Non-perturbative Brout-Englert-Higgs Physics

Axel Maas

December 2017 Odense Denmark





NAWI Graz Natural Sciences



Der Wissenschaftsfonds.

Brout-Englert-Higgs physics as a gauge theory

- Brout-Englert-Higgs physics as a gauge theory
 - What is Brout-Englert-Higgs "physics"?
 - Non-trivial implications for observables
 - What is the Higgs, the W/Z and the electron, anyway?

- Brout-Englert-Higgs physics as a gauge theory
 - What is Brout-Englert-Higgs "physics"?
 - Non-trivial implications for observables
 - What is the Higgs, the W/Z and the electron, anyway?
- Implications for BSM physics

- Brout-Englert-Higgs physics as a gauge theory
 - What is Brout-Englert-Higgs "physics"?
 - Non-trivial implications for observables
 - What is the Higgs, the W/Z and the electron, anyway?
- Implications for BSM physics
 - Qualitative possibly even a game changer

- Brout-Englert-Higgs physics as a gauge theory
 - What is Brout-Englert-Higgs "physics"?
 - Non-trivial implications for observables
 - What is the Higgs, the W/Z and the electron, anyway?
- Implications for BSM physics
 - Qualitative possibly even a game changer
- Investigations on the lattice

- Brout-Englert-Higgs physics as a gauge theory
 - What is Brout-Englert-Higgs "physics"?
 - Non-trivial implications for observables
 - What is the Higgs, the W/Z and the electron, anyway?
- Implications for BSM physics
 - Qualitative possibly even a game changer
- Investigations on the lattice
 - Group theory
 - Symmetry "breaking" on the lattice

- The Higgs sector of the standard model
 - Observables, phase structure,...

- The Higgs sector of the standard model
 - Observables, phase structure,...
- Flavor in the standard model
 - ...and other additional effects

- The Higgs sector of the standard model
 - Observables, phase structure,...
- Flavor in the standard model
 - ...and other additional effects
- Implications beyond the standard model
 - 2HDM, GUTs, technicolor

- The Higgs sector of the standard model
 - Observables, phase structure,...
- Flavor in the standard model
 - ...and other additional effects
- Implications beyond the standard model
 - 2HDM, GUTs, technicolor

- Review article upcoming
- O'Raifeartaigh: Group structure of gauge theories

- Strong interactions are non-perturbative
 - Like QCD

- Strong interactions are non-perturbative
 - Like QCD
 - But not always: Asymptotic freedom

- Strong interactions are non-perturbative
 - Like QCD
 - But not always: Asymptotic freedom
- Weak interactions can be non-perturbative

- Strong interactions are non-perturbative
 - Like QCD
 - But not always: Asymptotic freedom
- Weak interactions can be non-perturbative
 - QED is weakly interacting, but has nonperturbative features like atoms, molecules, matter with phase structure,...
 - Bound states, phase transitions,...

- Strong interactions are non-perturbative
 - Like QCD
 - But not always: Asymptotic freedom
- Weak interactions can be non-perturbative
 - QED is weakly interacting, but has nonperturbative features like atoms, molecules, matter with phase structure,...
 - Bound states, phase transitions,...
- Are there (relevant) non-perturbative effects in the weak interactions and the Higgs?

- Strong interactions are non-perturbative
 - Like QCD
 - But not always: Asymptotic freedom
- Weak interactions can be non-perturbative
 - QED is weakly interacting, but has nonperturbative features like atoms, molecules, matter with phase structure,...
 - Bound states, phase transitions,...
- Are there (relevant) non-perturbative effects in the weak interactions and the Higgs?
 - How could there be not?

A precursor: Global symmetry breaking revisited

- Consider some scalar theory with Lagrangian invariant under some group
 - No anomalies

- Consider some scalar theory with Lagrangian invariant under some group
 - No anomalies
- Path integral formulation

$$\langle O \rangle = \int D \varphi O(\varphi) e^{iS(\varphi)}$$

- Consider some scalar theory with Lagrangian invariant under some group
 - No anomalies
- Path integral formulation

$$\langle O \rangle = \int D \varphi O(\varphi) e^{iS(\varphi)}$$

Integrates over all (!) field configurations

- Consider some scalar theory with Lagrangian invariant under some group
 - No anomalies
- Path integral formulation

$$\langle O \rangle = \int D \varphi O(\varphi) e^{iS(\varphi)}$$

- Integrates over all (!) field configurations
- Non-symmetric expectation values need to vanish

$$\langle \varphi \rangle = \sum_{\varphi} (\varphi + g \varphi + g' \varphi + ...) e^{iS} = \sum_{\varphi} e^{iS} \sum_{g} \varphi = \langle G \varphi \rangle = 0$$

- Consider some scalar theory with Lagrangian invariant under some group
 - No anomalies
- Path integral formulation

$$\langle O \rangle = \int D \varphi O(\varphi) e^{iS(\varphi)}$$

- Integrates over all (!) field configurations
- Non-symmetric expectation values need to vanish

$$\langle \varphi \rangle = \sum_{\varphi} (\varphi + g \varphi + g' \varphi + ...) e^{iS} = \sum_{\varphi} e^{iS} \sum_{g} \varphi = \langle G \varphi \rangle = 0$$

- Does not depend on the parameters: Always
- Symmetry is always manifest

Introduce a source

$$\langle O \rangle = \int D \varphi O(\varphi) e^{iS(\varphi) + j\varphi}$$

Introduce a source

$$\langle O \rangle = \int D \varphi O(\varphi) e^{iS(\varphi) + j\varphi}$$

Preferred direction – result biased

$$\langle \varphi \rangle (j) = \sum_{\varphi} (\varphi e^{ij\varphi} + g\varphi e^{ijg\varphi} + ...) e^{iS} = \sum_{\varphi} e^{iS} \sum_{g} \varphi e^{ij\varphi} \neq 0$$

Introduce a source

$$\langle O \rangle = \int D \varphi O(\varphi) e^{iS(\varphi) + j\varphi}$$

Preferred direction – result biased

$$\langle \varphi \rangle (j) = \sum_{\varphi} (\varphi e^{ij\varphi} + g\varphi e^{ijg\varphi} + ...) e^{iS} = \sum_{\varphi} e^{iS} \sum_{g} \varphi e^{ij\varphi} \neq 0$$

• Take limit of vanishing source. If

 $\lim_{j \to 0} \langle \varphi \rangle(j) = v \neq \langle \varphi \rangle(0)$ the theory experiences SSB

t

Introduce a source

$$\langle O \rangle = \int D \varphi O(\varphi) e^{iS(\varphi) + j\varphi}$$

Preferred direction – result biased

$$\langle \varphi \rangle (j) = \sum_{\varphi} (\varphi e^{ij\varphi} + g\varphi e^{ijg\varphi} + ...) e^{iS} = \sum_{\varphi} e^{iS} \sum_{g} \varphi e^{ij\varphi} \neq 0$$

• Take limit of vanishing source. If

$$\lim_{j \to 0} \langle \varphi \rangle(j) = v \neq \langle \varphi \rangle(0)$$

he theory experiences SSB

• Goldstone's theorem only at $j \rightarrow 0$, not at 0

Source is external to the theory

- Source is external to the theory
- What is external in particle physics?

- Source is external to the theory
- What is external in particle physics?
 - Theory is a low-energy effective theory
 - Source encodes effect from the UV theory
 - Then the symmetry is not a symmetry of the UV theory
 - E.g. QCD and chiral symmetry

- Source is external to the theory
- What is external in particle physics?
 - Theory is a low-energy effective theory
 - Source encodes effect from the UV theory
 - Then the symmetry is not a symmetry of the UV theory
 - E.g. QCD and chiral symmetry
 - Theory is part of a system
 - Source encodes influence of the remainder
 - E.g. extrauniverse influences

- Source is external to the theory
- What is external in particle physics?
 - Theory is a low-energy effective theory
 - Source encodes effect from the UV theory
 - Then the symmetry is not a symmetry of the UV theory
 - E.g. QCD and chiral symmetry
 - Theory is part of a system
 - Source encodes influence of the remainder
 - E.g. extrauniverse influences
- No other possibilities

Metastability

- Theory without external influence fully symmetric
 - No Goldstone bosons, full degeneracy pattern

Metastability

- Theory without external influence fully symmetric
 - No Goldstone bosons, full degeneracy pattern
- Still: There is a difference between SSB occurs with external influence or not

Metastability

- Theory without external influence fully symmetric
 - No Goldstone bosons, full degeneracy pattern
- Still: There is a difference between SSB occurs with external influence or not
 - Metastability
Metastability

- Theory without external influence fully symmetric
 - No Goldstone bosons, full degeneracy pattern
- Still: There is a difference between SSB occurs with external influence or not
 - Metastability
 - Detectable by $\langle \left(\int \phi\right)^2 \rangle \neq 0$
 - Measures relative orientation, but invariant

Metastability

- Theory without external influence fully symmetric
 - No Goldstone bosons, full degeneracy pattern
- Still: There is a difference between SSB occurs with external influence or not
 - Metastability
 - Detectable by $\langle \left(\int \phi\right)^2 \rangle \neq 0$
 - Measures relative orientation, but invariant
 - Note: $\left< \phi^2 \right>$ does not work

Intact

- Intact
- Metastable

- Intact
- Metastable
- Spontaenously broken

- Intact
- Metastable
- Spontaenously broken
- Explicitly broken by Lagrangian or anomaly

- Intact
- Metastable
- Spontaenously broken
- Explicitly broken by Lagrangian or anomaly
- A split like $\varphi \rightarrow \psi + v$ does not break a symmetry
 - Only hides it
 - Symmetry is realized in a non-trivial way

The Higgs sector of the standard model

Consider the Higgs sector of the standard model

- Consider the Higgs sector of the standard model
- . The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$

• Ws W^a_{μ} W

• Coupling g and some numbers f^{abc}

- Consider the Higgs sector of the standard model
- . The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + gf^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Ws W^a_{μ} W
- Higgs h_i
- Coupling g and some numbers f^{abc} and t_a^{ij}

- Consider the Higgs sector of the standard model
- The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Ws W^a_{μ} W
- Higgs h_i (h)

• Couplings g, v, λ and some numbers f^{abc} and t_{j}^{ij}

- Consider the Higgs sector of the standard model
- . The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Ws W^a_{μ} W
- Higgs h_i
- No QED: Ws and Zs are degenerate
- Couplings g, v, λ and some numbers f^{abc} and t_a^{ij}

- Consider the Higgs sector of the standard model
- . The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

- Consider the Higgs sector of the standard model
- . The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

• Local SU(2) gauge symmetry $W^{a}_{\mu} \rightarrow W^{a}_{\mu} + (\delta^{a}_{b}\partial_{\mu} - gf^{a}_{bc}W^{c}_{\mu})\phi^{b}$ $h_{i} \rightarrow h_{i} + gt^{ij}_{a}\phi^{a}h_{j}$

- Consider the Higgs sector of the standard model
- . The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - \nu^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

• Local SU(2) gauge symmetry $W^a_{\mu} \rightarrow W^a_{\mu} + (\delta^a_b \partial_{\mu} - g f^a_{bc} W^c_{\mu}) \phi^b$

$$h_i \rightarrow h_i + g t_a^{ij} \phi^a h_j$$

• Global SU(2) Higgs custodial (flavor) symmetry

• Acts as right-transformation on the Higgs field only $W^a_\mu
ightarrow W^a_\mu
ightarrow M^a_\mu
ightarrow h_i
ightarrow h_i + a^{ij} h_j + b^{ij} h_j^*$

- Consider the Higgs sector of the standard model
- . The Higgs sector is a gauge theory

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j})^{+} D^{\mu}_{ik} h_{k} + \lambda (h^{a} h_{a}^{+} - v^{2})^{2}$$
$$W^{a}_{\mu\nu} = \partial_{\mu} W^{a}_{\nu} - \partial_{\nu} W^{a}_{\mu} + g f^{a}_{bc} W^{b}_{\mu} W^{c}_{\nu}$$
$$D^{ij}_{\mu} = \delta^{ij} \partial_{\mu} - ig W^{a}_{\mu} t^{ij}_{a}$$

Local SU(2) gauge symmetry

$$W^{a}_{\mu} \rightarrow W^{a}_{\mu} + (\delta^{a}_{b}\partial_{\mu} - gf^{a}_{bc}W^{c}_{\mu})\phi^{b} \qquad h_{i} \rightarrow h_{i} + gt^{ij}_{a}\phi^{a}h_{j}$$

• Global SU(2) Higgs custodial (flavor) symmetry

$$X = \begin{vmatrix} h_1 & -h_2^+ \\ h_2 & h_1^+ \end{vmatrix} \Rightarrow X' = G_{gauge} X G_{custodial}$$

$$L = -\frac{1}{4} W^{a}_{\mu\nu} W^{\mu\nu}_{a} + (D^{ij}_{\mu} h^{j}) + D^{\mu}_{ik} h_{k} + \lambda (h^{a} h^{+}_{a} - v^{2})^{2}$$

[Bohm et al. 2001]

$$L = \lambda (h^a h_a^+ - v^2)^2$$

$$L = \lambda (h^a h_a^+ - v^2)^2$$



$$L = \lambda (h^a h_a^+ - v^2)^2$$



$$L = \lambda (h^a h_a^+ - v^2)^2$$



$$L = \lambda (h^a h_a^+ - v^2)^2$$



$$L = \lambda (h^a h_a^+ - v^2)^2$$



$$L = \lambda (h^a h_a^+ - v^2)^2$$



- Classical analysis of the Higgs sector
- Non-zero condensate shifts Higgs mass to an ordinary mass

$$L = \lambda (h^a h_a^+ - v^2)^2$$



- Classical analysis of the Higgs sector
- Non-zero condensate shifts Higgs mass to an ordinary mass
- Perform perturbative expansion around the classical vacuum

- Minimize action classically
 - Yields $hh^+ = v^2$ Higgs vev

- Minimize action classically
 - Yields $hh^+ = v^2$ Higgs vev
 - Assume quantum corrections to this are small

- Minimize action classically
 - Yields $hh^+ = v^2$ Higgs vev
 - Assume quantum corrections to this are small
- Perform global gauge transformation such that

$$h(x) = \begin{vmatrix} \varphi^{1}(x) + i\varphi^{2}(x) \\ v + \eta(x) + i\varphi^{3}(x) \end{vmatrix} \Rightarrow \langle h \rangle = \begin{vmatrix} 0 \\ v \end{vmatrix}$$

- Minimize action classically
 - Yields $hh^+ = v^2$ Higgs vev
 - Assume quantum corrections to this are small
- Perform global gauge transformation such that

$$h(x) = \begin{vmatrix} \varphi^{1}(x) + i\varphi^{2}(x) \\ v + \eta(x) + i\varphi^{3}(x) \end{vmatrix} \Rightarrow \langle h \rangle = \begin{vmatrix} 0 \\ v \end{vmatrix}$$

• ηmass depends at tree-level on v

- Minimize action classically
 - Yields $hh^+ = v^2$ Higgs vev
 - Assume quantum corrections to this are small
- Perform global gauge transformation such that

$$h(x) = \begin{vmatrix} \varphi^{1}(x) + i\varphi^{2}(x) \\ v + \eta(x) + i\varphi^{3}(x) \end{vmatrix} \Rightarrow \langle h \rangle = \begin{vmatrix} 0 \\ v \end{vmatrix}$$

- ηmass depends at tree-level on v
- Perform perturbation theory

Not all charge directions equal

- Not all charge directions equal
 - This is not physical, but merely a choice of gauge

- Not all charge directions equal
 - This is not physical, but merely a choice of gauge
 - "Spontaneous gauge symmetry breaking"
 - Broken by the gauge choice, not by the dynamics
 - Dynamics only affect the length of the Higgs field

- Not all charge directions equal
 - This is not physical, but merely a choice of gauge
 - "Spontaneous gauge symmetry breaking"
 - Broken by the gauge choice, not by the dynamics
 - Dynamics only affect the length of the Higgs field
 - Local symmetry intact and cannot be broken [Elitzur PR'75]

- Not all charge directions equal
 - This is not physical, but merely a choice of gauge
 - "Spontaneous gauge symmetry breaking"
 - Broken by the gauge choice, not by the dynamics
 - Dynamics only affect the length of the Higgs field
 - Local symmetry intact and cannot be broken [Elitzur PR'75]
- Gauge symmetry no longer manifest
 - Symmetry expressed in STIs/WTIs
Implications of global transformation

- Not all charge directions equal
 - This is not physical, but merely a choice of gauge
 - "Spontaneous gauge symmetry breaking"
 - Broken by the gauge choice, not by the dynamics
 - Dynamics only affect the length of the Higgs field
 - Local symmetry intact and cannot be broken [Elitzur PR'75]
- Gauge symmetry no longer manifest
 - Symmetry expressed in STIs/WTIs
- But only way to get a working perturbation theory!

Implications of global transformation

- Not all charge directions equal
 - This is not physical, but merely a choice of gauge
 - "Spontaneous gauge symmetry breaking"
 - Broken by the gauge choice, not by the dynamics
 - Dynamics only affect the length of the Higgs field
 - Local symmetry intact and cannot be broken [Elitzur PR'75]
- Gauge symmetry no longer manifest
 - Symmetry expressed in STIs/WTIs
- But only way to get a working perturbation theory! [Lee et al.'72]
 - Otherwise W/Z massless to all orders

- Possible to have gauges without vev [Maas'13]
 - Do not select a global direction



- Possible to have gauges without vev [Maas'13]
 - Do not select a global direction



- Possible to have gauges without vev [Maas'13]
 - Do not select a global direction



- Possible to have gauges without vev [Maas'13]
 - Do not select a global direction

The physics of the Higgs sector a: The phase diagram

 Lattice simulations have an intrinsic cutoff – the lattice spacing a

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
 - Calculate for several a with all independent observables fixed
 - "Lines of constant physics"

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
 - Calculate for several a with all independent observables fixed
 - "Lines of constant physics"

Coupling(s)



- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
 - Calculate for several a with all independent observables fixed
 - "Lines of constant physics"



Mass(es)

- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
 - Calculate for several a with all independent observables fixed
 - "Lines of constant physics"



- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
 - Calculate for several a with all independent observables fixed
 - "Lines of constant physics"



- Lattice simulations have an intrinsic cutoff – the lattice spacing a
 - Full theory reached at zero lattice spacing
 - If it exists: Triviality problem
- Masses, couplings, and actions are specified at this scale
 - Numerical procedure:
 - Calculate for several a with all independent observables fixed
 - "Lines of constant physics"
 - Different starting points yield different physics



[Osterwalder et al.'78, Fradkin et al.'79 Caudy et al.'07]

• (Lattice-regularized)

f(Classical Higgs mass)

[Osterwalder et al.'78, Fradkin et al.'79 Caudy et al.'07]

• (Lattice-regularized)

f(Classical Higgs mass)



[Osterwalder et al.'78, Fradkin et al.'79 Caudy et al.'07]

• (Lattice-regularized)

f(Classical Higgs mass)



[Osterwalder et al.'78, Fradkin et al.'79 Caudy et al.'07]

• (Lattice-regularized) (sou phase diagram



[Osterwalder et al.'78, Fradkin et al.'79 Caudy et al.'07]



[Osterwalder et al.'78, Fradkin et al.'79 Caudy et al.'07]

- (Lattice-regularized) ²
 phase diagram
 continuous
 - Free energy can be shown to be analytic
 - On a finite lattice
 - Continuum limit
 - Validity for unregularized theory?
 - Works (likely) only for the standard-model case
 - Osterwalder-Seiler/Fradkin-Shenker construction



[Osterwalder et al.'78, Fradkin et al.'79 Caudy et al.'07]

- (Lattice-regularized) phase diagram continuous
 - Separation only in fixed gauges



[Osterwalder et al.'78, Fradkin et al.'79 Caudv et al.'071

- (Lattice-regularized) f(Classical Higgs mass) phase diagram continuous
 - Separation only in fixed gauges
 - But different in different in different gauges!



[Osterwalder et al.'78, Fradkin et al.'79 Caudy et al.'07]

- (Lattice-regularized)
 phase diagram
 continuous
 - Separation only in fixed gauges
- Same asymptotic states in confinement and Higgs pseudo-phases



g(Classical gauge coupling)

 Same asymptotic states irrespective of coupling strengths

[Osterwalder et al.'78, Fradkin et al.'79 Caudy et al.'07]

- (Lattice-regularized)
 phase diagram
 continuous
 - Separation only in fixed gauges
- Same asymptotic states in confinement and Higgs pseudo-phases



- Same asymptotic states irrespective of coupling strengths
- Other states than 'Higgs' and 'W'?





- Quantum effects remove BEH effect
 - Opposite does not happen



- Quantum effects remove BEH effect
 - Opposite does not happen
- Interacting continuum limit? [Gies & Zambelli'15]



- Quantum effects remove BEH effect
 - Opposite does not happen
- Interacting continuum limit? [Gies & Zambelli'15]



- Quantum effects remove BEH effect
 - Opposite does not happen
- Interacting continuum limit? [Gies & Zambelli'15]
 - LCP: 0⁺, 1⁻ mass, $\alpha(200\,GeV)$ (miniMOM scheme)

Selection criterion for candidate new physics

- Selection criterion for candidate new physics
 - Usual test: Perturbative running of couplings



- Selection criterion for candidate new physics
 - Usual test: Perturbative running of couplings
 - "Better" behavior than the standard model (1)No Landau poles (small coupling)
 (2)No triviality problem



- Selection criterion for candidate new physics
 - Usual test: Perturbative running of couplings
 - "Better" behavior than the standard model
 - (1)No Landau poles (small coupling)
 - (2)No triviality problem
 - (3)No or little fine-tuning



- Selection criterion for candidate new physics
 - Usual test: Perturbative running of couplings
 - "Better" behavior than the standard model
 - (1)No Landau poles (small coupling)
 - (2)No triviality problem
 - (3)No or little fine-tuning
 - Violated by QED (1,2), Yukawa (1,2), Higgs (1-3)
Ultraviolet structure



- Selection criterion for candidate new physics
 - Usual test: Perturbative running of couplings
 - "Better" behavior than the standard model

(1)No Landau poles (small coupling)

- (2)No triviality problem
- (3)No or little fine-tuning

• Violated by QED (1,2), Yukawa (1,2), Higgs (1-3)

Sufficient?

Ultraviolet structure



- Selection criterion for candidate new physics
 - Usual test: Perturbative running of couplings
 - "Better" behavior than the standard model

(1)No Landau poles (small coupling)

- (2)No triviality problem
- (3)No or little fine-tuning

• Violated by QED (1,2), Yukawa (1,2), Higgs (1-3)

• Sufficient? No.



- Asymptotic safety [Weinberg'79, Gies et al.'13-'17, Litim et al.'14-'17]
 - Sufficient to have a (small) finite coupling
 - . From (non-)perturbative cancellations in $\beta\mbox{-functions}$



- Asymptotic safety [Weinberg'79, Gies et al.'13-'17, Litim et al.'14-'17]
 - Sufficient to have a (small) finite coupling
 - . From (non-)perturbative cancellations in $\beta\mbox{-functions}$
- Quantum gravity can backcouple [Wetterich et al.'09, Eichhorn et al.'13-'17]
 - May solve all of these problems



- Asymptotic safety [Weinberg'79, Gies et al.'13-'17, Litim et al.'14-'17]
 - Sufficient to have a (small) finite coupling
 - . From (non-)perturbative cancellations in $\beta\mbox{-functions}$
- Quantum gravity can backcouple [Wetterich et al.'09, Eichhorn et al.'13-'17]
 - May solve all of these problems
- Fine-tuned special trajectories [Callaway'88,Litim et al.'14-'17]
 - All order cancellations solve problem



- Asymptotic safety [Weinberg'79, Gies et al.'13-'17, Litim et al.'14-'17]
 - Sufficient to have a (small) finite coupling
 - . From (non-)perturbative cancellations in $\beta\mbox{-functions}$
- Quantum gravity can backcouple [Wetterich et al.'09, Eichhorn et al.'13-'17]
 - May solve all of these problems
- Fine-tuned special trajectories [Callaway'88,Litim et al.'14-'17,Gies et al.'15,'16]
 - All order cancellations solve problem



- Experimentally hard to find
 - Energy-dependence of running couplings
 - Tiny deviations at accessible energies: Precision tests



- Experimentally hard to find
 - Energy-dependence of running couplings
 - Tiny deviations at accessible energies: Precision tests



- Experimentally hard to find
 - Energy-dependence of running couplings
 - Tiny deviations at accessible energies: Precision tests
 - Particle content constrained



- Experimentally hard to find
 - Energy-dependence of running couplings
 - Tiny deviations at accessible energies: Precision tests
 - Particle content constrained
- Quantum gravity has implications for cosmology
 - Cosmological constant becomes running
 - Tests against astrophysical data

The physics of the Higgs sector b: Observable particles





What is seen in experiment?



What did we see here?

What is seen in experiment?



What did we see here? Is this **really** the Higgs?

What is seen in experiment?



What did we see here? Is this **really** the Higgs? And what do we mean by Higgs, anyway?

- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections

- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge

- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable

- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable
- Gauge-invariant states are composite

- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable
- Gauge-invariant states are composite
 - Not asymptotic states in perturbation theory

- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable
- Gauge-invariant states are composite
 - Not asymptotic states in perturbation theory
 - Higgs-Higgs



- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable
- Gauge-invariant states are composite
 - Not asymptotic states in perturbation theory
 - Higgs-Higgs, W-W



- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable
- Gauge-invariant states are composite
 - Not asymptotic states in perturbation theory
 - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable
- Gauge-invariant states are composite
 - Not asymptotic states in perturbation theory
 - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Physical spectrum: Observable particles
 - Experiments measure peaks in cross-sections
- Elementary fields depend on the gauge
 - Cannot be observable
- Gauge-invariant states are composite
 - Not asymptotic states in perturbation theory
 - Higgs-Higgs, W-W, Higgs-Higgs-W etc.



- Why does perturbation theory work?
- Mass spectrum?

Why it does not matter in the standard model

Introducing gauge-invariant perturbation theory

 Masses are determined by poles of propagators

- Masses are determined by poles of propagators
- 2 propagators

- Masses are determined by poles of propagators
- 2 propagators
 - W/Z $D^{ab}_{\mu\nu}(x-y) = \langle W^a_{\mu}(x) W^b_{\nu}(y) \rangle$
 - Degenerate without QED

- Masses are determined by poles of propagators
- 2 propagators
 - W/Z $D^{ab}_{\mu\nu}(x-y) = \langle W^a_{\mu}(x) W^b_{\nu}(y) \rangle$
 - Degenerate without QED
 - Scalar $D_{H}^{ij}(x-y) = <\eta^{i}(x)\eta^{j+1}(y) >$

- Masses are determined by poles of propagators
- 2 propagators
 - W/Z $D^{ab}_{\mu\nu}(x-y) = \langle W^a_{\mu}(x) W^b_{\nu}(y) \rangle$
 - Degenerate without QED
 - Scalar $D_{H}^{ij}(x-y) = <\eta^{i}(x)\eta^{j+1}(y) >$
- (Tree-level/perturbative) poles of Higgs and W

- Masses are determined by poles of propagators
- 2 propagators
 - W/Z $D^{ab}_{\mu\nu}(x-y) = \langle W^a_{\mu}(x)W^b_{\nu}(y) \rangle$
 - Degenerate without QED
 - Scalar $D_{H}^{ij}(x-y) = <\eta^{i}(x)\eta^{j+1}(y) >$
- (Tree-level/perturbative) poles of Higgs and W
 - But only in a fixed gauge
 - Elementary fields are gauge-dependent
 - Without gauge fixing propagators are $\sim \delta(x-y)$

- Masses are determined by poles of propagators
- 2 propagators
 - W/Z $D^{ab}_{\mu\nu}(x-y) = \langle W^a_{\mu}(x)W^b_{\nu}(y) \rangle$
 - Degenerate without QED
 - Scalar $D_{H}^{ij}(x-y) = <\eta^{i}(x)\eta^{j+}(y) >$
- (Tree-level/perturbative) poles of Higgs and W
 - But only in a fixed gauge
 - Elementary fields are gauge-dependent
 - Without gauge fixing propagators are $\sim \delta(x-y)$
- Can be calculated in gauge-fixed lattice simulations

Effective mass



W for different lattice parameters

[Maas & Mufti'13]

Masses of the Higgs and the W/Z [Maas & Mufti'13]

Effective mass



W for different lattice parameters
Effective mass



W for different lattice parameters

Effective mass Effective mass E⁴ (1) 1.2 [/a_0 Mm Finite volume effect 120 100 0.8 80 0.6 60 0.4 40 0.2 20 0[.] 0^L 0.002 0.004 0.006 0.008 0.01 0.012 0.001 0.005 0.006 0.002 0.003 0.004 0.0 t [GeV-1] t [GeV⁻¹] Perturbative effect Higgs for (Oehme-Zimmermann) different lattice parameters

W for different lattice parameters





Intermission The W mass and QCD

[Quigg & Shrock'09]

 In the same way as technicolor QCD 'breaks' the electroweak gauge symmetry

- In the same way as technicolor QCD 'breaks' the electroweak gauge symmetry
 - Origin: Dynamically chiral symmetry breaking
 - Purely non-perturbative effect
 - Quark-Antiquark condensate

- In the same way as technicolor QCD 'breaks' the electroweak gauge symmetry
 - Origin: Dynamically chiral symmetry breaking
 - Purely non-perturbative effect
 - Quark-Antiquark condensate
- Acts exactly like the Higgs condensate

- In the same way as technicolor QCD 'breaks' the electroweak gauge symmetry
 - Origin: Dynamically chiral symmetry breaking
 - Purely non-perturbative effect
 - Quark-Antiquark condensate
- Acts exactly like the Higgs condensate
- $\ensuremath{\bullet}$ Will create (additional) mass for the W/Z







• Is it like this?



• Is it like this? No!



• Is it like this? No! Cannot create mass.



- Is it like this? No! Cannot create mass.
- Acts like an additional contribution to the condensate

 $M_W^2 \sim g^2_{weak} v^2_{Higgs} \rightarrow g^2_{weak} (v^2_{Higgs} + N_f < \overline{q}q > 3/2)$

• Essentially quark condensate



- Is it like this? No! Cannot create mass.
- Acts like an additional contribution to the condensate

- Essentially quark condensate
- Expected size : Typical effect: 30-50 MeV



- Is it like this? No! Cannot create mass.
- Acts like an additional contribution to the condensate

- Essentially quark condensate
- Expected size : Typical effect: 30-50 MeV
- Larger as current experimental error of ~20 MeV



- Is it like this? No! Cannot create mass.
- Acts like an additional contribution to the condensate

- Essentially quark condensate
- Expected size : Typical effect: 30-50 MeV
- Larger as current experimental error of ~20 MeV
- Acts like a static mass when added at tree-level



- Is it like this? No! Cannot create mass.
- Acts like an additional contribution to the condensate

- Essentially quark condensate
- Expected size : Typical effect: 30-50 MeV
- Larger as current experimental error of ~20 MeV
- Acts like a static mass when added at tree-level
 - Unitarity violation is canceled non-perturbatively

Needs to be accounted for

- Needs to be accounted for
- Same order as new physics effects
 - E.g. in 2HDM models
 - Could lead to 'false' new physics claims

- Needs to be accounted for
- Same order as new physics effects
 - E.g. in 2HDM models
 - Could lead to 'false' new physics claims
- Other non-perturbative QCD corrections exist
 - 300 MeV mass for the top (and bottom) quark
 - Higgs mixes with (heavy) mesons

- Needs to be accounted for
- Same order as new physics effects
 - E.g. in 2HDM models
 - Could lead to 'false' new physics claims
- Other non-perturbative QCD corrections exist
 - 300 MeV mass for the top (and bottom) quark
 - Higgs mixes with (heavy) mesons
- New particle with color affected
- New non-perturbative condensates contribute

Back on track Physical Mass Spectrum

 Determination of bound-state masses as in QCD or other theories

- Determination of bound-state masses as in QCD or other theories
 - Write down operator basis for quantum numbers

- Determination of bound-state masses as in QCD or other theories
 - Write down operator basis for quantum numbers
 - Only spin and custodial quantum numbers
 - Others are not gauge-invariant!

- Determination of bound-state masses as in QCD or other theories
 - Write down operator basis for quantum numbers
 - Only spin and custodial quantum numbers
 - Others are not gauge-invariant!
 - Including smeared operators

- Determination of bound-state masses as in QCD or other theories
 - Write down operator basis for quantum numbers
 - Only spin and custodial quantum numbers
 - Others are not gauge-invariant!
 - Including smeared operators
 - Correlators with variational analysis
 - Determine exponential behavior

- Determination of bound-state masses as in QCD or other theories
 - Write down operator basis for quantum numbers
 - Only spin and custodial quantum numbers
 - Others are not gauge-invariant!
 - Including smeared operators
 - Correlators with variational analysis
 - Determine exponential behavior
- Caveats
 - MUCH more noisy than QCD
 - >50k configurations

- Determination of bound-state masses as in QCD or other theories
 - Write down operator basis for quantum numbers
 - Only spin and custodial quantum numbers
 - Others are not gauge-invariant!
 - Including smeared operators
 - Correlators with variational analysis
 - Determine exponential behavior
- Caveats
 - MUCH more noisy than QCD
 - >50k configurations
 - Scalar has disconnected contributions



• Simpelst 0⁺ bound state $h^+(x)h(x)$



- Simpelst 0⁺ bound state $h^+(x)h(x)$
 - Same quantum numbers as the Higgs
 No weak or flavor charge



- Simpelst 0⁺ bound state $h^+(x)h(x)$
 - Same quantum numbers as the Higgs
 No weak or flavor charge



- Simpelst 0⁺ bound state $h^+(x)h(x)$
 - Same quantum numbers as the Higgs
 No weak or flavor charge
Higgsonium



• Simplist 0⁺ bound state $h^+(x)h(x)$

Same quantum numbers as the Higgs

• No weak or flavor charge

Mass is about 120 GeV

Higgsonium



• Simpelst 0⁺ bound state $h^+(x)h(x)$

Same quantum numbers as the Higgs

• No weak or flavor charge

Mass is about 120 GeV

Higgsonium



• Simplist 0⁺ bound state $h^+(x)h(x)$

Same quantum numbers as the Higgs

• No weak or flavor charge

Mass is about 120 GeV

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV
- Coincidence?

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV
- Coincidence? No.

[Fröhlich et al. PLB 80 Maas'12, Törek & Maas'16]

1) Formulate gauge-invariant operator

[Fröhlich et al. PLB 80 Maas'12, Törek & Maas'16]

- 1) Formulate gauge-invariant operator
 - 0⁺ singlet: $\langle (h^+ h)(x)(h^+ h)(y) \rangle$

[Fröhlich et al. PLB 80 Maas'12, Törek & Maas'16]

1) Formulate gauge-invariant operator

0⁺ singlet: $\langle (h^+ h)(x)(h^+ h)(y) \rangle$

2) Expand Higgs field around fluctuations

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = c + v^2 \langle \eta^+ (x)\eta(y) \rangle \\ + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

[Fröhlich et al. PLB 80 Maas'12, Törek & Maas'16]

1) Formulate gauge-invariant operator

0⁺ singlet: $\langle (h^+ h)(x)(h^+ h)(y) \rangle$

2) Expand Higgs field around fluctuations

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = c + v^2 \langle \eta^+ (x)\eta(y) \rangle + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

3) Standard perturbation theory

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = c + v^2 \langle \eta^+ (x)\eta(y) \rangle + \langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$$

[Fröhlich et al. PLB 80 Maas'12, Törek & Maas'16]

1) Formulate gauge-invariant operator

0⁺ singlet: $\langle (h^+ h)(x)(h^+ h)(y) \rangle$

2) Expand Higgs field around fluctuations

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = c + v^2 \langle \eta^+ (x)\eta(y) \rangle + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

3) Standard perturbation theory

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = c + v^2 \langle \eta^+ (x)\eta(y) \rangle + \langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$$

[Fröhlich et al. PLB 80 Maas'12, Törek & Maas'16]

1) Formulate gauge-invariant operator

0⁺ singlet: $\langle (h^+ h)(x)(h^+ h)(y) \rangle$

2) Expand Higgs field around fluctuations

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = c + v^2 \langle \eta^+ (x)\eta(y) \rangle + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

3) Standard perturbation theory

Bound state $\langle (h^+ h)(x)(h^+ h)(y) = c + v^2 \langle \eta^+ (x)\eta(y) \rangle$ mass $+ \langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$

[Fröhlich et al. PLB 80 Maas'12, Törek & Maas'16]

1) Formulate gauge-invariant operator

0⁺ singlet: $\langle (h^+ h)(x)(h^+ h)(y) \rangle$

2) Expand Higgs field around fluctuations

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = c + v^2 \langle \eta^+ (x)\eta(y) \rangle + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

3) Standard perturbation theory

Bound state $\langle (h^+ h)(x)(h^+ h)(y) \rangle = c + v^2 \langle \eta^+ (x)\eta(y) \rangle$ mass $+ \langle \eta^+ (x)\eta(y) \rangle \langle \eta^+ (x)\eta(y) \rangle + O(g,\lambda)$ $2 \times Higgs m$

2 x Higgs mass: Scattering state

[Fröhlich et al. PLB 80 Maas'12, Törek & Maas'16]

1) Formulate gauge-invariant operator

0⁺ singlet: $\langle (h^+ h)(x)(h^+ h)(y) \rangle$

2) Expand Higgs field around fluctuations

$$\langle (h^+ h)(x)(h^+ h)(y) \rangle = c + v^2 \langle \eta^+ (x)\eta(y) \rangle + v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle$$

 $\langle \chi \rangle \eta(y) \rangle \langle \eta^+(x) \eta(y) \rangle + O(g,\lambda)$

3) Standard perturbation theory

Bound

state

mass

 $(h^+ h)(x)(h^+ h)(y) = c + v (\eta^+ (x)\eta(y))$ Higgs mass

> 2 x Higgs mass: Scattering state

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV
- Coincidence? No.
 - Duality between elementary states and bound states
- $\langle (h^+ h)(x)(h^+ h)(y) \rangle \overset{h=\nu+\eta}{\approx} const. + \langle \eta^+ (x)\eta(y) \rangle + O(\eta^3)$
 - Same poles to leading order

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV
- Coincidence? No.
 - Duality between elementary states and bound states
- $\langle (h^+ h)(x)(h^+ h)(y) \rangle \overset{h=\nu+\eta}{\approx} const. + \langle \eta^+ (x)\eta(y) \rangle + O(\eta^3)$
 - Same poles to leading order
- Fröhlich-Morchio-Strocchi (FMS) mechanism

- Higgsonium: 120 GeV, Higgs at tree-level: 120 GeV
 - Scheme exists to shift Higgs mass always to 120 GeV
- Coincidence? No.
 - Duality between elementary states and bound states
- $\langle (h^+ h)(x)(h^+ h)(y) \rangle \overset{h=v+\eta}{\approx} const. + \langle \eta^+ (x)\eta(y) \rangle + O(\eta^3)$
 - Same poles to leading order
- Fröhlich-Morchio-Strocchi (FMS) mechanism
- Deeply-bound relativistic state
 - Mass defect~constituent mass
 - Cannot describe with quantum mechanics
 - Very different from QCD bound states



• Vector state 1⁻ with operator $tr t^a \frac{h^+}{\sqrt{h^+ h}} D_{\mu} \frac{h}{\sqrt{h^+ h}}$



- Vector state 1⁻ with operator $tr t^a \frac{h^+}{\sqrt{h^+ h}} D_{\mu} \frac{h}{\sqrt{h^+ h}}$
 - Only in a Higgs phase close to a simple particle



- Vector state 1⁻ with operator $tr Q_{\sqrt{h^+ h}}^{h^+} D_{\mu} \frac{h}{\sqrt{h^+ h}}$
 - Only in a Higgs phase close to a simple particle
 - Custodial triplet, instead of gauge triplet



- Vector state 1⁻ with operator $tr t^a \frac{h^+}{\sqrt{h^+ h}} D_{\mu} \frac{h}{\sqrt{h^+ h}}$
 - Only in a Higgs phase close to a simple particle
 - Custodial triplet, instead of gauge triplet



• Vector state 1⁻ with operator $tr t^a \frac{h}{\sqrt{h+h}} D_{\mu} \frac{h}{\sqrt{h+h}}$

- Only in a Higgs phase close to a simple particle
- Custodial triplet, instead of gauge triplet
- Mass about 80 GeV

- Vector state: 80 GeV
- W at tree-level: 80 GeV
 - W not scale or scheme dependent

[Fröhlich et al.'80 Maas'12]

- Vector state: 80 GeV
- W at tree-level: 80 GeV
 - W not scale or scheme dependent
- Same mechanism

 $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y) \rangle$

- Vector state: 80 GeV
- W at tree-level: 80 GeV
 - W not scale or scheme dependent
- Same mechanism

$$\langle (h^{+} D_{\mu}h)(x)(h^{+} D_{\mu}h)(y) \rangle$$

$$h = v + \eta$$

$$\approx const. + \langle W_{\mu}(x)W_{\mu}(y) \rangle + O(\eta^{3})$$

$$\partial v = 0$$

- Vector state: 80 GeV
- W at tree-level: 80 GeV
 - W not scale or scheme dependent
- Same mechanism

$$\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y) \rangle$$

$$h = v + \eta$$

$$\approx const. + \langle W_{\mu}(x)W_{\mu}(y) \rangle + O(\eta^3)$$

$$\partial v = 0$$

mo polos at loading order

- Same poles at leading order
 - Remains true beyond leading order

- Vector state: 80 GeV
- W at tree-level: 80 GeV
 - W not scale or scheme dependent
- Same mechanism

$$\langle (h^+ D_{\mu} h)(x)(h^+ D_{\mu} h)(y) \rangle$$

$$h = v + \eta$$

$$\approx const. + \langle W_{\mu}(x) W_{\mu}(y) \rangle + O(\eta^3)$$

$$\partial v = 0$$

- Same poles at leading order
 - Remains true beyond leading order
 - Exchanges a gauge for a custodial triplet

The remainder of the standard model

Is there an experimental lever to see all this?

[Fröhlich et al.'80, Egger, Maas, Sondenheimer'17]

• QED

Also requires gauge-invariant description

- QED
 - Also requires gauge-invariant description
 - But can be done using phase factors (Dirac phases)
 - 'Almost' local non-trivial on the lattice

- QED
 - Also requires gauge-invariant description
 - But can be done using phase factors (Dirac phases)
 - 'Almost' local non-trivial on the lattice
- Right-handed Dirac neutrinos trivial
 - Not gauged under anything

- QED
 - Also requires gauge-invariant description
 - But can be done using phase factors (Dirac phases)
 - 'Almost' local non-trivial on the lattice
- Right-handed Dirac neutrinos trivial
 - Not gauged under anything
- Quarks: Bound by confinement in hadrons

- QED
 - Also requires gauge-invariant description
 - But can be done using phase factors (Dirac phases)
 - 'Almost' local non-trivial on the lattice
- Right-handed Dirac neutrinos trivial
 - Not gauged under anything
- Quarks: Bound by confinement in hadrons?

- QED
 - Also requires gauge-invariant description
 - But can be done using phase factors (Dirac phases)
 - 'Almost' local non-trivial on the lattice
- Right-handed Dirac neutrinos trivial
 - Not gauged under anything
- Quarks: Bound by confinement in hadrons?
- Leptons
 - Flavor is actually weak gauge interaction

Flavor

[Fröhlich et al.'80, Egger, Maas, Sondenheimer'17]
[Fröhlich et al.'80, Egger, Maas, Sondenheimer'17]

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
 Replaced by bound state EMS applicable
- Replaced by bound state FMS applicable

 $\langle (h_{ia}^{+} f_{a})(x)^{+} (h_{ib}^{+} f_{b})(y) \rangle \overset{h=\nu+\eta}{\approx} \langle f_{a}^{+} (x) f_{a}(y) \rangle + O(\eta)$

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
- Replaced by bound state FMS applicable

 $\langle (h_{a}^{+} f_{a})(x)^{+} (h_{b}^{+} f_{b})(y) \rangle \overset{h=\nu+\eta}{\approx} \langle f_{a}^{+} (x) f_{a}(y) \rangle + O(\eta)$

Gauge-invariant state, but custodial doublet

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
 Replaced by bound state FMS applicable

 $\langle (h_{a}^{+} f_{a})(x)^{+} (h_{b}^{+} f_{b})(y) \rangle \overset{h=\nu+\eta}{\approx} \langle f_{a}^{+} (x) f_{a}(y) \rangle + O(\eta)$

- Gauge-invariant state, but custodial doublet
- Yukawa terms break custodial symmetry
 - Different masses for doublet members

- Flavor has two components
 - Global SU(3) generation
 - Local SU(2) weak gauge (up/down distinction)
- Same argument: Weak gauge not observable
 Replaced by bound state FMS applicable

 $\langle (h_{a}^{+} f_{a})(x)^{+} (h_{b}^{+} f_{b})(y) \rangle \overset{h=\nu+\eta}{\approx} \langle f_{a}^{+} (x) f_{a}(y) \rangle + O(\eta)$

- Gauge-invariant state, but custodial doublet
- Yukawa terms break custodial symmetry
 - Different masses for doublet members
- Hard to test but maybe even more possibilities

- Flavor is replaced by custodial symmetry
- Straightforward for leptons

- Flavor is replaced by custodial symmetry
- Straightforward for leptons
- Implications for hadrons?

- Flavor is replaced by custodial symmetry
- Straightforward for leptons
- Implications for hadrons?
- Open flavor must be replaced by custodial symmetry

- Flavor is replaced by custodial symmetry
- Straightforward for leptons
- Implications for hadrons?
- Open flavor must be replaced by custodial symmetry
- Requires Higgs component

- Flavor is replaced by custodial symmetry
- Straightforward for leptons
- Implications for hadrons?
- Open flavor must be replaced by custodial symmetry
- Requires Higgs component
- Consider nucleon
- qqq open flavor, cannot be gauge invariant
 - Impossible to build a gauge-invariant 3-quark state

- Flavor is replaced by custodial symmetry
- Straightforward for leptons
- Implications for hadrons?
- Open flavor must be replaced by custodial symmetry
- Requires Higgs component
- Consider nucleon
- qqq open flavor, cannot be gauge invariant
 - Impossible to build a gauge-invariant 3-quark state
- Replacement: qqqh

- Flavor is replaced by custodial symmetry
- Straightforward for leptons
- Implications for hadrons?
- Open flavor must be replaced by custodial symmetry
- Requires Higgs component
- Consider nucleon
- qqq open flavor, cannot be gauge invariant
 - Impossible to build a gauge-invariant 3-quark state
- Replacement: qqqh
 - FMS mechanism as usual yields QCD

- Flavor is replaced by custodial symmetry
- Straightforward for leptons
- Implications for hadrons?
- Open flavor must be replaced by custodial symmetry
- Requires Higgs component
- Consider nucleon
- qqq open flavor, cannot be gauge invariant
 - Impossible to build a gauge-invariant 3-quark state
- Replacement: qqqh
 - FMS mechanism as usual yields QCD
 - Detectable at LHC?

- Flavor is replaced by custodial symmetry
- Straightforward for leptons
- Implications for hadrons?
- Open flavor must be replaced by custodial symmetry
- Requires Higgs component
- Consider nucleon
- qqq open flavor, cannot be gauge invariant
 - Impossible to build a gauge-invariant 3-quark state
- Replacement: qqqh
 - FMS mechanism as usual yields QCD
 - Detectable at LHC? Large QCD background

- Flavor is replaced by custodial symmetry
- Straightforward for leptons
- Implications for hadrons?
- Open flavor must be replaced by custodial symmetry
- Requires Higgs component
- Consider nucleon
- qqq open flavor, cannot be gauge invariant
 - Impossible to build a gauge-invariant 3-quark state
- Replacement: qqqh
 - FMS mechanism as usual yields QCD
 - Detectable at LHC? Large QCD background. Test leptons

[Maas'12]



Collision of bound states

[Maas'12]



Collision of bound states - 'constituent' particles



- Collision of bound states 'constituent' particles
- Higgs partners just spectators
 - Similar to pp collisions

[Maas'12]



- Collision of bound states 'constituent' particles
- Higgs partners just spectators
 - Similar to pp collisions
- Sub-leading contributions

[Maas'12]



- Collision of bound states 'constituent' particles
- Higgs partners just spectators
 - Similar to pp collisions
- Sub-leading contributions
 - Ordinary ones: Large and detected

[Maas'12]



- Collision of bound states 'constituent' particles
- Higgs partners just spectators
 - Similar to pp collisions
- Sub-leading contributions
 - Ordinary ones: Large and detected
 - New ones: Small, require more sensitivity



Description of impact?





 $\langle hehe | h\mu h\mu \rangle$



 $\langle hehe|h\mu h\mu \rangle = \langle ee|\mu\mu \rangle$

Ordinary contribution



 $\langle hehe|h\mu h\mu \rangle = \langle ee|\mu\mu \rangle + \langle \eta\eta \rangle \langle ee|\mu\mu \rangle$

- Ordinary contribution
- Modification of ordinary contribution



 $\langle hehe|h\mu h\mu \rangle = \langle ee|\mu\mu \rangle + \langle \eta\eta \rangle \langle ee|\mu\mu \rangle + \langle ee \rangle \langle \eta\eta|\mu\mu \rangle$

- Ordinary contribution
- Modification of ordinary contribution
- Higgs as initial state



 $\langle hehe|h\mu h\mu \rangle = \langle ee|\mu\mu \rangle + \langle \eta\eta \rangle \langle ee|\mu\mu \rangle + \langle ee \rangle \langle \eta\eta|\mu\mu \rangle + \dots$

- Ordinary contribution
- Modification of ordinary contribution
- Higgs as initial state
- More contributions...



 $\langle hehe|h\mu h\mu \rangle = \langle ee|\mu\mu \rangle + \langle \eta\eta \rangle \langle ee|\mu\mu \rangle + \langle ee \rangle \langle \eta\eta|\mu\mu \rangle + \dots$

- Ordinary contribution
- Modification of ordinary contribution
- Higgs as initial state
- More contributions...complicated



Description of impact? PDF-type language!



- Description of impact? PDF-type language!
- Interacting particles either electrons



- Description of impact? PDF-type language!
- Interacting particles either electrons or Higgs





- Description of impact? PDF-type language!
- Interacting particles either electrons or Higgs
- Fragmentation 100% efficient like for quarks




















[Egger et al.'17]

• Why three generations?

- Why three generations?
- Leptons/Quarks are bound states
 - Could other generations be internal excitations?

- Why three generations?
- Leptons/Quarks are bound states
 - Could other generations be internal excitations?
- Consider a one-flavor standard model with such excitations

- Why three generations?
- Leptons/Quarks are bound states
 - Could other generations be internal excitations?
- Consider a one-flavor standard model with such excitations
 - The low-energy effective theory is exactly the ordinary three-generation standard model

- Why three generations?
- Leptons/Quarks are bound states
 - Could other generations be internal excitations?
- Consider a one-flavor standard model with such excitations
 - The low-energy effective theory is exactly the ordinary three-generation standard model
 - CKM/PMNS matrices are decay matrix elements

- Why three generations?
- Leptons/Quarks are bound states
 - Could other generations be internal excitations?
- Consider a one-flavor standard model with such excitations
 - The low-energy effective theory is exactly the ordinary three-generation standard model
 - CKM/PMNS matrices are decay matrix elements
 - CP violation is dynamically generated

- Why three generations?
- Leptons/Quarks are bound states
 - Could other generations be internal excitations?
- Consider a one-flavor standard model with such excitations
 - The low-energy effective theory is exactly the ordinary three-generation standard model
 - CKM/PMNS matrices are decay matrix elements
 - CP violation is dynamically generated
 - Far-reaching consequences for BSM searches and cosmology

- Why three generations?
- Leptons/Quarks are bound states
 - Could other generations be internal excitations?
- Consider a one-flavor standard model with such excitations
 - The low-energy effective theory is exactly the ordinary three-generation standard model
 - CKM/PMNS matrices are decay matrix elements
 - CP violation is dynamically generated
 - Far-reaching consequences for BSM searches and cosmology
 - Almost impossible to check with current methods

Why it can matter beyond the standard model

And when this can be dealt with using gauge-invariant perturbation theory

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true?

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]
 - Fluctuations can invalidate it





Effective mass

Gauge-dependent quantity



FMS prediction



FMS prediction



Too low: Finite volume effect

FMS prediction

Too low: Finite volume effect



Elastic decay threshold Higgs as resonance Expensive, signal very bad



Too low: Finite volume effect



Elastic decay threshold Higgs as resonance Expensive, signal very bad

Higgs and W mass agrees FMS stops working



Too low: Finite volume effect



Elastic decay threshold Higgs as resonance Expensive, signal very bad

Higgs and W mass agrees FMS stops working So does Brout-Englert-Higgs!







Does not coincide with weak/strong coupling transitions!



Does not coincide with weak/strong coupling transitions! Why?

Higgs mass



No strong dependence of mass range on cutoff - expected

Perturbative predictivity: Mass ratios


Perturbative predictivity: Mass ratios



Perturbative predictivity: Mass ratios



Perturbative predictivity: Mass ratios



Comparability to the standard model

2 correct masses only fix two parameters, but
 3 parameters needed

Comparability to the standard model

- 2 correct masses only fix two parameters, but
 3 parameters needed
- Comparison to standard model complicated
 - States stable, no W/Z splitting

Comparability to the standard model

- 2 correct masses only fix two parameters, but
 3 parameters needed
- Comparison to standard model complicated
 - States stable, no W/Z splitting
- Couplings run differently proceed with caution





[Maas, unpublished]



[[]Maas, unpublished]



[[]Maas, unpublished]



[Maas, unpublished]



[Maas, unpublished]

Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]
 - Fluctuations can invalidate it
- So what could happen?

[Maas et al. Unpublished]

 (Singlet) quartic gauge coupling and resonance formation in the same channel



[Maas et al. Unpublished]

 (Singlet) quartic gauge coupling and resonance formation in the same channel



[Maas et al. Unpublished]

 (Singlet) quartic gauge coupling and resonance formation in the same channel



[Maas et al. Unpublished]

 (Singlet) quartic gauge coupling and resonance formation in the same channel



Resonance peak in final state invariant mass?









SPECULATIVE



[Low-energy effective Lagrangian, MC by Sherpa 1.4.2]

[Maas et al. Unpublished]

SPECULATIVE



[Low-energy effective Lagrangian, MC by Sherpa 1.4.2]

[Maas et al. Unpublished]

SPECULATIVE



[Low-energy effective Lagrangian, MC by Sherpa 1.4.2]

- E.g. excited Higgs: Decay channel: 2W
- Does it happen in the standard model?

[Maas et al. Unpublished]

SPECULATIVE



[Low-energy effective Lagrangian, MC by Sherpa 1.4.2]

- E.g. excited Higgs: Decay channel: 2W
- Does it happen in the standard model?
- If yes, this would fake new physics!



[Maas et al. Unpublished]

[Low-energy effective Lagrangian, MC by Sherpa 1.4.2]

- E.g. excited Higgs: Decay channel: 2W
- Does it happen in the standard model?
- If yes, this would fake new physics!
 - And would be hard to detect

- Each quantum number channel has a spectrum
 - Discreet in a finite volume



- Each quantum
 number channel
 has a spectrum
 - Discreet in a finite volume
- States can be either stable, excited states,



- Each quantum
 number channel
 has a spectrum
 - Discreet in a finite volume
- States can be either stable, excited states, resonances



- Each quantum
 number channel
 has a spectrum
 - Discreet in a finite volume
- States can be either stable, excited states, resonances or scattering states









Spectrum



Spectrum



Spectrum


Excited states on the lattice

Spectrum Scattering states Inelastic threshold: H->2H 250 200 Elastic threshold: H->2W 150 [VaD] m Ground state 100 50 0 0.1 0.05 0

[Luescher'85,'86,'90,'91]

Excited states on the lattice

250

200

<u>1</u>50[′] [Лар] ш

100

50

0

0.05

Spectrum Scattering states Inelastic threshold: H->2H Avoided level crossing Identification and widths from phase shifts Elastic threshold: H->2W Ground state

[Luescher'85,'86,'90,'91]

0.1

Search: Excited Higgs







Typical spectra

Spectrum Higgs-like



Typical spectra



Typical spectra



Typical spectra



Typical spectra

Spectrum Higgs-like

Spectrum QCD-like



- Generically different low-lying spectra
 - 0⁺ lighter in QCD-like region
 - 1⁻ lighter in Higgs-like region

Typical spectra

Spectrum Higgs-like

Spectrum QCD-like



- Generically different low-lying spectra
 - 0⁺ lighter in QCD-like region
 - 1⁻ lighter in Higgs-like region
- Coincides with gauge-dependent definitions

Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]
 - Fluctuations can invalidate it
- So what could happen? Probably nothing.

Status of the standard model

- Physical states are bound states
 - Observed in experiment
 - Described using gauge-invariant perturbation theory based on the FMS mechanism
 - Mostly the same as ordinary perturbation theory
- Is this always true? No. [Maas'15, Maas & Mufti'14]
 - Fluctuations can invalidate it
- So what could happen? Probably nothing.
- What happens for a different structure
 - Different gauge group or more Higgs

Example 1: 2HDM

Like the standard model Gauge-invariant and ordinary perturbation theory coincide

[Maas'15, Maas & Pedro'16]

- Additional Higgs doublet
- Enlarged custodial group

- Additional Higgs doublet
- Enlarged custodial group
- BEH Effect FMS mechanism applicable
 - In a suitable basis, all condensates contained in a single doublet

- FMS states for maximal custodial group:
 - Scalar sector Singlet

 $\langle (h^+ h)(x)(h^+ h)(y) \rangle \approx const. + \langle \eta_h^+ (x)\eta_h(y) \rangle + O(\eta_h^3)$

- FMS states for maximal custodial group:
 - Scalar sector Singlet

 $\langle (h^+ h)(x)(h^+ h)(y) \rangle \approx const. + \langle \eta_h^+ (x)\eta_h(y) \rangle + O(\eta_h^3)$

Scalar Sector Quadruplet

 $\langle (a + \Gamma a)(x)(a + \Gamma a)(y) \rangle \approx const. + \langle \eta_a + (x) \Gamma \eta_a(y) \rangle + O(\eta_a^3)$

Splitted into 1+3 states for broken group

- FMS states for maximal custodial group:
 - Scalar sector Singlet

 $\langle (h^+ h)(x)(h^+ h)(y) \rangle \approx const. + \langle \eta_h^+ (x) \eta_h(y) \rangle + O(\eta_h^3)$

- Scalar Sector Quadruplet
- $\langle (a + \Gamma a)(x)(a + \Gamma a)(y) \rangle \approx const. + \langle \eta_a + (x) \Gamma \eta_a(y) \rangle + O(\eta_a^3)$

Splitted into 1+3 states for broken group

• Vector triplet

 $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y) \rangle \approx const. + \langle W_{\mu}(x)W_{\mu}(y) \rangle + O(\eta_h^3)$

All other states expand to scattering states

- FMS states for maximal custodial group:
 - Scalar sector Singlet

 $\langle (h^+ h)(x)(h^+ h)(y) \rangle \approx const. + \langle \eta_h^+ (x)\eta_h(y) \rangle + O(\eta_h^3)$

- Scalar Sector Quadruplet
- $\langle (a + \Gamma a)(x)(a + \Gamma a)(y) \rangle \approx const. + \langle \eta_a + (x)\Gamma \eta_a(y) \rangle + O(\eta_a^3)$

Splitted into 1+3 states for broken group

• Vector triplet

 $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y) \rangle \approx const. + \langle W_{\mu}(x)W_{\mu}(y) \rangle + O(\eta_h^3)$

- All other states expand to scattering states
- Validity: Requires non-perturbative check

- FMS states for maximal custodial group:
 - Scalar sector Singlet

 $\langle (h^+ h)(x)(h^+ h)(y) \rangle \approx const. + \langle \eta_h^+ (x)\eta_h(y) \rangle + O(\eta_h^3)$

Scalar Sector Quadruplet

 $\langle (a + \Gamma a)(x)(a + \Gamma a)(y) \rangle \approx const. + \langle \eta_a + (x)\Gamma \eta_a(y) \rangle + O(\eta_a^3)$

Splitted into 1+3 states for broken group

• Vector triplet

 $\langle (h^+ D_{\mu}h)(x)(h^+ D_{\mu}h)(y) \rangle \approx const. + \langle W_{\mu}(x)W_{\mu}(y) \rangle + O(\eta_h^3)$

- All other states expand to scattering states
- Validity: Requires non-perturbative check
- Discrete factor groups could yield doubling

- Additional Higgs doublet
- Enlarged custodial group
- BEH Effect FMS mechanism applicable
 - In a suitable basis, all condensates contained in a single doublet
 - Yields again perturbative spectrum

• Discrete factor groups may be a problem

• Key: Global multiplet structure diverse

- Additional Higgs doublet
- Enlarged custodial group
- BEH Effect FMS mechanism applicable
 - In a suitable basis, all condensates contained in a single doublet
 - Yields again perturbative spectrum

• Discrete factor groups may be a problem

- Key: Global multiplet structure diverse
- Size of fluctuations needs to be checked non-perturbatively!

Example 2: GUT-like structure

Gauge-invariant perturbation theory correct and different from ordinary perturbation theory

[Maas'15 Törek & Maas '15, '16]

 GUTs: Large gauge group, small custodial group

- GUTs: Large gauge group, small custodial group
 - Standard model structure: diagonal subgroup not gauge-invariant

- GUTs: Large gauge group, small custodial group
 - Standard model structure: diagonal subgroup not gauge-invariant
- Toy-GUT: SU(3) broken to SU(2)

Test for GUTs



Separation into Higgs-like and QCD-like

- GUTs: Large gauge group, small custodial group
 - Standard model structure: diagonal subgroup not gauge-invariant
- Toy-GUT: SU(3) broken to SU(2)
 - Perturbative spectrum
 - 1 massive Higgs, 3 massless and
 - 5 (1 (heavier) + 4 (lighter)) massive vectors

- GUTs: Large gauge group, small custodial group
 - Standard model structure: diagonal subgroup not gauge-invariant
- Toy-GUT: SU(3) broken to SU(2)
 - Perturbative spectrum
 - 1 massive Higgs, 3 massless and
 - 5 (1 (heavier) + 4 (lighter)) massive vectors
 - FMS spectrum
 - 1 massive scalar, 1 massive vector
 - Same masses as Higgs and heaviest gauge boson

- GUTs: Large gauge group, small custodial group
 - Standard model structure: diagonal subgroup not gauge-invariant
- Toy-GUT: SU(3) broken to SU(2)
 - Perturbative spectrum
 - 1 massive Higgs, 3 massless and
 - 5 (1 (heavier) + 4 (lighter)) massive vectors
 - FMS spectrum
 - 1 massive scalar, 1 massive vector
 - Same masses as Higgs and heaviest gauge boson
 - ...or something else?

[Maas & Törek'16 Maas, Sondenheimer & Törek'17]

Perturbation theory Gauge-dependent





[Maas & Törek'16 Maas, Sondenheimer & Törek'17]

Perturbation theory Gauge-dependent

FMS expansion U(1) singlets



[Maas & Törek'16 Maas, Sondenheimer & Törek'17]

Perturbation theory Gauge-dependent FMS expansion U(1) singlets



Qualitatively different spectrum

[Maas & Törek'16 Maas, Sondenheimer & Törek'17]



Qualitatively different spectrum

[Maas & Törek'16 Maas, Sondenheimer & Törek'17]



Qualitatively different spectrum


Qualitatively different spectrum



Qualitatively different spectrum

Scalars similar but no confirmation yet (statistics!)



Theory weakly interacting



Theory weakly interacting



[Maas & Törek'16+unpublished Maas, Sondenheimer & Törek'17]



[Maas, Sondenheimer & Törek'17]

 SU(3) case is generic for SU(N) with a single fundamental Higgs

- SU(3) case is generic for SU(N) with a single fundamental Higgs
 - Elementary spectrum: More particles

- SU(3) case is generic for SU(N) with a single fundamental Higgs
 - Elementary spectrum: More particles
 - Physical spectrum: Independent of N
 - Reason: Same custodial symmetry

- SU(3) case is generic for SU(N) with a single fundamental Higgs
 - Elementary spectrum: More particles
 - Physical spectrum: Independent of N
 - Reason: Same custodial symmetry
 - Always a mismatch to perturbation theory
 - No lattice tests yet

- SU(3) case is generic for SU(N) with a single fundamental Higgs
 - Elementary spectrum: More particles
 - Physical spectrum: Independent of N
 - Reason: Same custodial symmetry
 - Always a mismatch to perturbation theory
 - No lattice tests yet
- Multiple Higgs fields: Larger custodial symmetry

- SU(3) case is generic for SU(N) with a single fundamental Higgs
 - Elementary spectrum: More particles
 - Physical spectrum: Independent of N
 - Reason: Same custodial symmetry
 - Always a mismatch to perturbation theory
 - No lattice tests yet
- Multiple Higgs fields: Larger custodial symmetry
- Note: Predicts vector lighter than scalar!

[Maas, Sondenheimer & Törek'17]

• SU(N) with adjoint Higgs

[Maas, Sondenheimer & Törek'17]

• SU(N) with adjoint Higgs subtle

- SU(N) with adjoint Higgs subtle
- N=2

- SU(N) with adjoint Higgs subtle
- N=2
 - Perturbatively breaks to U(1): QED?

- SU(N) with adjoint Higgs subtle
- N=2
 - Perturbatively breaks to U(1): QED?
 - Not gauge-invariant no U(1) of SU(2) special!

- SU(N) with adjoint Higgs subtle
- N=2
 - Perturbatively breaks to U(1): QED?
 - Not gauge-invariant no U(1) of SU(2) special!
 - But: Massless vector state predicted

- SU(N) with adjoint Higgs subtle
- N=2
 - Perturbatively breaks to U(1): QED?
 - Not gauge-invariant no U(1) of SU(2) special!
 - But: Massless vector state predicted
 - Standard gauge-invariant perturbation theory

- SU(N) with adjoint Higgs subtle
- N=2
 - Perturbatively breaks to U(1): QED?
 - Not gauge-invariant no U(1) of SU(2) special!
 - But: Massless vector state predicted
 - Standard gauge-invariant perturbation theory
 - If interactions are right could be an effective U(1)
 - Not yet tested

- SU(N) with adjoint Higgs subtle
- N=2
 - Perturbatively breaks to U(1): QED?
 - Not gauge-invariant no U(1) of SU(2) special!
 - But: Massless vector state predicted
 - Standard gauge-invariant perturbation theory
 - If interactions are right could be an effective U(1)
 - Not yet tested
 - Still mismatch in other channels

- SU(N) with adjoint Higgs subtle
- N>2

- SU(N) with adjoint Higgs subtle
- N>2
 - Multiple breaking patterns

- SU(N) with adjoint Higgs subtle
- N>2
 - Multiple breaking patterns
 - E.g. SU(3): SU(2)xU(1) and U(1)xU(1)

- SU(N) with adjoint Higgs subtle
- N>2
 - Multiple breaking patterns
 - E.g. SU(3): SU(2)xU(1) and U(1)xU(1)
 - Gauge-invariant distinction per configuration

- SU(N) with adjoint Higgs subtle
- N>2
 - Multiple breaking patterns
 - E.g. SU(3): SU(2)xU(1) and U(1)xU(1)
 - Gauge-invariant distinction per configuration
 - But: Average over all configurations!

- SU(N) with adjoint Higgs subtle
- N>2
 - Multiple breaking patterns
 - E.g. SU(3): SU(2)xU(1) and U(1)xU(1)
 - Gauge-invariant distinction per configuration
 - But: Average over all configurations!
 - Is there a dominance?
 - Different phases or gradual distinction?
 - Mixing of features or distinct physics?

- SU(N) with adjoint Higgs subtle
- N>2
 - Multiple breaking patterns
 - E.g. SU(3): SU(2)xU(1) and U(1)xU(1)
 - Gauge-invariant distinction per configuration
 - But: Average over all configurations!
 - Is there a dominance?
 - Different phases or gradual distinction?
 - Mixing of features or distinct physics?
 - Spectrum still different in every pattern
 - Too many massless states, too few massive

Example 3: Technicolor

No gauge-invariant perturbation theory but interesting implications

 Higgs replaced by bound state of new fermions (techniquarks) and new gauge interaction (technicolor)

- Higgs replaced by bound state of new fermions (techniquarks) and new gauge interaction (technicolor)
 - No BEH effect: FMS cannot work

- Higgs replaced by bound state of new fermions (techniquarks) and new gauge interaction (technicolor)
 - No BEH effect: FMS cannot work
- Observable states must still be gaugeinvariant
 - Needs to create Higgs and W/Z(!) signals by (new) bound states

- Higgs replaced by bound state of new fermions (techniquarks) and new gauge interaction (technicolor)
 - No BEH effect: FMS cannot work
- Observable states must still be gaugeinvariant
 - Needs to create Higgs and W/Z(!) signals by (new) bound states
 - Vectors must be lighter
 - Behavior not yet seen for strong interactions
 - Usually: Scalars and pseudoscalars

Summary

UPCOMING REVIEW IN MID DECEMBER

Brief introduction: Törek & Maas'16

• Observable (particles) must be gauge-invariant

Summary

- Observable (particles) must be gauge-invariant
- In non-Abelian gauge theories: Bound states

Summary

- Observable (particles) must be gauge-invariant
- In non-Abelian gauge theories: Bound states
- Gauge-invariant perturbation theory as a tool
- Observable (particles) must be gauge-invariant
- In non-Abelian gauge theories: Bound states
- Gauge-invariant perturbation theory as a tool
 - Requires a Brout-Englert-Higgs effect

- Observable (particles) must be gauge-invariant
- In non-Abelian gauge theories: Bound states
- Gauge-invariant perturbation theory as a tool
 - Requires a Brout-Englert-Higgs effect
 - Yields the same results for the standard model
 - More robust
 - Mostly not much more complicated

- Observable (particles) must be gauge-invariant
- In non-Abelian gauge theories: Bound states
- Gauge-invariant perturbation theory as a tool
 - Requires a Brout-Englert-Higgs effect
 - Yields the same results for the standard model
 - More robust
 - Mostly not much more complicated
- Applicable to beyond-the standard model

- Observable (particles) must be gauge-invariant
- In non-Abelian gauge theories: Bound states
- Gauge-invariant perturbation theory as a tool
 - Requires a Brout-Englert-Higgs effect
 - Yields the same results for the standard model
 - More robust
 - Mostly not much more complicated
- Applicable to beyond-the standard model
 - Structural requirement: Multiplets must match

- Observable (particles) must be gauge-invariant
- In non-Abelian gauge theories: Bound states
- Gauge-invariant perturbation theory as a tool
 - Requires a Brout-Englert-Higgs effect
 - Yields the same results for the standard model
 - More robust
 - Mostly not much more complicated
- Applicable to beyond-the standard model
 - Structural requirement: Multiplets must match
 - Dynamical requirement: Small fluctuations

- Observable (particles) must be gauge-invariant
- In non-Abelian gauge theories: Bound states
- Gauge-invariant perturbation theory as a tool
 - Requires a Brout-Englert-Higgs effect
 - Yields the same results for the standard model
 - More robust
 - Mostly not much more complicated
- Applicable to beyond-the standard model
 - Structural requirement: Multiplets must match
 - Dynamical requirement: Small fluctuations
 - Often at odds with standard phenomenology

Advertisment



Bound States in Strongly Coupled Systems A Galileo Galilei Workshop 12th-16th of March 2018 Firenze, Italy https://www.ggi.infn.it/showevent.pl?id=286

Alps 2018 An Alpine Physics Summit 15th-20th of April 2018 Obergurgl, Austria alps2018.hephy.at