Gauge fixing and the Ghost DSE

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  - Gribov-Singer ambiguity
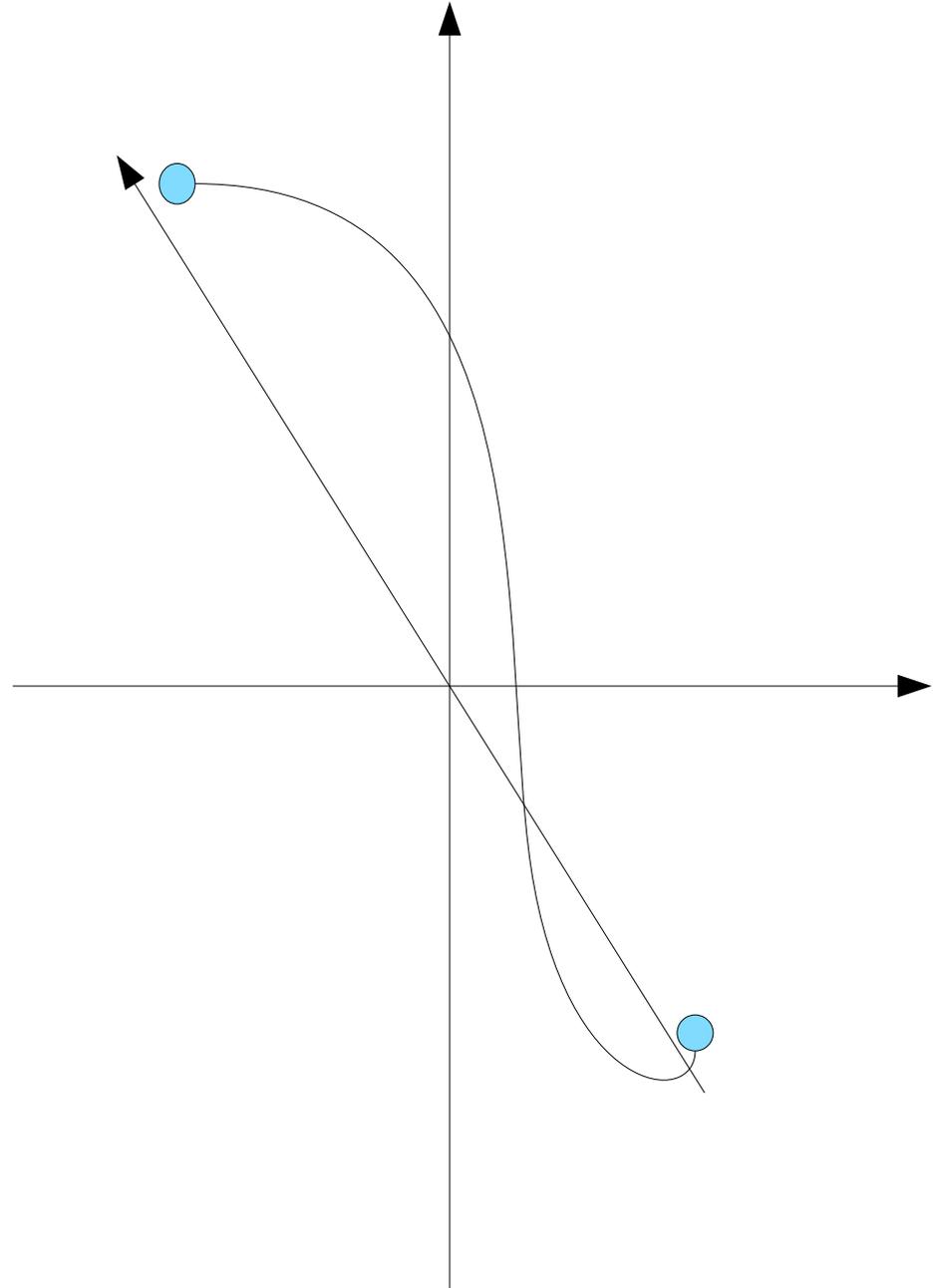
Configuration space (artist's view)
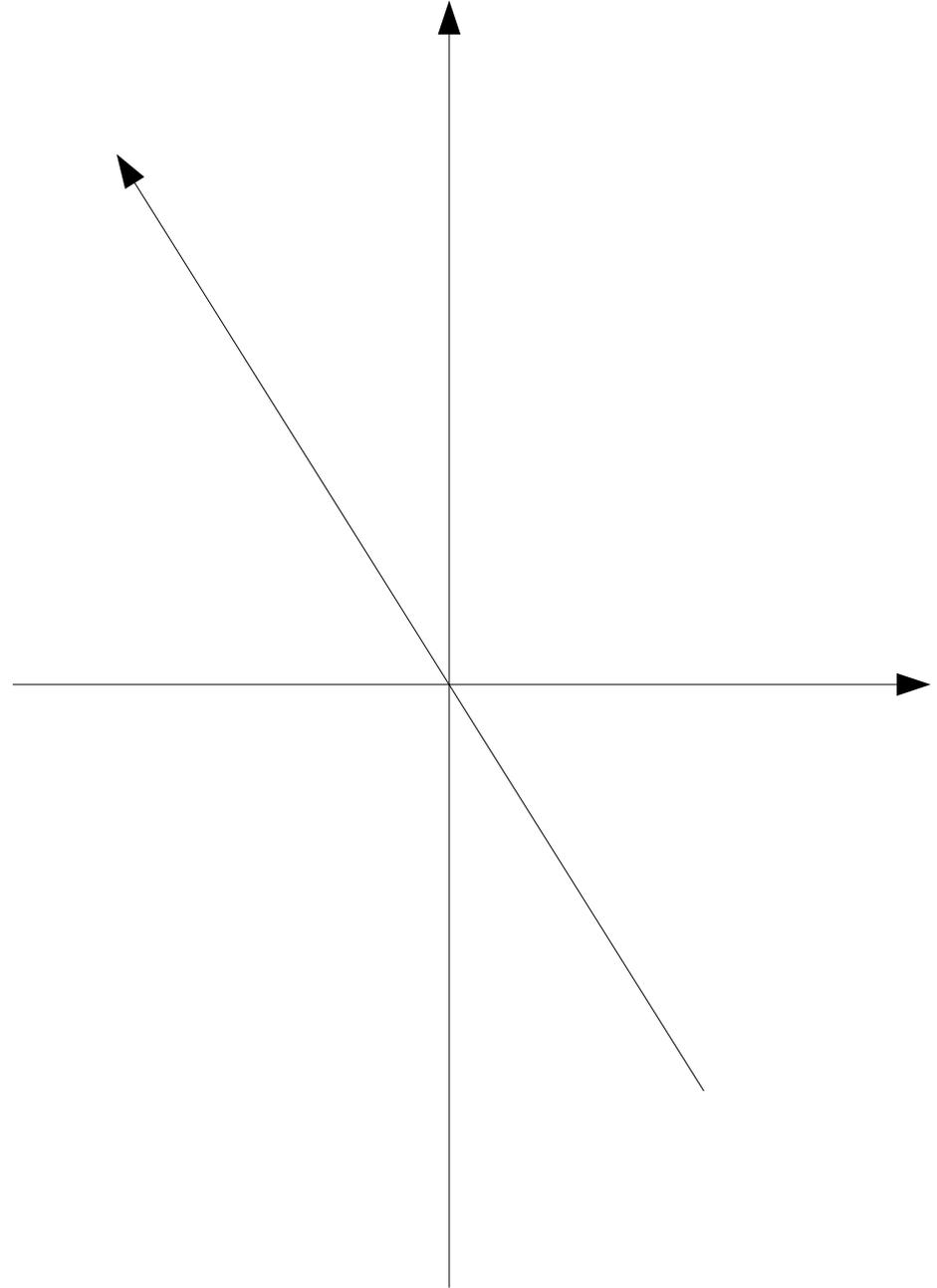
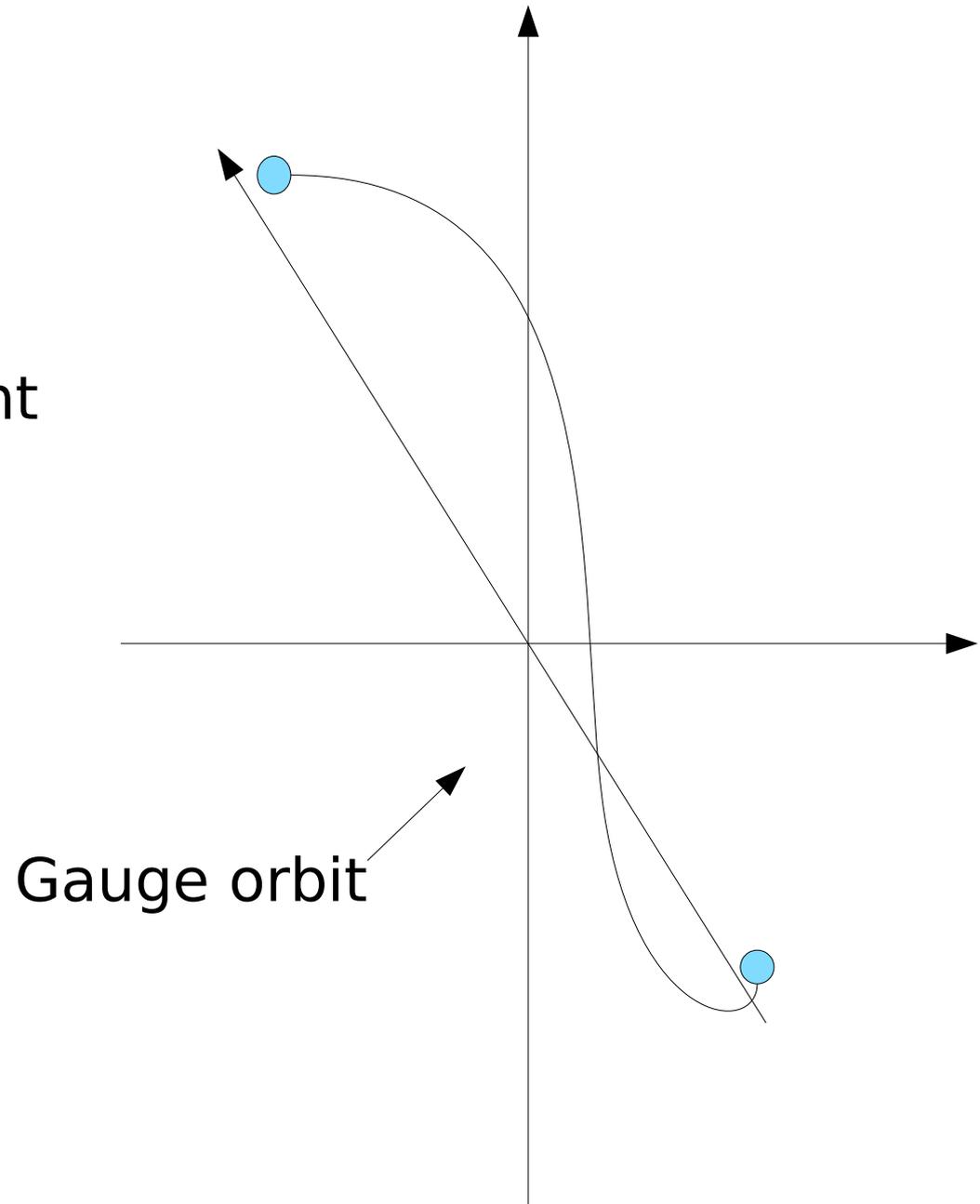# Configuration space (artist's view)
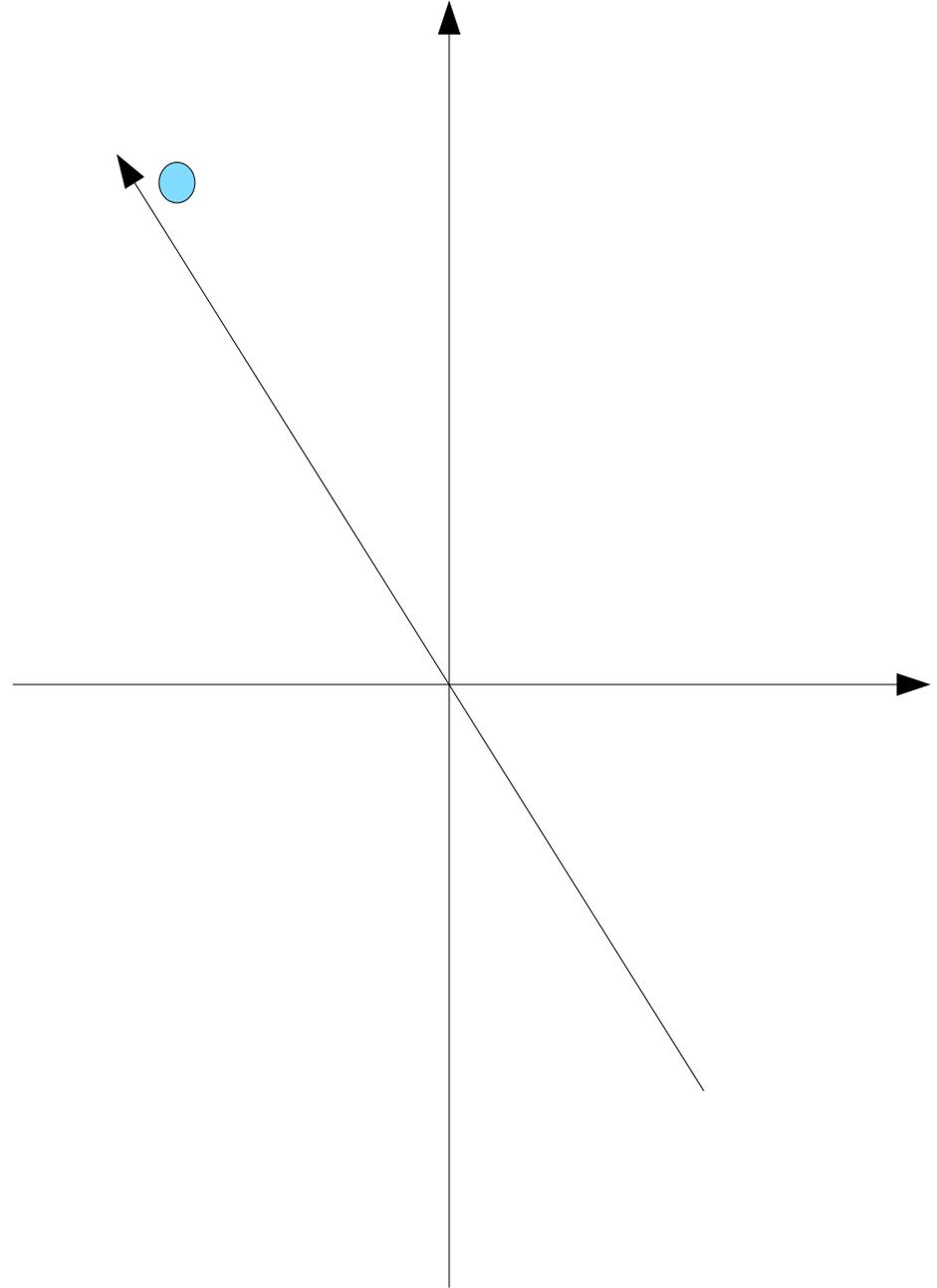
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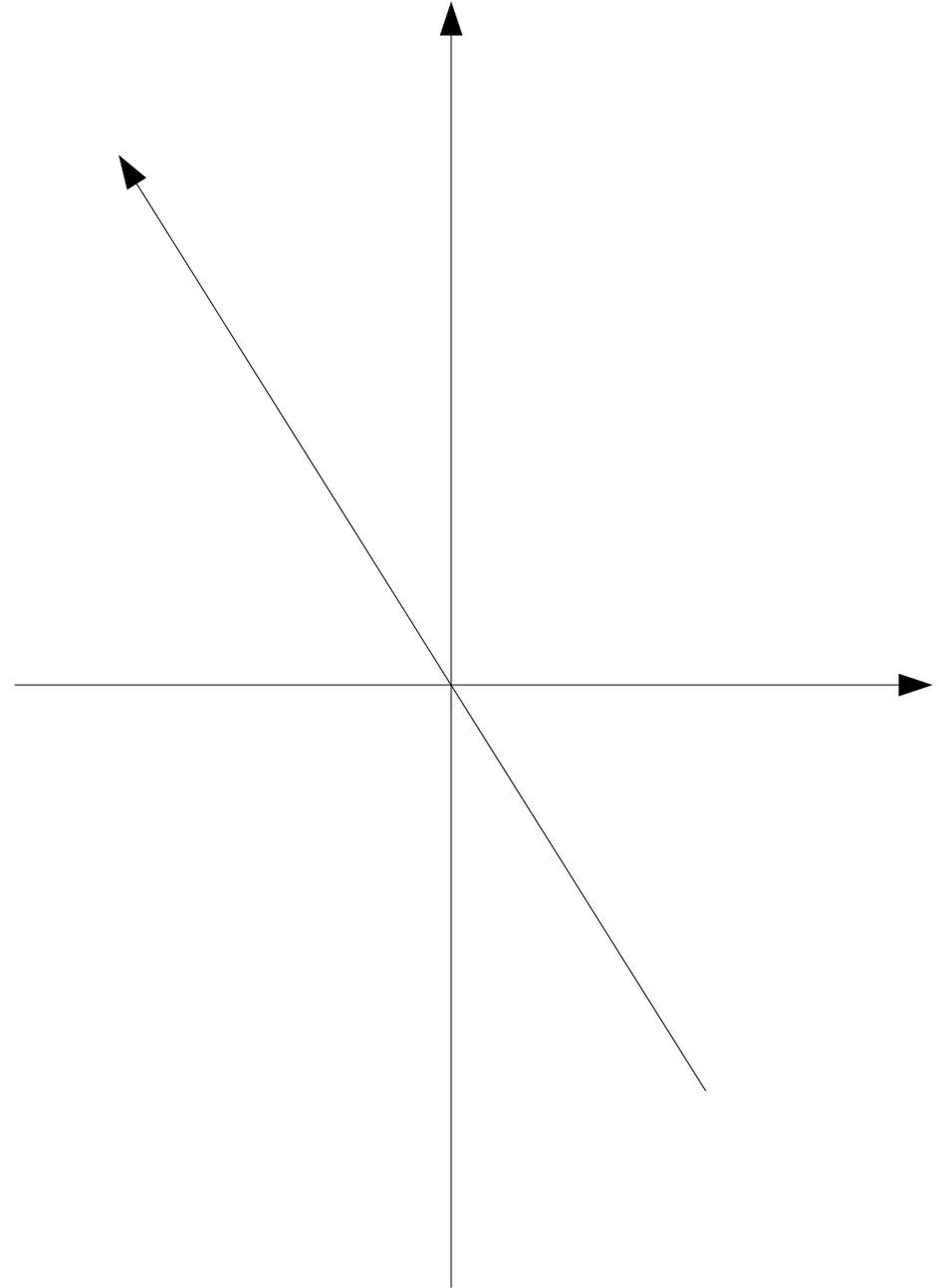 - Set of all gauge-equivalent copies is a gauge orbit

Gauge orbit

Configuration space (artist's view)

- Gauge freedom: Choice of internal coordinate system

- Fields and most correlation functions change under gauge transformations

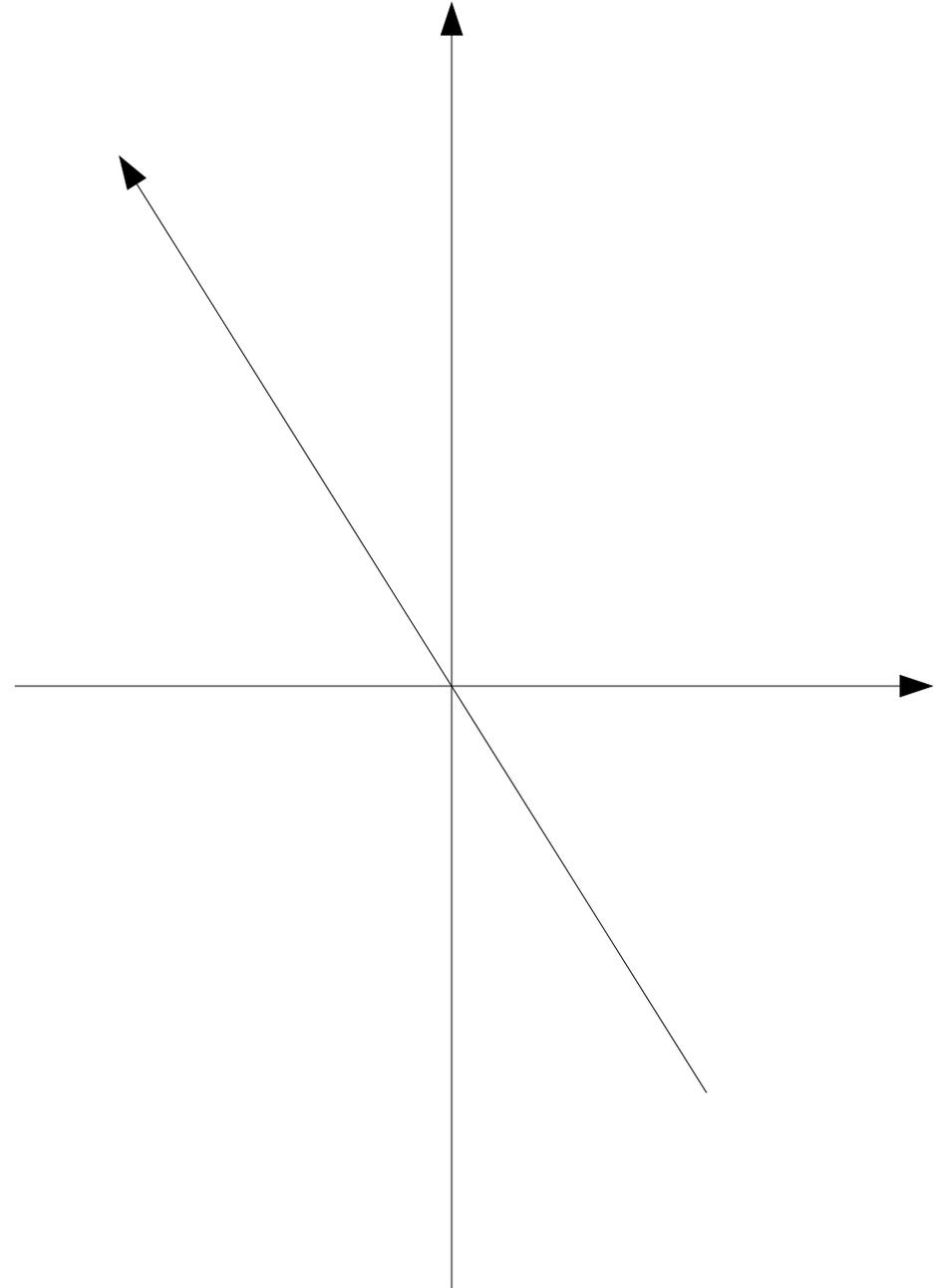- Requires a choice of gauge/coordinate system
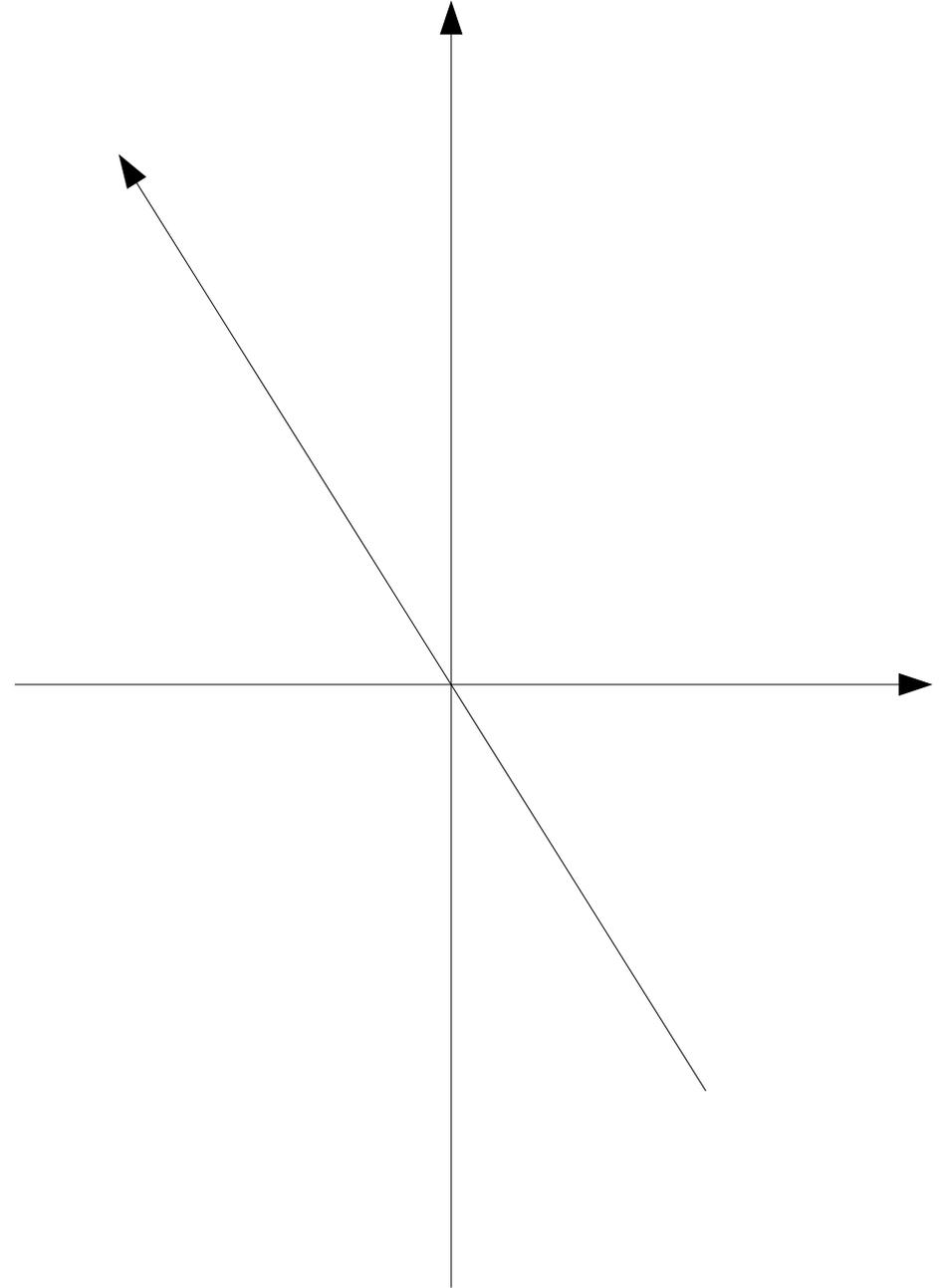
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  - Reduces configuration space to a hypersurface
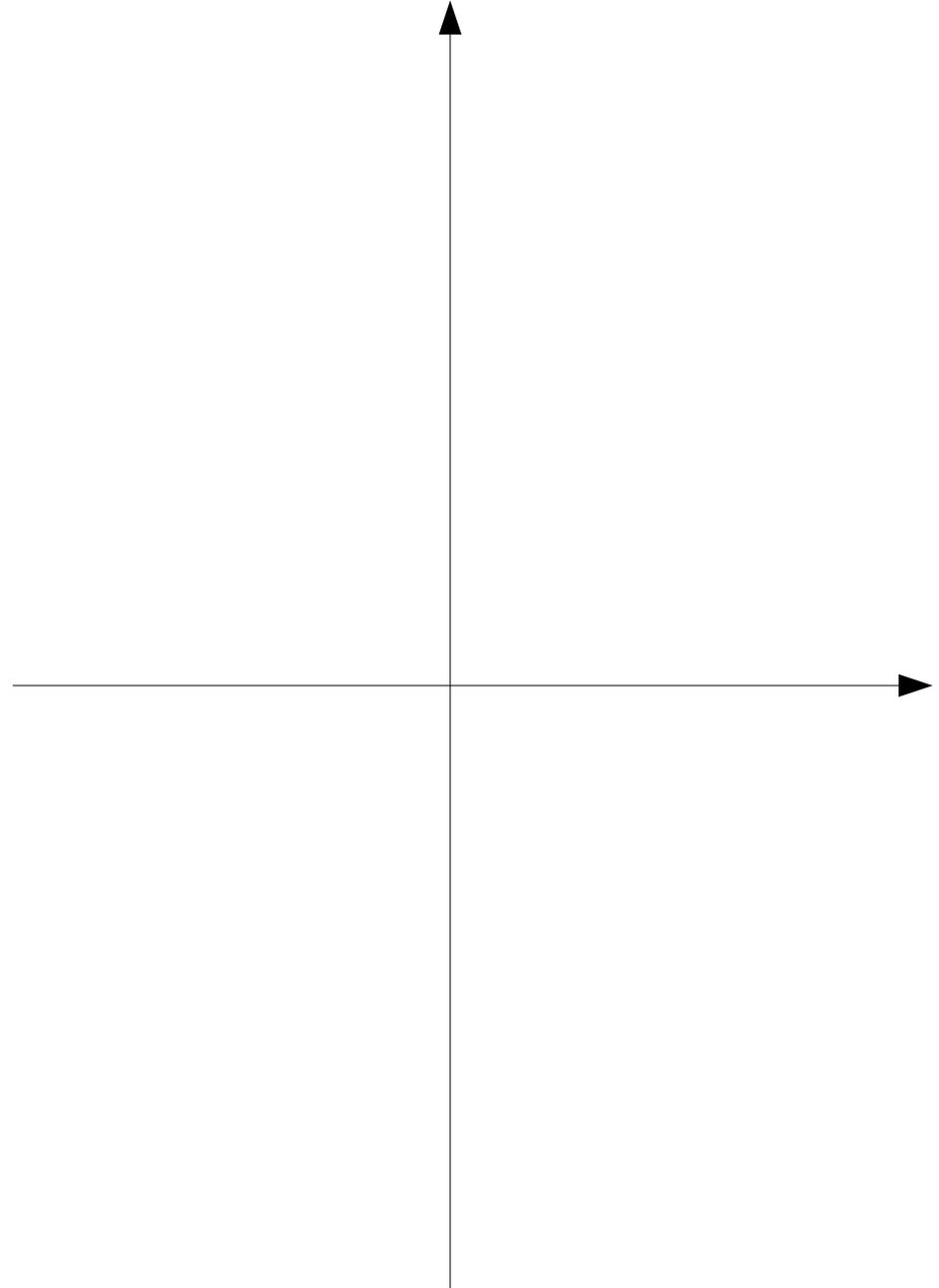
(Perturbative) Landau gauge

- Condition can be implemented using auxiliary fields, the ghost fields

- Lagrangian:

$$L = -\frac{1}{4} F^a_{\mu\nu} F^a_{\mu\nu} - \bar{c}^a \partial_\mu D^{ab}_\mu c^b$$

$$D^{ab}_\mu = \delta^{ab} \partial_\mu - g f^{abc} A^c_\mu$$

$$\lim_{\zeta \to 0} \int DA c \bar{c} \exp\left(\int dx \left(-\frac{1}{4} F^2 - \bar{c} \partial D c - \frac{1}{2\zeta} (\partial A)^2 \right) \right)$$

- Degrees of freedom: Gluons A^a_μ Ghosts $\bar{c}^a, c^a$
  - Not physical objects
  - Pure mathematical convenience

Unique gauge-fixing [Sobreiro & Sorella, 2005]

- Local gauge condition
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Unique gauge-fixing

- Local gauge condition
  - Landau gauge: $\partial_\mu A_\mu^a = 0$
- Sufficient for perturbation theory
- Insufficient beyond perturbation theory
  - There are gauge-equivalent configurations: Gribov copies
- No local gauge conditions known to select a unique gauge copy: Gribov-Singer ambiguity

Residual freedom

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  - Set of Gribov copies, in general not continous
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- Construct non-local conditions
  - Choice: Leave the global color symmetry unfixed
- Method: Lattice

First horizon gauges
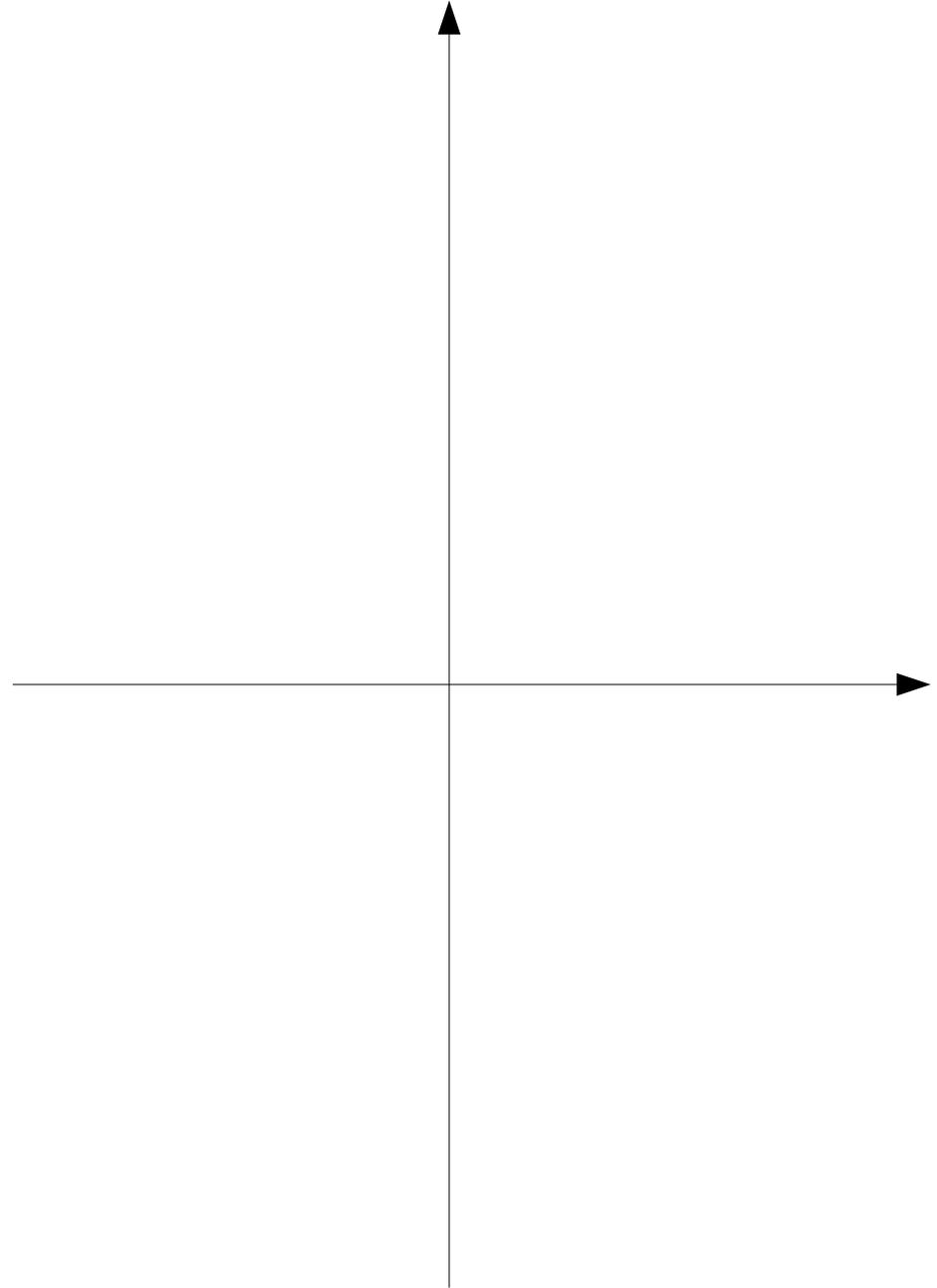
First horizon gauges

- Gribov horizon encloses all field configurations with non-negative Faddeev-Popov operator $-\partial_\mu D_\mu$

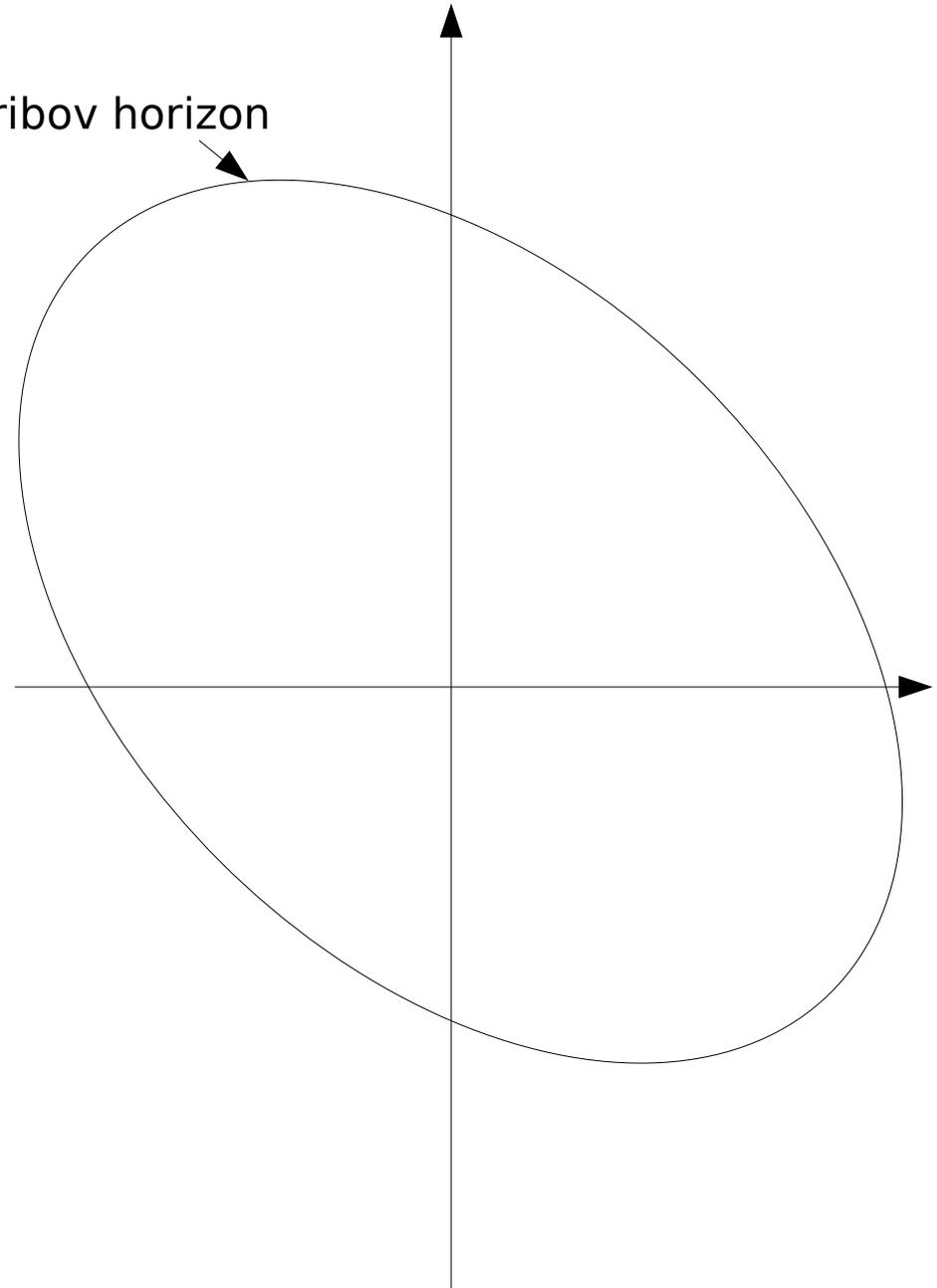- Convex and bounded

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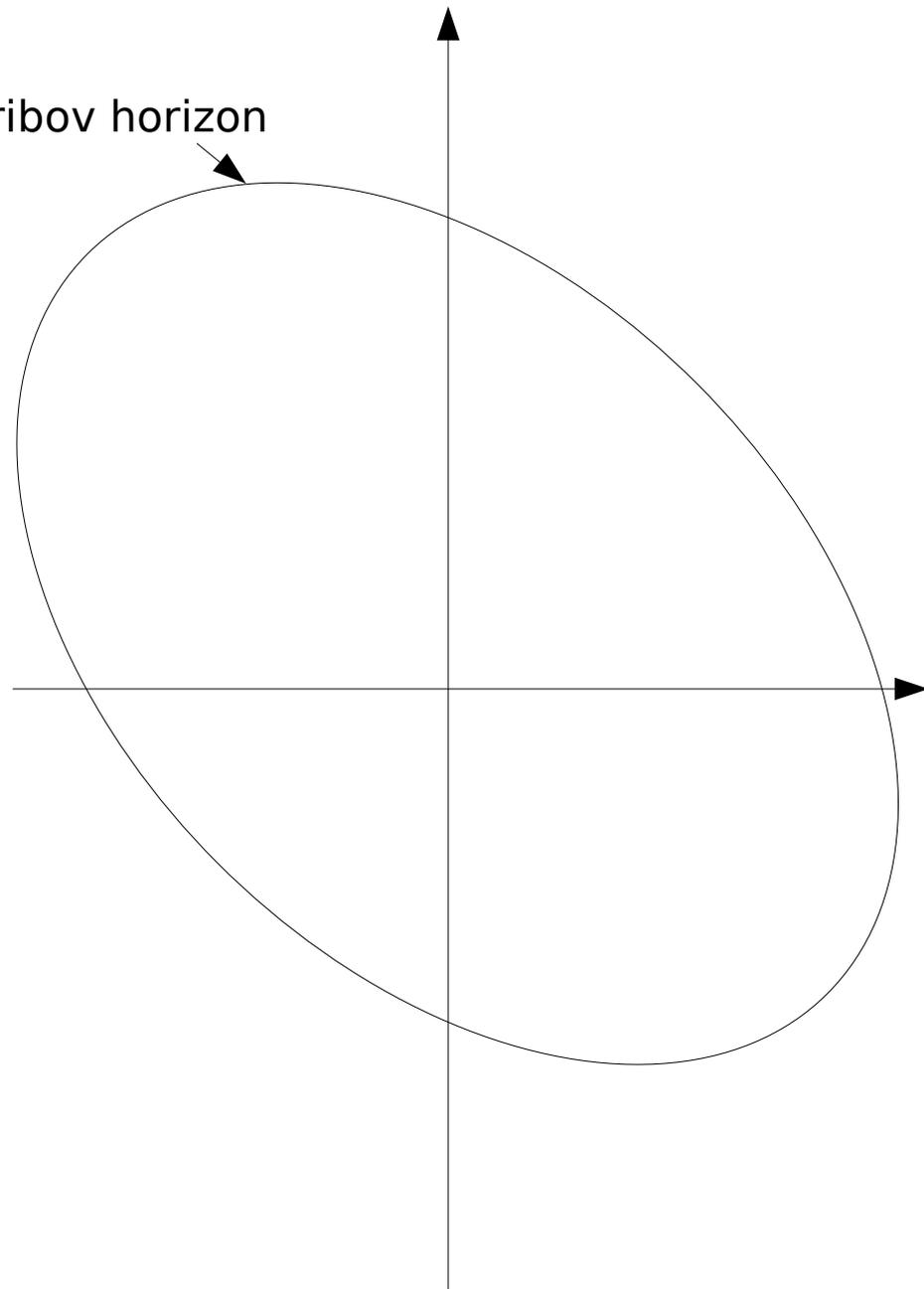
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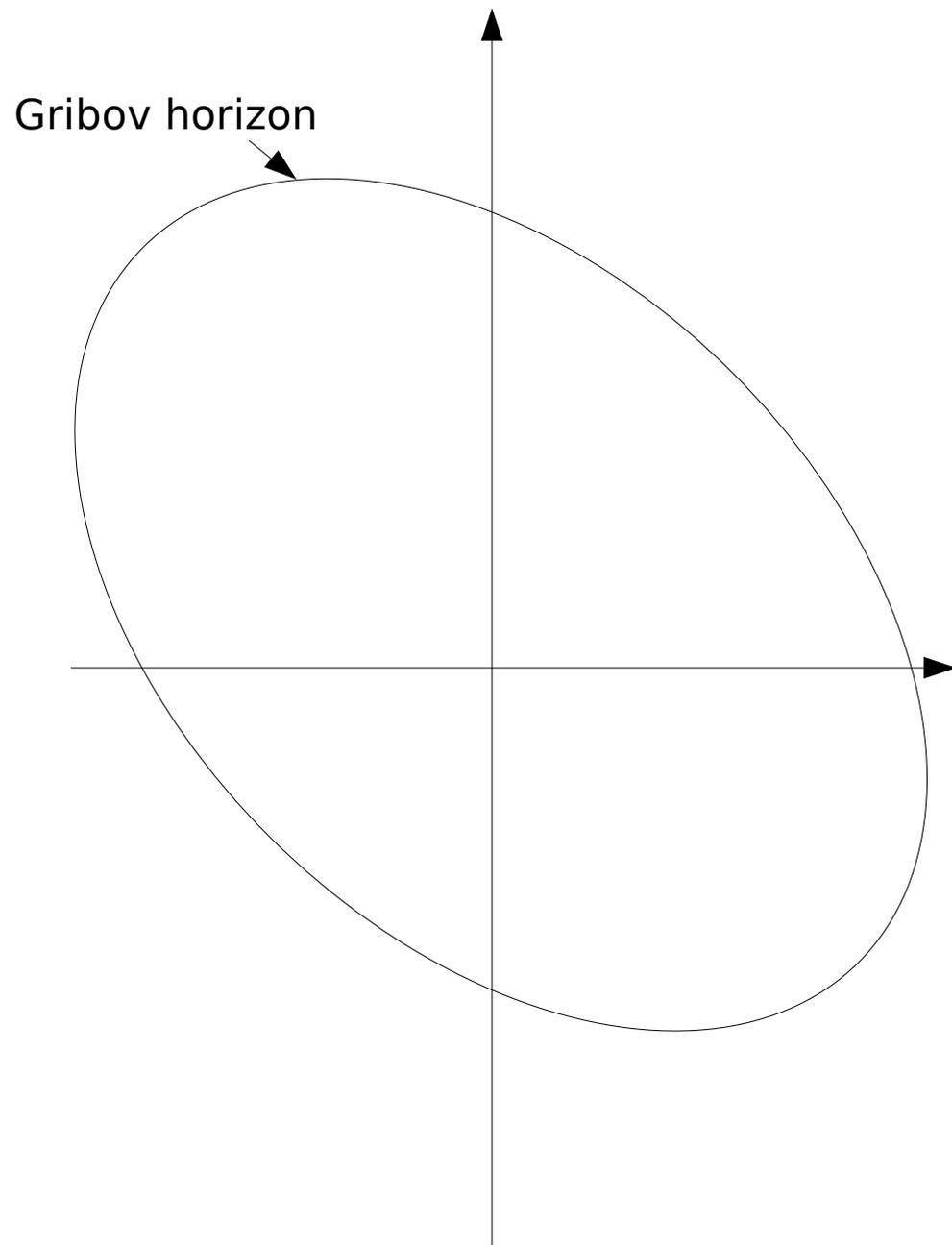- All gauge orbits pass through this region
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  - Literally on the lattice:

Gribov horizon
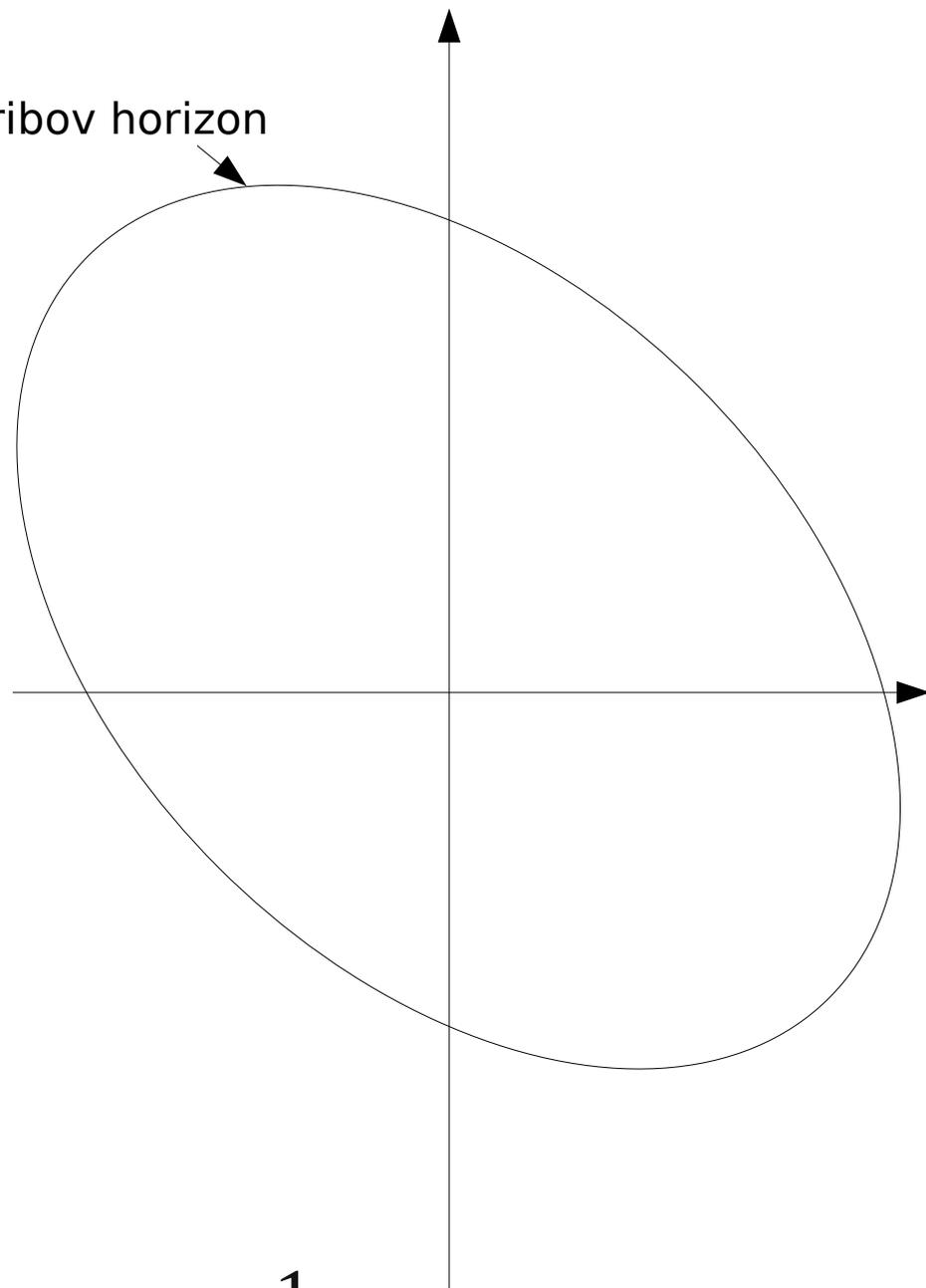
$$\lim_{\zeta \to 0} \int DAc\,\bar{c}\exp\left(\int dx\left(-\frac{1}{4}F^2 - \bar{c}\,\partial D\,c - \frac{1}{2\zeta}(\partial A)^2\right)\right)\Theta(-\partial D)$$

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- Two tests

Test 1: Quantum effective action

- Derivation of functional equations assume the reconstruction theorem
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  - Alternatively: Free energy is an analytic function of the sources

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  - The quantum effective action is an analytical function of the classical fields
  - Alternatively: Free energy is an analytic function of the sources
- Is this correct?
  - Unlikely: Implies necessarily a vanishing gluon propagator at zero momentum on the lattice [Zwanziger'13]
  - At odds with lattice results [Cucchieri et al.'07, von Smekal et al'07]

Lattice results

Lattice results
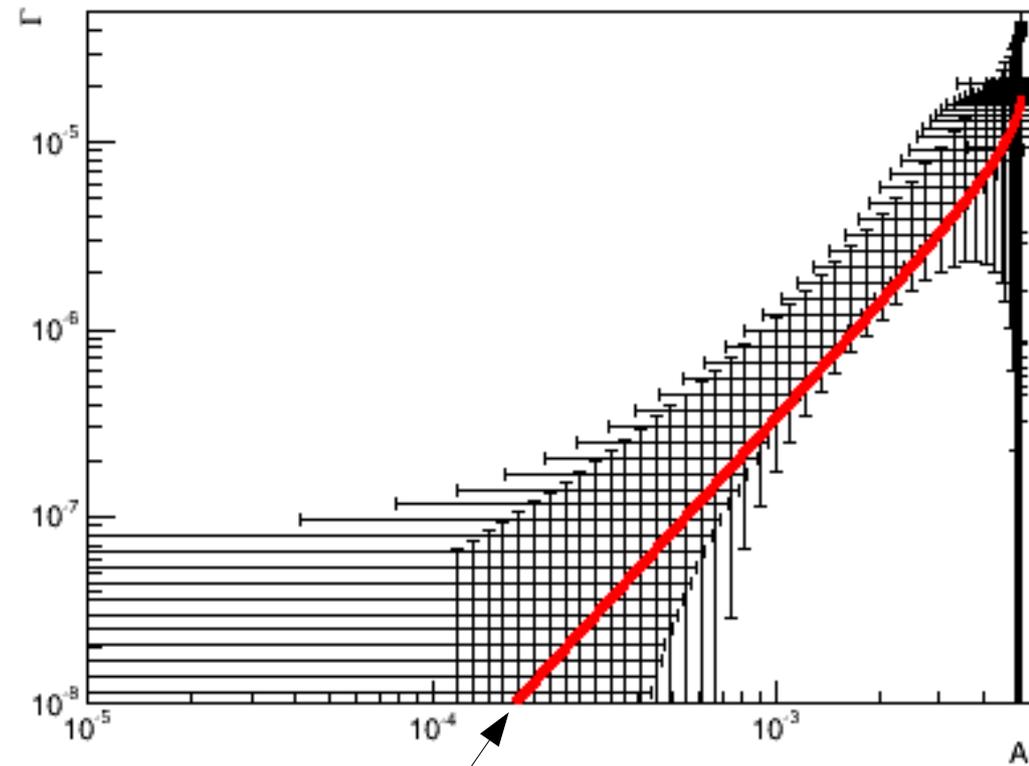
Free energy
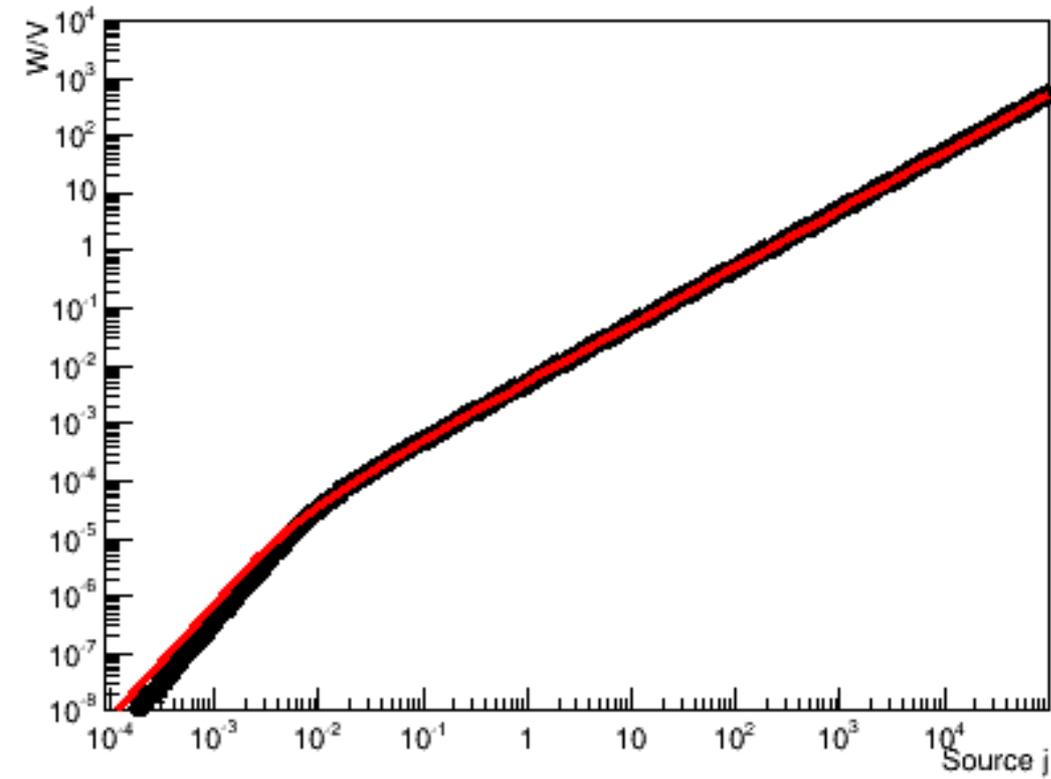
Lattice results

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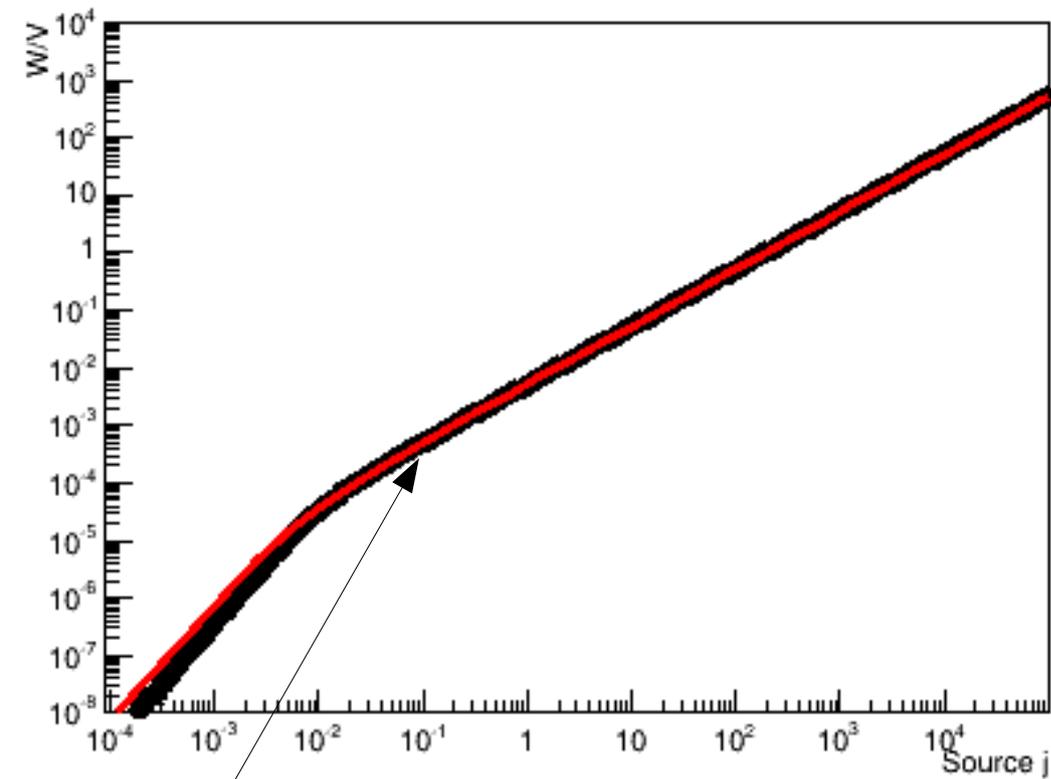
Fit: $c\left(\sqrt{1+(aj)^2}-1\right)$
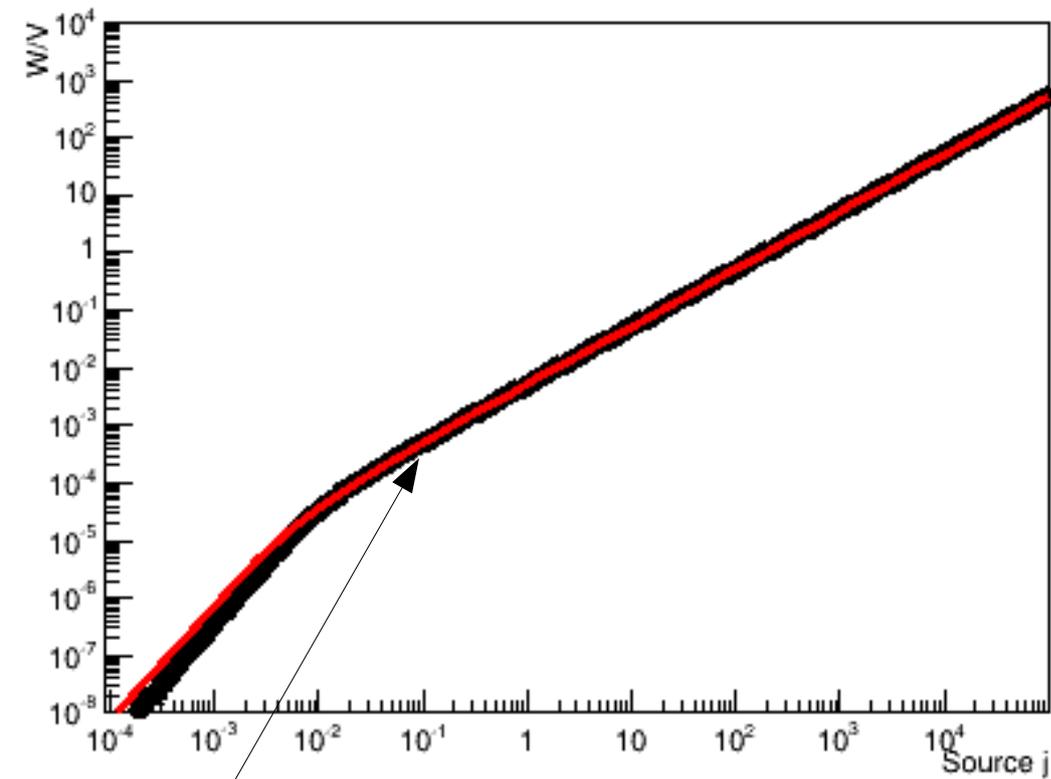
Free energy

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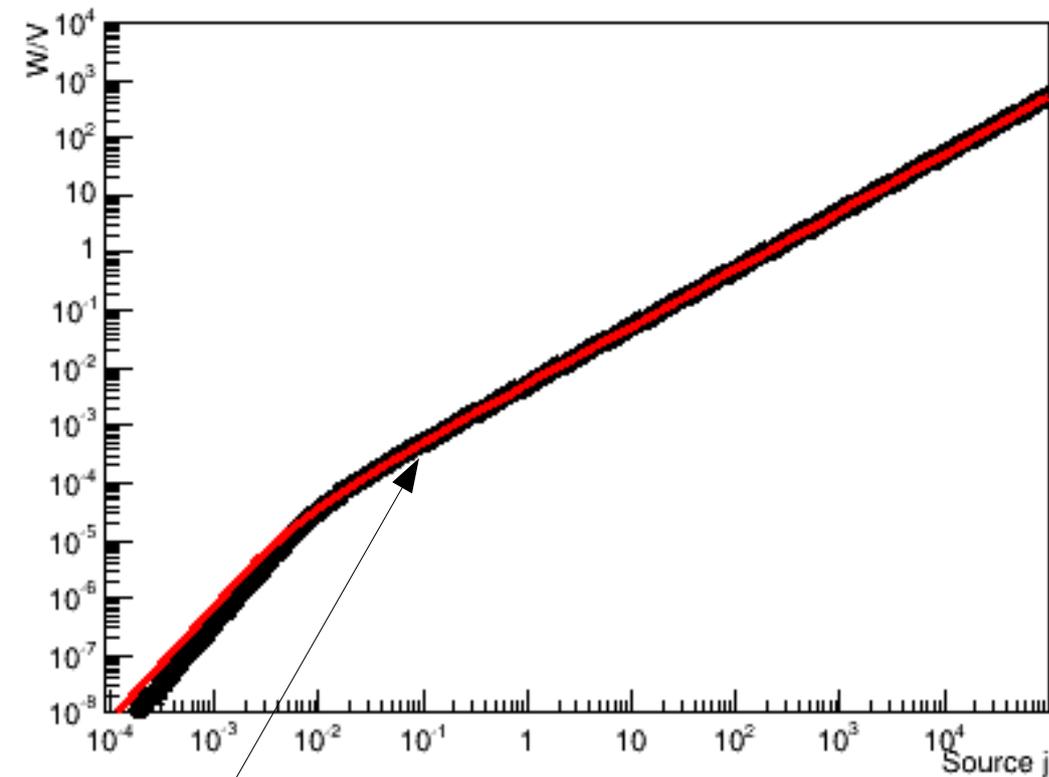
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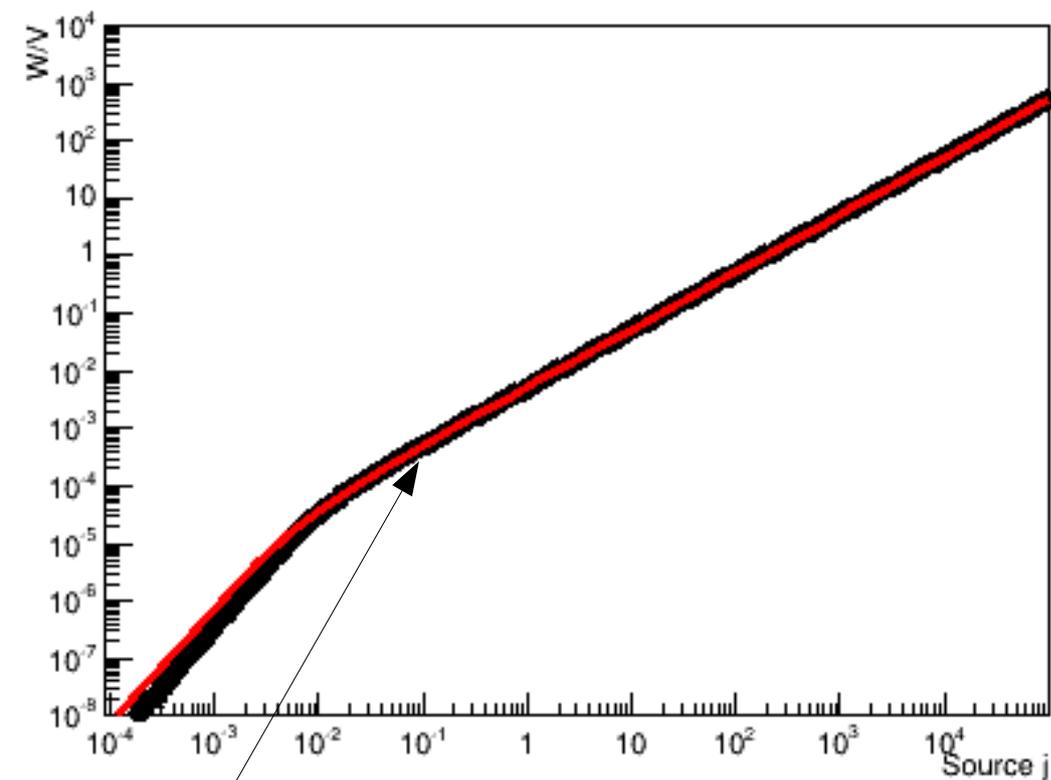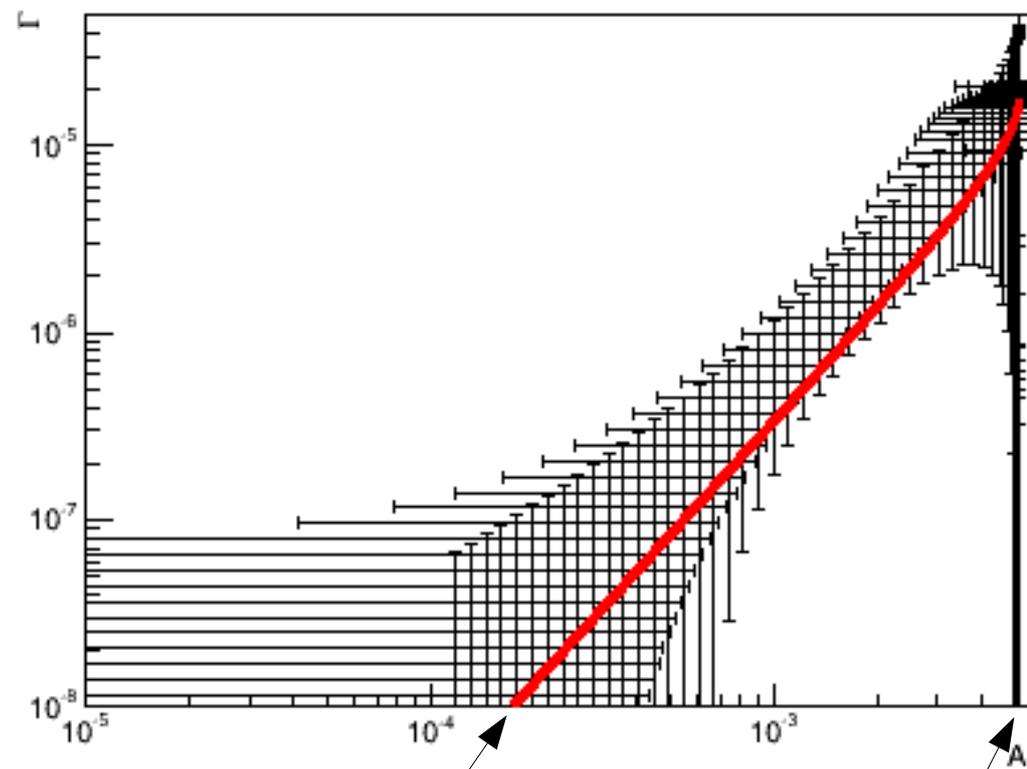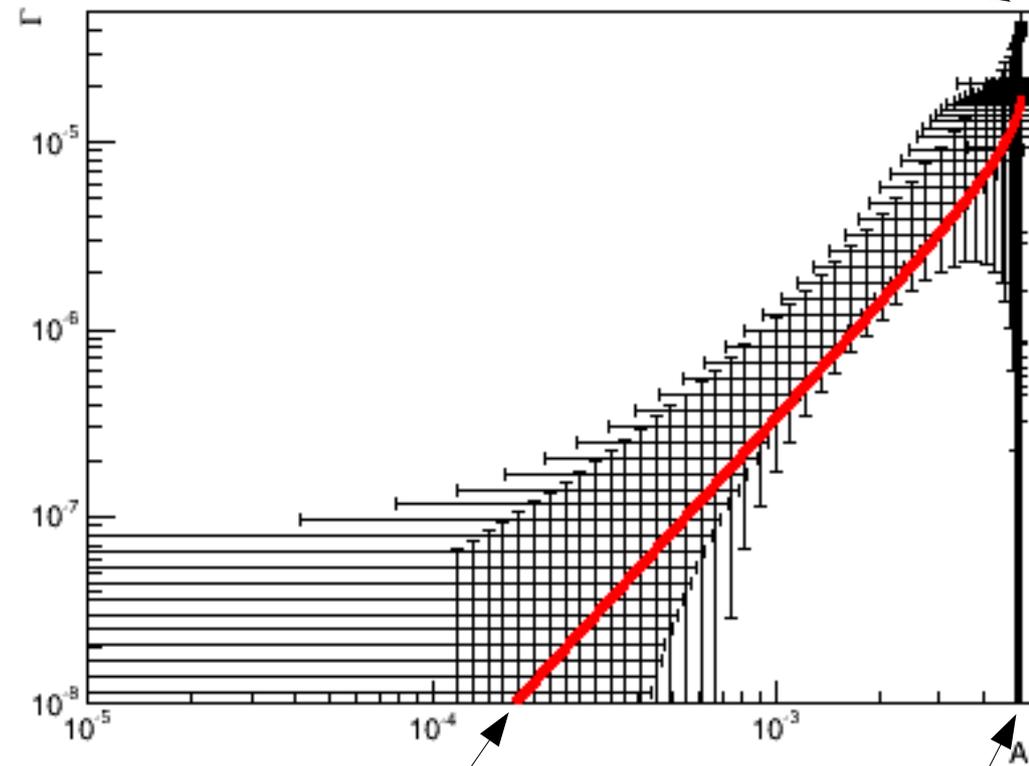
Singularity

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Maximum field amplitude:
Cut-off at the Gribov horizon

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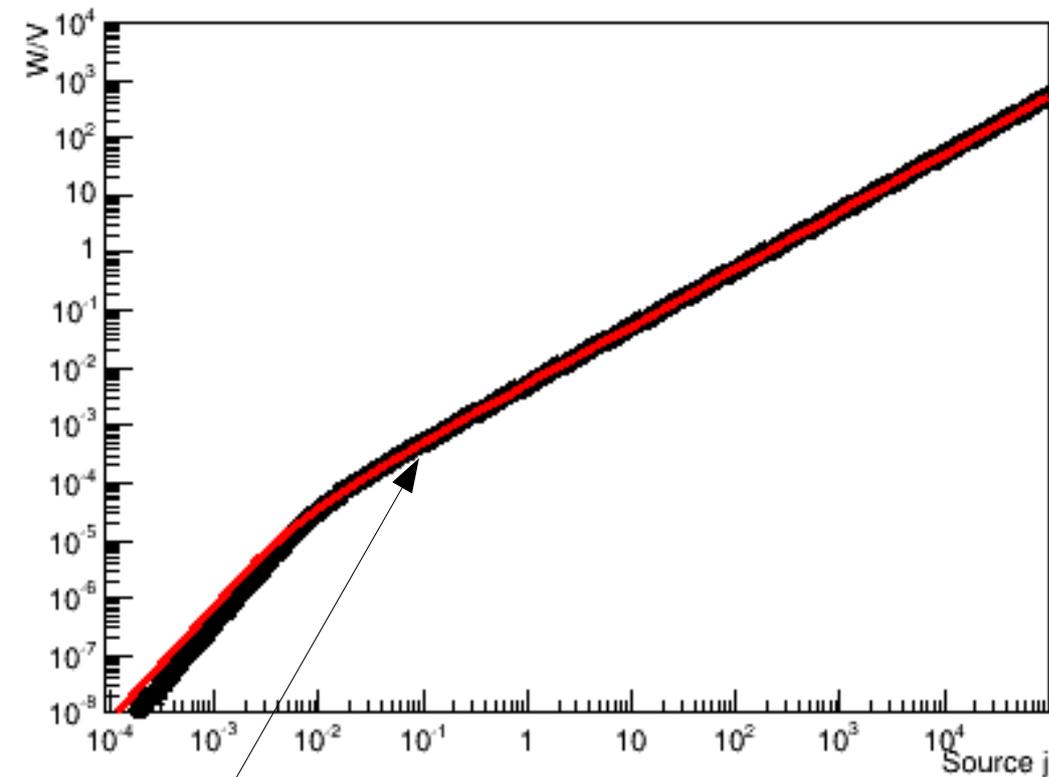Cut-off at the Gribov horizon

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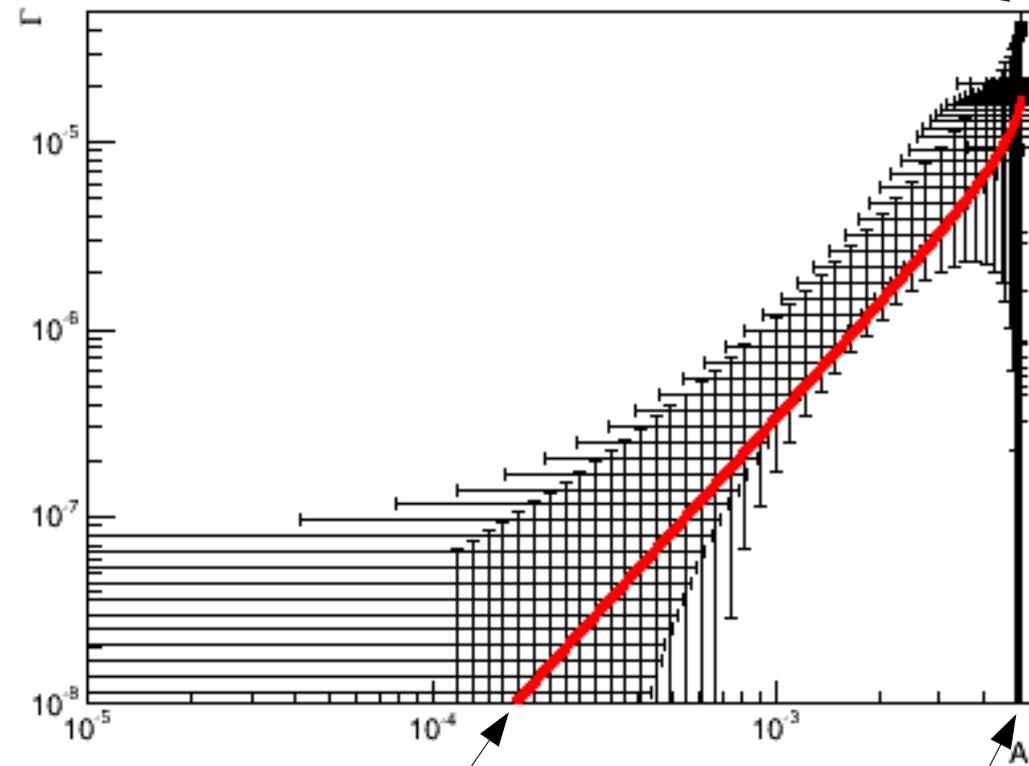
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The quantum effective action is not an analytic functional!

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- Need all elements from the lattice
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- Possible also for other equations
  - Link DSE shows similar results [Sternbeck et al.'15]

Determine a mis-function

Ghost DSE

$$0 = \frac{1}{\left(p^2 D(p)\right)} + 1 - \Pi(p)$$

$$\Pi(p) = \frac{c}{p^2} \int d^d q \, \Gamma_\mu^{tl}(p,q) D_{\mu\nu}(p+q) D(p) \Gamma_\nu(p,q)$$

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Ghost DSE   Tree-level

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Continuum structure, lattice input

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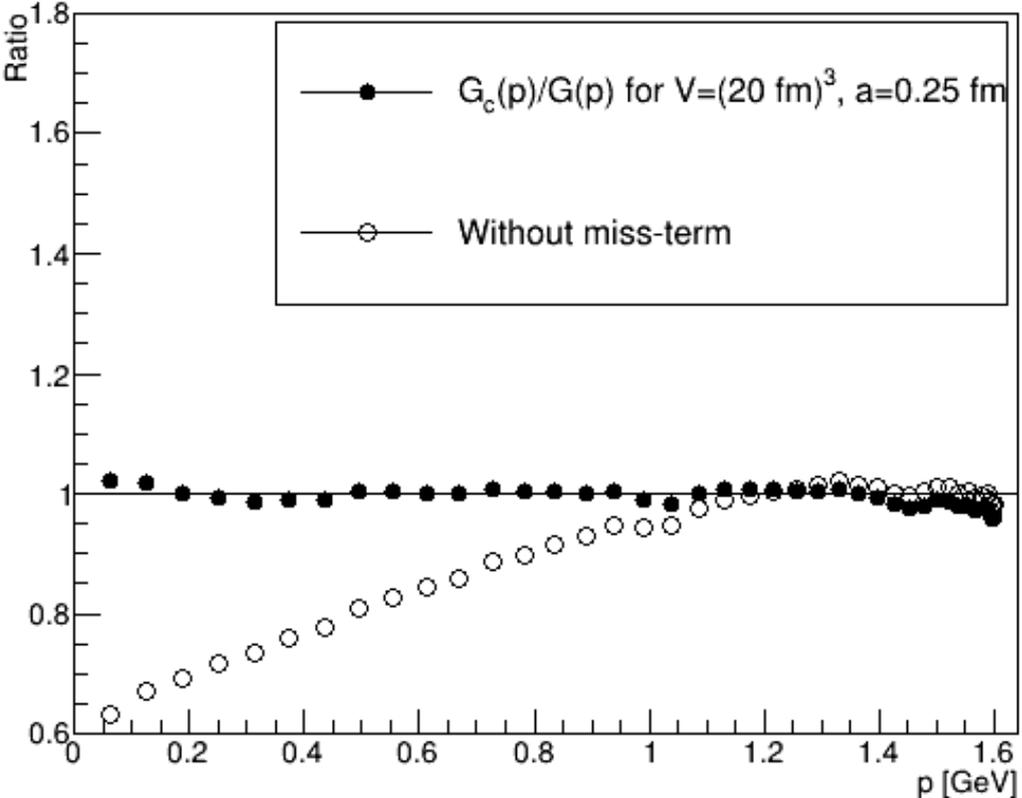
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Question can be
recast as ratio:
$$\frac{\dfrac{1}{1-\Pi(p)}}{p^2 D^{lattice}(p)} \overset{?}{=} 1 \overset{?}{=} \frac{\dfrac{1}{1-f(p)-\Pi(p)}}{p^2 D^{lattice}(p)}$$

Lattice results

Ratio of propagators

Legend: $G_c(p)/G(p)$ for $V=(20 \text{ fm})^3$, $a=0.25$ fm; Without miss-term
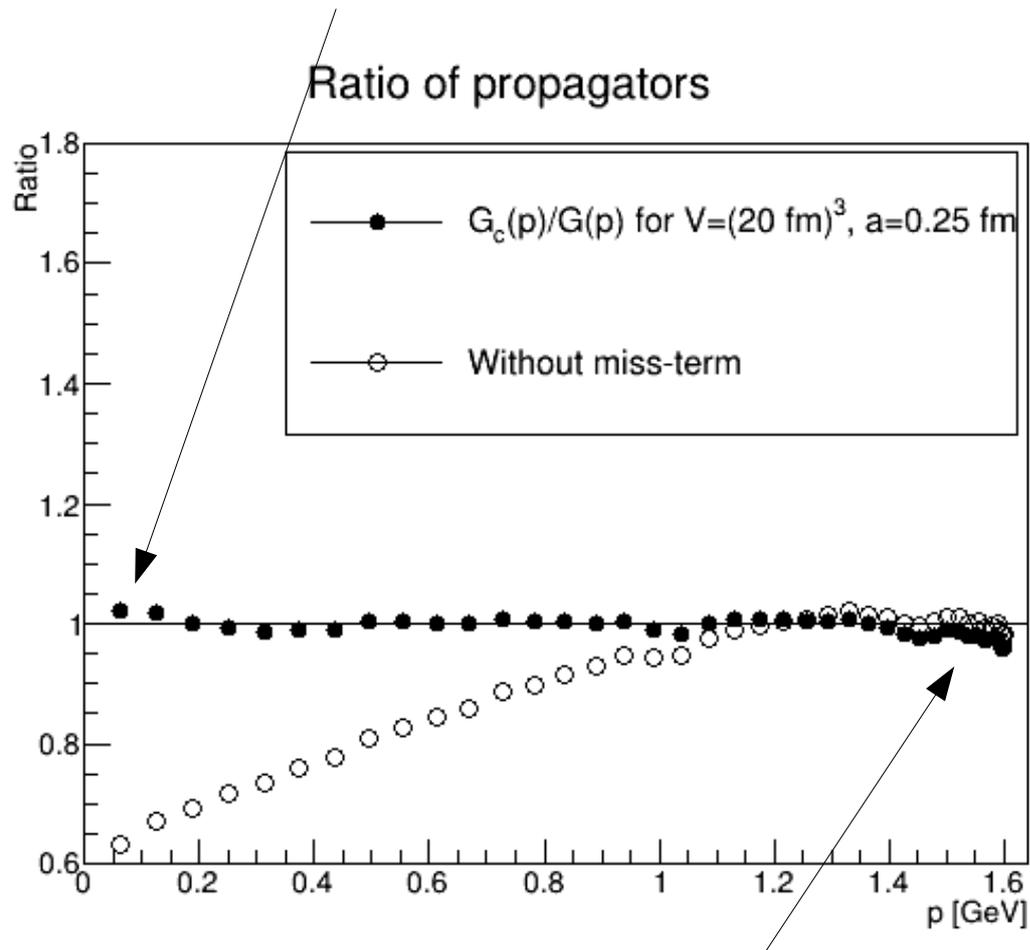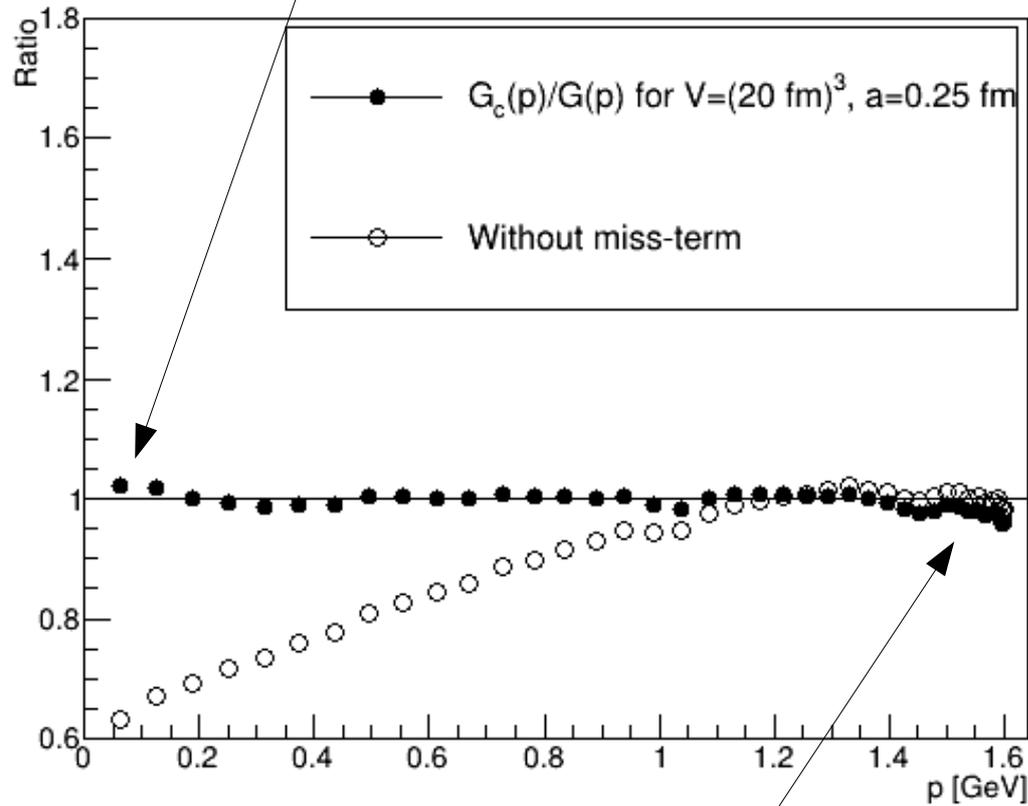
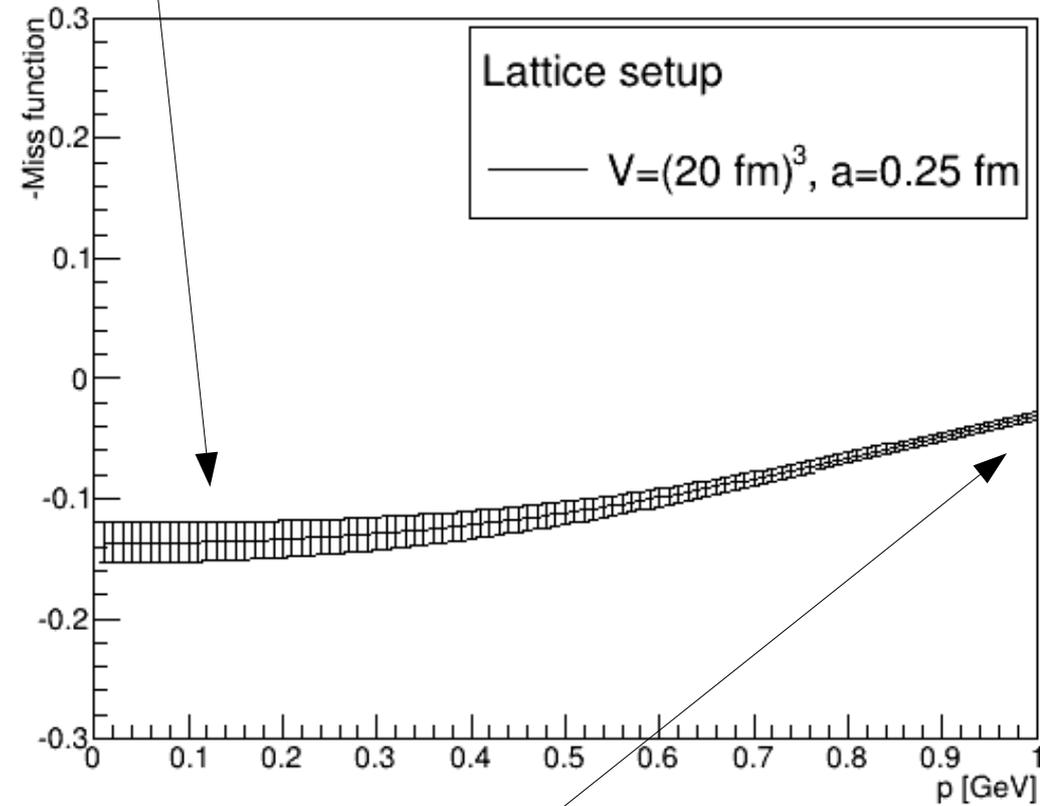# Lattice results

Substantial disagreement at low momenta

Agreement at large momenta

Lattice results

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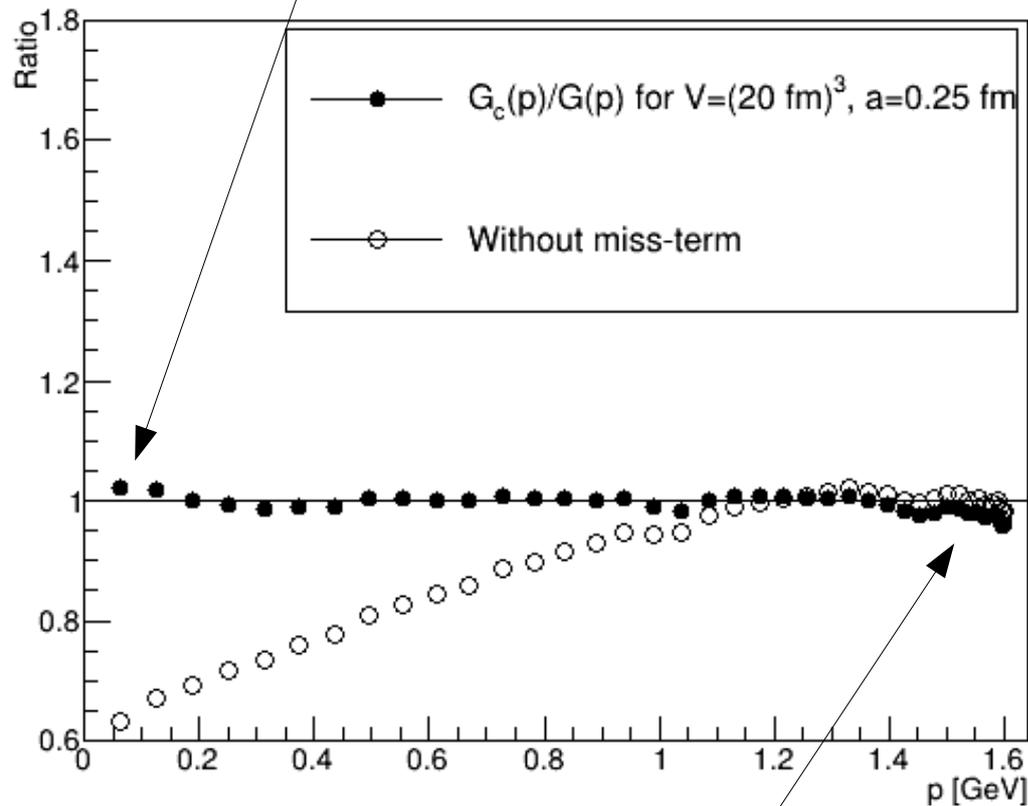
Miss function in three dimensions

$G_c(p)/G(p)$ for V=(20 fm)^3, a=0.25 fm

Without miss-term

Lattice setup

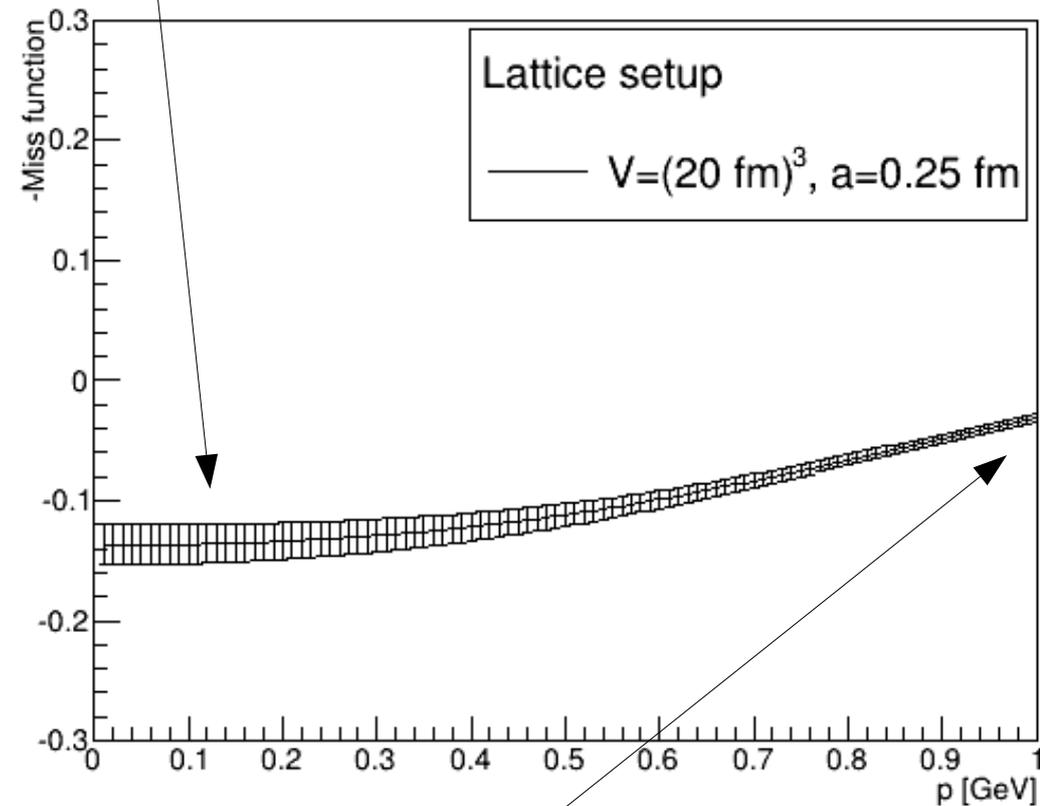V=(20 fm)^3, a=0.25 fm

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Same effect in 4 dimensions

Miss function in three dimensions
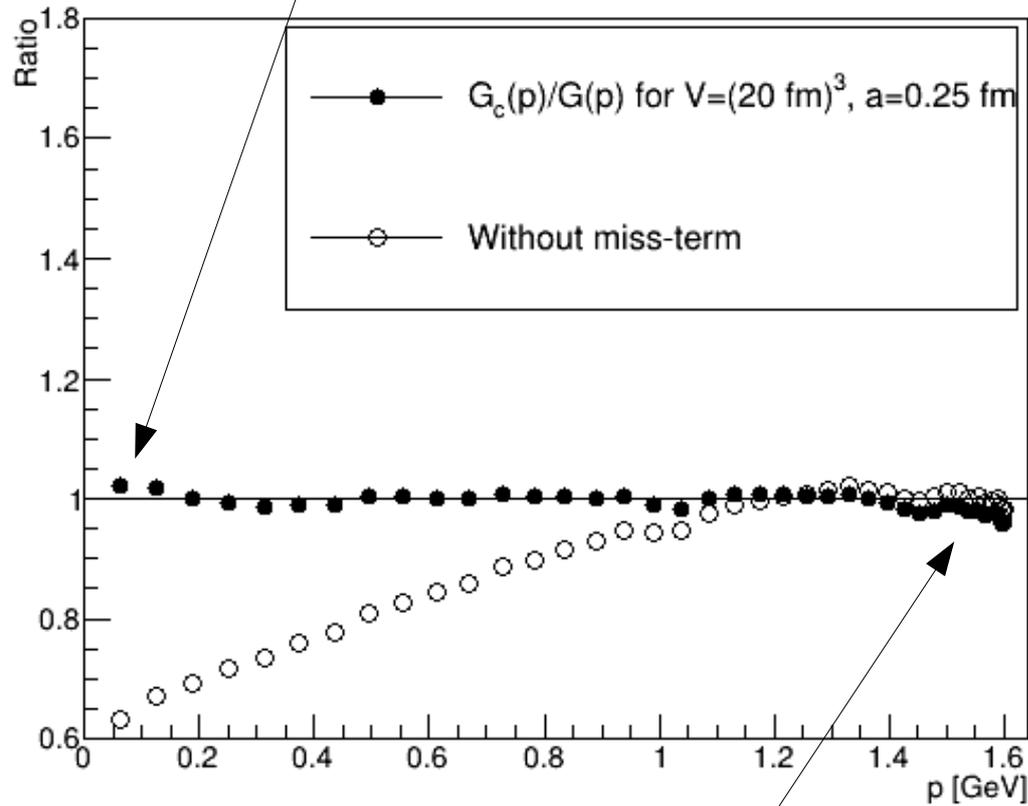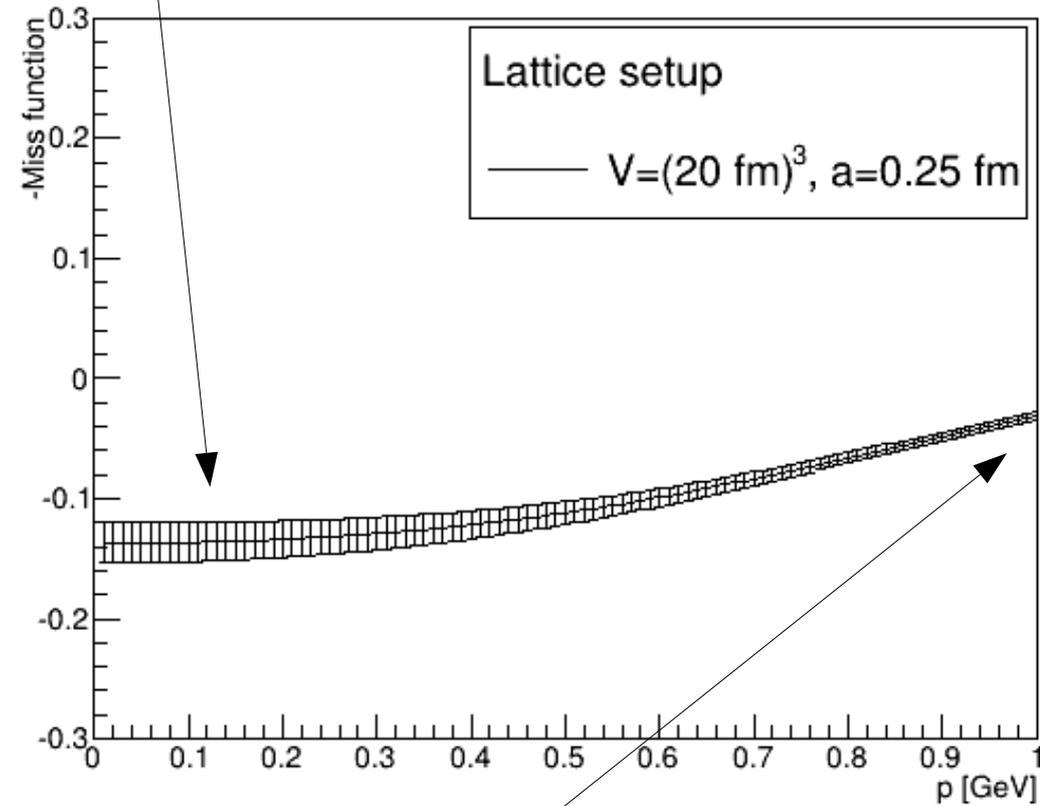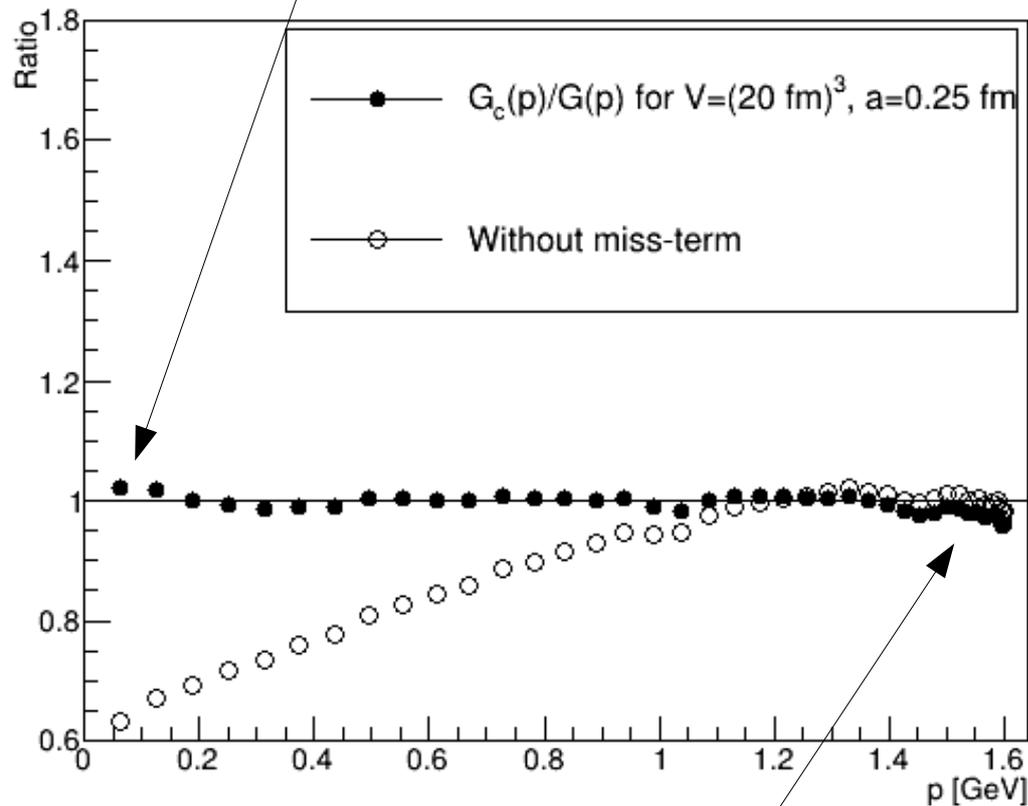
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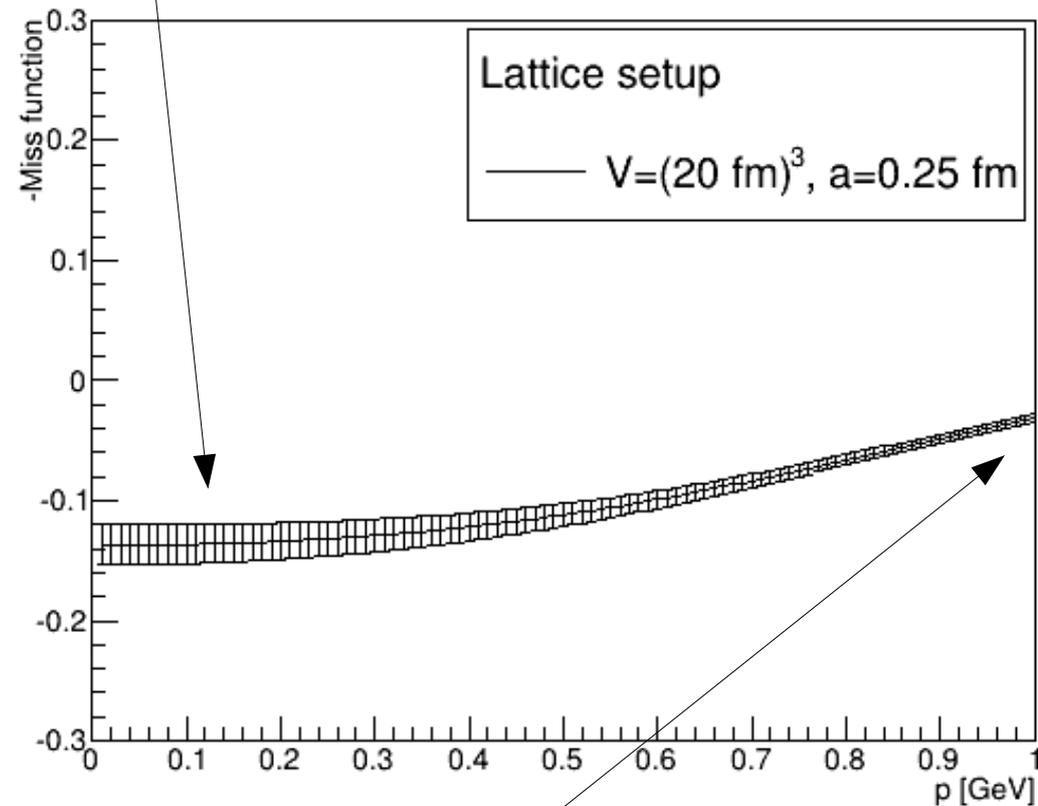$$f(p) = r \frac{1 - (p/b)^e}{1 + (p/c)^{2e}}$$

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Same effect in 4 dimensions
...and in 2 dimensions!

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- The misfunction needs to be a ghost/anti-ghost-dependent term in the gauge-fixing Lagrangian!
  - Should be reconstructable/derivable somehow