

# FMS mechanism in the standard model and beyond

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with

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University of Graz

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# Outline

[Törek and Maas, LATTICE2016, 1610.04188]

## Physical states in the standard model

Standard description

Gauge-invariant perturbation  
theory

Physical states in a grand-unified  
setting

# Physical states in the standard model

# The problem in the standard model

- Consider gauge-Higgs sector of the standard model:

$$\mathcal{L} = -\frac{1}{4} W_{\mu\nu}^a W^{a\mu\nu} + (D_\mu \phi)^\dagger (D^\mu \phi) - \lambda (\phi^\dagger \phi - v^2)^2$$

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- Global  $SU(2)$  Higgs flavor symmetry:  
Custodial symmetry

$$W_\mu^a \rightarrow W_\mu^a \quad \phi_i \rightarrow M_{ij} \phi_j + N_{ij} \phi_j^*$$

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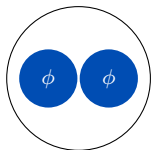
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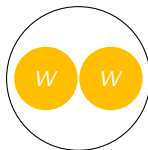
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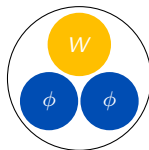
Higgs-Higgs



W-W



Higgs-Higgs-W



et cetera

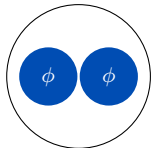
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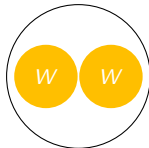
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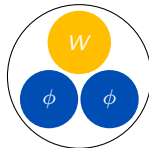
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- What is the mass spectrum?
- Why does perturbation theory work?

# Masses from propagators

- Poles of propagators  $\Rightarrow$  Masses

- Two propagators:

- W/Z:  $D_{\mu\nu}^{ab}(x - y) = \langle W_{\mu}^a(x) W_{\nu}^b(y) \rangle$

- Degenerate without QED

- Scalar:  $D_{ij}(x - y) = \langle \eta_i(x) \eta_j^{\dagger}(y) \rangle$

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  - Elementary fields are gauge dependent
  - No gauge fixing: Propagators  $\propto \delta(x - y)$

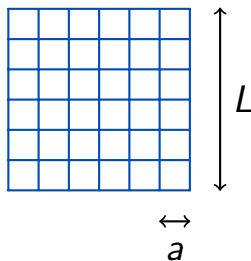


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- For gauge-invariant states: Non-perturbative method

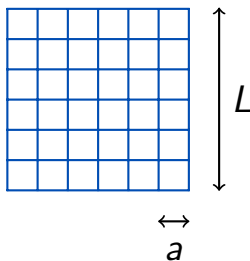
# Lattice calculations

- Finite volume (hypercube)
- Discretization  $\Rightarrow$  Finite hypercubic lattice



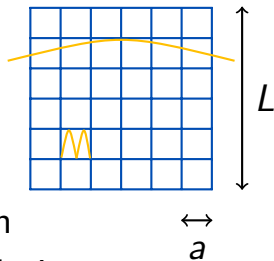
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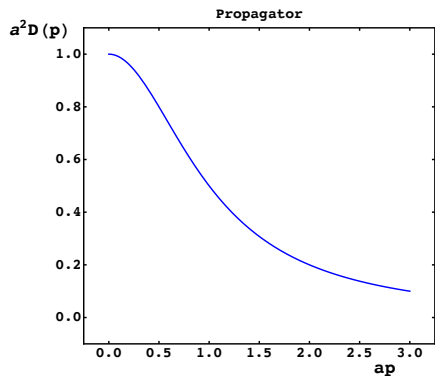


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- Artifacts
  - Finite volume and discretization
  - Masses vs. wave lengths: Resolution
  - Euclidean formulation

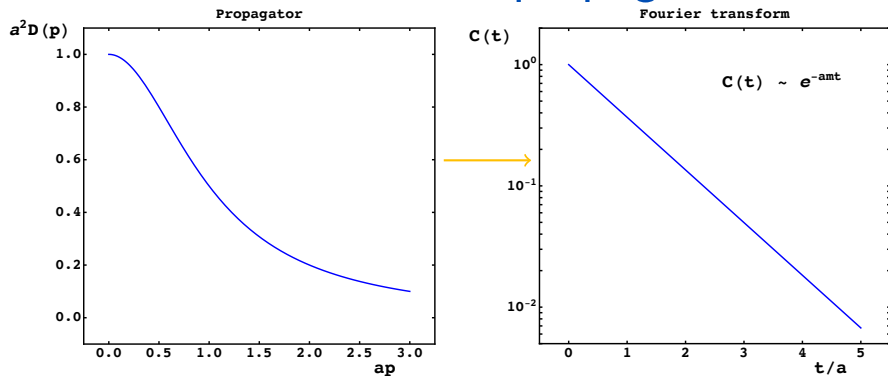


# Masses from Euclidean propagators



■ Propagator:  $D(p) = \langle \mathcal{O}(p) \mathcal{O}^\dagger(p) \rangle \propto \sum_i \frac{a_i}{p^2 + m_i^2}$

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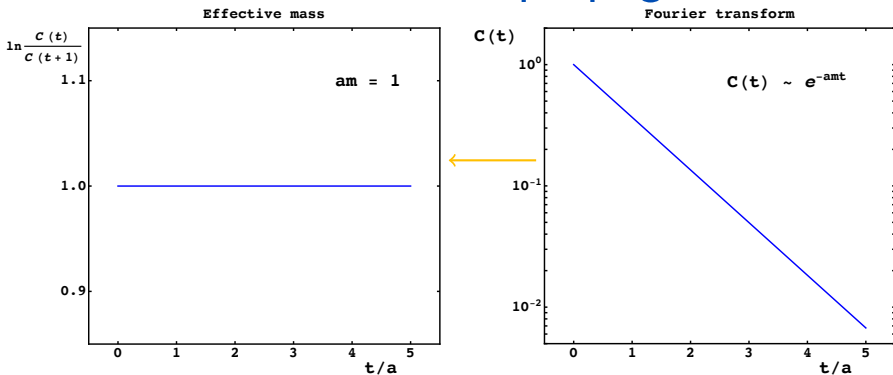


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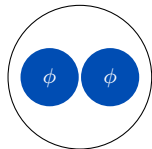
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■ Extract effective mass

# Mass spectrum: Higgs-Higgs

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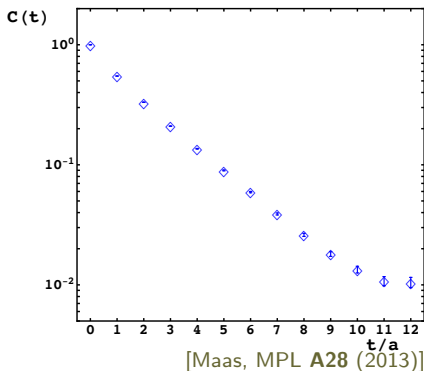


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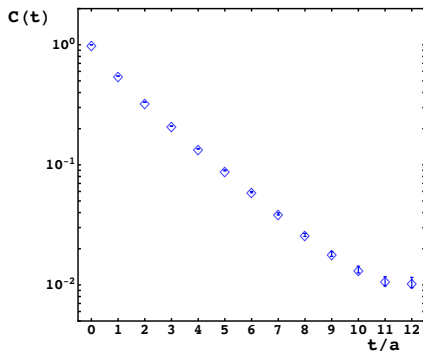
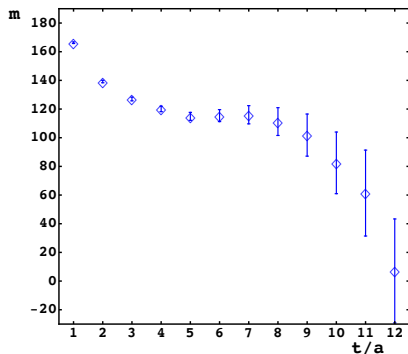
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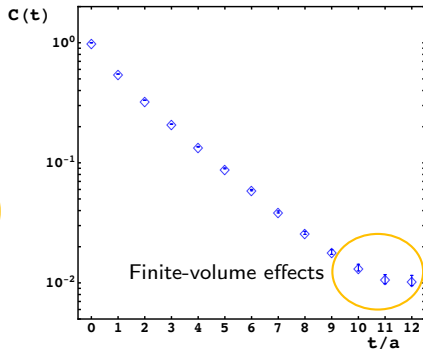
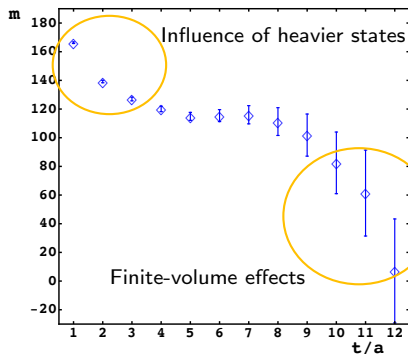
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[Maas, MPL **A28** (2013)]

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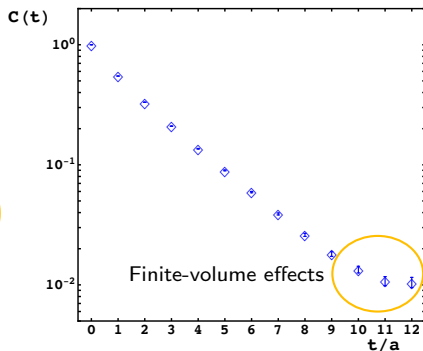
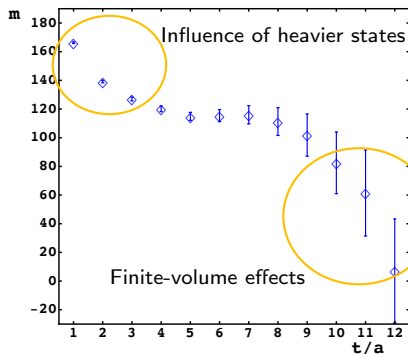
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[Maas, MPL **A28** (2013)]

- Mass is about 120 GeV: Same as Higgs

# Gauge-invariant perturbation theory

[Fröhlich *et al.*, PL **B97** (1980) and NP **B190** (1981) /

Törek and Maas, (LATTICE2016) 1610.04188]

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Bound state mass  $\rightarrow$   $\langle \mathcal{O}(x) \mathcal{O}^\dagger(y) \rangle$

Higgs mass  $\rightarrow$   $\langle \text{Re}[v^\dagger \eta](x) \text{Re}[v^\dagger \eta](y) \rangle_{\text{tl}}$

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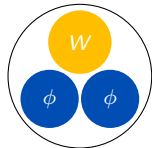
Bound state mass

Higgs mass

2 x Higgs mass scattering state

- Compare poles on both sides

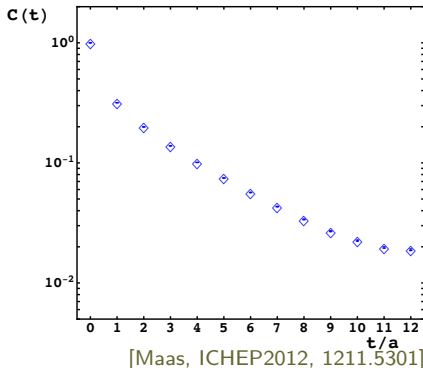
# Mass spectrum: Vector state



- Vector state  $1^-$ :  $\text{tr}[\tau^a \tilde{\phi}^\dagger(x) D_\mu \tilde{\phi}(x)]$
- $\tau^a$  generators of custodial group and  $\tilde{\phi} = \begin{pmatrix} \phi_1 & -\phi_2^* \\ \phi_2 & \phi_1^* \end{pmatrix}$
- Custodial triplet instead of gauge triplet

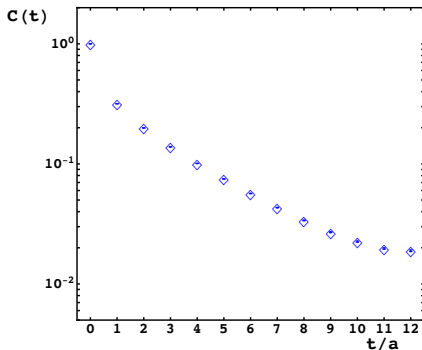
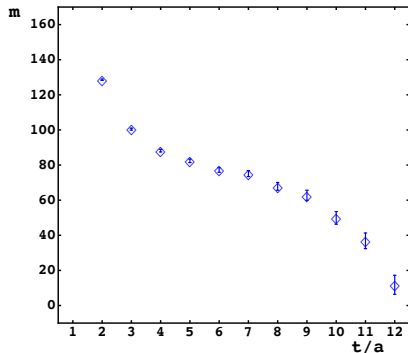
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[Maas, ICHEP2012, 1211.5301]

- Mass is about 80 GeV



# FMS mechanism for W

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- W at tree level: 80 GeV

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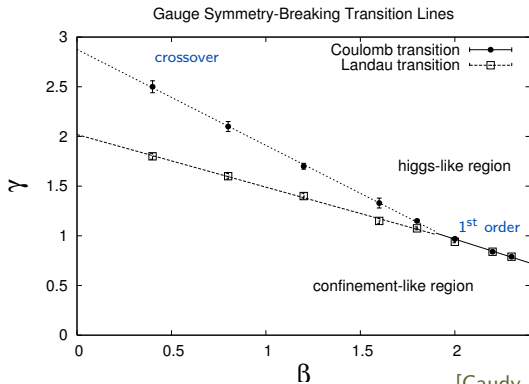
$$\langle \mathcal{O}_\mu^a(x) \mathcal{O}_\mu^{a\dagger}(y) \rangle \stackrel{\tilde{\phi}=v+\tilde{\eta}}{=} c + v^4 \langle W_\mu^a(x) W_\mu^a(y) \rangle + O(W\phi)$$

- Same poles to leading order

- Exchange of a gauge for a custodial triplet

# Phase diagram: $SU(2)$ gauge-Higgs model

- Depending on parameters (inverse gauge coupling  $\beta$ , classical Higgs mass  $\gamma$ ) different regions

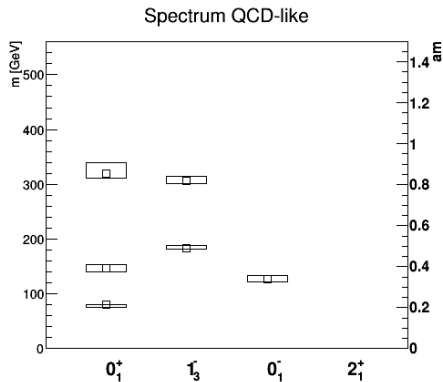
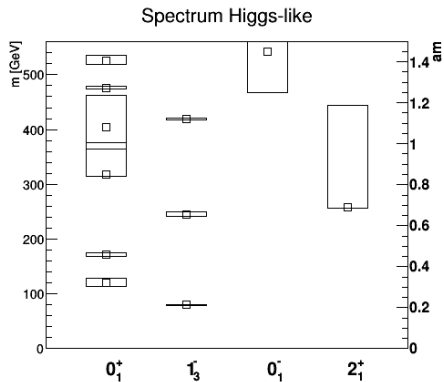


[Caudy and Greensite, PR **D78** (2008)]

- No unique transition line (depends on gauge)
- No phase transition in this model

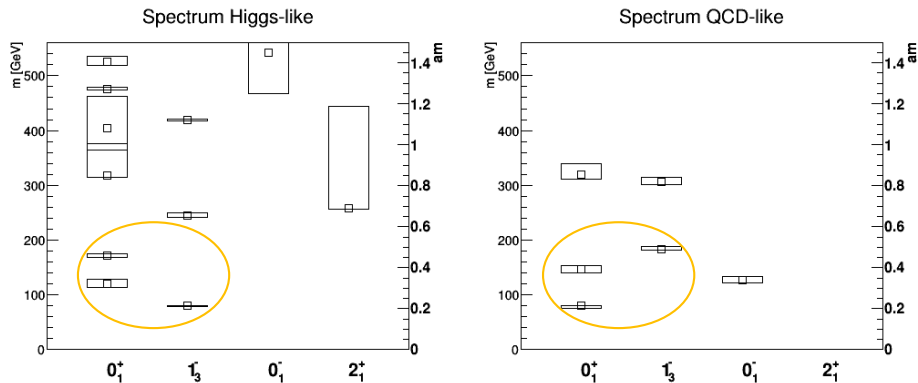
[Osterwalder and Seiler, AP**110** (1978); Fradkin and Shenker, PR **D19** (1979) ]

# Typical spectra



[Maas, MPL **A28** (2013) / Maas and Mufti, JHEP (2014) / Evertz *et al.*, PLB**171** (1986) /  
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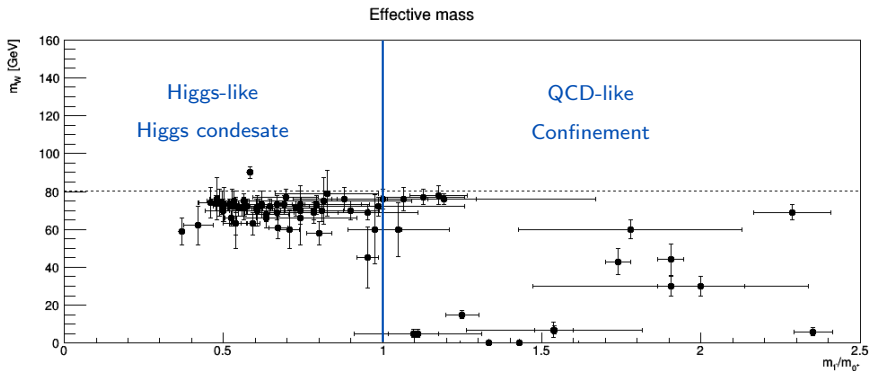


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■  $1^-$  lighter in Higgs-like region

■  $0^+$  lighter in QCD-like region

# Limits



[Maas and Mufti, JHEP (2014)]

- At  $m_{1-} = m_{0+}$  FMS stops working
- So does BEH effect

# Rest of the standard model

[Fröhlich *et al.*, PL **B97** (1980) and NP **B190** (1981) / Egger *et al.*, 1701.02881]

## ■ Quarks and gluons

- Bound by confinement in bound states
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(enormous mass defect)
- Except for right-handed neutrinos

$$\mathcal{O} = \tilde{\phi}^\dagger \begin{pmatrix} \nu \\ e \end{pmatrix} = \begin{pmatrix} \phi_2 \nu - \phi_1 e \\ \phi_1^* \nu + \phi_2^* e \end{pmatrix} \underset{=}{\phi_i = v \delta_{i,2} + \eta_i} \nu \begin{pmatrix} \nu \\ e \end{pmatrix} + O(\eta)$$

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## ■ Photons

- Can also be included

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- Has to be checked for BSM theories

# Physical spectrum in a grand-unified setting

[Törek and Maas, **PRD95** (2017), 1607.05860]

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- Perturbative construction:  $SU(3) \xrightarrow{\langle \phi \rangle} SU(2)$

- Perturbative spectrum:
  - 4 + 1 massive and 3 massless gauge bosons
  - 1 massive Higgs boson

$J^P = 1^-$  singlet channel

- Conflict expected in vector channel

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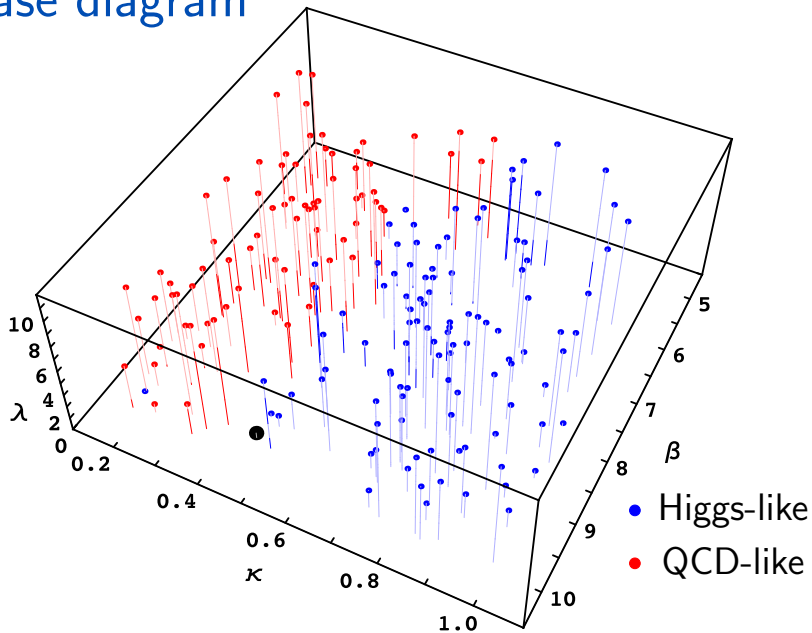
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- Only a single massive particle is predicted

$\Rightarrow$  Contradiction to perturbative spectrum

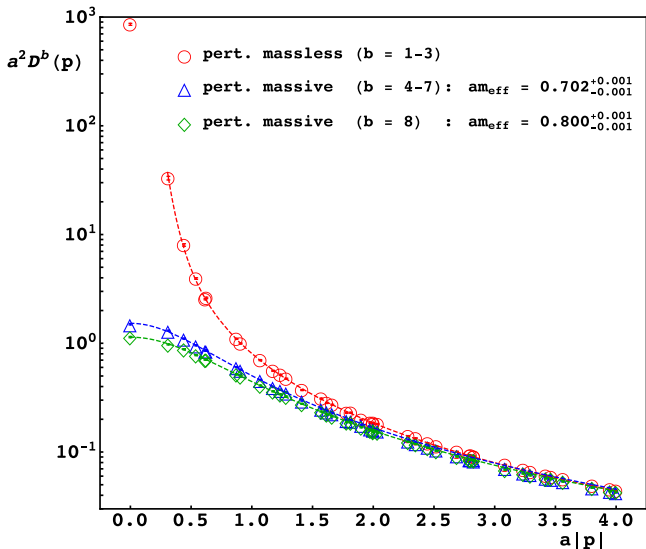


# Phase diagram

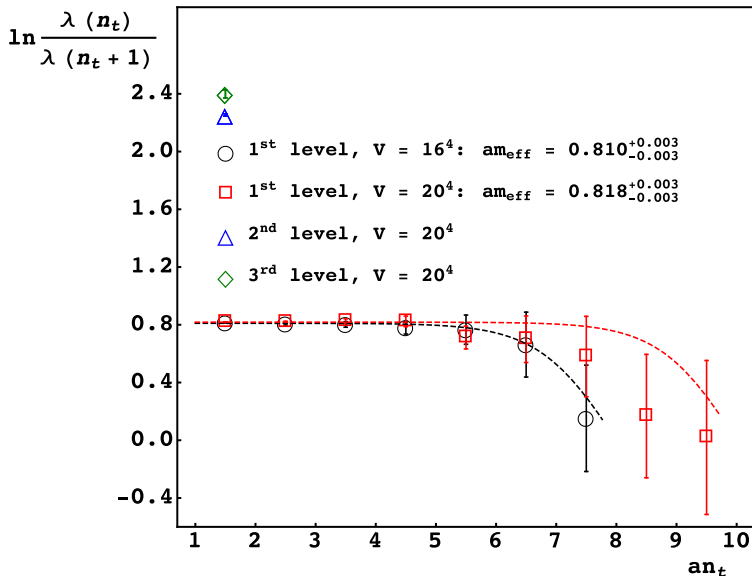


# Propagators

- Good agreement with tree-level perturbation theory

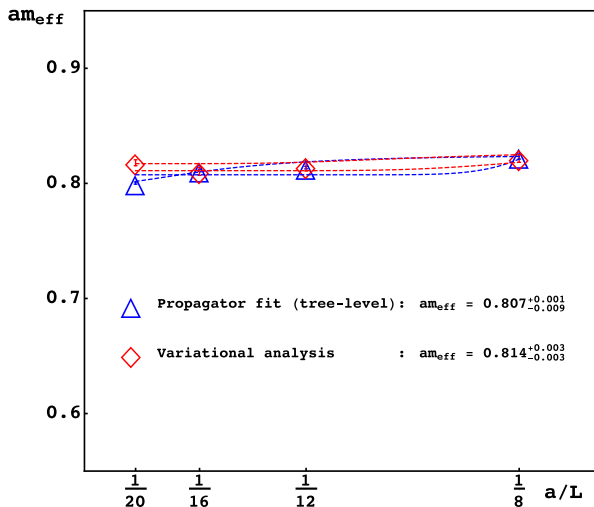


# Spectroscopy in $1^-$ channel - Results



# Volume dependency of $m_{\text{eff}}$

- Single massive ground state with mass of  $W^8$
- Exactly like FMS mechanism predicts



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  - Same light d.o.f. as in standard model: Good candidate

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[Maas, Sondenheimer and Törek, work in progress]

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- More realistic GUT:  $SU(5)$   
Higgs in adjoint and fundamental representation

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  - Yields same results for standard model
  - Mostly not much more complicated
- Applicable to BSM theories
  - Structural requirement: Multiplets must match
  - Dynamical requirement: Small fluctuations
  - Verification requires non-perturbative methods

**Thank you!**