Hehl et al.'76]
 - Arbitrary local choices of coordinates do not affect observables – pure passive formulation

Setup

Standard gravity

$$Z = \int_{\Omega} Dg_{\mu\nu} D\phi^a e^{iS[\phi, e] + iS_{EH}[e]}$$

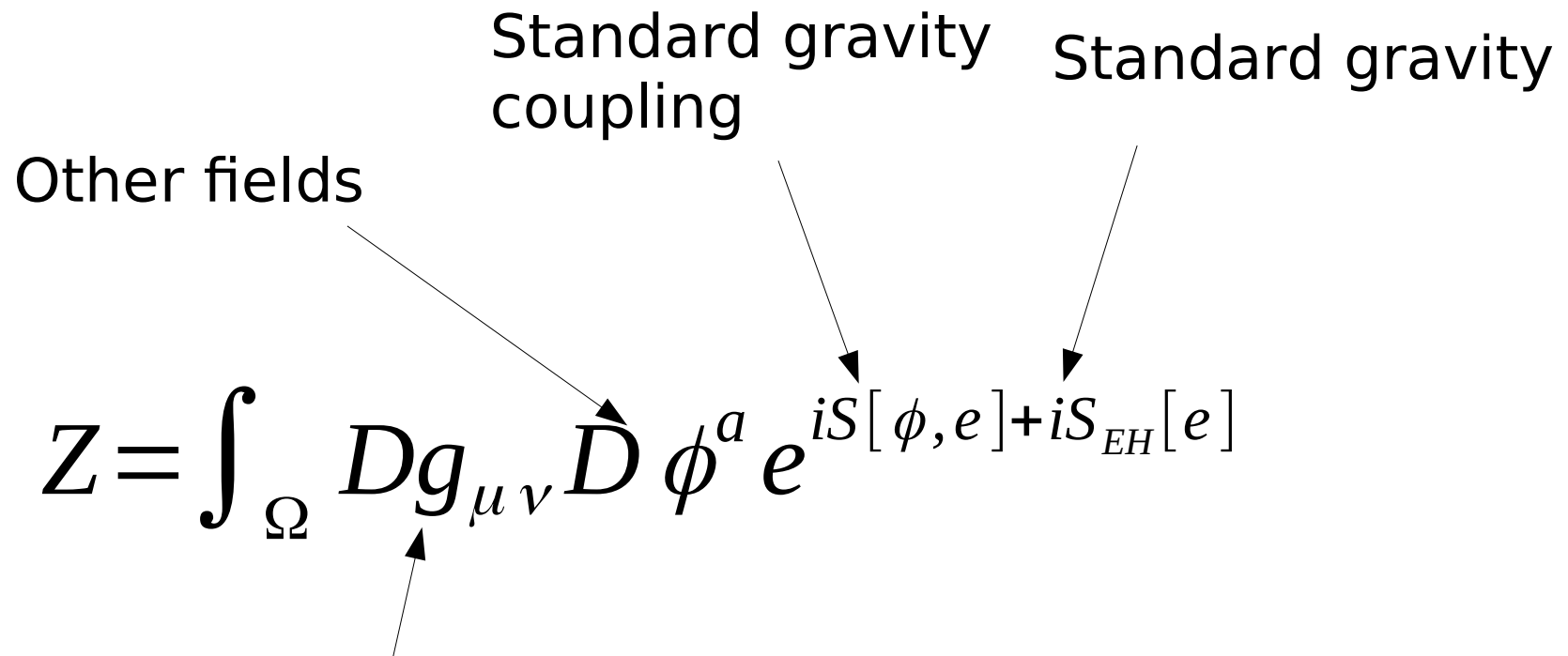
- Integration variable currently arbitrary choice
 - Manifold topologies when factoring out diffeomorphisms
 - Other choices (e.g. vierbein) possible when integrating over diffeomorphism orbits

Setup

Other fields

Standard gravity coupling

Standard gravity

$$Z = \int_{\Omega} Dg_{\mu\nu} \vec{D} \phi^a e^{iS[\phi, e] + iS_{EH}[e]}$$


- Integration variable currently arbitrary choice
 - Manifold topologies when factoring out diffeomorphisms
 - Other choices (e.g. vierbein) possible when integrating over diffeomorphism orbits

Setup

Standard gravity coupling Standard gravity

Other fields

$$Z = \int_{\Omega} Dg_{\mu\nu} \vec{D} \phi^a e^{iS[\phi, e] + iS_{EH}[e]}$$

- Integration variable currently arbitrary choice
 - Manifold topologies when factoring out diffeomorphisms
 - Other choices (e.g. vierbein) possible when integrating over diffeomorphism orbits
- Setup of FRG/Asymptotic safety, CDT, EDT

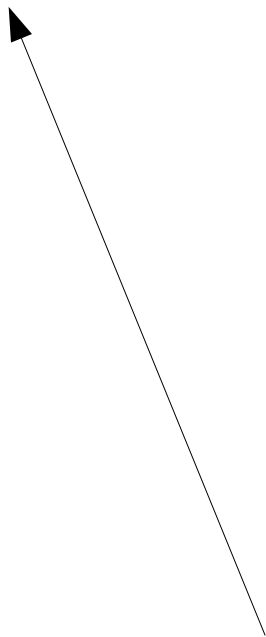
Dynamical formulation

$$\langle O \rangle = \int_{\Omega} Dg_{\mu\nu} D\phi^a O e^{iS[\phi, e] + iS_{EH}[e]}$$

Dynamical formulation

$$0 \neq \langle O \rangle = \int_{\Omega} Dg_{\mu\nu} D\phi^a O e^{iS[\phi, e] + iS_{EH}[e]}$$

non-zero



Dynamical formulation

$$0 \neq \langle O \rangle = \int_{\Omega} Dg_{\mu\nu} D\phi^a O e^{iS[\phi, e] + iS_{EH}[e]}$$

Needs to be invariant



to be non-zero



Dynamical formulation

$$0 \neq \langle O \rangle = \int_{\Omega} Dg_{\mu\nu} D\phi^a O e^{iS[\phi, e] + iS_{EH}[e]}$$

Needs to be invariant

- Locally under Diffeomorphism
- Locally under Lorentz transformation

to be non-zero

Dynamical formulation

$$0 \neq \langle O \rangle = \int_{\Omega} Dg_{\mu\nu} D\phi^a O e^{iS[\phi, e] + iS_{EH}[e]}$$

Needs to be invariant

- Locally under Diffeomorphism
- Locally under Lorentz transformation
- Locally under gauge transformation
- Globally under custodial,... transformation

to be non-zero

- Average metric vanishes: $\langle g_{\mu\nu}(x) \rangle = 0$
 - No preferred events, maximally symmetric

Observables

- Average metric vanishes: $\langle g_{\mu\nu}(x) \rangle = 0$
 - No preferred events, maximally symmetric
- Observables need to be diffeomorphism-invariant
 - Average space-time is an observation from invariants like the curvature scalars

Observables

- Average metric vanishes: $\langle g_{\mu\nu}(x) \rangle = 0$
 - No preferred events, maximally symmetric
- Observables need to be diffeomorphism-invariant
 - Average space-time is an observation from invariants like the curvature scalars
 - Local observables need to be composite
 - E.g. again curvature scalars

Observables

- Average metric vanishes: $\langle g_{\mu\nu}(x) \rangle = 0$
 - No preferred events, maximally symmetric
- Observables need to be diffeomorphism-invariant
 - Average space-time is an observation from invariants like the curvature scalars
 - Local observables need to be composite
 - E.g. again curvature scalars
 - Arguments/distances need to be invariants
 - E.g. geodesic distances [Ambjorn et al.'12, Schaden '15, Maas'19]

Causal Dynamical Triangulation [Ambjorn et al.'12,19]

$$\langle O \rangle = \int_{\Omega} Dg_{\mu\nu} D\phi^a O e^{iS[\phi,e] + iS_{EH}[e]}$$

Causal Dynamical Triangulation

[Ambjorn et al.'12,19]

$$\langle O \rangle = \int_{\Omega} Dg_{\mu\nu} O e^{iS_{EH}[e]}$$

Causal Dynamical Triangulation [Ambjorn et al.'12,19]

Topology of manifold


$$\langle O \rangle = \int_{\Omega} Dd(X, Y) O e^{iS_{EH}[e]}$$

Causal Dynamical Triangulation

[Ambjorn et al.'12,19]

Topology of manifold
– diffeomorphism invariant!


$$\langle O \rangle = \int_{\Omega} Dd(X, Y) O e^{iS_{EH}[e]}$$

Causal Dynamical Triangulation

[Ambjorn et al.'12,19]

Topology of manifold
– diffeomorphism invariant!

$$\langle O \rangle = \int_{\Omega} Dd(X, Y) O e^{iS_{EH}[e]}$$

Restricted to foliable,
pseudo-Riemannian manifolds
with fixed global structure

Causal Dynamical Triangulation [Ambjorn et al.'12,19]

$$\langle O \rangle = \sum_{\Omega} D d(X_i, Y_j) O e^{iS_{EH}[e]}$$

Replaced with a finite, discrete triangulation
by simplices – basically tetrads
(Regge calculus)

Causal Dynamical Triangulation [Ambjorn et al.'12,19]

Wick rotation allows use
of standard Monte-Carlo
(lattice) techniques

$$\langle O \rangle = \sum_{\Omega} D d(X_i, Y_j) O e^{-S_{EH}[e]}$$

Replaced with a finite, discrete triangulation
by simplices – basically tetrads
(Regge calculus)

Causal Dynamical Triangulation [Ambjorn et al.'12,19]

Wick rotation allows use of standard Monte-Carlo (lattice) techniques

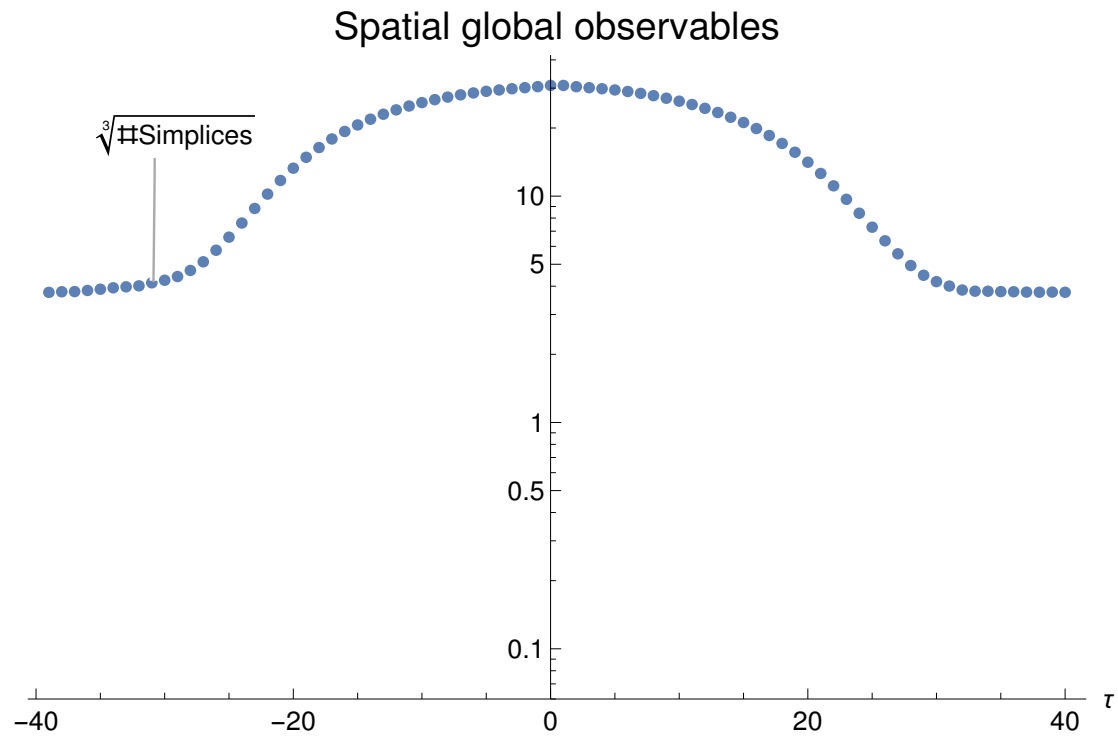
$$\langle O \rangle = \sum_{\Omega} Dd(X_i, Y_j) O e^{-S_{EH}[e]}$$

Replaced with a finite, discrete triangulation by simplices – basically tetrads (Regge calculus)

Systematic errors due to statistics, finite volume and discretization

Space-time in CDT

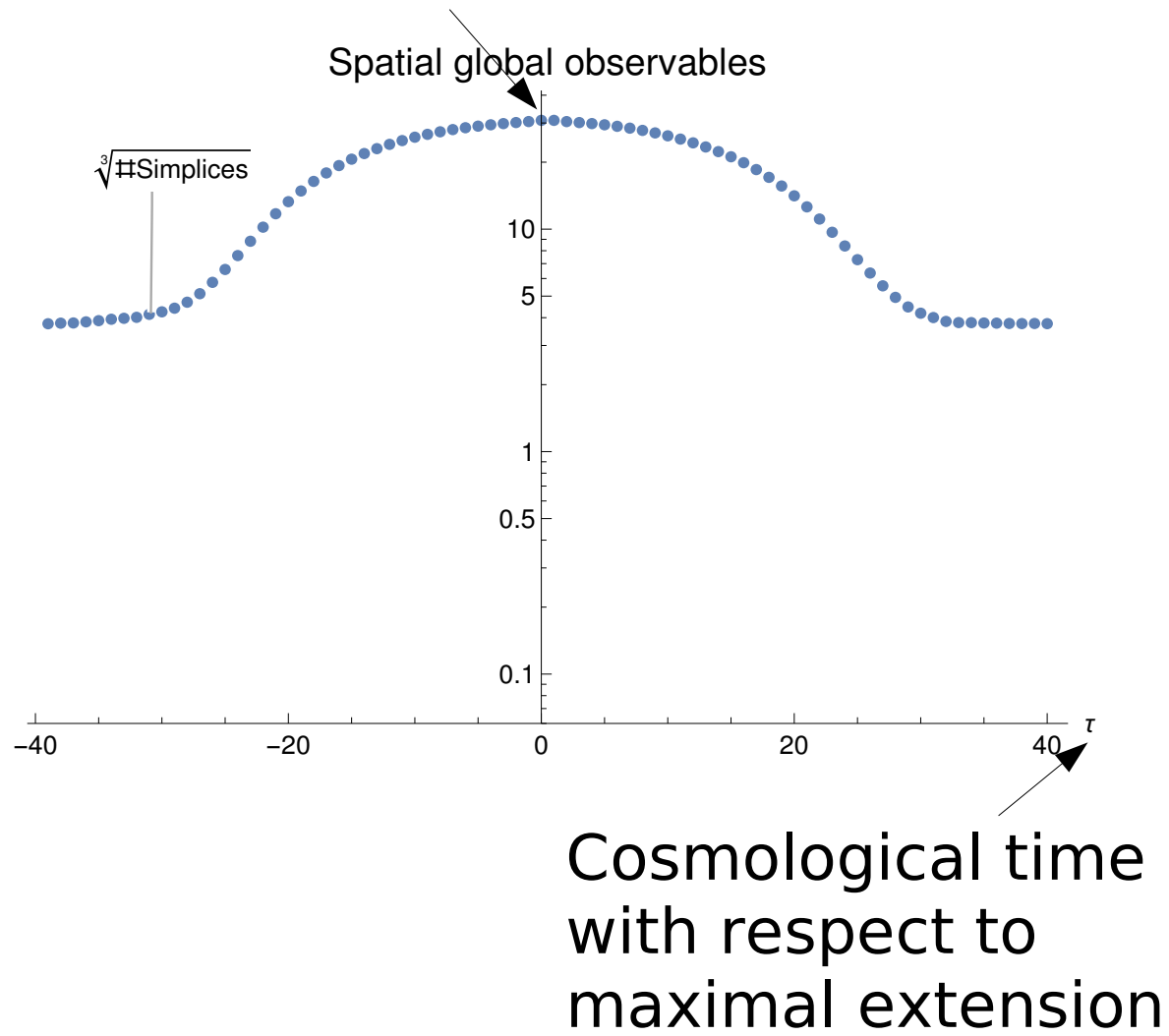
[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]

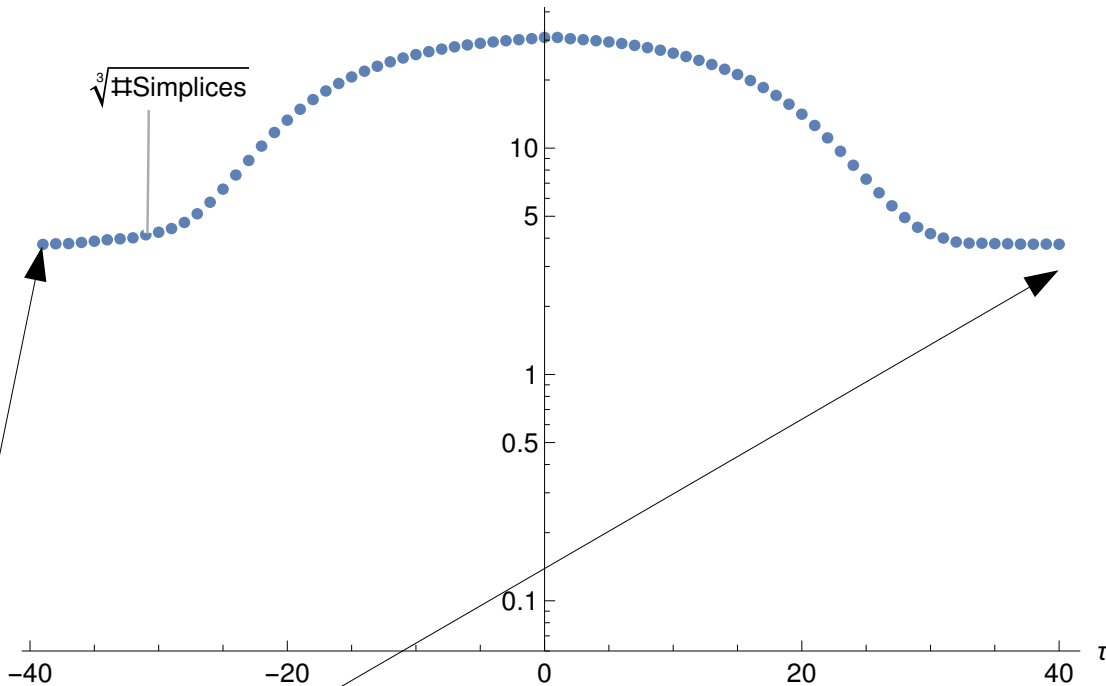
Maximum volume
constrained by finite size



Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]

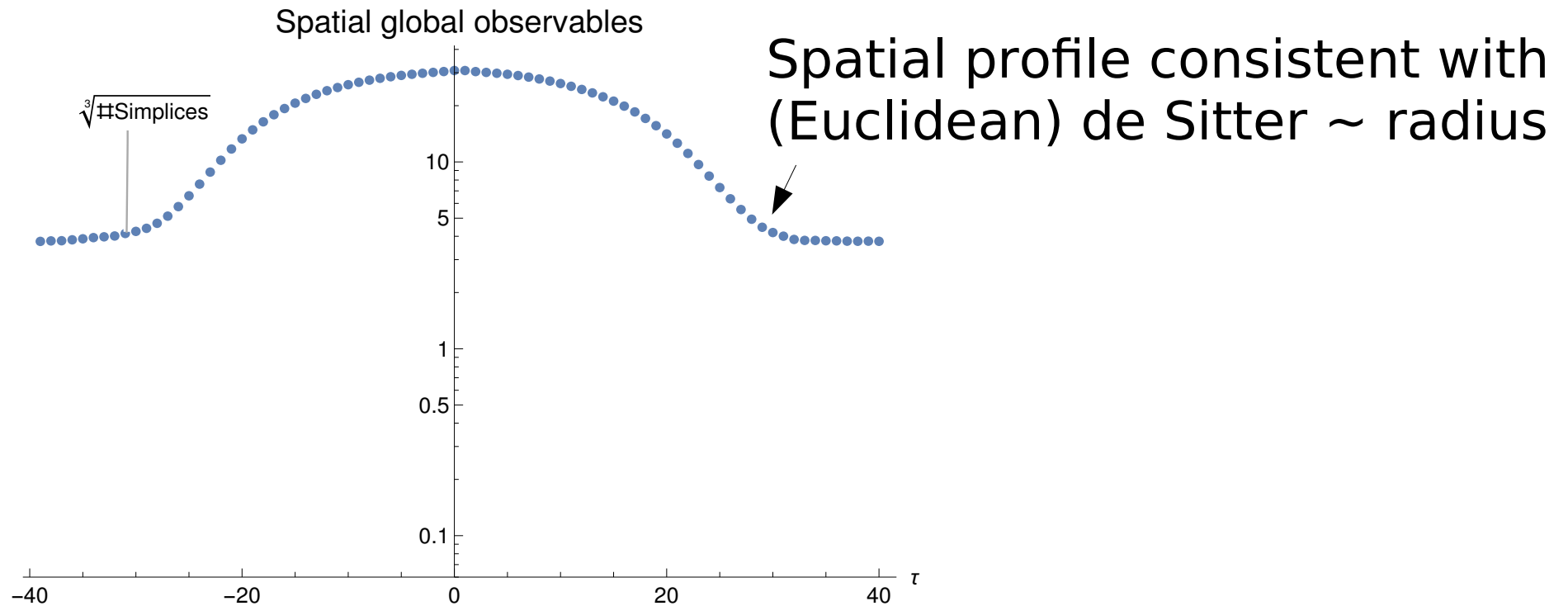
Spatial global observables



Periodicity
by boundary conditions

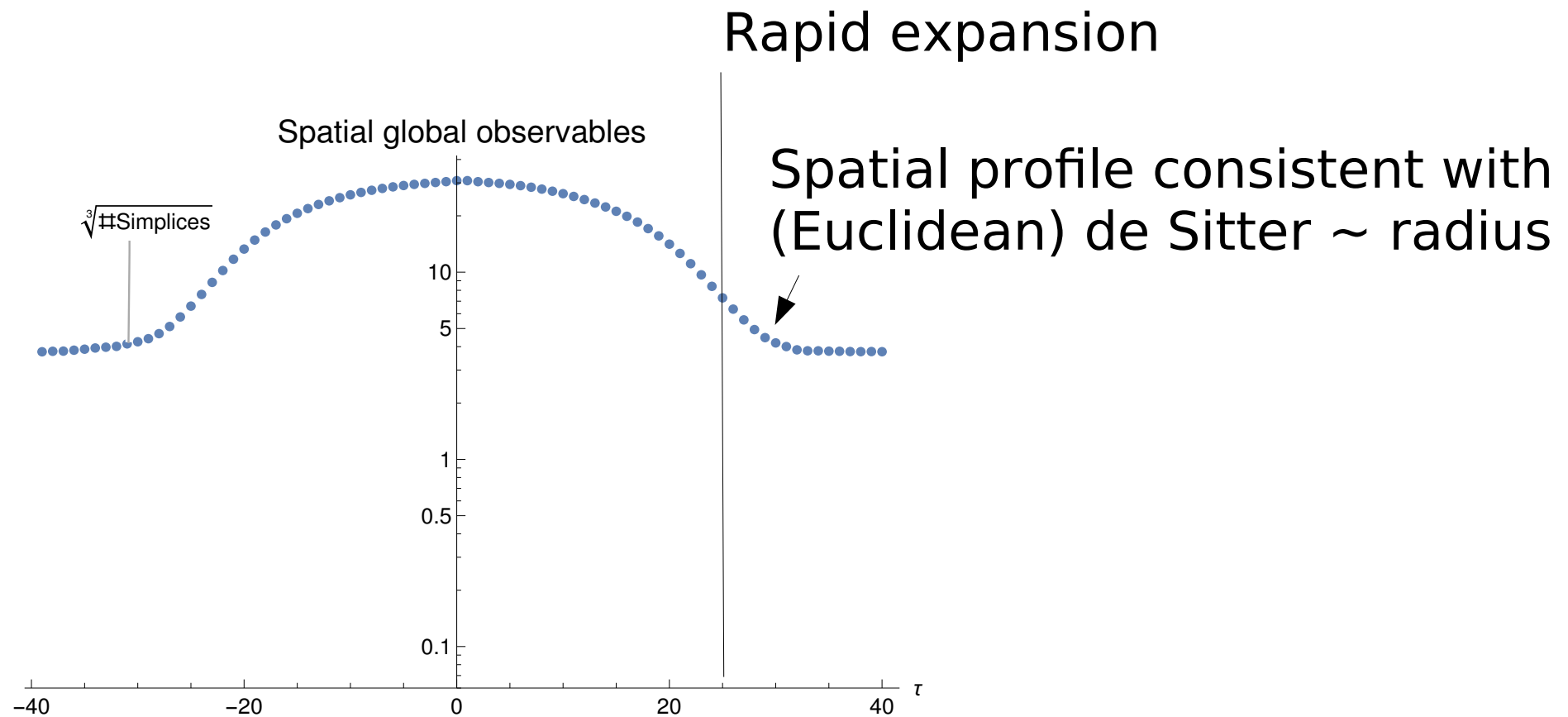
Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



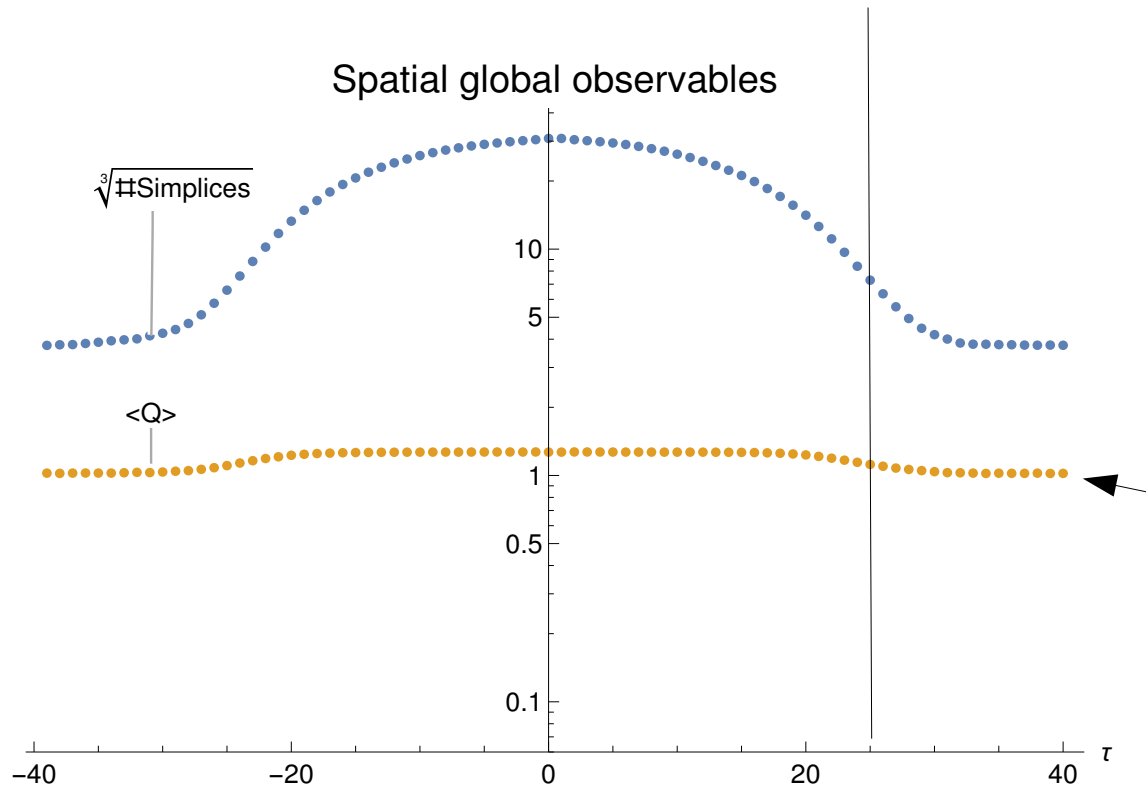
Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]

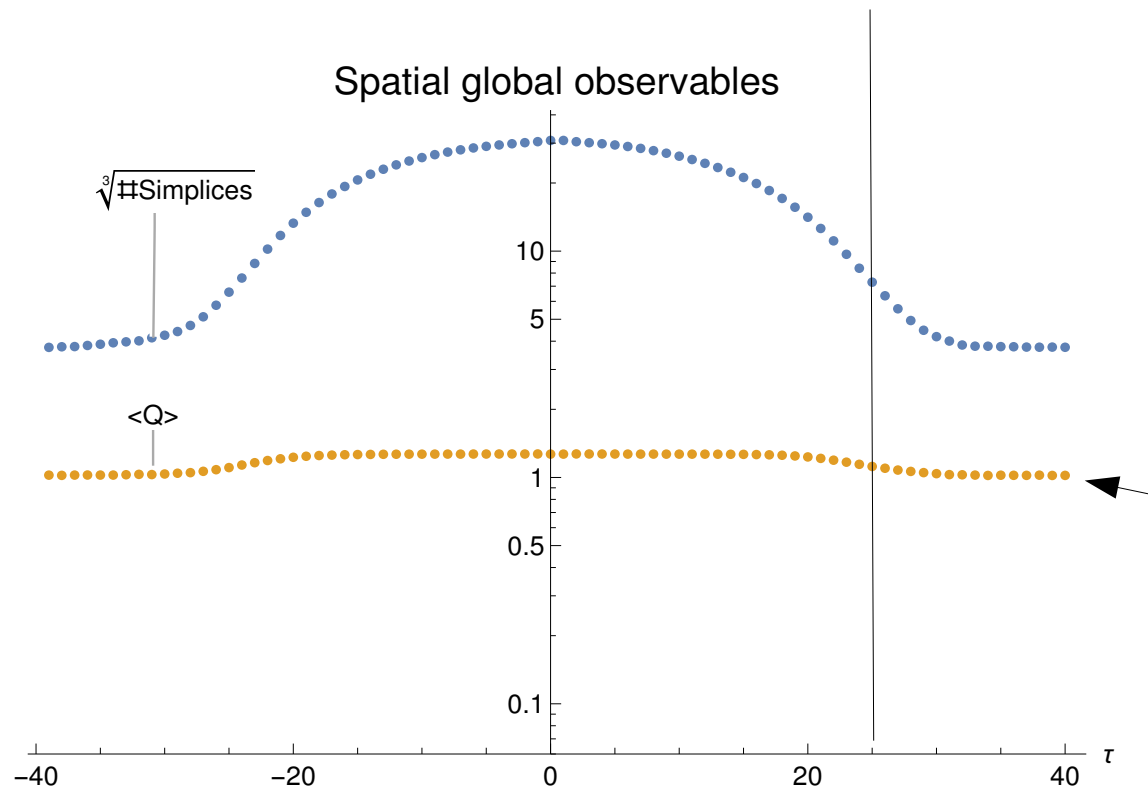


Quantum curvature Q

Basically distances of
surface points on a
sphere

Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



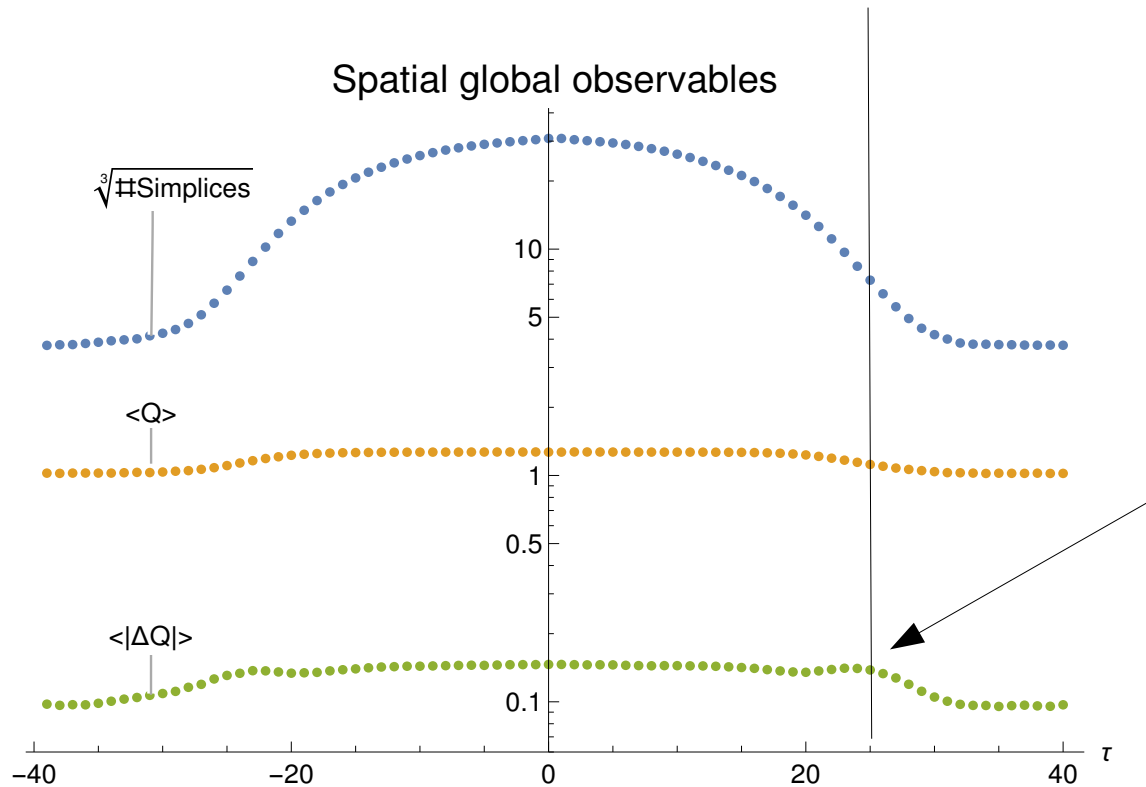
Quantum curvature Q

Basically distances of
surface points on a
sphere

Curvature scalar R can
be numerically
Estimated from Q

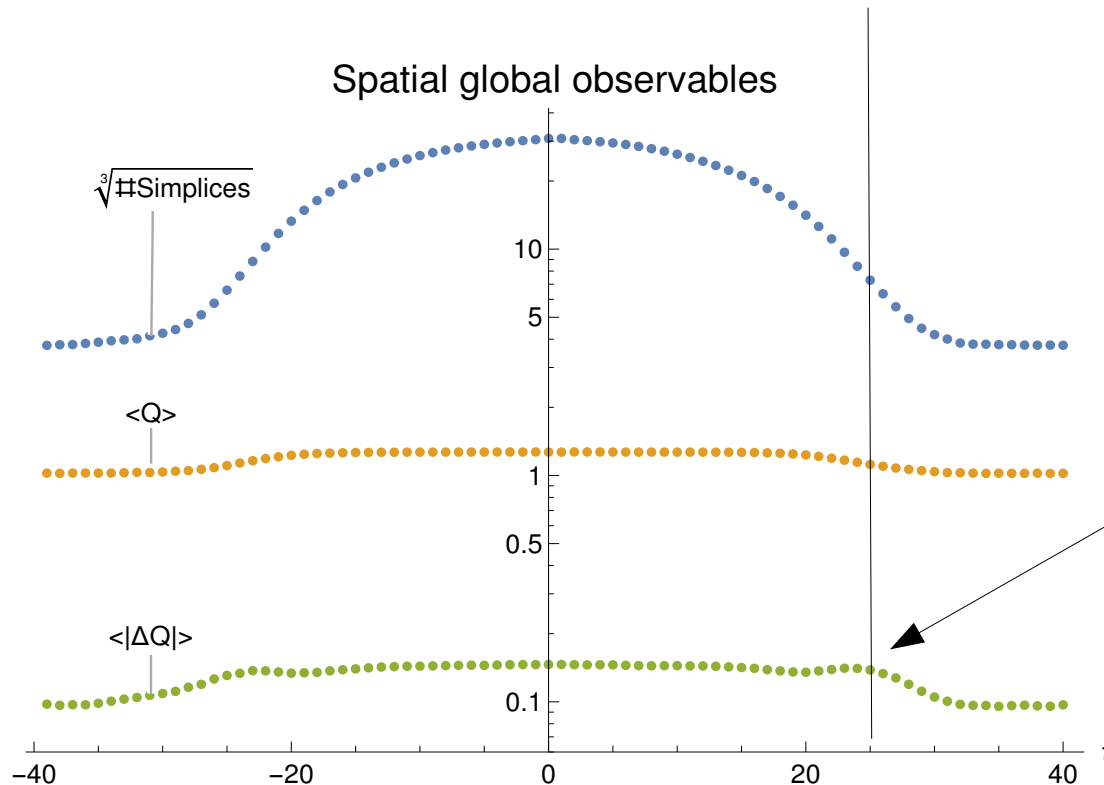
Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]

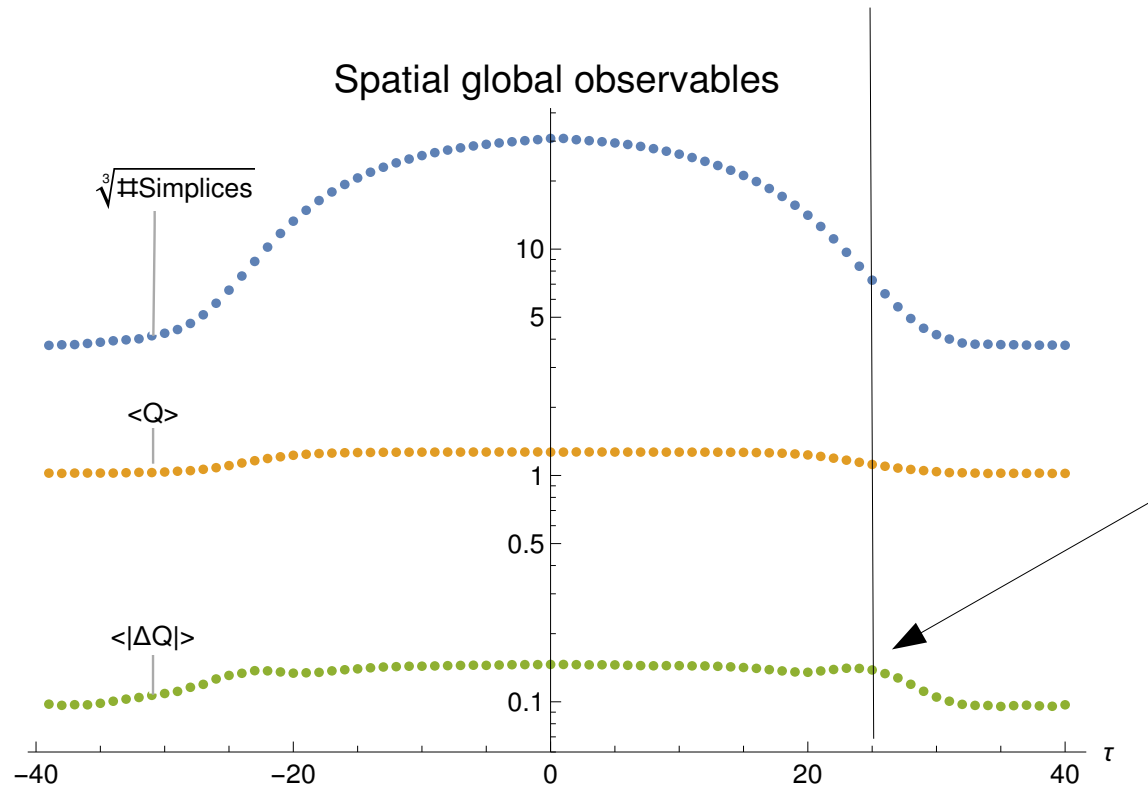


Quantum curvature
fluctuations peak
around fastest
expansion

Only cosmological
constant, no inflation!

Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]

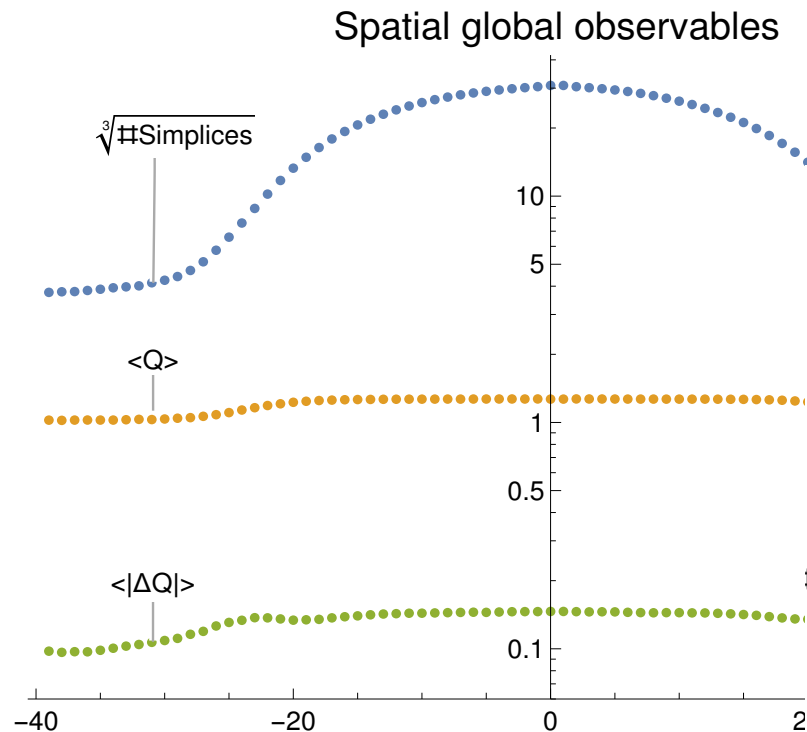


An average gauge-fixed metric
would be close to the de Sitter-like

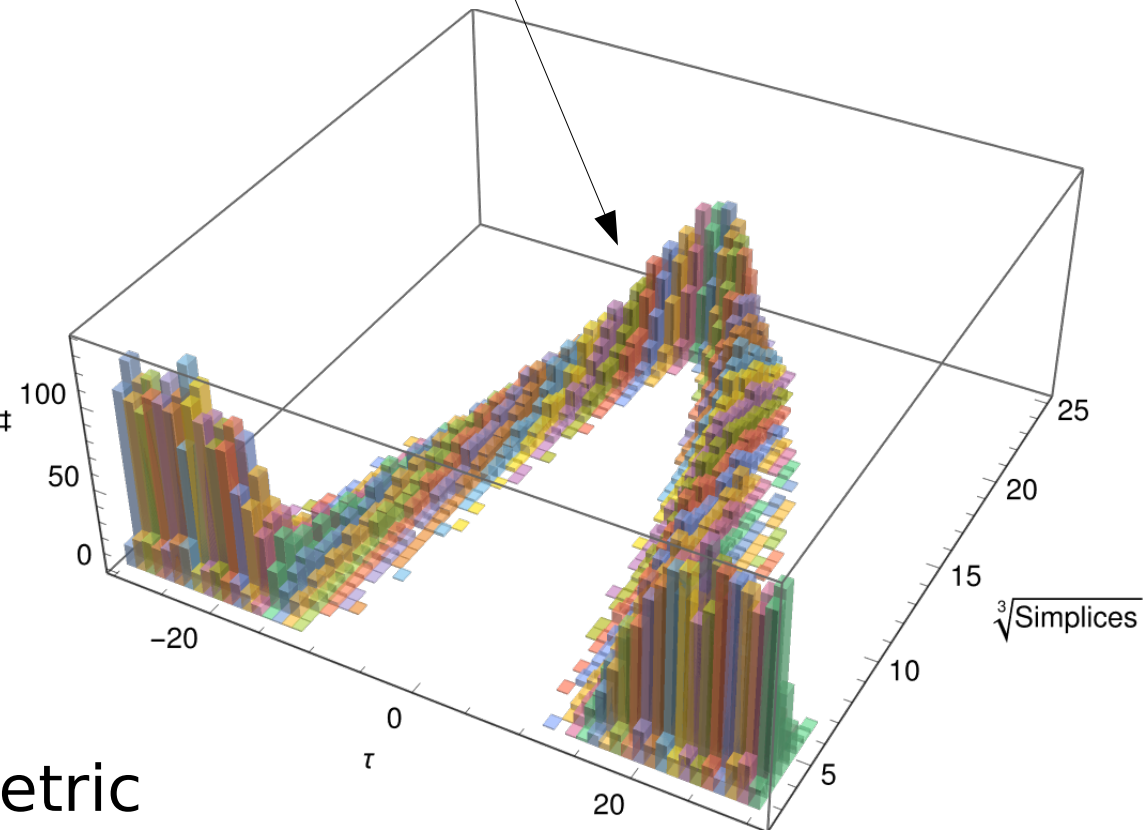
Like in classical general relativity

Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



Per configuration
histogram of the size
- relatively similar, no
large fluctuations

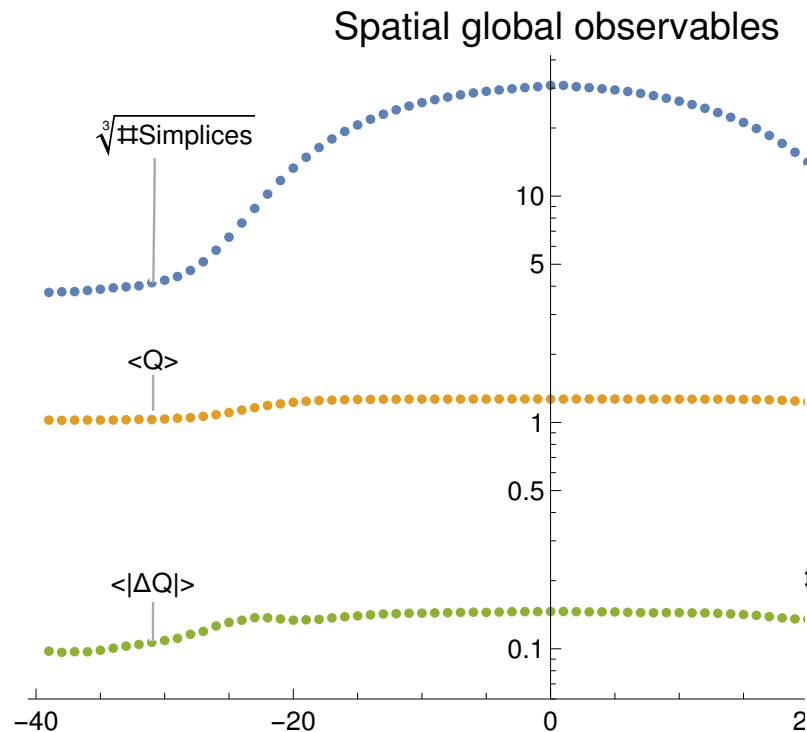


An average gauge-fixed metric
would be close to the de Sitter-like

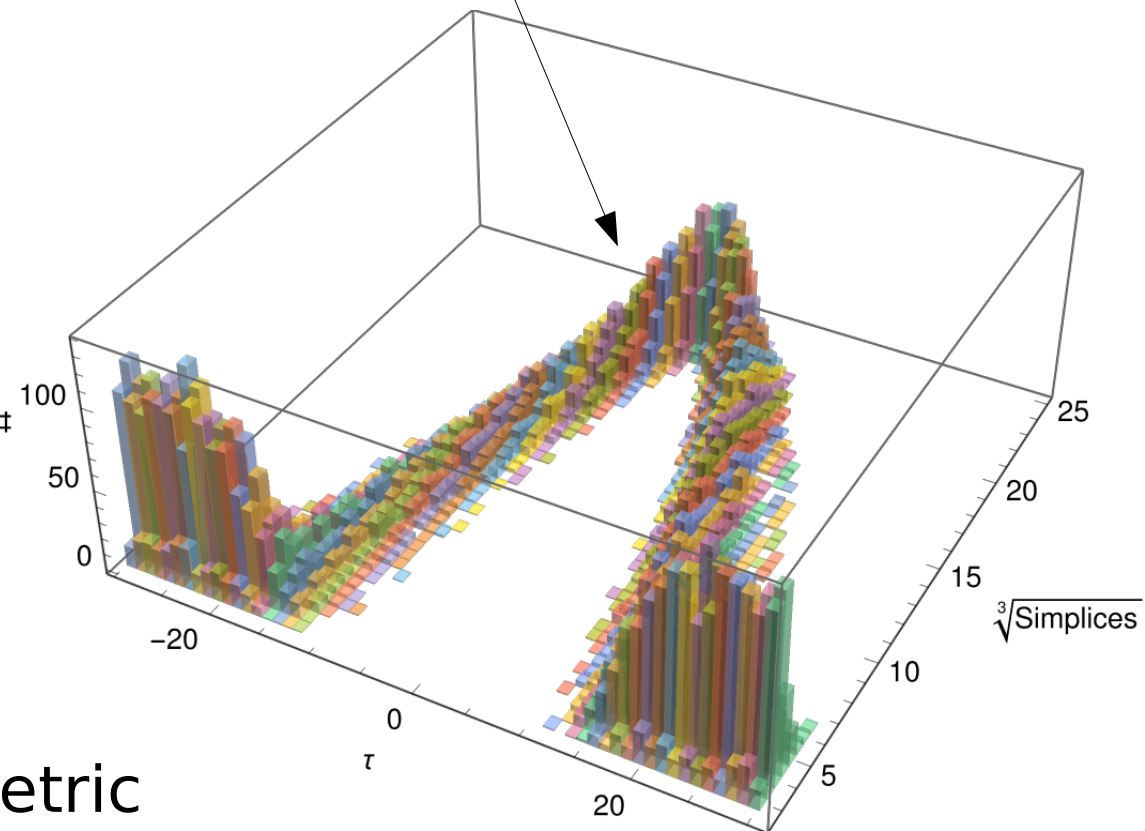
Like in classical general relativity

Space-time in CDT

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



Per configuration
histogram of the size
- relatively similar, no
large fluctuations



An average gauge-fixed metric
would be close to the de Sitter-like

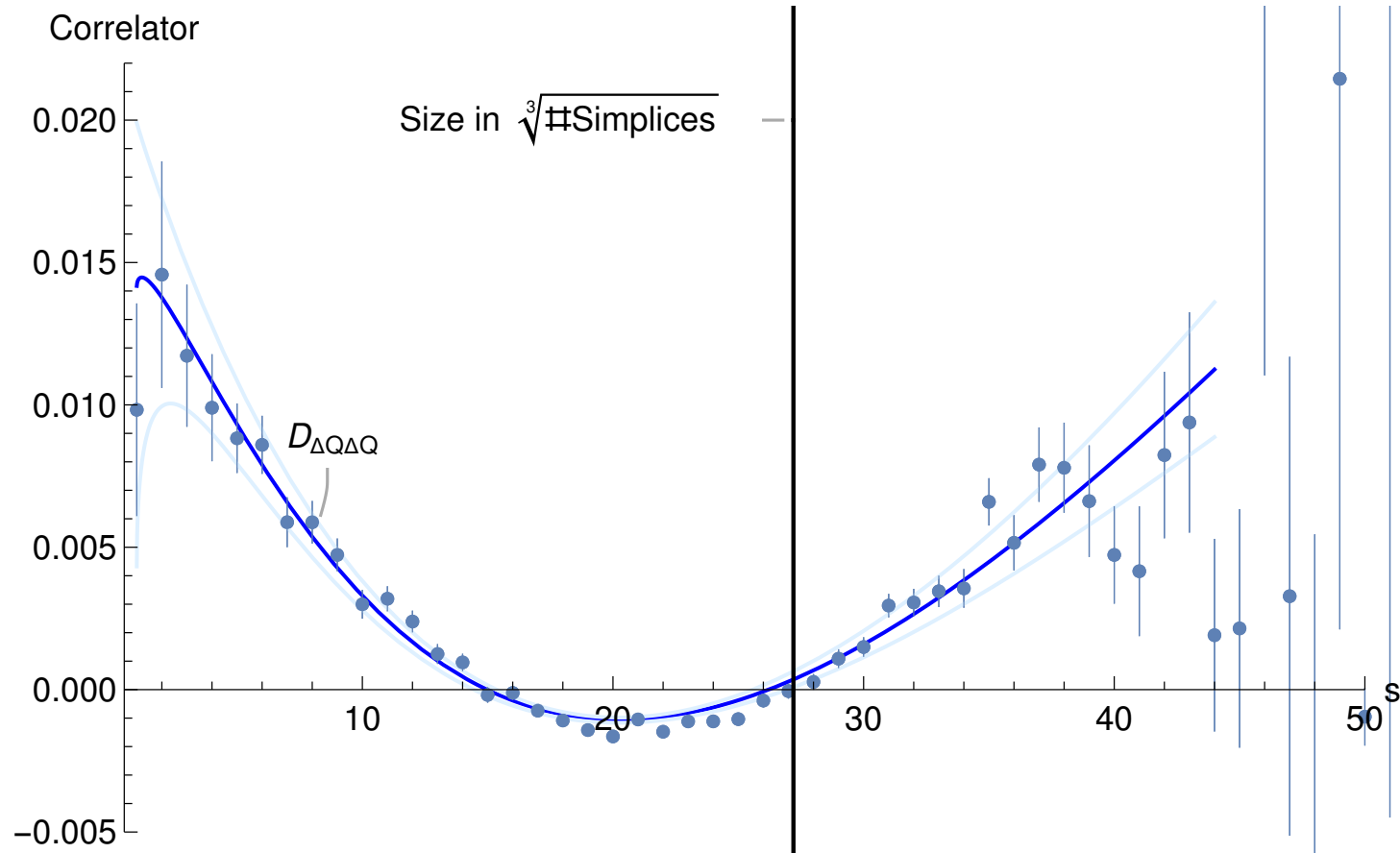
A Brout-Englert-Higgs effect in quantum gravity! [Maas'19]

- Geon: A diff-invariant composite scalar operator

- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$

- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$

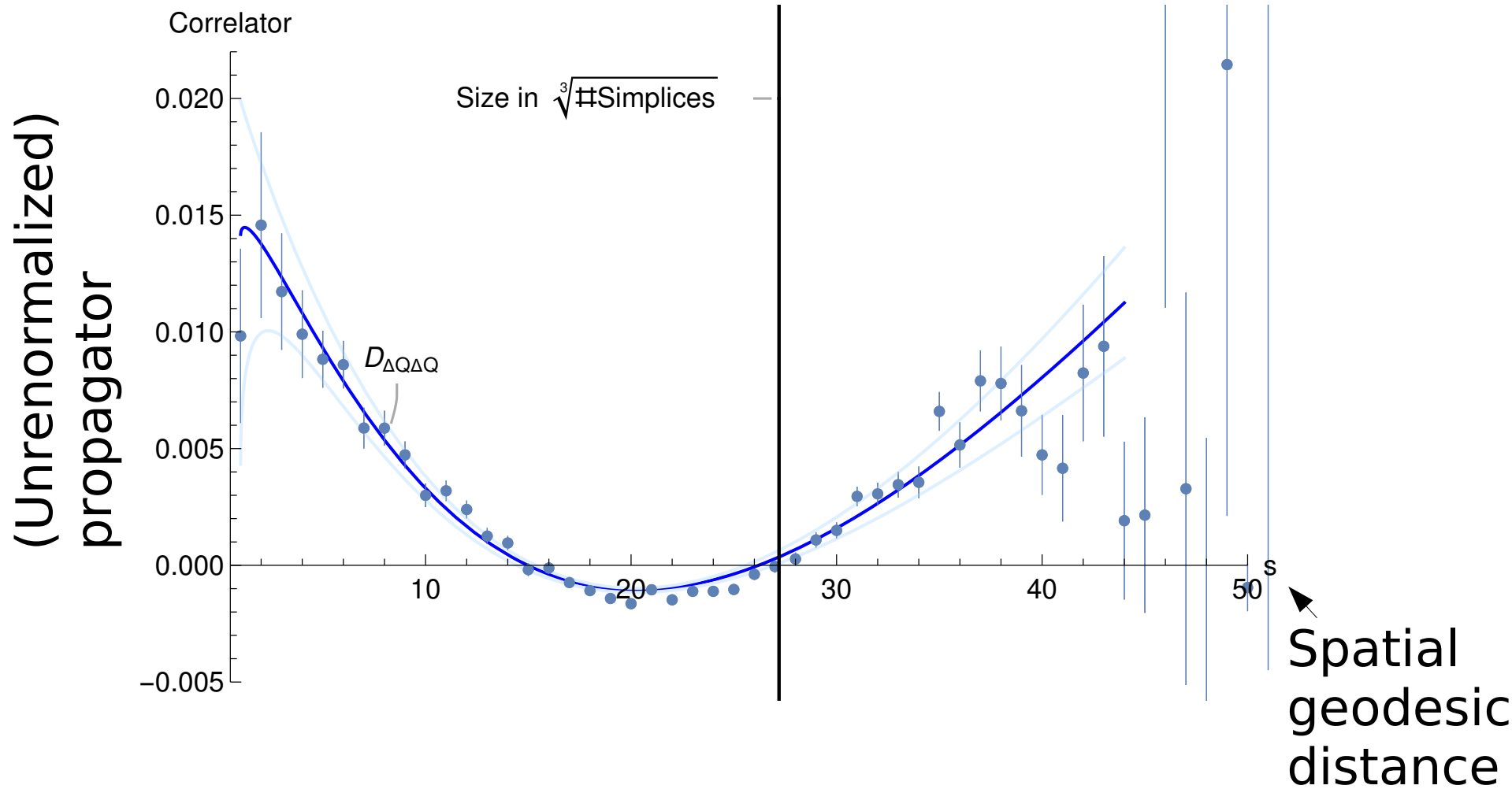
Geon correlator



- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$

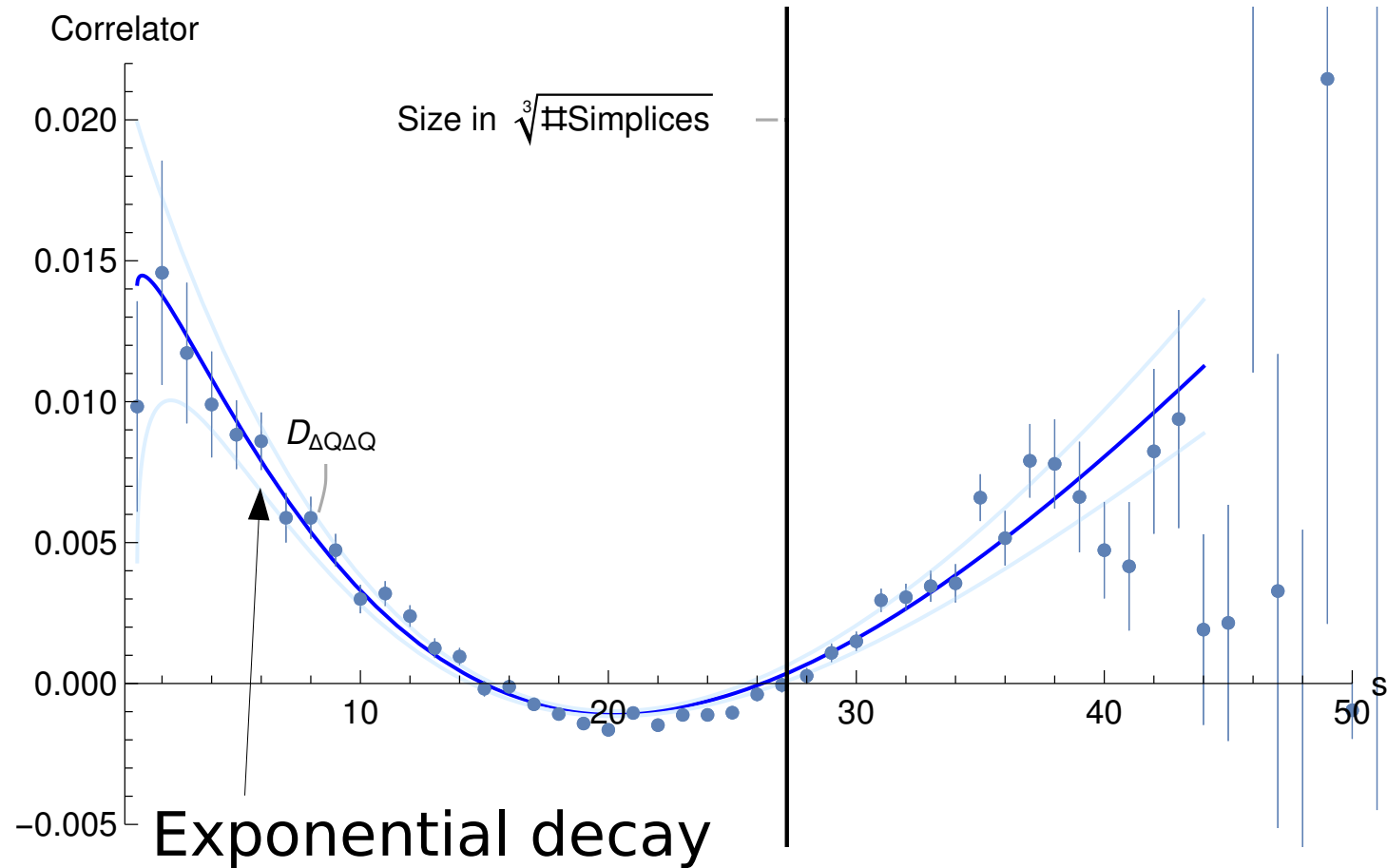
Geon correlator

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



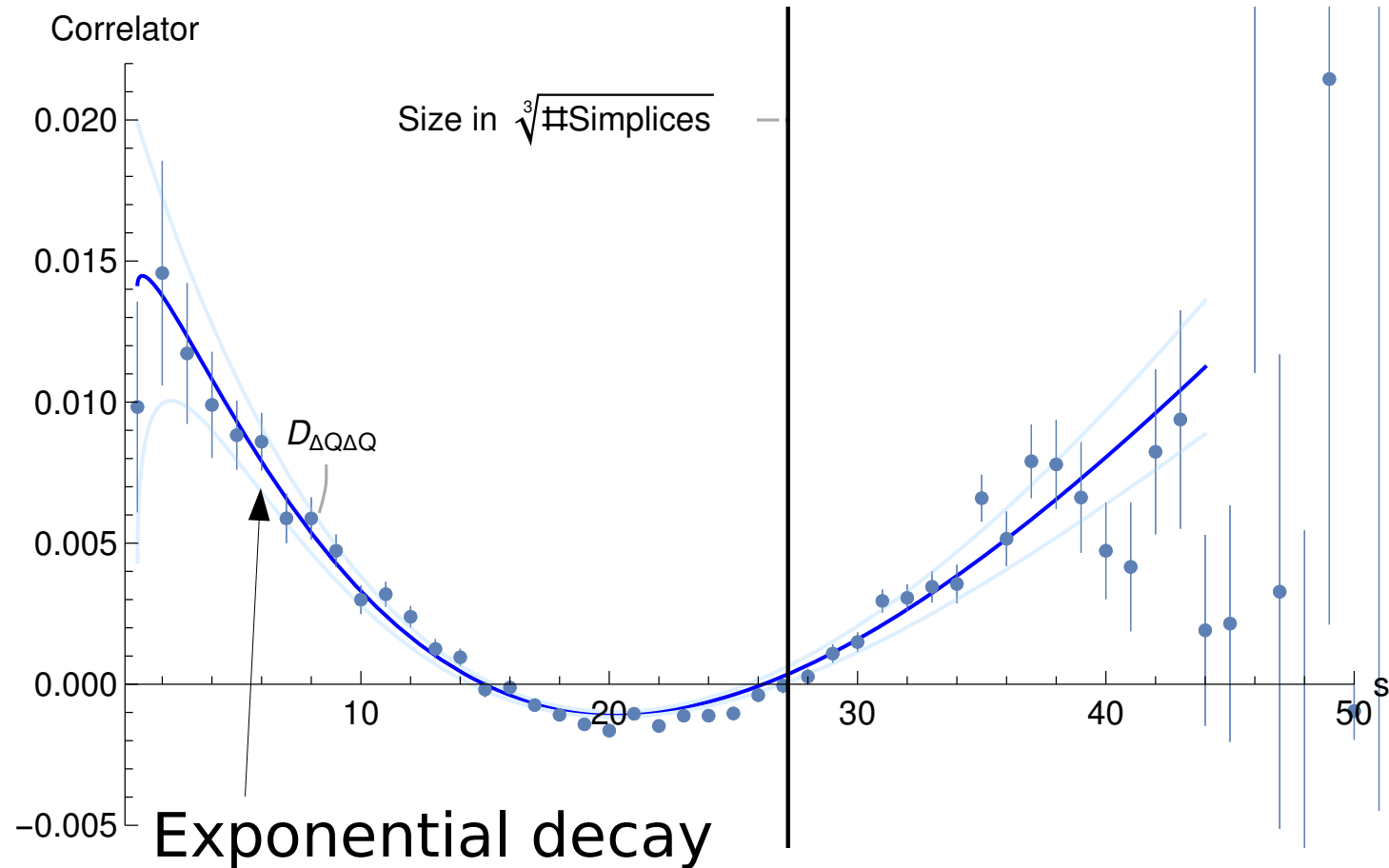
- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$

Geon correlator



- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$

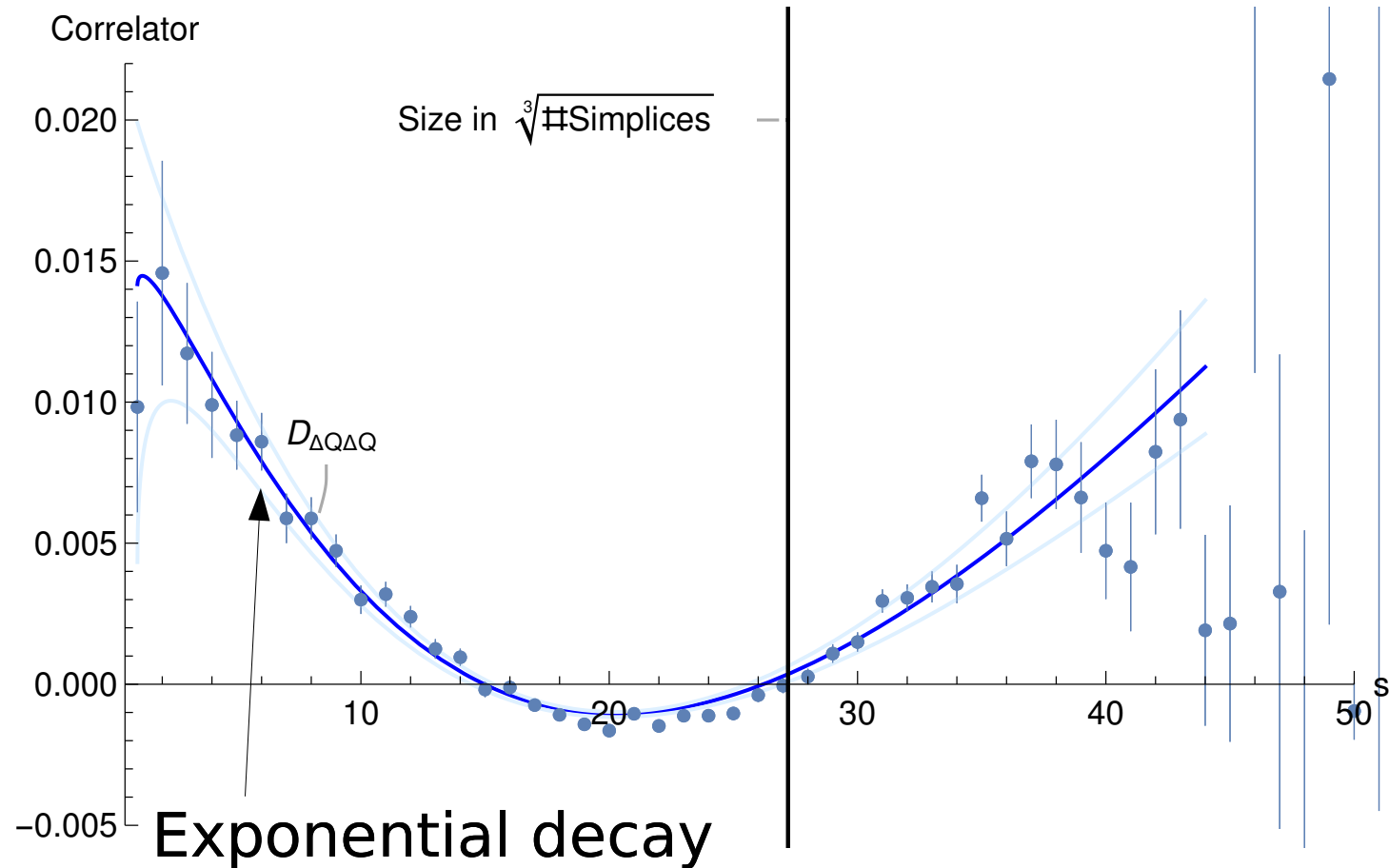
Geon correlator



- as in flat space-time for a massive particle

- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$

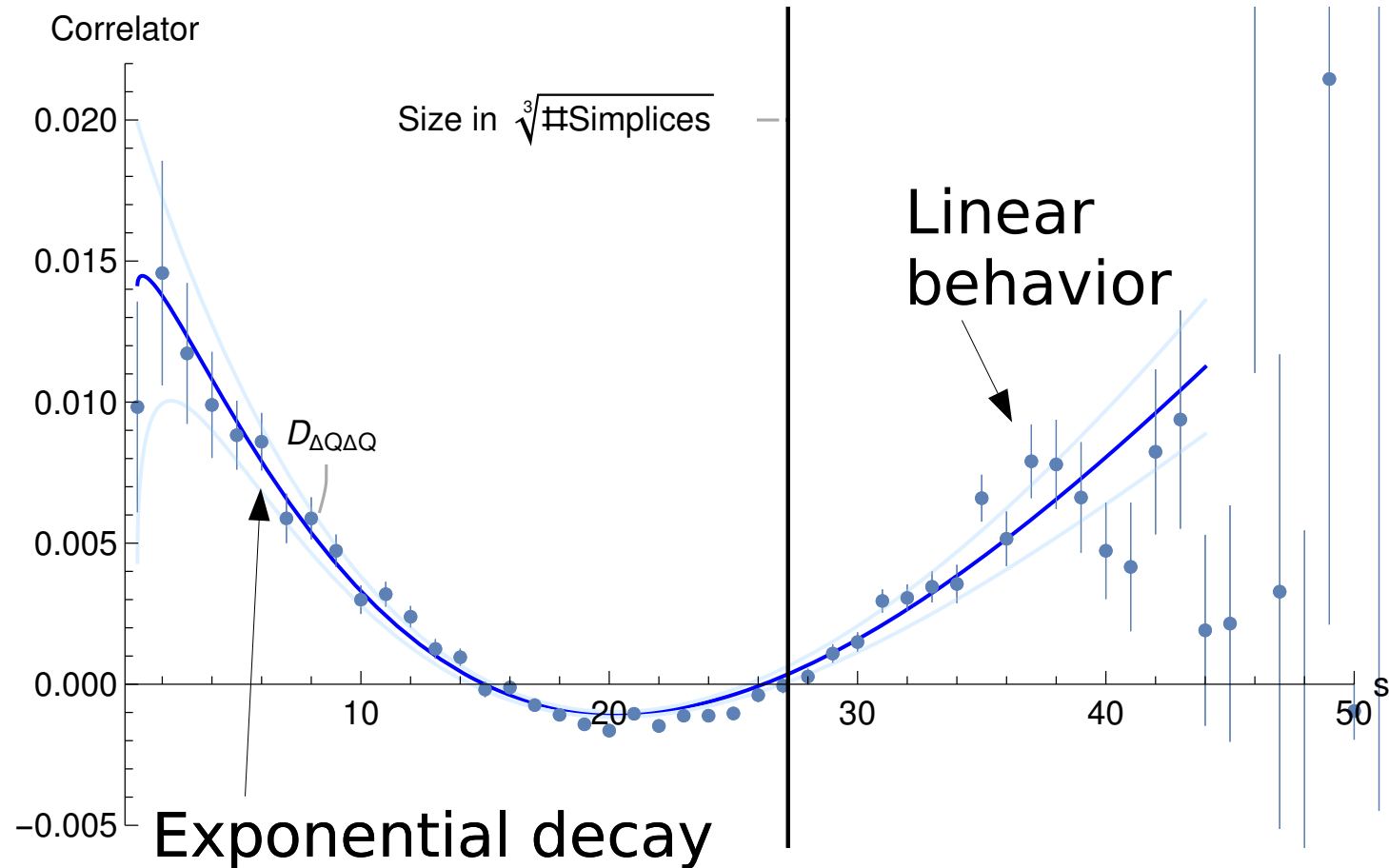
Geon correlator



- as in flat space-time for a massive particle

- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$
 - Self-bound gravitons, akin to a glueball

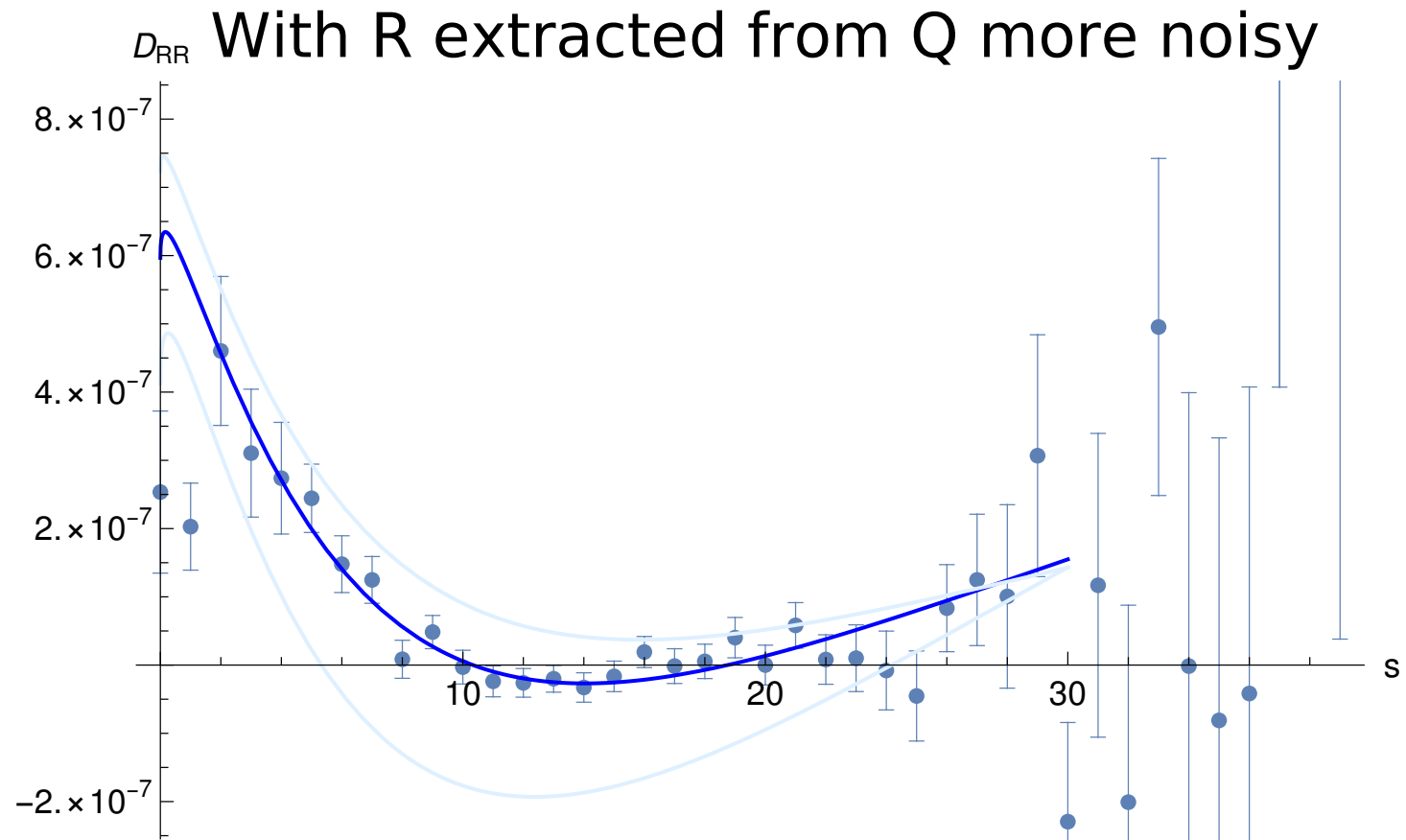
Geon correlator



- as in flat space-time for a massive particle

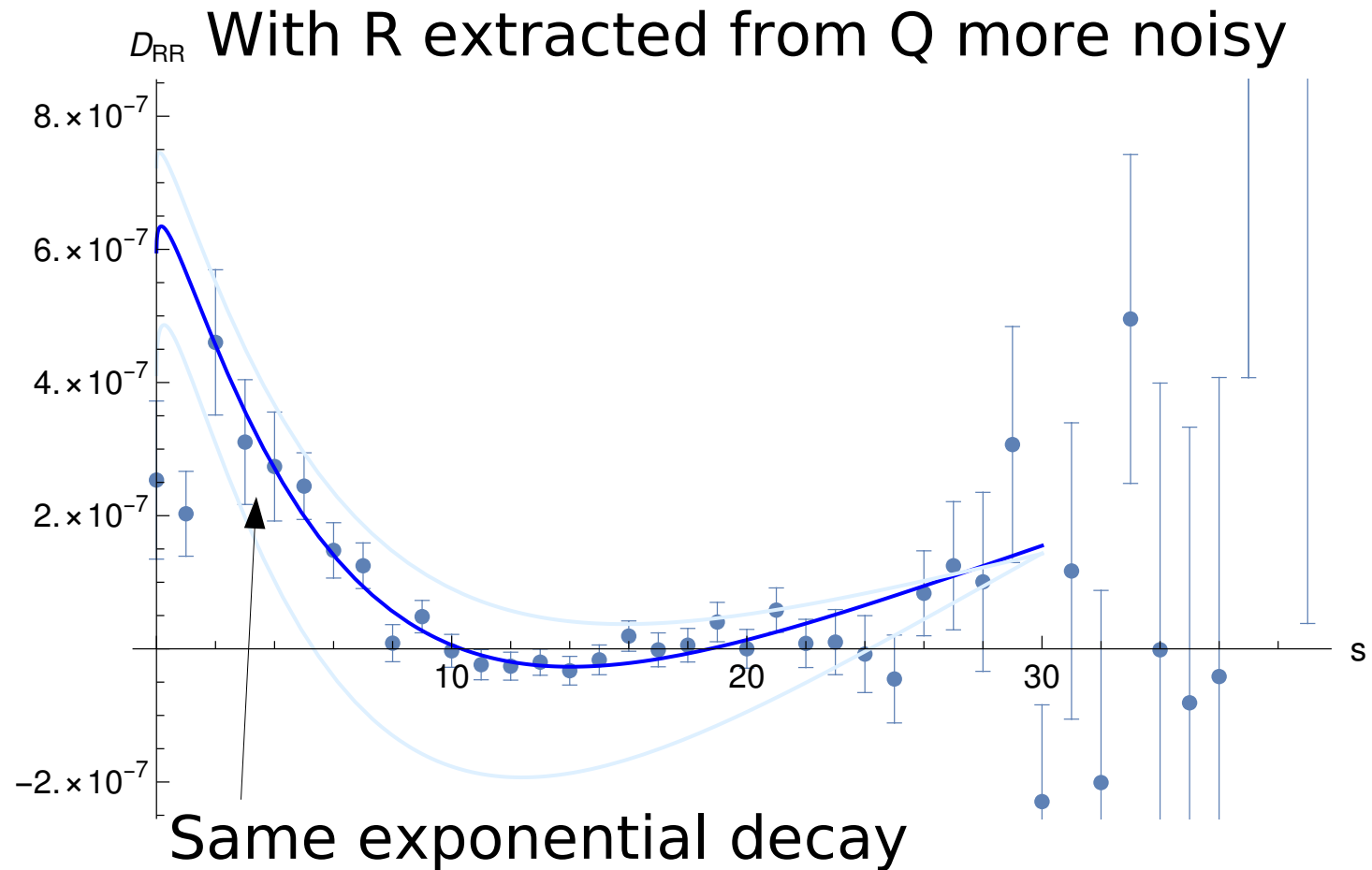
- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$
 - Self-bound gravitons, akin to a glueball

Geon correlator



- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$
 - Self-bound gravitons, akin to a glueball

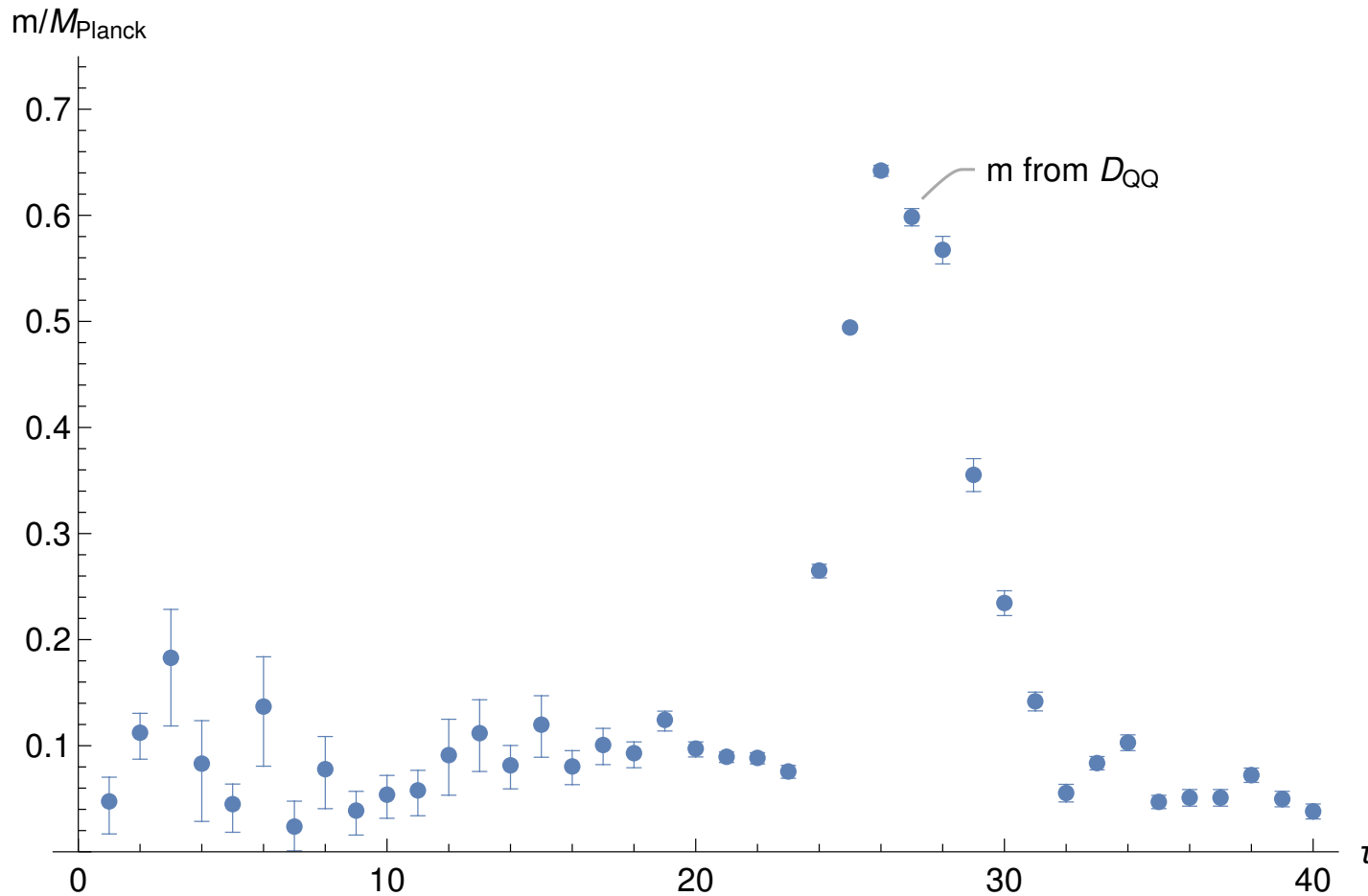
Geon correlator



- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$
 - Self-bound gravitons, akin to a glueball

Geon correlator

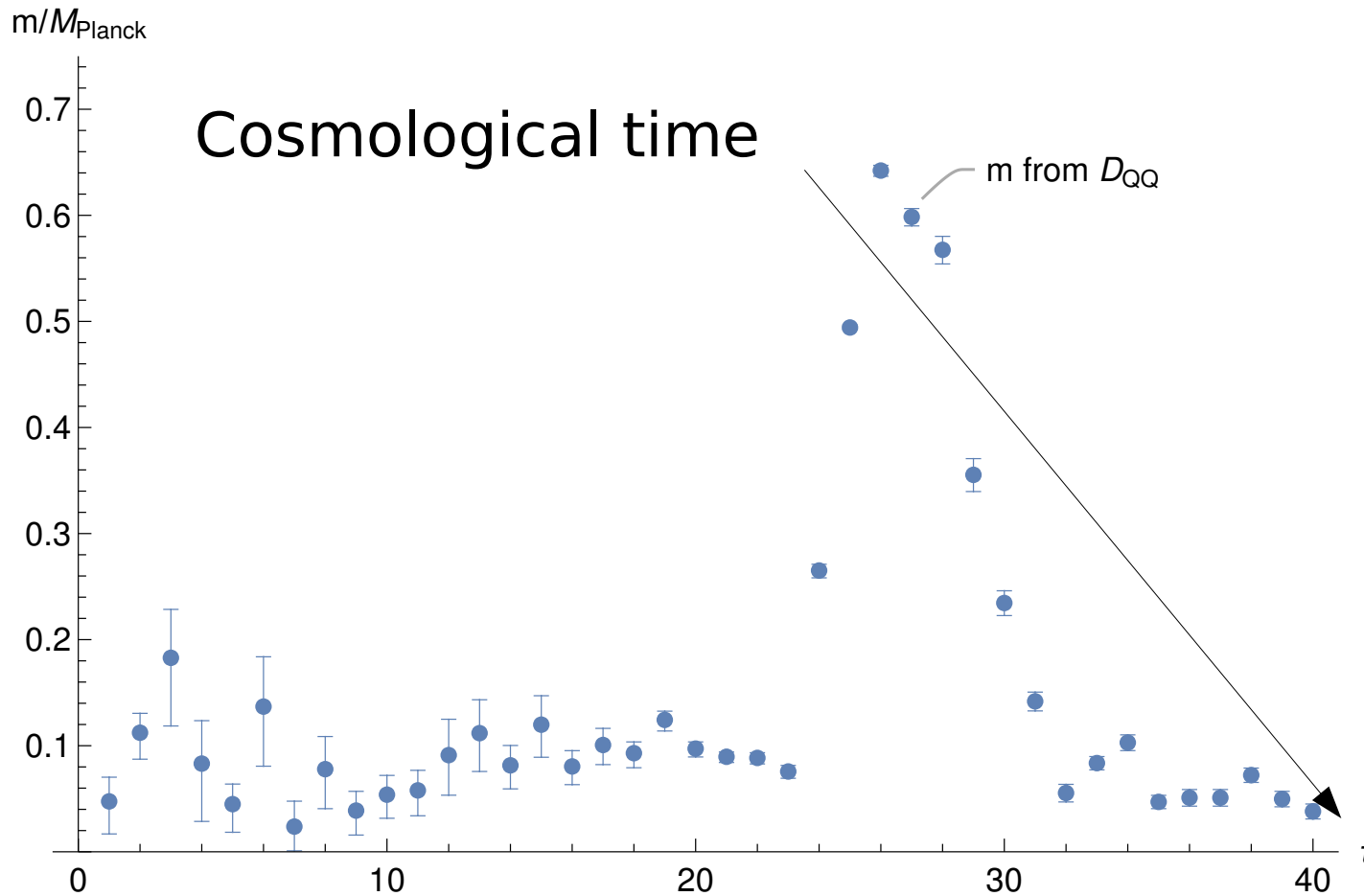
[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$
 - Self-bound gravitons, akin to a glueball

Geon correlator

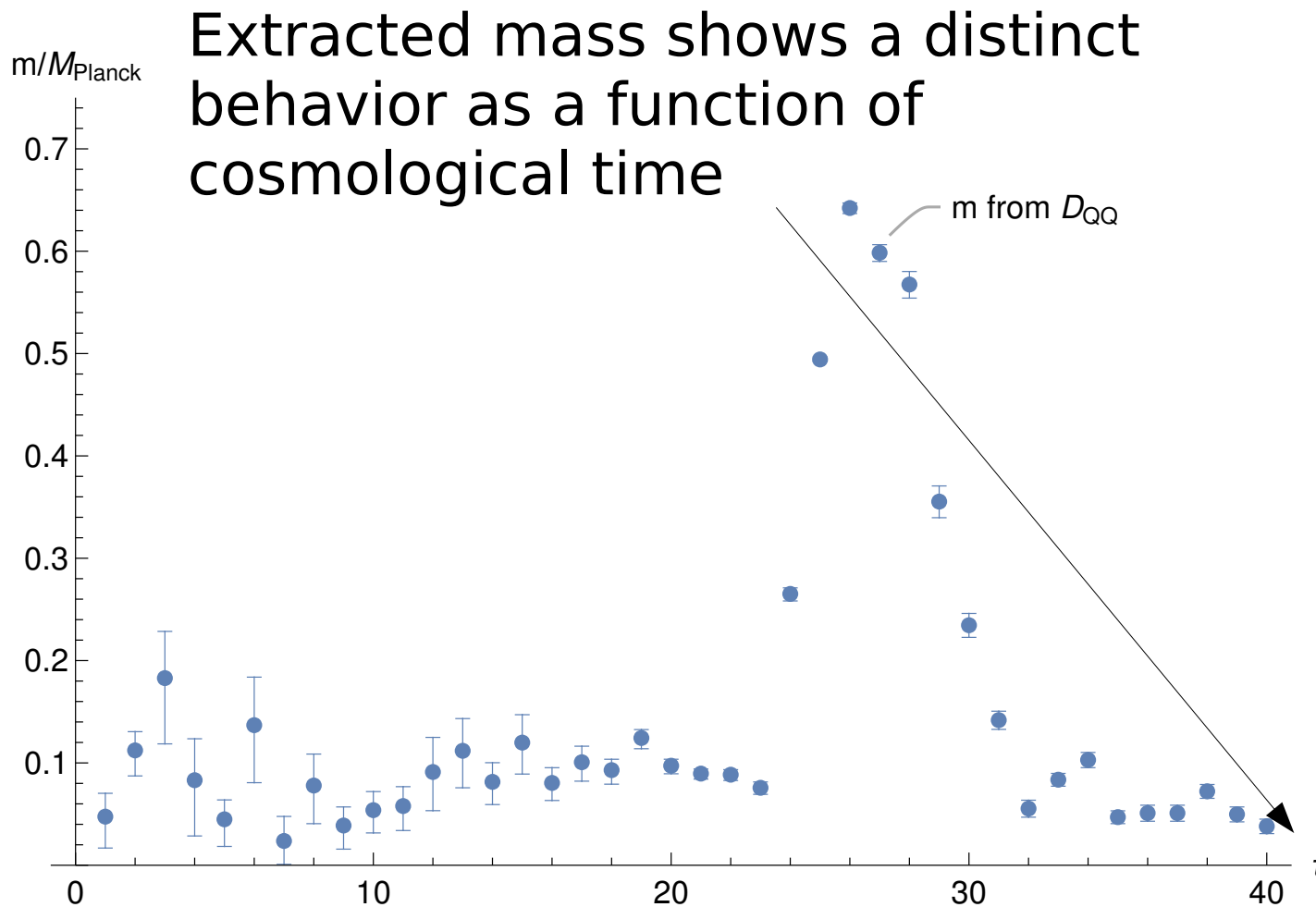
[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$
 - Self-bound gravitons, akin to a glueball

Geon correlator

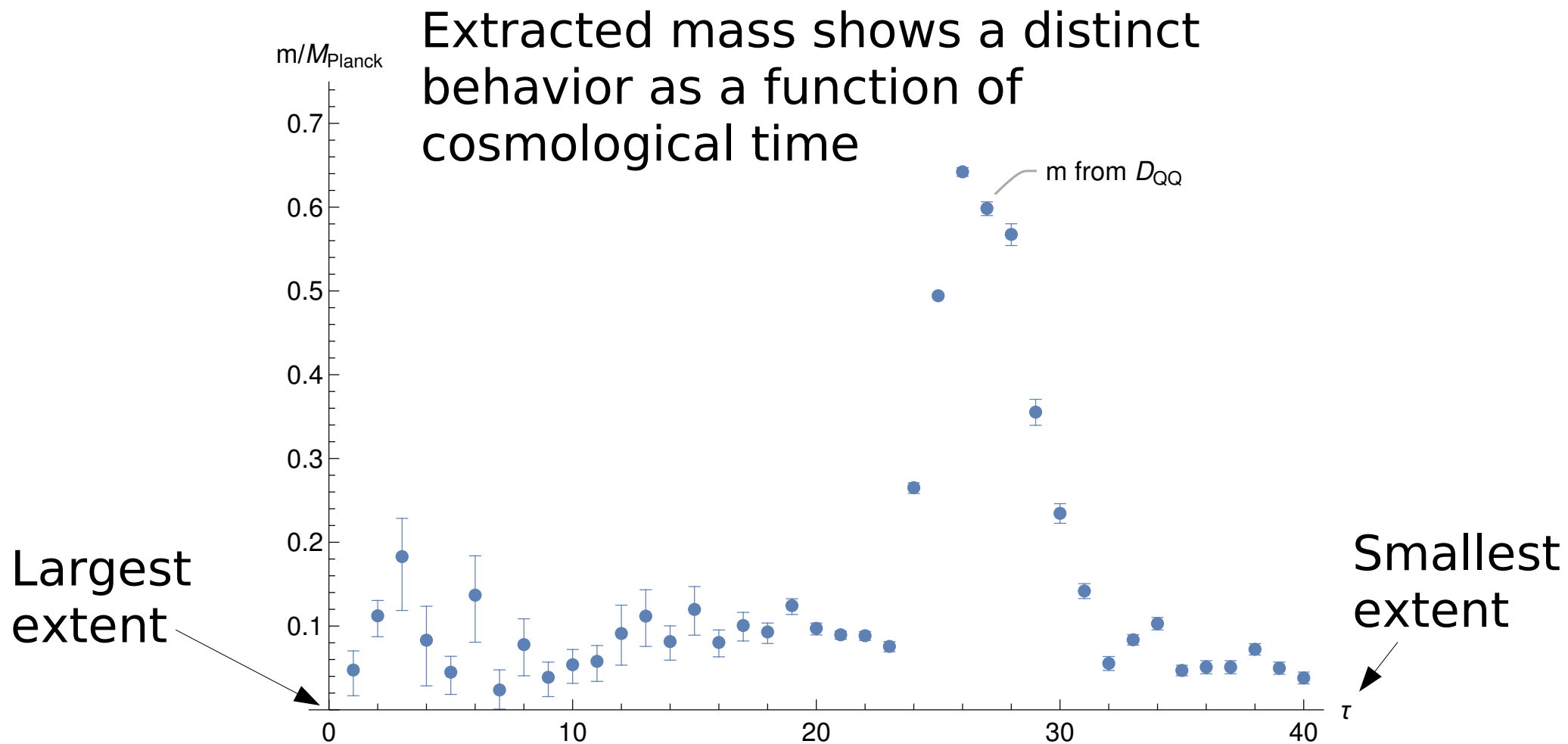
[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$
 - Self-bound gravitons, akin to a glueball

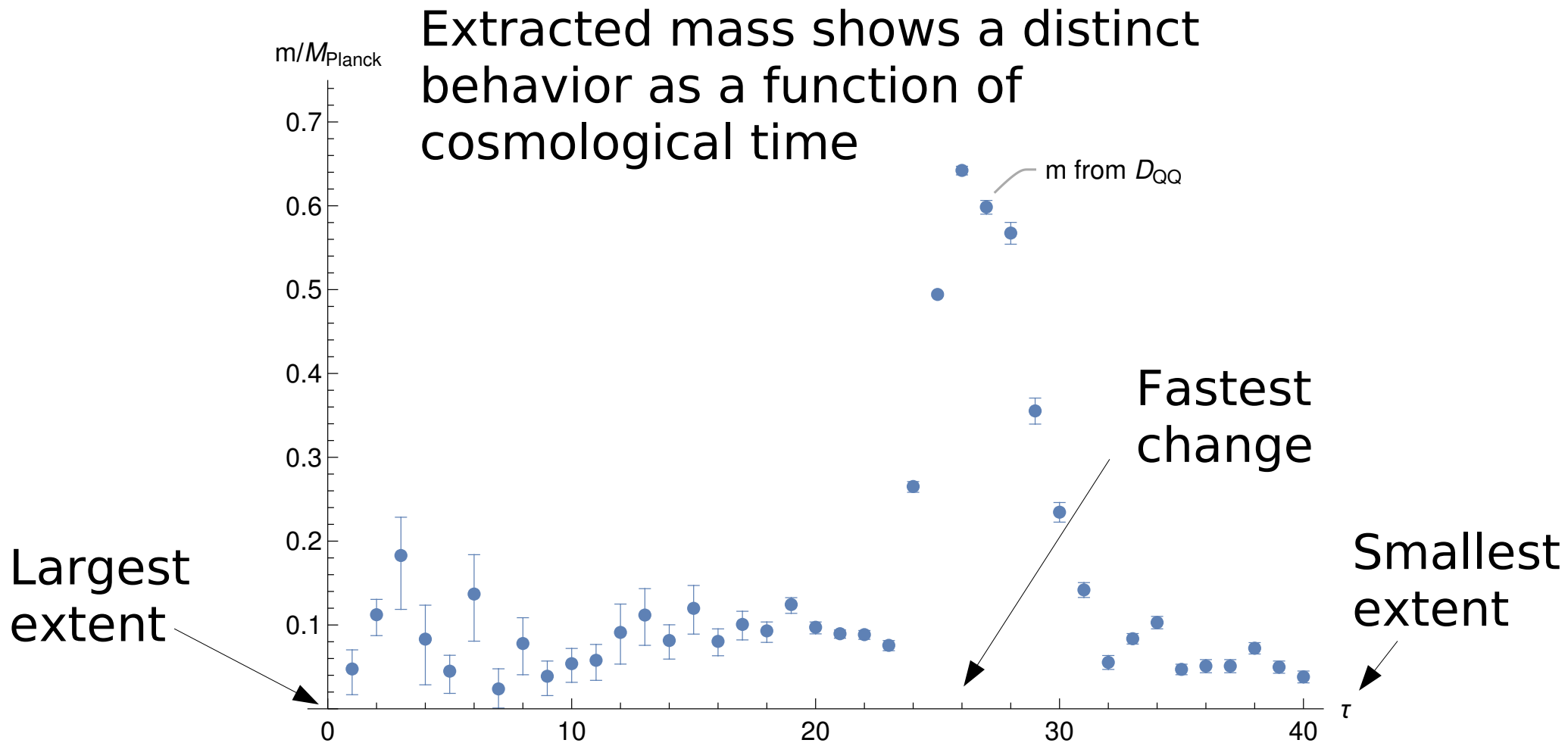
Geon correlator

[Ambjorn et al.'12,'19
Maas, Plätzer, Pressler'25]



- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$
 - Self-bound gravitons, akin to a glueball

Geon correlator



- Geon: A diff-invariant composite scalar operator
 - Continuum: Curvature fluctuations: $\langle R(x)R(y) \rangle / \langle 11 \rangle$
 - CDT: $\langle \Delta Q(x) \Delta Q(y) \rangle / \langle 11 \rangle$
 - Self-bound gravitons, akin to a glueball

Summary

- Quantum gravity (CDT) appears to exhibit a Brout-Englert-Higgs effect, explaining semiclassical behavior

Summary

- Quantum gravity (CDT) appears to exhibit a Brout-Englert-Higgs effect, explaining semiclassical behavior
- Still, only composite objects can be observables in quantum gravity

Summary

- Quantum gravity (CDT) appears to exhibit a Brout-Englert-Higgs effect, explaining semiclassical behavior
- Still, only composite objects can be observables in quantum gravity
- CDT simulations hints that these create reasonable physical degrees of freedom
 - Self-bound gravitons with a mass: Geons