Observable Spectrum in theories with a Brout-Englert-Higgs effect

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- Brout-Englert-Higgs Physics
	- The issue of gauge invariance

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- Describing the physical spectrum

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	- The issue of gauge invariance
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- Fröhlich-Morchio-Strocchi mechanism
- Reformulations and effective theories
- Impact beyond the Standard Model
	- Qualitative changes
	- Changed phenomenology

The issue

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Clashing formal theory with effective phenomenology

A toy model

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-

• Consider an SU(2) with a fundamental scalar

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- Essentially the standard model Higgs

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L = -\frac{1}{4} W_{\mu\nu}^a W_{\nu}^{\mu\nu}
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W_{\mu\nu}^a = \partial_{\mu} W_{\nu}^a - \partial_{\nu} W_{\mu}^a + gf_{bc}^a W_{\mu}^b W_{\nu}^c
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- Coupling q and some numbers f^{abc}

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- Ws W^a_μ *W*
- Higgs h_i (h
- Coupling *g* and some numbers *f abc* and *t a ij*

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L = -\frac{1}{4} W^a_{\mu\nu} W^a_{a} + (D^{\mathbf{i}j}_{\mu} h^j)^+ D^{\mu}_{ik} h_k + \lambda (h^a h_a^+ - v^2)^2
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- Couplings *g, v, λ* and some numbers *f abc* and *t a ij*
- Parameters selected for a BEH effect

A toy model: Symmetries

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• Local SU(2) gauge symmetry W^a_μ \int_{μ}^{a} → W_{μ}^{a} + $(\delta_{b}^{a} \partial_{\mu} - g f_{bc}^{a} W_{\mu}^{c})$ \int_a^c) ϕ ^b $h_i \rightarrow h_i + g t_a^{ij} \phi^a h_j$

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- Global SU(2) custodial (flavor) symmetry
	- Acts as (right-)transformation on the scalar field only $W^a_\mu \rightarrow W^a_\mu$ *^a h*→*h*Ω

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- Choose a suitable gauge and obtain 'spontaneous gauge symmetry breaking': $SU(2) \rightarrow 1$
- Get masses and degeneracies at treelevel
- Perform perturbation theory

Physical spectrum

Perturbation theory

 \bigcap

Physical spectrum

Perturbation theory Scalar fixed charge

 \bigcirc Custodial singlet

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Physical spectrum

 \bigcirc Both custodial singlets

[Fröhlich et al.'80, Banks et al.'79]

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- Physics has to be expressed in terms of manifestly gauge-invariant quantities
	- And this includes non-perturbative aspects...
	- ...even at weak coupling [Gribov'78,Singer'78,Fujikawa'82]

Physical states

[Fröhlich et al.'80, Banks et al.'79]

• Need physical, gauge-invariant particles

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- Has nothing to do with weak coupling
	- Think QED (hydrogen atom!)
- Can this matter?

Remember: Experiment tells that somehow the left is correct!

Experiment tells that somehow the left is correct Theory say the right is correct

There must exist a relation that both are correct

Physical particles Maas & Törek'16,'18, production and properties and properties and all'80,'81, property and property and property and property and property and property $\frac{1}{2}$

[Fröhlich et al.'80,'81,

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	- Bound state structure non-perturbative methods! - Lattice
		- Standard lattice spectroscopy problem
		- Standard methods
			- Smearing, variational analysis, systematic error analysis etc.
		- Very large statistics ($>$ 10⁵ configurations)

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Gauge-invariant

Scalar singlet

$\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf$ Both custodial singlets

$$
h(x)^{+} h(x) \quad (
$$

Custodial singlet $\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf$ Both custodial singlets

 $\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf$

$$
tr t^{a} \frac{h^{+}}{\sqrt{h^{+} h}} D_{\mu} \frac{h}{\sqrt{h^{+} h}}
$$

$$
tr \frac{\omega h^+}{\sqrt{h^+ h}} D_\mu \frac{h}{\sqrt{h^+ h}}
$$

Custodial singlet Triplet $\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf$ Both custodial singlets

Custodial singlet Triplet $\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf\bf$ Both custodial singlets

Why?

A microscopic origin - Fröhlich-Morchio-Strocchi mechanism

How to make predictions

[Fröhlich et al.'80,'81, Maas & Törek'16,'18, Maas, Sondenheimer & Törek'17 Maas & Sondenheimer '20]

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- Formulate gauge-invariant, composite operators
	- Bound state structure non-perturbative methods?
	- But coupling is still weak and there is a BEH
	- Perform double expansion [Fröhlich et al.'80, Maas'12]
		- Vacuum expectation value (FMS mechanism)
		- Standard expansion in couplings
		- Together: Augmented perturbation theory

[Fröhlich et al.'80,'81 Maas'12,'17]

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Higgs field

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+ v \langle \eta^* \eta^2 + \eta^*^2 \eta \rangle + \langle \eta^* \eta^2 \rangle

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[Fröhlich et al.'80,'81 Maas'12,'17]

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state \rightarrow (*h*⁺ *h*)(*x*)(*h*⁺ *h*)(*y*))=*v*²(η⁺ (*x*)η(*y*)) $+\langle \eta^+(x)\eta(y)\rangle \langle \eta^+(x)\eta(y)\rangle + O(g,\lambda)$ Bound mass

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tandard Perturbation Theory

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\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+(x) \eta(y) \rangle
$$
\nWhat about $\sqrt{\frac{+v \langle \eta^+ \eta^2 + \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle}{+v \langle \eta^+ \eta^{+2} \eta \rangle + \langle \eta^{+2} \eta^2 \rangle}}$, this? \rightarrow Duiife van Egmond's talk $\langle (h^+ h)(x)(h^+ h)(y) \rangle = v^2 \langle \eta^+(x) \eta(y) \rangle$

What about the vector?
Massilal About the vector?

Maas'12]

1) Formulate gauge-invariant operator 1 triplet: $\langle (\tau^i h^+ D_\mu h)(x) (\tau^j h^+ D_\mu h)(y) \rangle$

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	- $\langle (\tau^i h^{\; +} D_\mu h)(x) (\tau^j h^{\; +} D_\mu h)(y) \rangle = v^2 c_{ij}^{ab} \langle W_\mu^a(x) W^b(y)^\mu \rangle + ...$

1) Formulate gauge-invariant operator

1-triplet:
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Matrix from group structure

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Matrix from group structure

c projects custodial states to gauge states

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= $v^2 \langle W_\mu^i W_\mu^j \rangle + ...$
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Exactly one gauge boson for every physical state

Phenomenological Implications

Can we measure this?

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-
-
-
- -
	-
- -

• Two possibilities to measure extension

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	- Form factor
		- Difficult
			- Higgs and Z need to be both produced in the same process

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		- Standard vector boson scattering process at low energies
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- Elastic region: 160/180*GeV* ⩽√*s*⩽250*GeV*
	- s is the CMS energy in the initial/final ZZ/WW system

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\nScattering length-"size" Phase shift

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\n
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a_0 = \tan\left(\delta_J\right) / \sqrt{s - 4m_W^2}
$$
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\nPlase shift

\nAlattice Lüscher analysis

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	- Parameters slightly different
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Generic behavior

Elimination

-

What are the true degrees of freedom?

$$
L = -\frac{1}{4} W_{\mu\nu}^a W_a^{\mu\nu} + (D_\mu^{ij} h^j)^+ D_{ik}^\mu h_k + \lambda (h^a h_a^+ - v^2)^2
$$

$$
W_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + gf_{bc}^a W_\mu^b W_\nu^c
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+ $(D_{\mu}^{ij} h)^+ D_{ij}^{\mu} h + \lambda (h^2 - v^2)^2 + 3 \ln(h)$
 $Z_{\mu\nu}^a = \partial_{\mu} Z_{\nu}^a - \partial_{\nu} Z_{\mu}^a + gf_{bc}^a Z_{\mu}^b Z_{\nu}^c$
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Global SU(2) symmetry

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• Perform a variable transformation *L*=− 1 4 Z^a_μ _v $Z^{\mu \nu}_a$ + *m*³ Z^a_μ Z^{μ}_a $+(\frac{D_{\mu}^{ij}h}{\partial t}) + D_{ij}^{\mu}h + \lambda(h^2 - v^2)^2 + 3\ln(h)$ $Z_{\mu\nu}^{a} = \partial_{\mu} Z_{\nu}^{a} - \partial_{\nu} Z_{\mu}^{a} + gf_{bc}^{a} Z_{\mu}^{b} Z_{\nu}^{c}$ D_{μ}^{ij} $=$ $\delta^{ij}\partial_{\mu}$ igZ_{μ}^a t_a^{ij} *ij* Massive Vector field Global SU(2) symmetry $h = \sqrt{\phi^+ \phi}$ $Z^a = tr \tau^a$ ϕ ⁺ $D\phi$

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Non-trivial tree-level structure defects or large λ

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Well-defined theory, can be simulated on the lattice

[Jersak et al.'85, Evertz et al.'86]

The structure of the particles

- Elementary particles appear point-like
	- But equivalent theories: At the quantum level extended objects!
	- Others: Looks like having a substructure
		- Peaks at 1/2, 1/3, 1/4,...

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	- But not in perturbation theory
- Properties of particles do not uniquely determine if elementary or composites of a gauge theory

Generalizing

• Rewriting of a gauge theory as an ungauged theory

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[Berghofer et al.'21]

- Possible for QED
	- Including the Aharanov-Bohm effect [Strocchi et al.'74]
- Yang-Mills theory induces an infinite number of variables, Wilson loops of all sizes [Gambini et al. '96]
- (Quantum) gravity as dynamical triangulation \rightarrow Talk by Renate Loll [Regge '61, Ambjorn et al. '12]
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	- (Quantum) gravity as dynamical triangulation \rightarrow Talk by Renate Loll [Regge '61, Ambjorn et al. '12]
	- Generalization: Dressing field method \rightarrow Talk by Jordan Francois
- Conversely: Many (all?) ungauged theories can be written explicitly as a gauge theory
	- (Generalized) Kretschmann hypothesis: Always possible [Kretschmann '17, Einstein '18, Kibble '67, Pitts '09, Francois '18]

Physical states

- Need physical, gauge-invariant particles
	- **Cannot** be the elementary particles
	- Non-Abelian nature is relevant
- Need more than one particle: Composite particles
	- Higgs-Higgs, W-W, Higgs-Higgs-W etc.

- Has nothing to do with weak coupling
	- Think QED (hydrogen atom!)
- Can this matter?

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[Fröhlich et al.'80, Egger, Maas, Sondenheimer'17]

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	- Global SU(3) generation
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$$
\left\| \left(\left| h_2 - h_1 \right|_{L} \right| V_L \right\|_{L} (x) + \left\| \left(\left| h_2 - h_1 \right| V_L \right|_{L} \right) (y) \right\|
$$

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	- Different masses for doublet members
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- Yukawa terms break custodial symmetry
	- Different masses for doublet members
- Extends non-trivially to hadrons

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	- Compressed mass scales
	- One generation
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	- Dirac fermions: left/righthanded non-degenerate
	- Quenched

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- Supports FMS prediction

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- Supports FMS prediction grant for unquenching '24-'28

Standard Model

3 Generations

Explicit CP violation

Standard Model

- 1 Generation
- C and P violation

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Composite fermions

Standard Model

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Explicit CP violation

Interesting option **Interesting option**

Standard Model

- 1 Generation
- C and P violation

Composite fermions

+ Assumption: each composite fermion has 3+ internal excitations Standard Model

3 Generations

Explicit CP violation

has 3+ internal excitations

Standard Model 3 Generations Explicit CP violation Low-energy effective theory

CKM matrix

...but no method available (yet) to check this? [Greensite '21]

New physics - Qualitative changes

[Maas'15 Maas, Sondenheimer, Törek'17]

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	- Mapping of custodial symmetry to gauge symmetry
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- Is this generally true?
	- No: Depends on gauge group, representations, and custodial groups
	- Can work sometimes (2HDM, MSSM) [Maas,Pedro'16, Maas,Schreiner'23]
	- Generally qualitative differences
		- Most dramatic consequences: GUTs

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	- No coupling universality for Abelian charges
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- Why are gauge symmetries except the strong are linked to the Higgs?
	- \cdot Gauging of the SU(2)xU(1) subgroup of the O(4) Higgs global symmetry
- Why do all gauge couplings become similarly strong at about the same energy of \sim 10¹⁵ GeV?

Grand-unified theories

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- Could all be emergent features?
- Scenario: Only one gauge interaction
	- Grand unification
	- Standard model low-energy effective theory
- Standard scenario of double breaking
	- Two Brout-Englert-Higgs effect
	- \cdot One breaks at 10^{15} GeV
	- The other at the electroweak scale
	- Requires at least one more Higgs
		- Other particle content scenario-dependent

- Requires leptoquarks
	- Surplus gauge bosons
	- Connects necessarily quarks and leptons
		- One gauge group only!
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- Many (simple) scenarios ruled out...
- ...but many more alive (w or w/o SUSY)

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- **Ws** W^a_μ W
- Coupling q and some numbers f^{abc}

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$$

- Ws W^a_μ *W*
- Higgs h_i (h
- Coupling *g* and some numbers *f abc* and *t a ij*

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- Global U(1) custodial (flavor) symmetry

• Acts as (right-)transformation on the scalar field only $W^a_\mu \rightarrow W^a_\mu$ $h \rightarrow exp(ia)h$

Vector Gauge-dependent

 $'SU(3) \rightarrow SU(2)'$

Confirmed in gauge-fixed lattice calculations [Maas et al.'16]

[Maas & Törek'16,'18 Maas, Sondenheimer & Törek'17]

What about the vector? **Example 2018**

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group structure

c ab projects out only one field

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Only one state remains in the spectrum at mass of gauge boson 8 (heavy singlet)

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- Now: States without elementary analouge
	- Gauge-invariant states from 3 Higgs fields
	- \cdot Baryon analogue U(1) acts as baryon number
	- Lightest must exist and be absolutely stable

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	- All channels: \leq 3
	- Aim: Ground state for each channel
		- Characterization through scattering states

[Dobson et al.'22]

Typical spectrum

PRELIMINARY Typical spectrum [Dobson et al.'22]

Typical spectrum

[Dobson et al.'22 Maas '17 Jenny et al.'22]

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		- Without associated symmetry
		- Not a Goldstone of dilation symmetry breaking, like SM photon

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- Massless state: Autocorrelation bad!

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- FMS: Massless state [Maas et al.'17]

Result External External External External External Extends (Afferrante et al.'20)

No massive states seen yet – but no suitable methods available

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Summary

• Full invariance necessary for physical observables in path integrals

> Review: 1712.04721 Update: 2305.01960 Philosophical aspects: 2110.00616

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• A new perspective on particle physics

Review: 1712.04721 Update: 2305.01960 Philosophical aspects: 2110.00616